

What is the appropriate field theory?

Kay Wiese

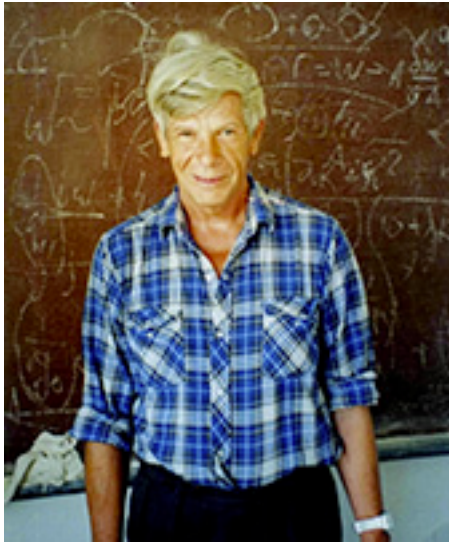
Ecole Normale Supérieure (ENS), Paris
with Cathelijne ter Burg, Gauthier Mukerjee

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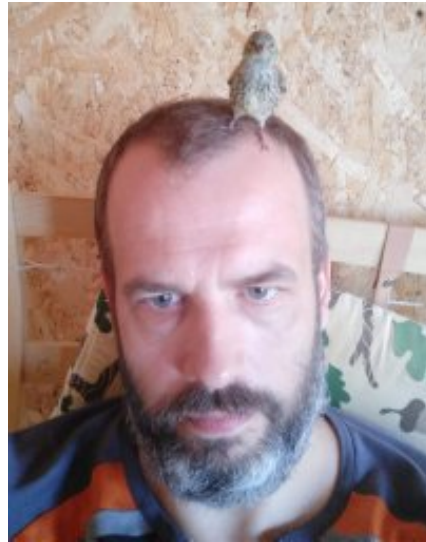
<http://www.phys.ens.fr/~wiese/>

Review: arXiv:2102.01215

...dedicated to my Russian and Ukrainian friends and colleagues...



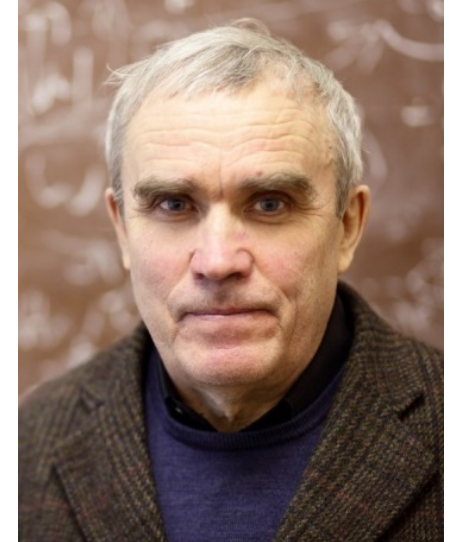
Sasha



Misha



Andrei



Juri



Nicolai



Kolya



Boris



Mykola

What is the appropriate field theory?

- Ising model: ϕ^4 -theory ✓
- quantum gravity ?
- turbulence ?
- KPZ equation ?

we believe that for KPZ

$$\mathcal{S} = \int_{x,t} \tilde{h}(x,t) \left[\partial_t h(x,t) + \nabla^2 h(x,t) + \lambda \left[\nabla h(x,t) \right]^2 + 2D\tilde{h}(x,t) \right]$$

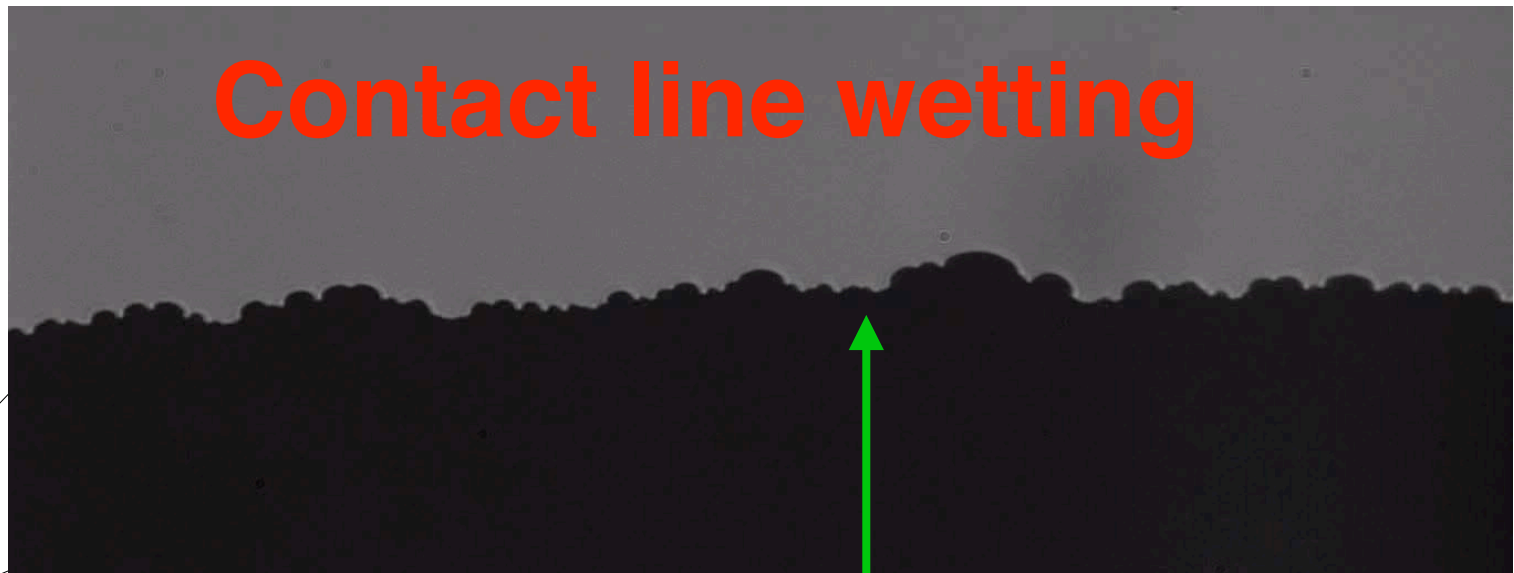
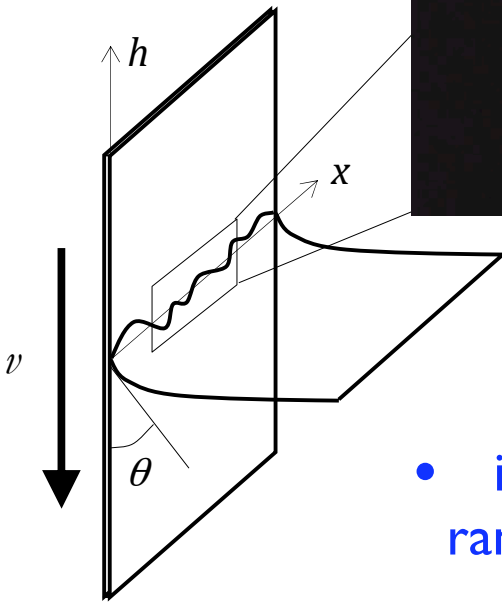
- but the perturbative expansion describes weak-strong coupling crossover in $d = 2 + \epsilon$ dimensions. Something is missing here...
- strong-coupling regime accessible via the directed polymer, but in completely different variables!

