

# Probing proton intrinsic charm in photon or Z boson production accompanied by heavy jets at the LHC

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We consider an observable very sensitive to the nonzero intrinsic charm (IC) contribution to the proton density. It is the ratio between the differential cross sections of the photon or Z-boson and  $c$ -jet production in the  $pp$  collision,  $\gamma(Z) + c$ , and the  $\gamma(Z)$  and the  $b$ -jet production. It is shown that this ratio can be approximately flat or increasing at large  $\gamma(Z)$  transverse momenta  $p_T$ , and their pseudorapidities  $1.5 < \eta < 2.4$  if the IC contribution is taken into account. On the contrary, in the absence of the IC, this ratio decreases as  $p_T$  grows. We also present the ratios of the cross sections integrated over  $p_T$  as a function of the IC probability  $w$ . It is shown that these ratios are mostly independent on the theoretical uncertainties, and such predictions could therefore be much more promising for the search for the intrinsic charm signal at the LHC compared to the predictions for  $p_T$  spectra, which significantly depend on these uncertainties.

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## I. INTRODUCTION

The hypothesis of the *intrinsic* (or valence-like) heavy quark component, the quark Fock state  $|uudQ\bar{Q}\rangle$  [1–4] in a proton, suggested by Brodsky with coauthors [1,2] (BHPS model), is intensively discussed in connection with an opportunity to verify it experimentally [5–15]. Up to now, there has been a long-standing debate about the possible existence of the *intrinsic* charm (IC) and *intrinsic* strange (IS) quarks in a proton [7,15–17]. Thorough theoretical and experimental studies of these intrinsic heavy quark components would be very important for the experiments performed at the LHC.

Recently it was shown that the possible existence of the intrinsic heavy quark components in the proton can be seen not only in the inclusive heavy flavor production at high energies [8], but also in the semi-inclusive production of prompt photons or vector bosons accompanied by heavy quark jets [9,13]. An experimental hint on the possible existence of the IC contribution was observed in the Tevatron experiment on the prompt photon production in the association of the  $c$  and  $b$  jets in the  $p\bar{p}$  annihilation at  $\sqrt{s} = 1.98$  TeV [18,19]. It was shown that the description of the Tevatron data within the perturbative QCD (pQCD) could be significantly improved if the IC contributions were taken into account. The photon transverse momentum ( $p_T$ ) spectrum in the  $\gamma + c$  production and the ratio of the spectra in the  $\gamma + c$  and  $\gamma + b$  productions measured at the Tevatron [20] are better described within the BHPS model [1,2], which includes the IC contributions. According to the pQCD calculations [21], in the absence of the IC contribution, this ratio decreases as  $p_T$  grows, while the Tevatron data show its flat behavior at large  $p_T \geq 100$  GeV [20].

The possible IC signal can also be observed in the hard  $pp$  production of the gauge bosons  $Z$  or  $W$  accompanied by

heavy flavors. As was shown in Ref. [13], the ratio of the  $Z + c$  and  $W +$  heavy jet production cross sections maximizes the sensitivity to the IC component of the proton. Our early predictions about a possible intrinsic charm signal in the production of prompt photons or gauge bosons accompanied by heavy flavor jets concerned their transverse momenta distributions in the midrapidity region of  $pp$  collisions at the LHC energies [9,13]. They were obtained with the IC probability about  $w = 3.5\%$ , which is the upper limit due to constraints from the HERA data on the deep inelastic scattering. However, the upper limit of the IC probability in a proton is still very actively debated [7,15–17]. Therefore, in the present paper we focus mainly on the predictions for searching at any  $w$  for the IC signal in the observables, which are poorly sensitive to the theoretical uncertainties—namely, the ratios between the  $\gamma(Z) + c$  and  $\gamma(Z) + b$  cross sections in  $pp$  collisions at the LHC energies. An important advantage of these observables is that many theoretical uncertainties—for example, heavy quark masses, the factorization and/or renormalization scales—are canceled, as will be demonstrated below. We show that the measurement of these ratios is much more promising for the search for the IC signal.

Below, we perform the calculations in two ways. First, we use the parton-level Monte Carlo event generator MCFM [22], which implements the NLO pQCD calculations of associated Z-boson and heavy flavor jet production. The detailed description of the MCFM routine is available [22]. To generate the prompt photon and heavy jet production cross sections, we apply the  $k_T$  factorization approach [23,24], which becomes a commonly recognized tool in the high-energy phenomenology. Our main motivation is that it gives a better description of the Tevatron data compared to the NLO pQCD calculations [21], as was claimed in

