# The spectrum and separability of 2-qubit mixed X-states

Arsen Khvedelidze, <u>Astghik Torosyan</u>

LIT, JINR, DUBNA, RUSSIA

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#### Qubit

A generic mixed state <sup>1</sup> of an *n*-level quantum system is described by an  $n \times n$  complex matrix - the density matrix  $\rho$ , satisfying the following conditions:

- 1 Hermicity:  $\rho = \rho^{\dagger}$ ,
- 2 finite trace:  $Tr(\rho) = 1$ ,
- 3 positive semidefiniteness:  $\rho \ge 0$ .

The state of a qubit is given by a density matrix:

$$\rho = \frac{1}{2}(1 + \boldsymbol{\alpha} \cdot \boldsymbol{\sigma}), \qquad \boldsymbol{\alpha}^2 \leq 1.$$
(1)

where  $\alpha = Tr(\sigma \rho)$  is the expectation and  $\sigma$  is the set of Pauli matrices.

<sup>1</sup>The special class of idempotent matrices, satisfying  $\rho^2 = \rho$ , corresponds to the so-called *pure states*. A mixed state is a mixture of pure states.  $\square \rightarrow \square \square \square \square \square \square$ 

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#### **Composite states**

The space of states of the system, obtained by joining two systems 1 and 2, is a subspace of the tensor product of their individual Hilbert spaces  $\mathcal{H}_1$  and  $\mathcal{H}_2$ :

$$\mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2. \tag{2}$$

The density matrix  $\rho$ , describing mixed states of system  $\mathcal{H}$ , is separable, if it allows the convex decomposition:

$$\rho = \sum_{k} \omega_{k} \rho_{1}^{k} \otimes \rho_{2}^{k}, \qquad \sum_{k} \omega_{k} = 1, \quad \omega_{k} \ge 0,$$
(3)

where  $\rho_1^k$  and  $\rho_2^k$  represent the density matrices, acting on the corresponding multiplier of  $\mathcal{H}$ . Otherwise it is entangled.

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#### **Two qubits**

Consider the density matrix of two qubits, parametrized in the Fano form:

$$\rho = \frac{1}{4} [\mathbb{I}_2 \otimes \mathbb{I}_2 + \mathbf{a} \cdot \boldsymbol{\sigma} \otimes \mathbb{I}_2 + \mathbf{b} \cdot \mathbb{I}_2 \otimes \boldsymbol{\sigma} + c_{ij} \sigma_i \otimes \sigma_j], \qquad (4)$$

#### where

- a = (a<sub>1</sub>, a<sub>2</sub>, a<sub>3</sub>) and b = (b<sub>1</sub>, b<sub>2</sub>, b<sub>3</sub>) are the Bloch vectors of the constituent qubits,
- $C = ||c_{ij}||$  is the so-called "correlation matrix",
- $\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$ ,  $\sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$ ,  $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$  are the Pauli matrices.

### Are there mixed states, which are separable for an arbitrary spectrum of $\rho$ ?

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#### 2-qubit X-states

The density matrices of the form:

$$\rho_X := \begin{pmatrix} \rho_{11} & 0 & 0 & \rho_{14} \\ 0 & \rho_{22} & \rho_{23} & 0 \\ 0 & \rho_{32} & \rho_{33} & 0 \\ \rho_{41} & 0 & 0 & \rho_{44} \end{pmatrix}$$

$$\begin{cases} \rho_{11}, \rho_{22}, \rho_{33}, \rho_{44} \in \mathbb{R}, \\ \rho_{14} = \overline{\rho}_{14}, \ \rho_{23} = \overline{\rho}_{32}, \\ \sum_{i=1}^{4} \rho_{ii} = 1, \end{cases}$$
(5)

are called the X-states. The matrix (5) is unitary equivalent to the diagonal matrix

$$\rho_{X} = KWP \operatorname{diag}\left(\lambda_{1}, \lambda_{2}, \lambda_{3}, \lambda_{4}\right) PW^{\dagger}K^{\dagger}, \qquad (6)$$

where  $K = \exp\left(i\frac{u}{2}\sigma_3\right) \otimes \exp\left(i\frac{v}{2}\sigma_3\right) \in SU(2) \otimes SU(2)$  and

$$W = \begin{pmatrix} \frac{i\frac{\phi_1}{2}\sigma_2}{e^{i\frac{\phi_2}{2}\sigma_2}} & \mathbf{0} \\ \hline \mathbf{0} & e^{i\frac{\phi_2}{2}\sigma_2} \end{pmatrix}, \ P = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix}.$$
(7)