How to find the appropriate field theory?

- how to get out of this dilemma?
- try to measure the effective theory!
- here: disordered elastic systems

(C) E. Rolley

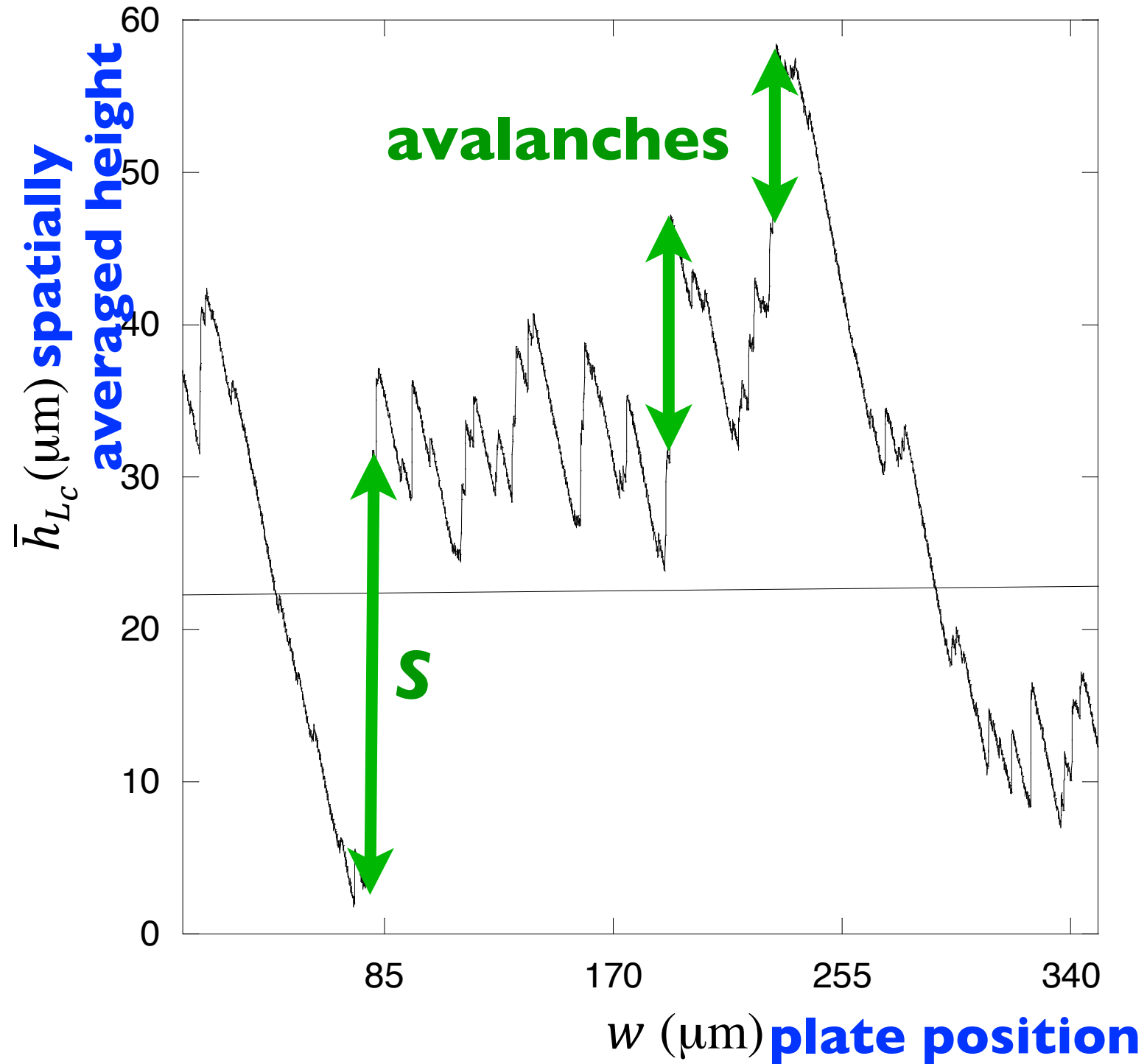
Contact line wetting



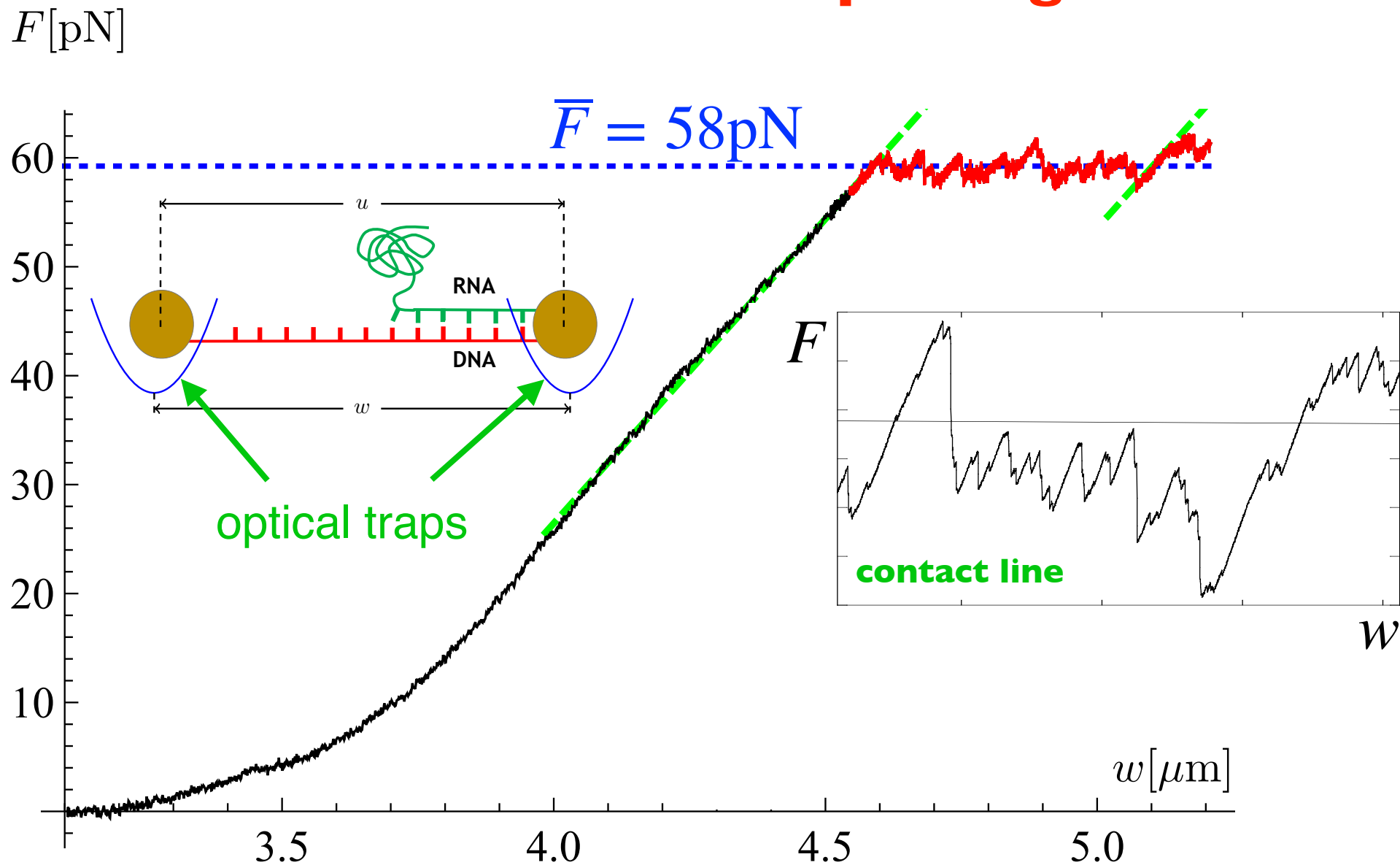
avalanche

- isobutanol on a randomly silanized

height jumps = avalanches

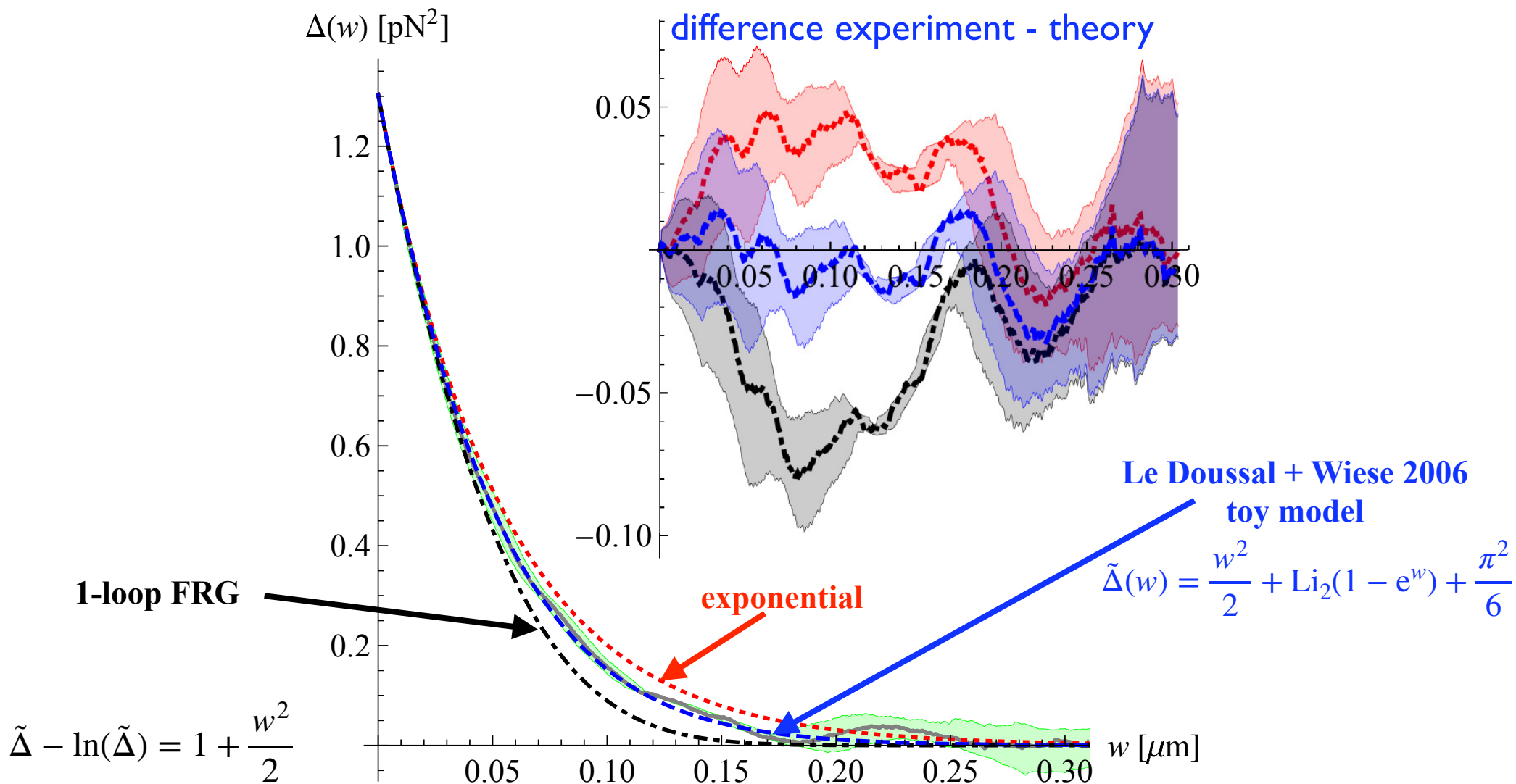


Force as a function of distance for RNA/DNA peeling



Force-force correlations

$$\Delta(w - w') := \overline{F_w F_{w'}}^c \equiv \overline{F_w F_{w'}} - \overline{F_w} \overline{F_{w'}}$$



Field theory

Equation of motion (for SR elasticity for simplicity)

height of the interface

$$\partial_t u(x, t) = \nabla^2 u(x, t) + m^2[w - u(x, t)] + F(x, u(x, t))$$

$w = vt$

Forces are drawn from a **Gaussian**, and have correlations

$$\overline{F(x, u)F(x', u')^c} = \delta^d(x - x')\Delta(u - u')$$

Field theory (MSR=classical limit $\hbar \rightarrow 0$ of Keldysh)

$$\mathcal{S}[\tilde{u}, u] = \int_{x,t} \tilde{u}(x, t) \left[\partial_t u(x, t) - \nabla^2 u(x, t) + m^2(u(x, t) - w) \right] \\ - \frac{1}{2} \int_{x,t,t'} \tilde{u}(x, t) \tilde{u}(x, t') \Delta(u(x, t) - u(x, t'))$$

was measured

Why did we measure Δ ?