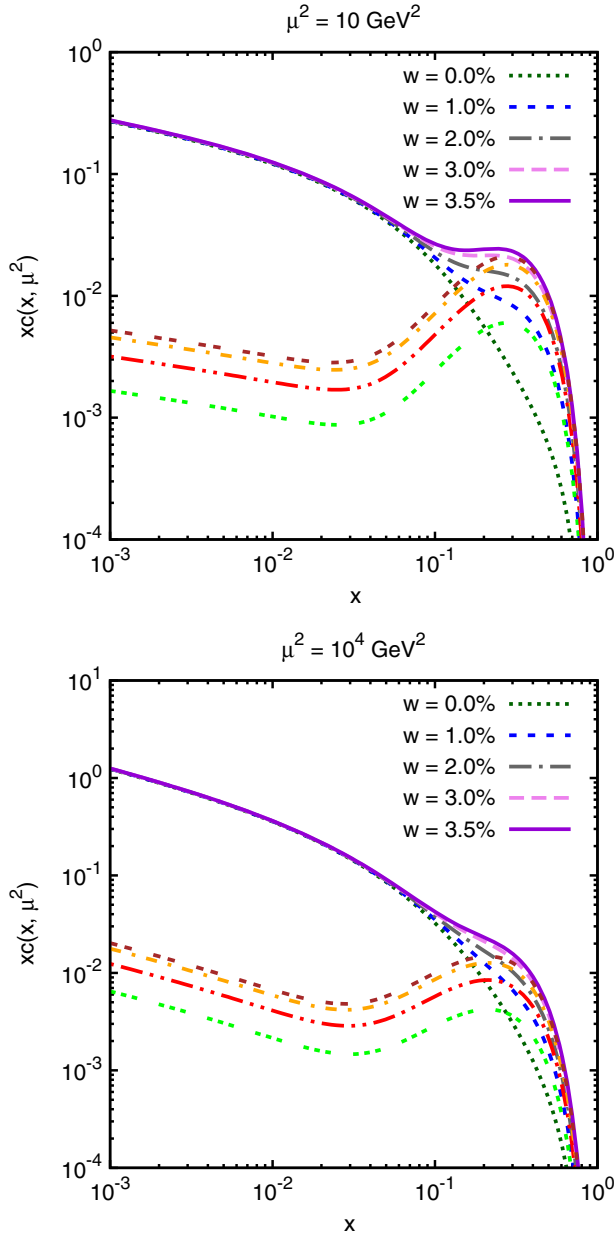


FIG. 1. The total charmed quark density  $xc(x, \mu^2)$  as a function of  $x$  at different values of  $w$  at  $\mu^2 = 10 \text{ GeV}^2$  (top) and  $\mu^2 = 10^4 \text{ GeV}^2$  (bottom). The triple-dashed line is the IC contribution at  $w = 1\%$ , the dash-double-dotted line corresponds to the IC at  $w = 2\%$ , the dash-dotted curve is the IC at  $w = 3\%$ , and the double-dashed line corresponds to the IC at  $w = 3.5\%$ .

Refs. [18,19]. We apply this approach to the associated  $Z$  and heavy jet production to perform an independent cross-check of our results.<sup>1</sup>

The outline of our paper is as follows: In Secs. II and III, we recall basic ideas with a brief review of calculation steps. In Sec. IV, we present the numerical results of our

<sup>1</sup>Unfortunately, the MCFM routine does not produce the prompt photon and heavy jet production cross sections.

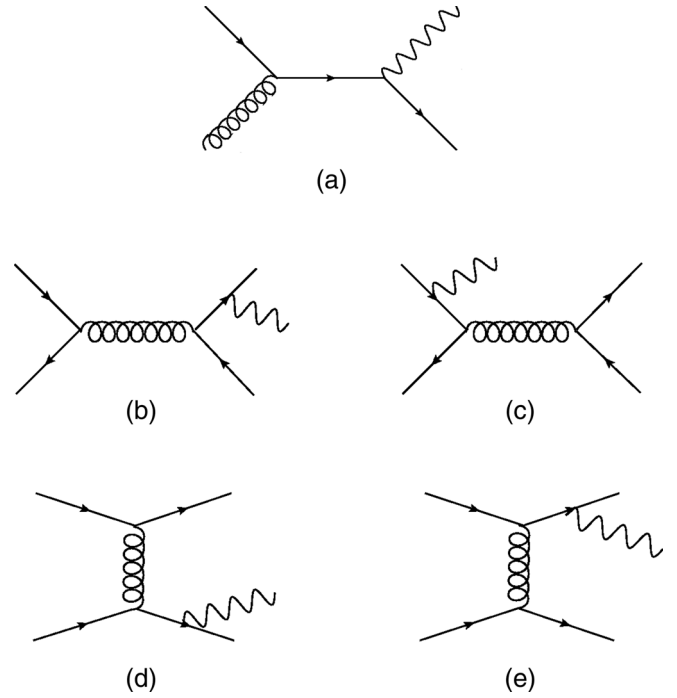


FIG. 2. (a) The  $\mathcal{O}(\alpha\alpha_s)$  and (b)–(e)  $\mathcal{O}(\alpha\alpha_s^2)$  contributions to the  $\gamma(Z) + Q$  production taken into account in the  $k_T$  factorization calculations.

calculations and a discussion. Finally, Sec. V contains our conclusions.

## II. INTRINSIC CHARM DENSITY IN A PROTON AS A FUNCTION OF IC PROBABILITY $w$

According to Refs. [6,12,25], the intrinsic charm distribution at the starting scale  $\mu_0^2$  as a function of  $x$  can be presented in the following approximated form:

$$c_{\text{int}}(x, \mu_0^2) = c_0 w x^2 [(1-x)(1+10x+x^2) + 6x(1+x)\ln(x)], \quad (1)$$

where  $w$  is the probability to find the Fock state  $|uudc\bar{c}\rangle$  in the proton,  $c_0$  is the normalization constant, and the masses of the light quarks and the nucleon are negligible compared to the charm quark mass. The inclusion of the nonzero nucleon mass leads to a more complicated analytic form [26]. According to the BHPS model [1,2], the charm density in a proton is the sum of the *extrinsic* and *intrinsic* charm densities,

$$xc(x, \mu_0^2) = xc_{\text{ext}}(x, \mu_0^2) + xc_{\text{int}}(x, \mu_0^2). \quad (2)$$

The *extrinsic*, or ordinary, quarks and gluons are generated on a short time scale associated with the large-transverse-momentum processes. Their distribution functions satisfy the standard QCD evolution equations. Contrariwise, the