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#### **Spectrum of** 2–**qubit** *X*–**states**

The spectrum  $\{\lambda_1, \lambda_2, \lambda_3, \lambda_4\}$  of the diagonal density matrix  $\rho_X^2$  forms the partially ordered simplex <sup>3</sup>  $\Delta_3$  (Fig. 1):

$$\begin{cases} \sum_{i=1}^{4} \lambda_i = 1, \\ 0 \le \lambda_2 \le \lambda_1 \le 1, \\ 0 \le \lambda_4 \le \lambda_3 \le 1. \end{cases}$$
(8)



 ${}^{2}\rho_{X} = KWP \operatorname{diag}(\lambda_{1}, \lambda_{2}, \lambda_{3}, \lambda_{4}) PW^{\dagger}K^{\dagger}.$ 

<sup>3</sup>Partially ordered simplex is the quotient of a standard simplex by action of transposition subgroup  $P_2 \times P_2 \subset P_4$ .

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## The separability as a function of density matrices eigenvalues $\{\lambda\}$

According to the Peres-Horodecki criterion, which is a necessary and sufficient condition of separability for  $2 \otimes 2$  and  $2 \otimes 3$  dimensional systems, a state  $\rho$  is separable iff its partial transposition is semi-positive as well.

The partial transposition  $\rho^{T_2}$  of a 2-qubit density matrix with respect to the ordinary transposition operation T in the second subsystem is defined as:

$$\rho^{T_2} = I \otimes T\rho, \qquad T(\sigma_1, \sigma_2, \sigma_3) \to (\sigma_1, -\sigma_2, \sigma_3).$$
(9)

Similarly, one can use the alternative action:  $\rho^{T_1} = T \otimes I\rho$ .

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#### The separability conditions

Applying the Peres-Horodecki separability criterion to the X-state density matrix  $\rho_X$ <sup>4</sup>, we conclude, that it is separable iff:

$$(\lambda_1 - \lambda_2)^2 \cos^2 \phi_1 + (\lambda_3 - \lambda_4)^2 \sin^2 \phi_2 \le (\lambda_1 + \lambda_2)^2, \quad (10)$$

$$(\lambda_3 - \lambda_4)^2 \cos^2 \phi_2 + (\lambda_1 - \lambda_2)^2 \sin^2 \phi_1 \le (\lambda_3 + \lambda_4)^2.$$
(11)

New variables (x, y) and parameters (a, b, c, d) as functions of the density matrix eigenvalues and angles  $\phi_1$  and  $\phi_2$ :

$$\begin{cases} x = (\lambda_1 - \lambda_2)^2 \cos^2 \phi_1, \\ y = (\lambda_3 - \lambda_4)^2 \cos^2 \phi_2, \end{cases}$$
(12) 
$$\begin{cases} a = (\lambda_1 + \lambda_2)^2 - (\lambda_3 - \lambda_4)^2, \\ b = -(\lambda_1 - \lambda_2)^2 + (\lambda_3 + \lambda_4)^2, \\ c = (\lambda_1 - \lambda_2)^2, \ d = (\lambda_3 - \lambda_4)^2. \end{cases}$$
(13)

 ${}^{4}\rho_{X} = \mathcal{K}WP \operatorname{diag}\left(\lambda_{1}, \lambda_{2}, \lambda_{3}, \lambda_{4}\right) PW^{\dagger}K^{\dagger}.$ 

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The parameters (a, b, c, d) obey the inequalities:

$$a + b \ge 0, \quad a + d \ge 0, \quad b + c \ge 0,$$
 (14)

Thus, the separability conditions in the form of two inequalities (10)<sup>5</sup> and (11) <sup>6</sup> linearize:

$$\begin{cases} x - y \le a, & 0 \le x \le c, \\ y - x \le b, & 0 \le y \le d. \end{cases}$$
(15)