action

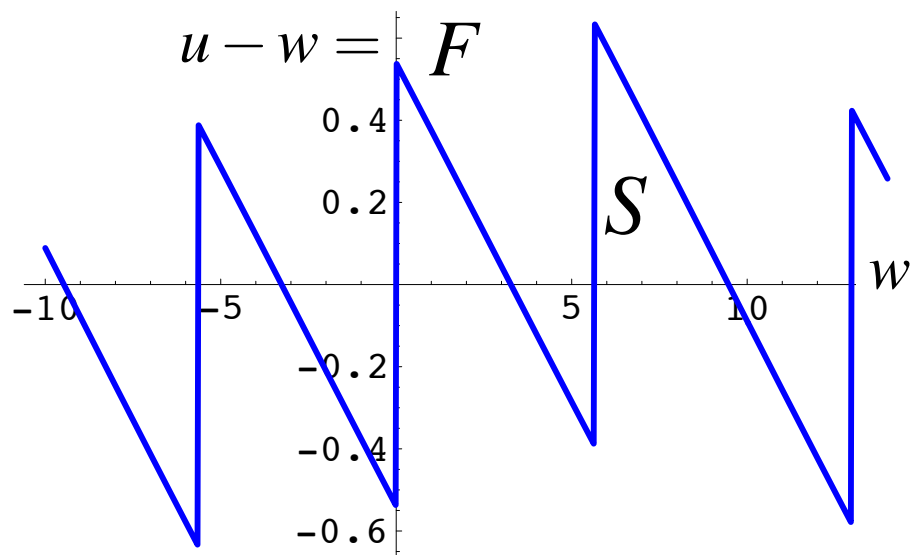
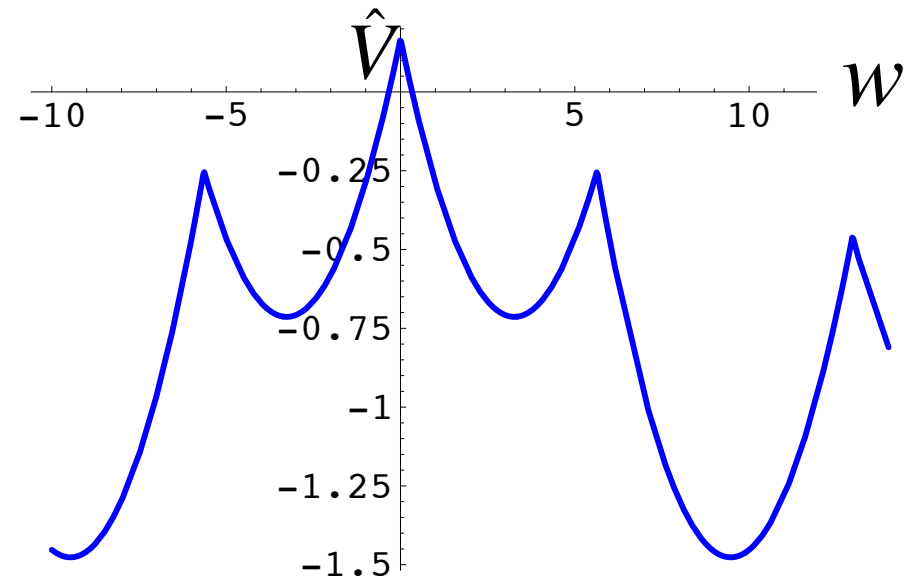
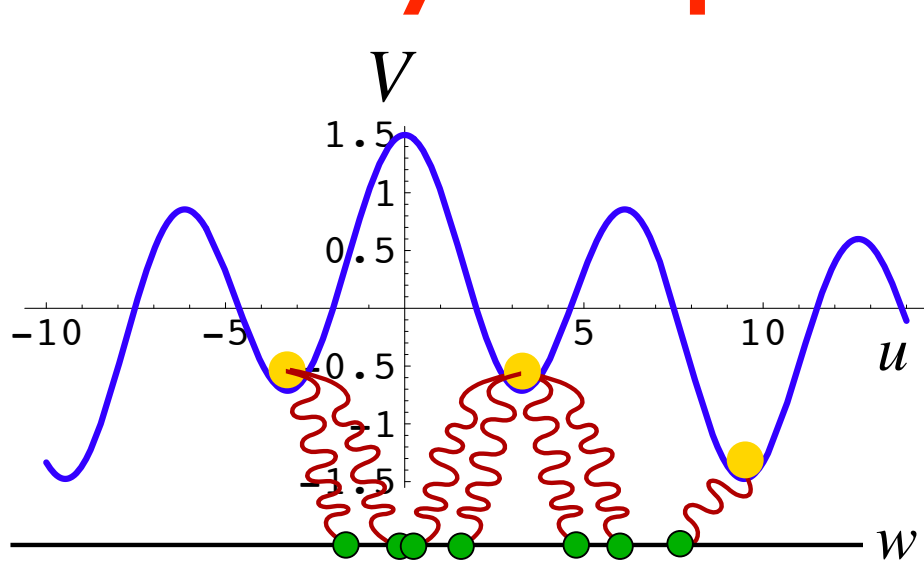
$$\mathcal{S}[\tilde{u}, u] = \int_{x,t} \tilde{u}(x, t) \left[\partial_t u(x, t) - \nabla^2 u(x, t) + m^2 \left(u(x, t) - w \right) \right]$$
$$- \frac{1}{2} \int_{x,t,t'} \tilde{u}(x, t) \tilde{u}(x, t') \Delta \left(u(x, t) - u(x, t') \right)$$

$$u_w := \lim_{t \rightarrow \infty} \frac{1}{L^d} \int_x u(x, t) \Big|_w$$

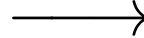
center of mass at large t , i.e. $\omega \rightarrow 0$

$$\Delta(w - w') \equiv \Gamma^{(2)} = \mathcal{L} \circ \overline{u_w u_{w'}} = [\mathcal{R}^{-1}]^2 \overline{u_w u_{w'}} = (m^2)^2 \overline{u_w u_{w'}}$$