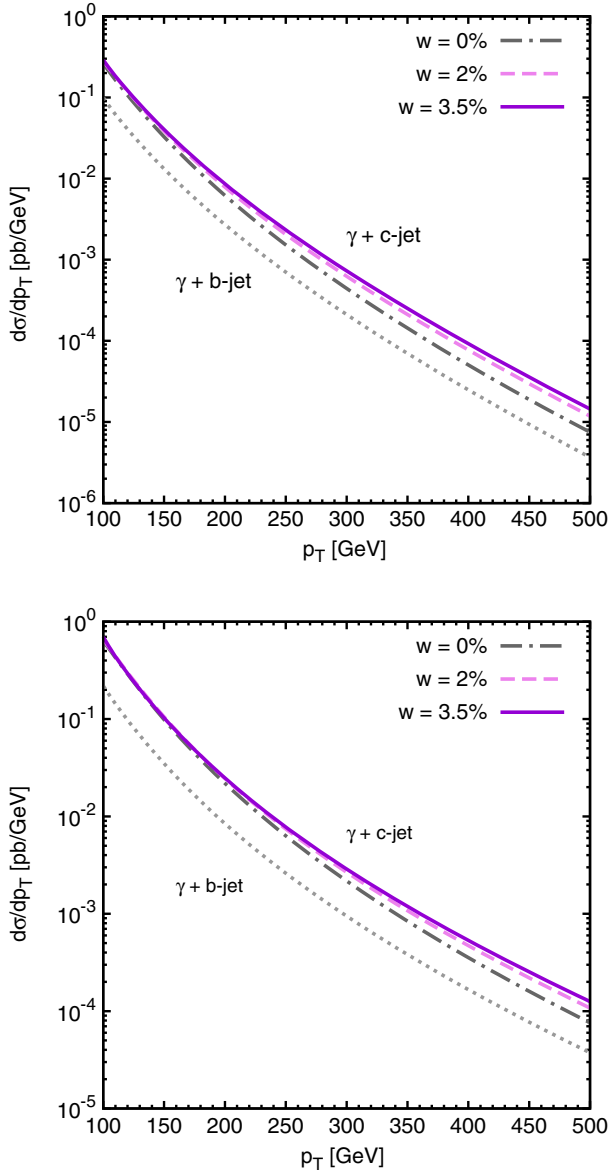


FIG. 3. The cross sections of the associated  $\gamma + c$  and  $\gamma + b$  production in the  $pp$  collision calculated as a function of the photon transverse momentum  $p_T$  at  $\sqrt{s} = 8$  TeV (top) and  $\sqrt{s} = 13$  TeV (bottom) within the  $k_T$  factorization approach. The kinematical conditions are described in the text.

*intrinsic* quarks and gluons can be associated with a bound-state hadron dynamics, and it is believed that they have a nonperturbative origin. The lifetime of this Fock state should be much more than the interaction time of the hard probe [26]. Some comments on this are presented below. It was argued [2] that the existence, for example, of intrinsic heavy quark pairs  $c\bar{c}$  and  $b\bar{b}$  within the proton state can be due to the gluon-exchange and vacuum-polarization graphs.

The charm density  $xc(x, \mu^2)$  at an arbitrary scale  $\mu^2$  is calculated using the Dokshitzer-Gribov-Lipatov-Altarelli-Parisi (DGLAP) equations [27]. Let us stress here that

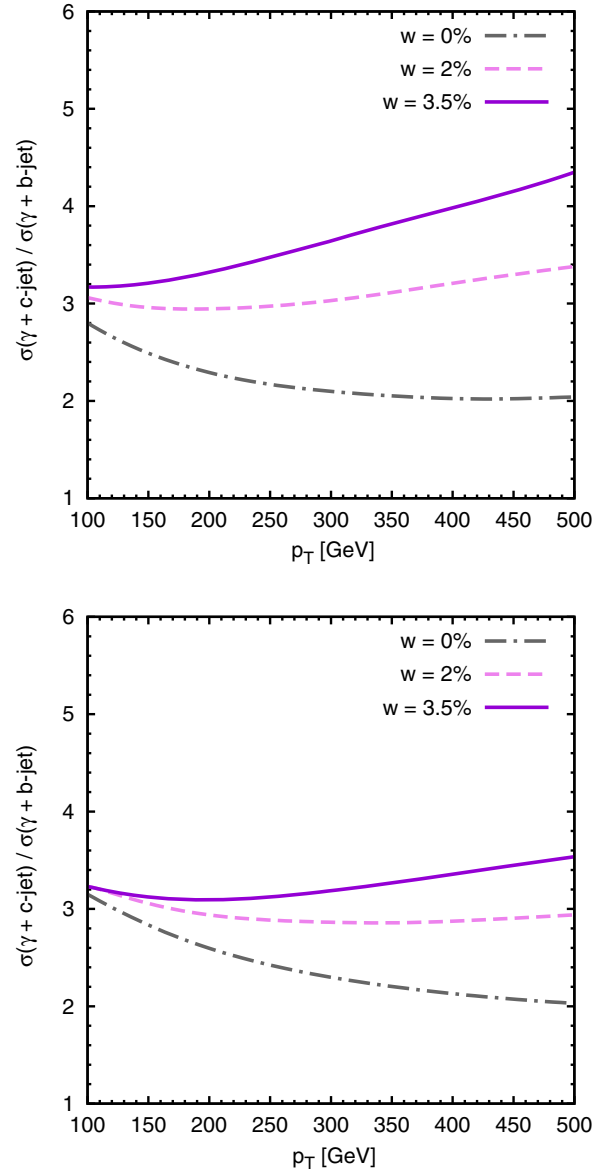


FIG. 4. The cross-section ratio of the  $\gamma + c$  production to the  $\gamma + b$  one in the  $pp$  collision calculated as a function of the photon transverse momentum  $p_T$  at  $\sqrt{s} = 8$  TeV (top) and  $\sqrt{s} = 13$  TeV (bottom) within the  $k_T$  factorization approach. The kinematical conditions are described in the text.

both the intrinsic part  $xc_{\text{int}}$  and the extrinsic one  $xc_{\text{ext}}$  depend on  $\mu^2$ . In the general case, there is some mixing between the two parts of Eq. (2) during the DGLAP evolution. However, such mixing is negligible [12,28], especially at large  $\mu^2$  and  $x$ . It can be seen from the comparison between our calculations of charmed quark densities presented in Fig. 1, where this mixing was included within the CTEQ [29] set; and Fig. 2 of Ref. [12], when the mixing between two parts of the charm density was neglected. Our results on the total charm density  $xc(x, \mu^2)$  are in good agreement with the calculations [12] in the whole kinematical region of  $x$ ,