Hence, the inequalities (15) have solutions for all possible values of parameters from the restrictions (14).

$$5(\lambda_1 - \lambda_2)^2 \cos^2 \phi_1 + (\lambda_3 - \lambda_4)^2 \sin^2 \phi_2 \le (\lambda_1 + \lambda_2)^2.$$

$$6(\lambda_3 - \lambda_4)^2 \cos^2 \phi_2 + (\lambda_1 - \lambda_2)^2 \sin^2 \phi_1 \le (\lambda_3 + \lambda_4)^2.$$

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For eigenvalues from the partially ordered simplex  $\Delta_3$  the separability conditions <sup>7</sup> of the *X*-state density matrix  $\rho_X$  <sup>8</sup> determine non empty domain (Fig. 2) for angles  $\phi_1$  and  $\phi_2$ .



Figure 2: Plots (I-V) - families of solutions: Domain (I) : a < 0, b = -a,  $c \ge 0$ ,  $d \ge b$ ; Domain (II): a < 0, b > -a,  $c \ge 0$ ,  $d \ge -a$ ; Domain (III): a = 0,  $b \ge 0$ ,  $c \ge 0$ ,  $d \ge 0$ ; Domain (IV): a > 0,  $-a \le b \le 0$ ,  $c \ge -b$ ,  $d \ge 0$ ; Domain (V): a > 0, b > 0,  $c \ge 0$ ,  $d \ge 0$ .

There exists 4-parametric family of separable mixed X-states of 2-qubits with an arbitrary spectrum of the density matrix.

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#### The absolute separability conditions

The states of *n*-dimensional quantum system, remaining separable under the adjoint action of the SU(n)-transformations of  $n \times n$  density matrices, are called absolute separable.

Are there X-states, which are separable for arbitrary angles  $\phi_1$  and  $\phi_2$ ? The inequalities in the eigenvalues of X-matrices, defining the absolutely separable X-states, read:





Figure 3: The absolute separability region.

#### On the generalization to an arbitrary 2-qubit states

The Peres-Horodecki separability criterion can be written in the form of polynomial inequalities in the SU(4) Casimir invariants  $\mathfrak{C}_2$ ,  $\mathfrak{C}_3$ ,  $\mathfrak{C}_4$  and two  $SU(2) \times SU(2)$ -invariant polynomials <sup>9</sup>. In general case,

- the determinants det(C), det(M) are analogues of angles  $\phi_1$ ,  $\phi_2$  of X-states,
- the Casimir invariants  $\mathfrak{C}_2$ ,  $\mathfrak{C}_3$ ,  $\mathfrak{C}_4$  are analogues of eigenvalues  $\lambda_1$ ,  $\lambda_2$ ,  $\lambda_3$ ,  $\lambda_4$  of X-states.

<sup>9</sup>Here the determinants of correlation  $C = ||c_{ij}||$  and Schlienz-Mahler matrix  $M = ||c_{ij} - a_i b_j||$  are the  $SU(2) \times SU(2)$ - polynomial invariants  $B \times A \equiv A = A$ 

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#### Conjecture

#### **CONJECTURE:** The inequalities

$$\begin{cases} 0 \leq 3\mathfrak{C}_2 - 2\mathfrak{C}_3 - 4 \det(\mathcal{C}) \leq 1, \\ 0 \leq (1 - 3\mathfrak{C}_2)^2 + 8\mathfrak{C}_3 - 12\mathfrak{C}_4 + 16 \det(\mathcal{M}) \leq 1, \end{cases}$$
(17)

have solutions for unknown det(C) and det(M) for all values of Casimir invariants  $\mathfrak{C}_2$ ,  $\mathfrak{C}_3$  and  $\mathfrak{C}_4$ , which are constrained by the inequalities:

$$\begin{cases} 0 \le \mathfrak{C}_2 \le 1, \\ 0 \le 3\mathfrak{C}_2 - 2\mathfrak{C}_3 \le 1, \\ 0 \le (1 - 3\mathfrak{C}_2)^2 + 8\mathfrak{C}_3 - 12\mathfrak{C}_4 \le 1. \end{cases}$$
(18)

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#### Thank you for attention



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