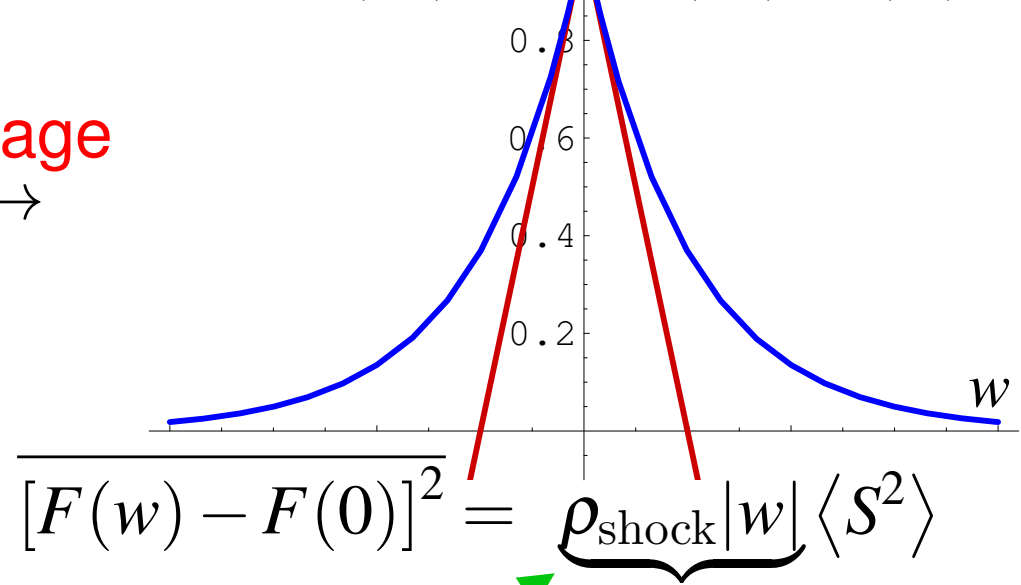
Why a cusp in the effective action?



average

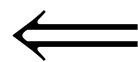


$$\Delta(w) = \overline{F(w)F(0)}^c$$

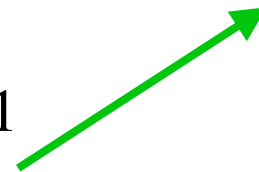


$$\overline{[F(w) - F(0)]^2} = \underbrace{\rho_{\text{shock}} |w|}_{p_{\text{shock}}} \langle S^2 \rangle$$