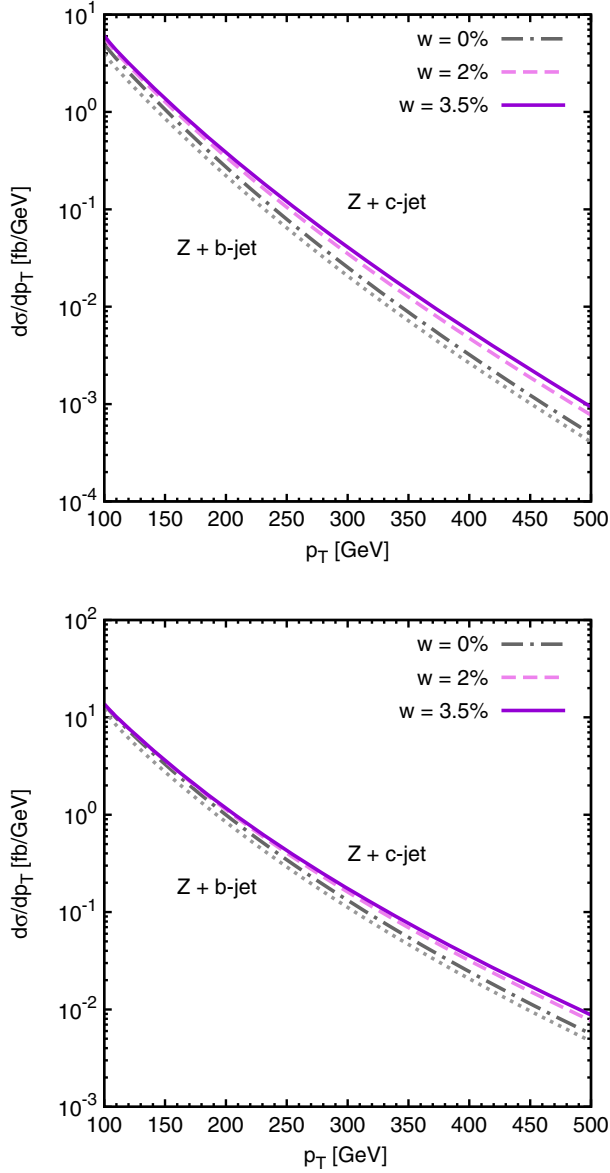


FIG. 5. The cross sections of the associated  $Z + c$  and  $Z + b$  production in the  $pp$  collision calculated as a function of the  $Z$ -boson transverse momentum  $p_T$  at  $\sqrt{s} = 8$  TeV (top) and  $\sqrt{s} = 13$  TeV (bottom) within the  $k_T$  factorization approach. The kinematical conditions are described in the text.

because at  $x < 0.1$  the IC contribution is much smaller than the extrinsic one. Therefore, one can apply the DGLAP evolution separately to the first part  $xc_{\text{ext}}(x, \mu_0^2)$  and the second part  $xc_{\text{int}}(x, \mu_0^2)$  of Eq. (2), as was done in Refs. [12,28]. Such calculations were done by the CTEQ [29] and CT14 [30] groups at some fixed values of the IC probability  $w$ . Namely, the CTEQ group used  $w = 1\%$  and  $w = 3.5\%$ , and CT14 used  $w = 1\%$  and  $w = 2\%$ .

Note that, according to a recent paper [26], the lifetime of the intrinsic charm should be more than the interaction time by a factor of at least 5, when the quark Fock state can be

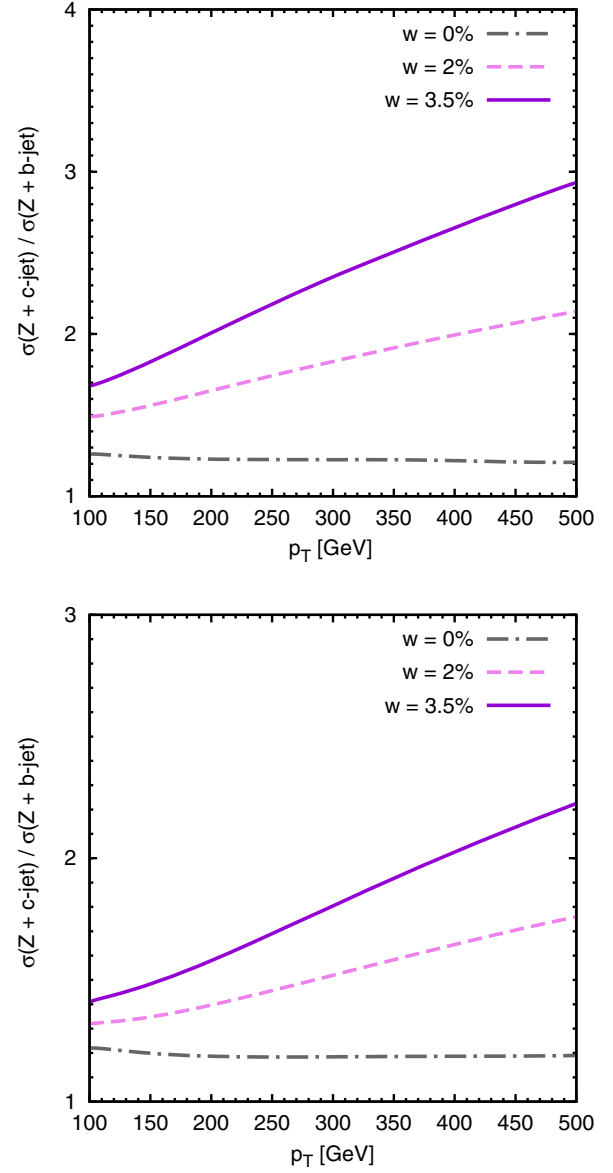


FIG. 6. The cross section ratio of the  $Z + c$  production to the  $Z + b$  one in the  $pp$  collision calculated as a function of the  $Z$ -boson transverse momentum  $p_T$  at  $\sqrt{s} = 8$  TeV (top) and  $\sqrt{s} = 13$  TeV (bottom) within the  $k_T$  factorization approach. The kinematical conditions are described in the text.

observed with the satisfactory accuracy. The ratio of these times is proportional to  $Q^2$  or  $p_T^2$  [26]. We will analyze the hard processes of  $\gamma(Z)$  production associated with heavy jets at LHC energies and  $p_T^2 \geq 10^4$  GeV<sup>2</sup>, when the lifetime of the intrinsic charm is much larger than the interaction time, where the intrinsic charm could be resolved.

Taking into account that the IC probability  $w$  enters into Eq. (2) as a constant in front of the function dependent on  $x$  and  $\mu^2$ , one can suggest a simple relation at any  $w \leq w_{\text{max}}$ :

$$xc_{\text{int}}(x, \mu^2) = \frac{w}{w_{\text{max}}} xc_{\text{int}}(x, \mu^2)|_{w=w_{\text{max}}}. \quad (3)$$

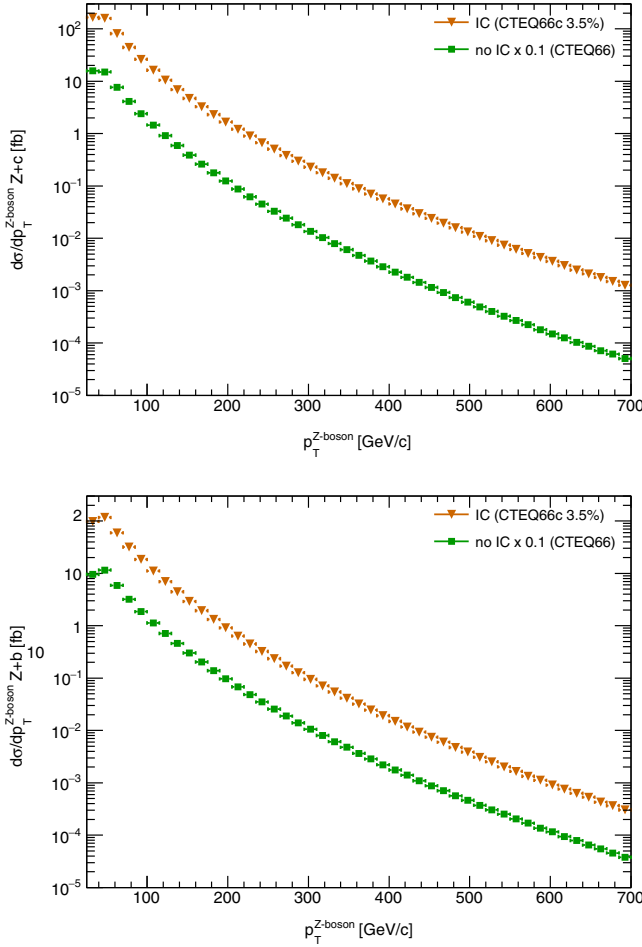


FIG. 7. The cross sections of the associated  $Z + c$  (top) and  $Z + b$  (bottom) production in  $pp$  collision calculated as a function of the  $Z$ -boson transverse momentum  $p_T$  at  $\sqrt{s} = 13$  TeV within the MCFM routine. The kinematical conditions are described in the text.

Actually, that is the linear interpolation between two charm densities at the scale  $\mu^2$ , obtained at  $w = w_{\max}$  and  $w = 0$ . Later, we adopt the charm distribution function from the CTEQ66M set [29]. We assume  $w_{\max} = 3.5\%$  everywhere, which corresponds to the CTEQ66c1 set [29]. Additionally, we perform the three-point interpolation of the charmed quark distributions (over  $w = 0$ ,  $w = 1\%$  and  $w = 3.5\%$ , which correspond to the CTEQ66M, CTEQ66c0 and CTEQ66c1 sets, respectively). These results differ from the ones based on Eq. (3) by no more than 0.5%, thus giving us confidence in our starting point.

Below, we apply the charmed quark density obtained by Eqs. (2) and (3) to calculate the total and differential cross sections of associated prompt photon or  $Z$ -boson and heavy flavor jet production,  $\gamma(Z) + Q$ , at the LHC conditions. The suggested procedure to calculate  $xc_{\text{int}}(x, \mu^2)$  at any  $w \leq w_{\max}$  allows us to reduce significantly the time for the calculation of these observables.

### III. THEORETICAL APPROACHES TO THE ASSOCIATED $\gamma(Z) + Q$ PRODUCTION

As was mentioned above, we perform the numerical calculations of the associated  $\gamma(Z) + Q$  production cross sections using the parton-level Monte Carlo event generator MCFM within the NLO pQCD, as well as the  $k_T$  factorization QCD approach. The MCFM is able to calculate the processes that involve the gauge bosons  $Z$  or  $W$  (see Ref. [22] for more information). In contrast to our early study of these processes [13] within the MCFM, we use this generator to calculate the differential and total cross sections of the  $Z + c$  and  $Z + b$  production in the  $pp$  collision and their ratio as a function  $w$ .

The  $k_T$  factorization approach [23,24] is based on the small- $x$  Balitsky-Fadin-Kuraev-Lipatov (BFKL) [31] gluon dynamics and provides solid theoretical grounds for the effects of the initial gluon radiation and the intrinsic parton transverse momentum.<sup>2</sup> Our main motivation to use the  $k_T$  factorization formalism here is that its predictions for the associated  $\gamma + Q$  production better agree with the Tevatron data compared to the NLO pQCD (see Refs. [18,19]). The consideration is mainly based on the  $\mathcal{O}(\alpha_s)$  off-shell (depending on the transverse momenta of initial quarks and gluons) quark-gluon Compton-like scattering subprocess; see Fig. 2(a). Within this approach, the transverse-momentum-dependent (TMD) parton densities include many high-order corrections, while the partonic amplitudes are calculated within the leading order (LO) of QCD. The off-shell quark-gluon Compton scattering amplitude is calculated within the Regge-ized parton approach [33–35] based on the effective action formalism [36], which ensures the gauge invariance of the obtained amplitudes despite the off-shell initial quarks and gluons.<sup>3</sup> The TMD parton densities are calculated using the Kimber-Martin-Ryskin (KMR) approach, currently developed within the NLO [38]. This approach is the formalism to construct the TMD quark and gluon densities from the known conventional parton distributions. The key assumption is that the  $k_T$  dependence appears at the last evolution step, so that the DGLAP evolution can be used up to this step. Numerically, for the input we used parton densities derived in Sec. II. Other details of these calculations are explained in Ref. [37].