$$-\Delta'(0^+) = \frac{\langle S^2 \rangle}{2 \langle S \rangle}$$



$$\rho_{\text{shock}} = \langle S \rangle^{-1}$$

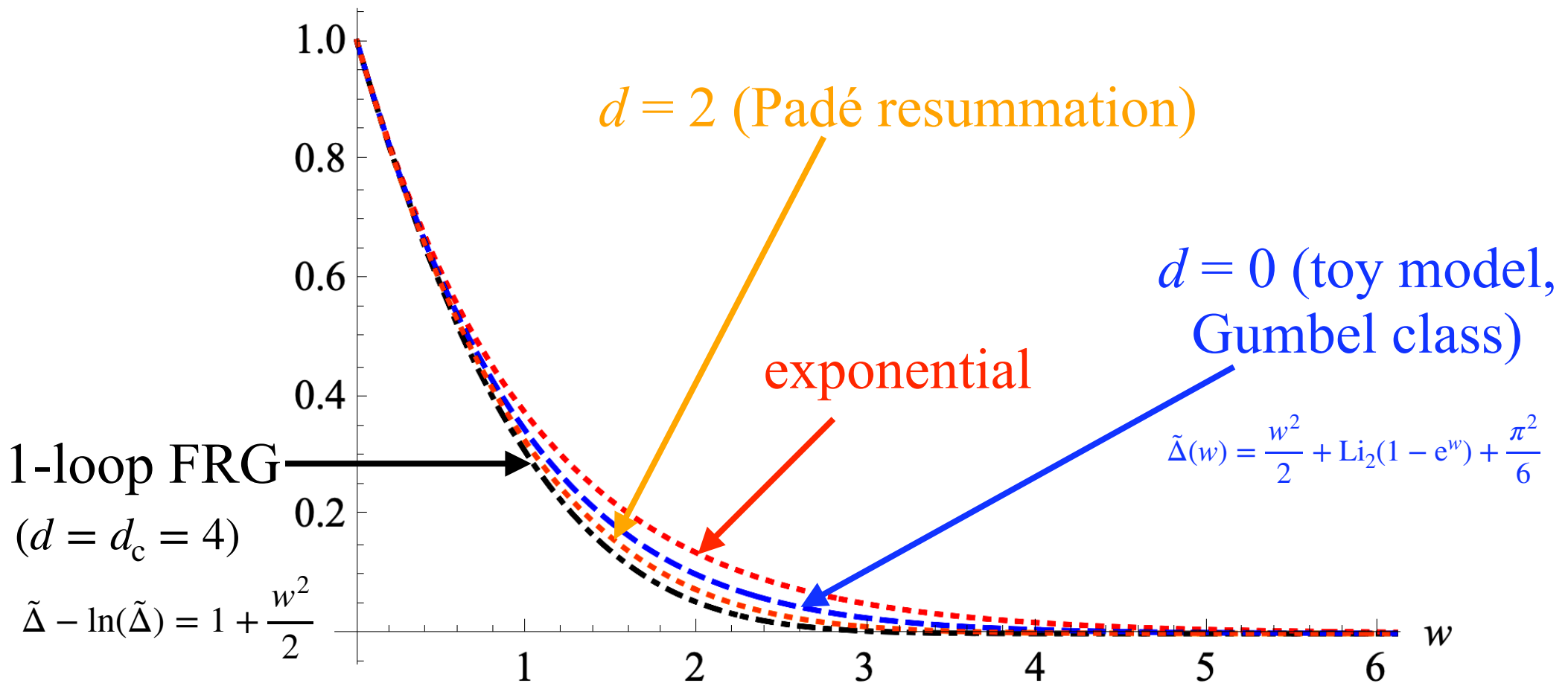


Renormalization of disorder

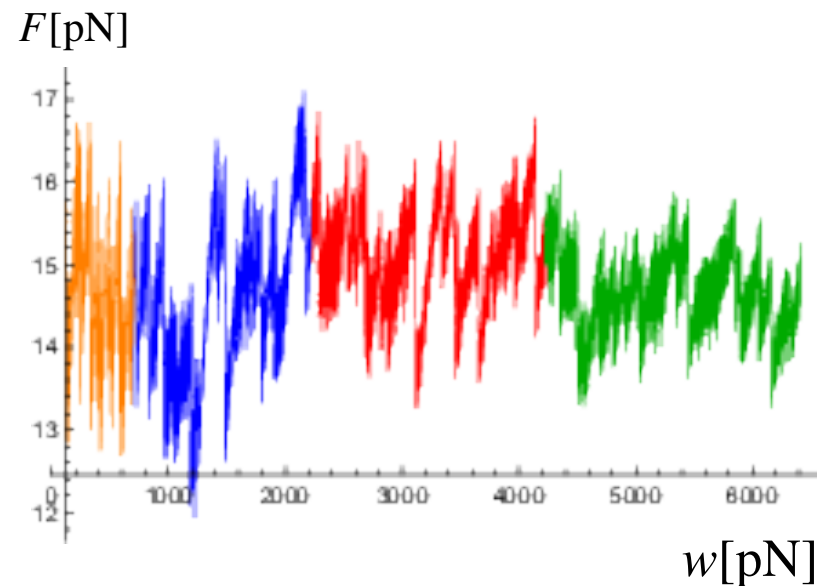
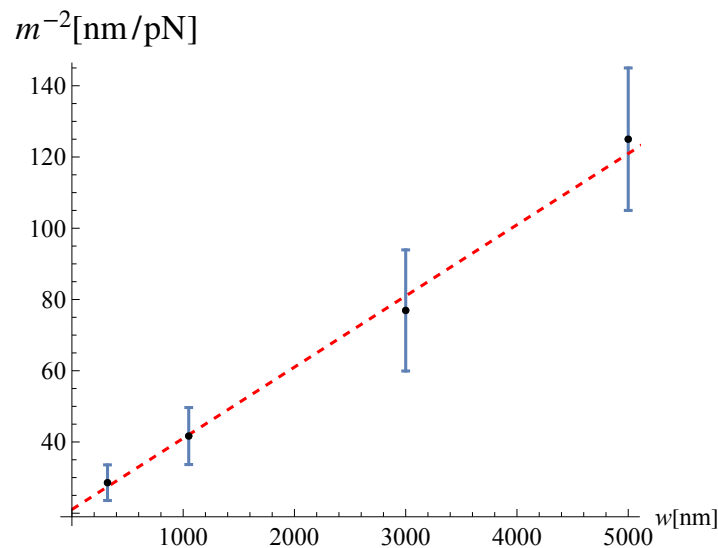
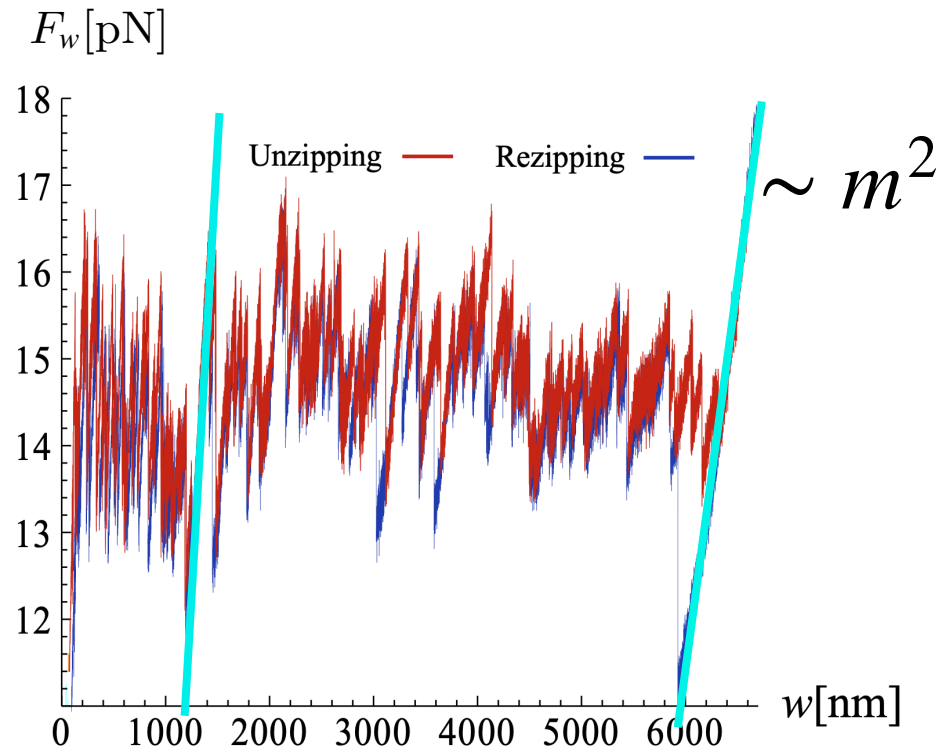
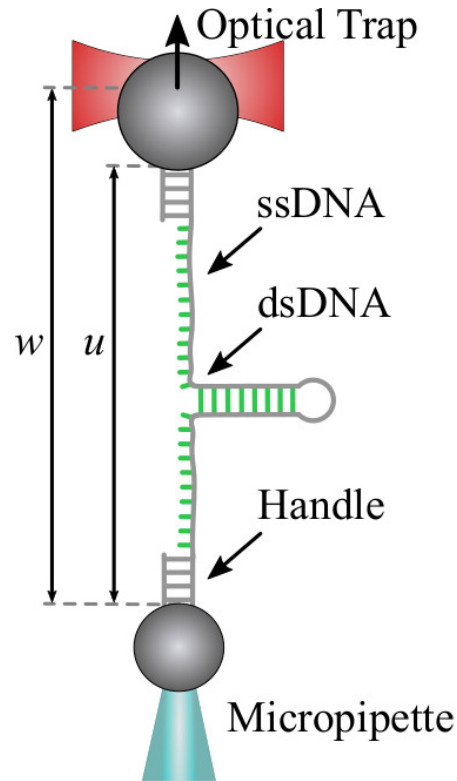
FRG

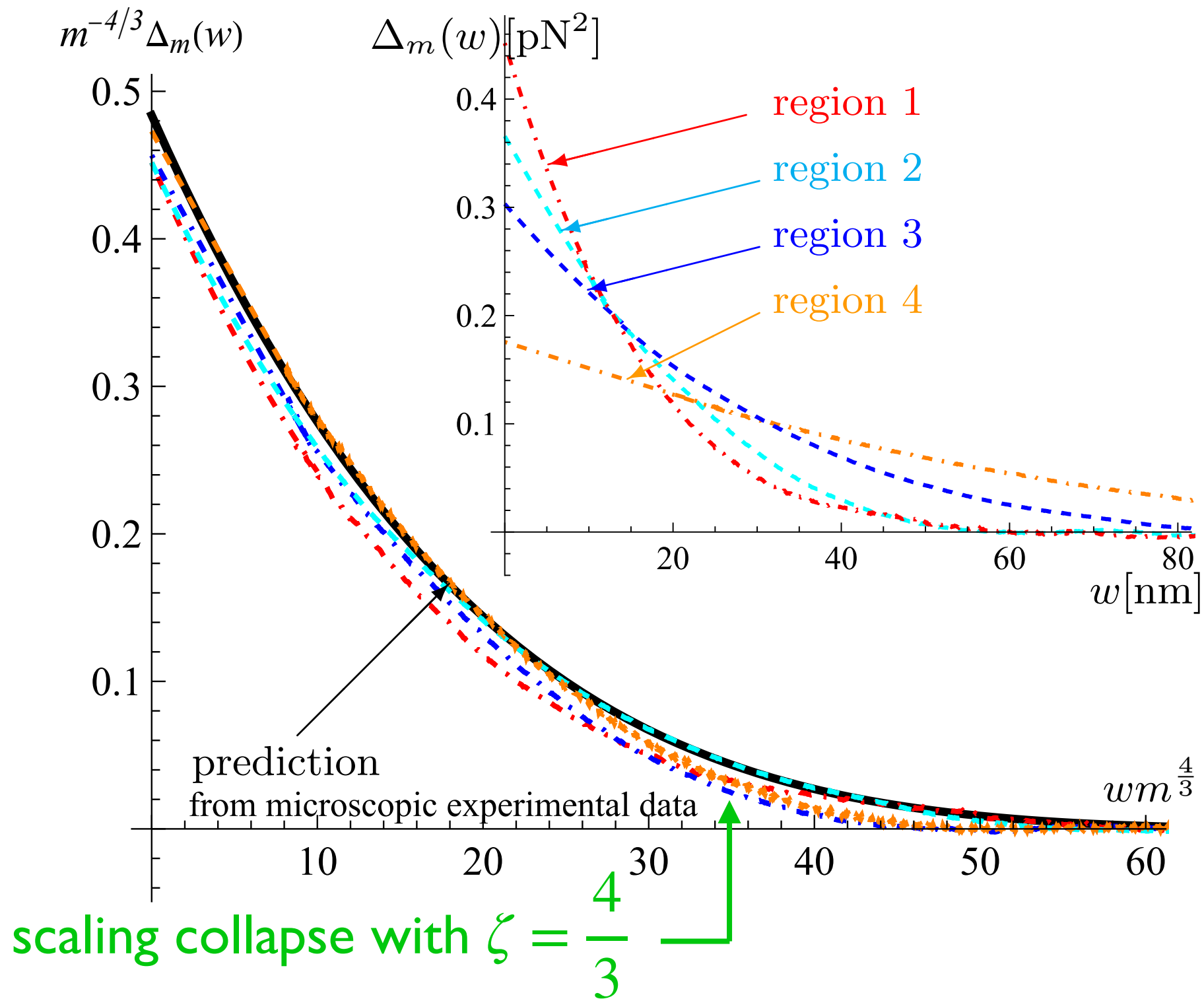
$$-\frac{md}{dm}\tilde{\Delta}(w) = (\epsilon - 2\zeta)\tilde{\Delta}(w) + \zeta w\tilde{\Delta}'(w) - \frac{1}{2}\partial_w^2 [\tilde{\Delta}(w) - \tilde{\Delta}(0)]^2 + \frac{1}{2}\partial_w^2 \left\{ [\tilde{\Delta}(w) - \tilde{\Delta}(0)] \tilde{\Delta}'(w)^2 + \tilde{\Delta}'(0^+)^2 \tilde{\Delta}(w) \right\}$$

Chauve, Le Doussal, Wiese 2004



Renormalization in DNA-unzipping

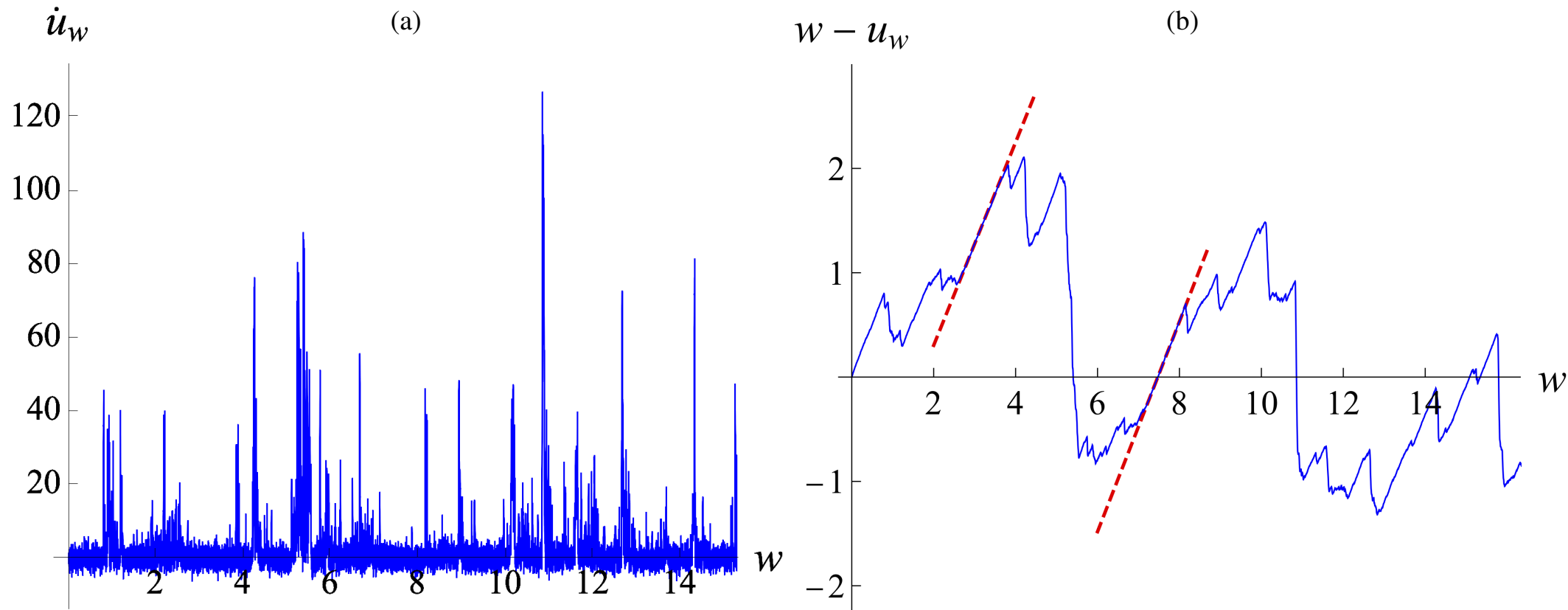




Magnetic domain walls ($d = 2$)

(data by F. Bohn, G. Durin, R.L. Sommer)