To improve the  $k_T$  factorization predictions at high transverse momenta, we take into account some  $\mathcal{O}(\alpha_s^2)$  contributions, namely  $q\bar{q} \rightarrow VQ\bar{Q}$  and  $qQ \rightarrow VqQ$  ones, where  $V$  denotes the photon or the  $Z$  boson; see Figs. 2(e). These contributions are significant at large  $x$  and therefore can be calculated in the usual collinear QCD factorization scheme. Thus, we rely on the combination of two techniques that is most suitable.

<sup>2</sup>A detailed description of the  $k_T$  factorization approach can be found, for example, in reviews [32].

<sup>3</sup>Here we use the expressions derived earlier [37].



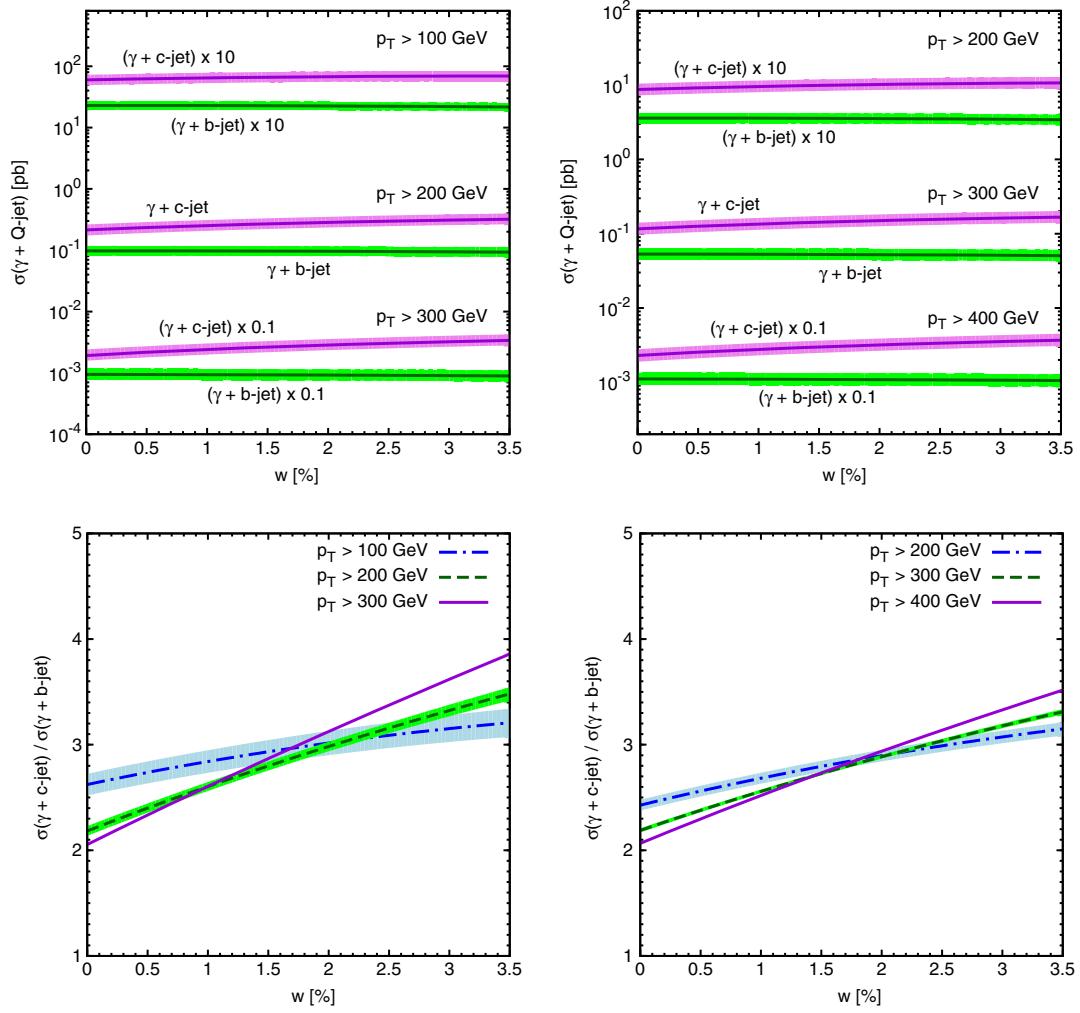


FIG. 8. Top: The cross sections of the associated  $\gamma + c$  and  $\gamma + b$  production in the  $pp$  collision as a function of  $w$  integrated over the photon transverse momenta  $p_T > p_T^{\min}$  for different  $p_T^{\min}$  at  $\sqrt{s} = 8$  TeV (left) and  $\sqrt{s} = 13$  TeV (right). The kinematical conditions are described in the text. Bottom: The corresponding ratios of these cross sections. The calculations were done using the  $k_T$  factorization approach. The bands correspond to the usual scale variation as it is described in the text.

#### IV. RESULTS AND DISCUSSION

Let us present the results of our calculations. First of all, we describe our numerical input. Following Ref. [39], we set the charmed and beauty quark masses  $m_c = 1.4$  GeV,  $m_b = 4.75$  GeV; the  $Z$ -boson mass  $m_Z = 91.1876$  GeV; and  $\sin^2 \theta_W = 0.23122$ . The chosen factorization and renormalization scales are  $\mu_R = \mu_F = \xi p_T$  or  $\mu_R = \mu_F = \xi m_T$ , where  $p_T$  is the produced photon transverse momentum and  $m_T$  is the  $Z$ -boson transverse mass. As usual, we vary the nonphysical parameter  $\xi$  between  $1/2$  and  $2$  about the default value  $\xi = 1$  in order to estimate the scale uncertainties of our calculations. We employ the two-loop formula for the strong coupling constant with active quark flavors  $n_f = 5$  at  $\Lambda_{\text{QCD}} = 226.2$  MeV and use the running QED coupling constant over a wide region of transverse momenta. The multidimensional integration in the  $k_T$

factorization calculations was performed by means of the Monte Carlo technique, using the VEGAS routine [40].

In our calculations, we also follow the conclusion obtained in our papers [9,13] that the IC signal in the hard processes discussed here can be detected at ATLAS or CMS of the LHC in the forward rapidity region  $1.5 < |\eta| < 2.4$ , and  $p_T > 50$  GeV. Additionally, we require  $|\eta(Q)| < 2.4$  and  $p_T(Q) > 25$  GeV, where  $\eta(Q)$  and  $p_T(Q)$  are the pseudorapidity and transverse momentum of the heavy quark jet in a final state, as was done in Refs. [9,13].

The results of our calculations are shown in Figs. 3–9. The transverse momentum distributions of photons and  $Z$  bosons accompanied by the  $c$  and  $b$  quarks are presented in Figs. 3, 5, and 7 at different IC probabilities  $w$  (namely,  $w = 0\%$ ,  $w = 2\%$ , and  $w = 3.5\%$ ) at  $\sqrt{s} = 8$  and  $13$  TeV. One can see in Figs. 5 and 7 that the MCFM and  $k_T$  factorization predictions for  $Z + Q$  production are very

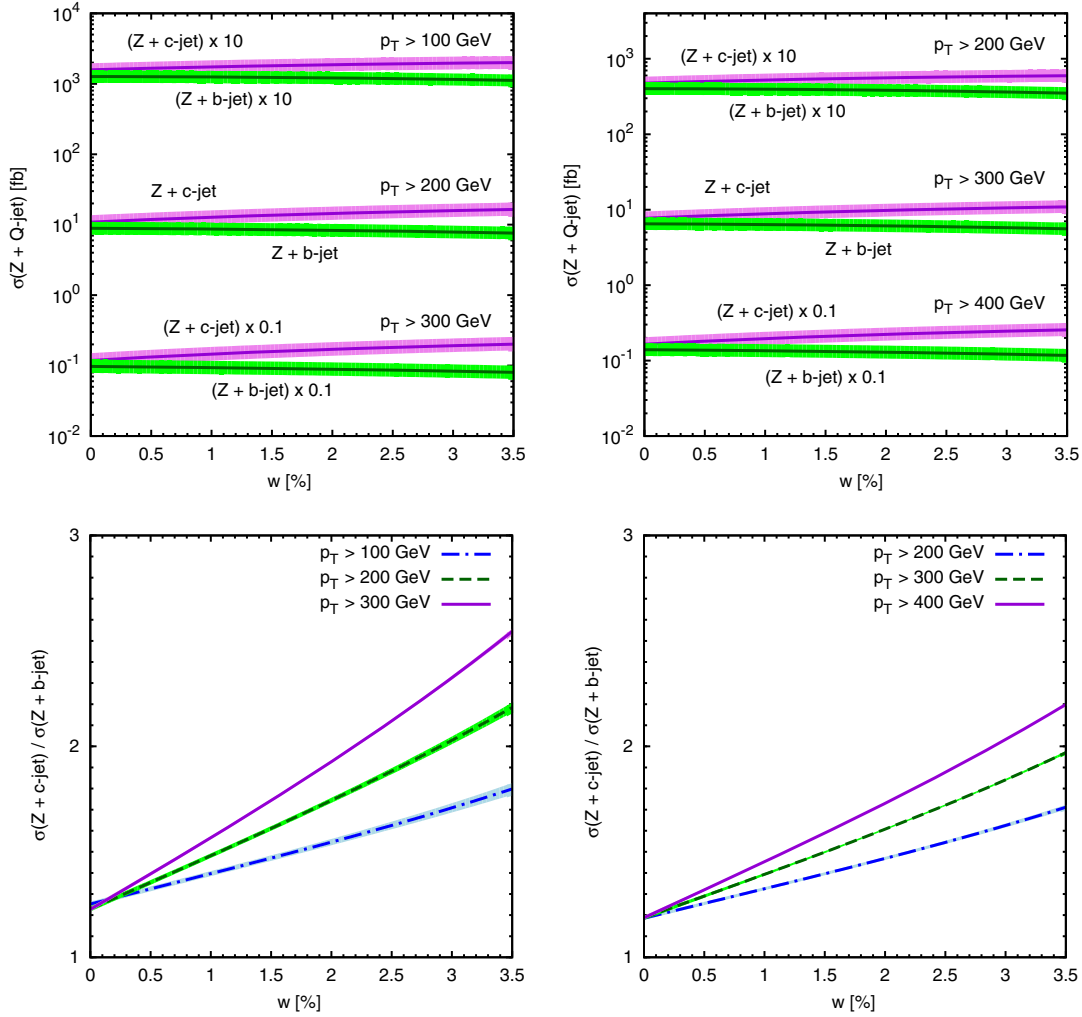


FIG. 9. Top: The cross sections of the associated  $Z + c$  and  $Z + b$  production in the  $pp$  collision as a function of  $w$  integrated over the  $Z$ -boson transverse momenta  $p_T > p_T^{\min}$  for different  $p_T^{\min}$  at  $\sqrt{s} = 8$  TeV (left) and  $\sqrt{s} = 13$  TeV (right). The kinematical conditions are described in the text. Bottom: The corresponding ratios of these cross sections. The calculations were done using the  $k_T$  factorization approach. The bands correspond to the usual scale variation as it is described in the text.

similar in the whole  $p_T$  region; therefore, below we will present the observables calculated within the  $k_T$  factorization approach only. The coincidence of these two calculations is due to effective allowance for the high-order corrections within the  $k_T$  factorization formalism (see, for example, Ref. [32] for more information). Both types of calculations predict a significant enhancement of  $p_T$  distributions due to the IC terms at  $p_T \geq 100$  GeV, which is in agreement with the previous studies [9,12,13].

The  $p_T$  spectrum ratios  $\sigma(\gamma + c)/\sigma(\gamma + b)$  and  $\sigma(Z + c)/\sigma(Z + b)$  versus  $p_T$  at different  $w$ 's are presented in Figs. 4 and 6. One can see that in the absence of the IC contribution, the ratio  $\sigma(\gamma + c)/\sigma(\gamma + b)$  is about 3 at  $p_T \sim 100$  GeV and decreases down to 2 at  $p_T \sim 500$  GeV. This behavior is the same for both energies  $\sqrt{s} = 8$  TeV and  $\sqrt{s} = 13$  TeV. If one takes into account the IC contributions, this ratio becomes approximately flat at  $w = 2\%$ , or

even increasing up to about 4 at  $w = 3.5\%$ . It is very close to the Tevatron data [20]: the constant ratio  $\sigma(\gamma + c)/\sigma(\gamma + b) \sim 3.5\text{--}4.5$  measured in the  $p\bar{p}$  collisions at  $110 < p_T < 300$  GeV and  $\sqrt{s} = 1.96$  TeV. However, this agreement cannot be treated as the IC indication due to huge experimental uncertainties (about 50%) and rather different kinematical conditions. If the IC contribution is included, the ratio  $\sigma(Z + c)/\sigma(Z + b)$  also increases by a factor of about 2 at  $w = 3.5\%$ , when the  $Z$ -boson transverse momentum grows from 100 GeV to 500 GeV (see Fig. 6). In the absence of the IC terms, this ratio slowly decreases.

One can consider other observables which could be useful to detect the IC signal—the cross sections discussed above but integrated over  $p_T > p_T^{\min}$ , where  $p_T^{\min} \geq 100$  GeV, and their ratios. Our predictions for such integrated cross sections versus the IC probability  $w$  at

$p_T^{\min} = 100, 200, \text{ and } 300 \text{ GeV}$  for  $\sqrt{s} = 8 \text{ TeV}$  and  $p_T^{\min} = 200, 300, \text{ and } 400 \text{ GeV}$  for  $\sqrt{s} = 13 \text{ TeV}$  are shown in Figs. 8 and 9.

All the  $p_T$  spectra have a significant scale uncertainty, as is shown in Ref. [13] (see also Figs. 8 and 9). According to Ref. [13], the ratio between the cross sections for the  $Z + Q$  and  $W + Q$  production in the  $pp$  collision is less sensitive to the scale variation calculated within the MCFM. Nevertheless, the uncertainty in this ratio at large  $p_T > 250 \text{ GeV}$  is about 40%–50%. In the present paper, we check these results for the ratios  $\sigma(\gamma + c)/\sigma(\gamma + b)$  and  $\sigma(Z + c)/\sigma(Z + b)$ . In Figs. 8 and 9 (bottom), we present these ratios versus the IC probability  $w$  calculated at different scales, when the cross sections of  $\gamma(Z) + Q$  production are integrated within the different intervals of transverse momentum. One can see a very small QCD scale uncertainty, especially at  $\sqrt{s} = 13 \text{ TeV}$  (bottom right), which is less than 1%. In contrast, the scale uncertainty for the integrated  $\gamma(Z) + Q$  cross sections (see Figs. 8 and 9, top) is significant and amounts to about 30%–40%. The sizable difference between the scale uncertainties for the ratios  $\sigma(Z + Q)/\sigma(W + Q)$  and  $\sigma(Z + c)/\sigma(Z + b)$  is due to the different matrix elements for the  $Z + Q$  and  $W + Q$  production in  $pp$  collisions, while the matrix elements for the  $Z + c$  and  $Z + b$  production are the same.

It is important that the calculated ratios  $\sigma(\gamma + c)/\sigma(\gamma + b)$  and  $\sigma(Z + c)/\sigma(Z + b)$  can be used to determine the IC probability  $w$  from the future LHC data. Moreover, these ratios are practically independent of the uncertainties of our calculations: actually, the curves corresponding to the usual scale variations as described above coincide with each other (see Figs. 8 and 9, bottom). Therefore, we can recommend these observables as a test for the hypothesis of the IC component inside the proton.

## V. CONCLUSION

The transverse momentum spectra of the prompt photons and  $Z$  bosons produced in association with the  $c$  or  $b$  jets in  $pp$  collisions are calculated using the MCFM (NLO pQCD)

and the  $k_T$  factorization approach at the LHC energies and pseudorapidities  $1.5 < \eta < 2.4$  using PDFs with and without the IC contribution. It is shown that these two approaches give similar results. We found that the contribution of the intrinsic charm can give a significant signal in the ratios  $\sigma(\gamma + c)/\sigma(\gamma + b)$  and  $\sigma(Z + c)/\sigma(Z + b)$  at forward pseudorapidities ( $1.5 < \eta < 2.4$ ) corresponding to the ATLAS and CMS facilities. If the IC contributions are taken into account, the ratio  $\sigma(\gamma + c)/\sigma(\gamma + b)$  as a function of the photon transverse momentum is approximately flat or increases at  $p_T > 100 \text{ GeV}$ . The similar flat behavior of this ratio was observed in the  $p\bar{p}$  annihilation at the Tevatron. In the absence of the IC contributions, this ratio decreases. Similarly, the ratio  $\sigma(Z + c)/\sigma(Z + b)$  increases when the  $Z$ -boson transverse momentum grows if the IC contribution is included and slowly decreases in the absence of the IC terms. We argued that the ratio of the cross sections  $\gamma(Z) + c$  and  $\gamma(Z) + b$  integrated over  $p_T > p_T^{\min}$  with  $p_T^{\min} \geq 100 \text{ GeV}$  can be used to determine the IC probability from the future LHC data. The advantage of the proposed ratios is that the theoretical uncertainties are very small, while the uncertainties for the  $p_T$  spectra of photons or  $Z$  bosons produced in association with the  $c$  or  $b$  jets are large. Therefore, the search for the IC signal by analyzing the ratio  $\sigma(\gamma/Z + c)/\sigma(\gamma/Z + b)$  can be more promising.

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