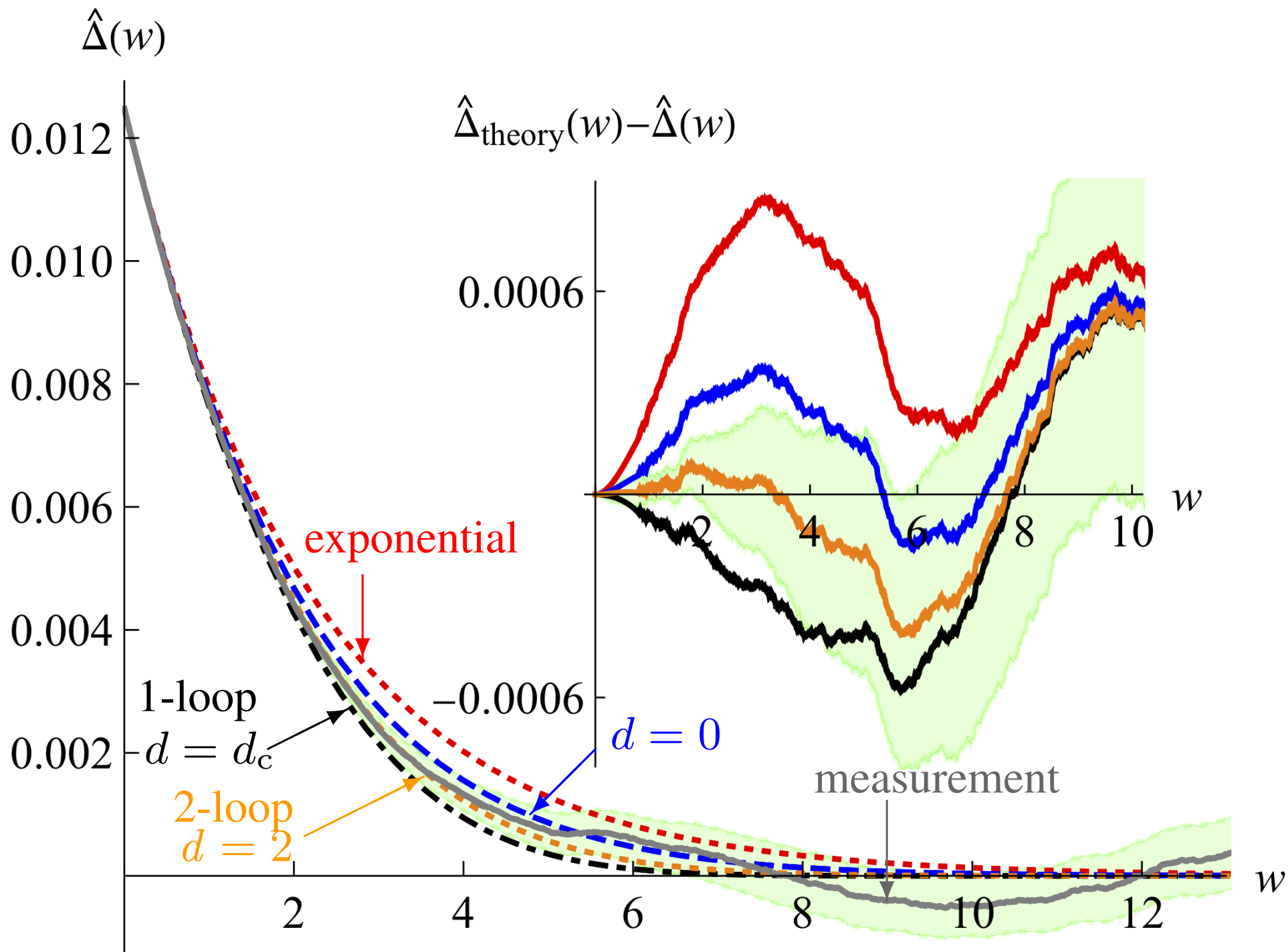
current in a pickup coil allows to construct :



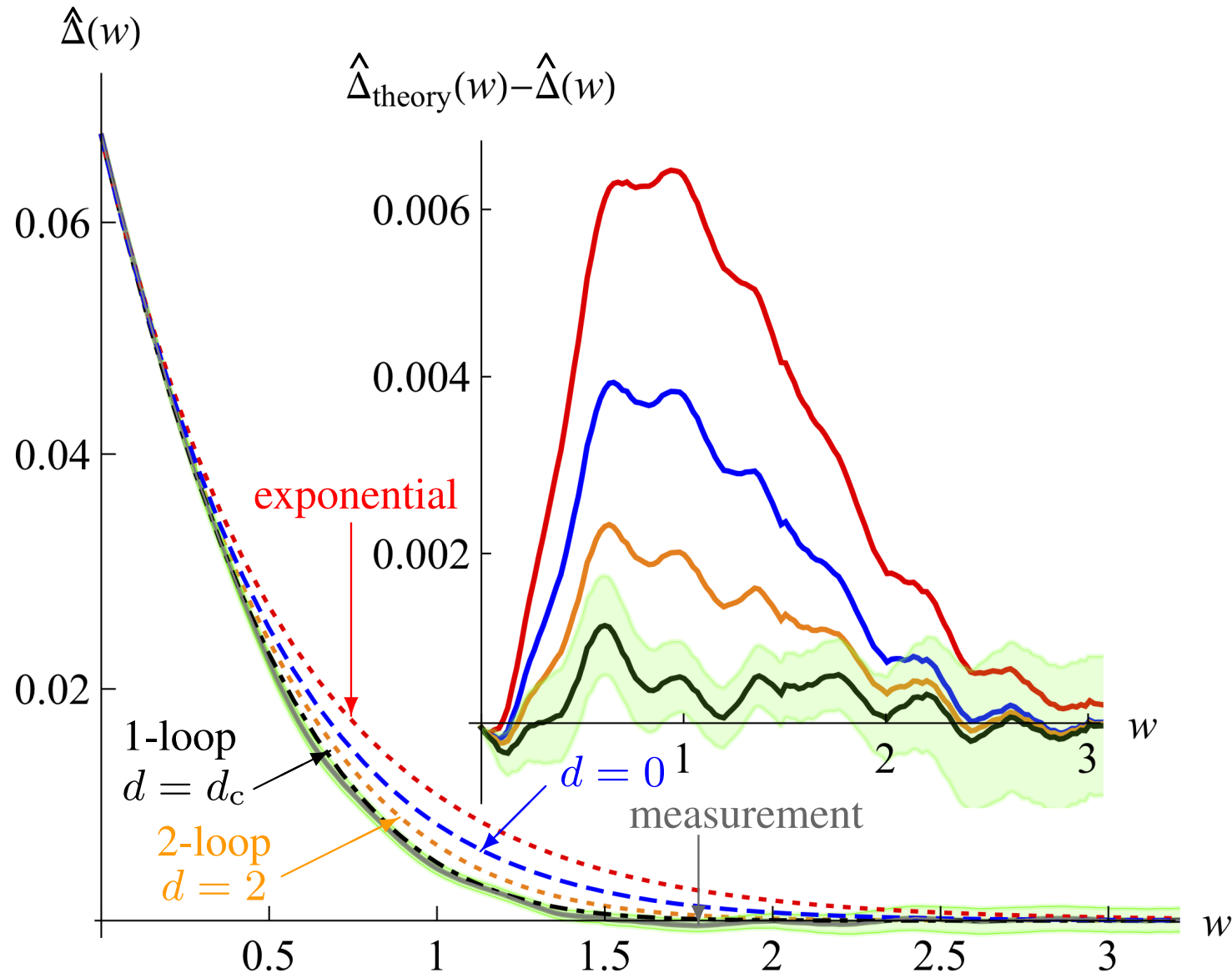
eliminate one **unknown** scale by the definition

$$\hat{\Delta}_v(w - w') := \overline{[w - u_w] [w' - u_{w'}]}^c = \frac{1}{m^4} \overline{F_w F_{w'}}^c$$

Magnetic domain walls SR elasticity ($d = 2$)

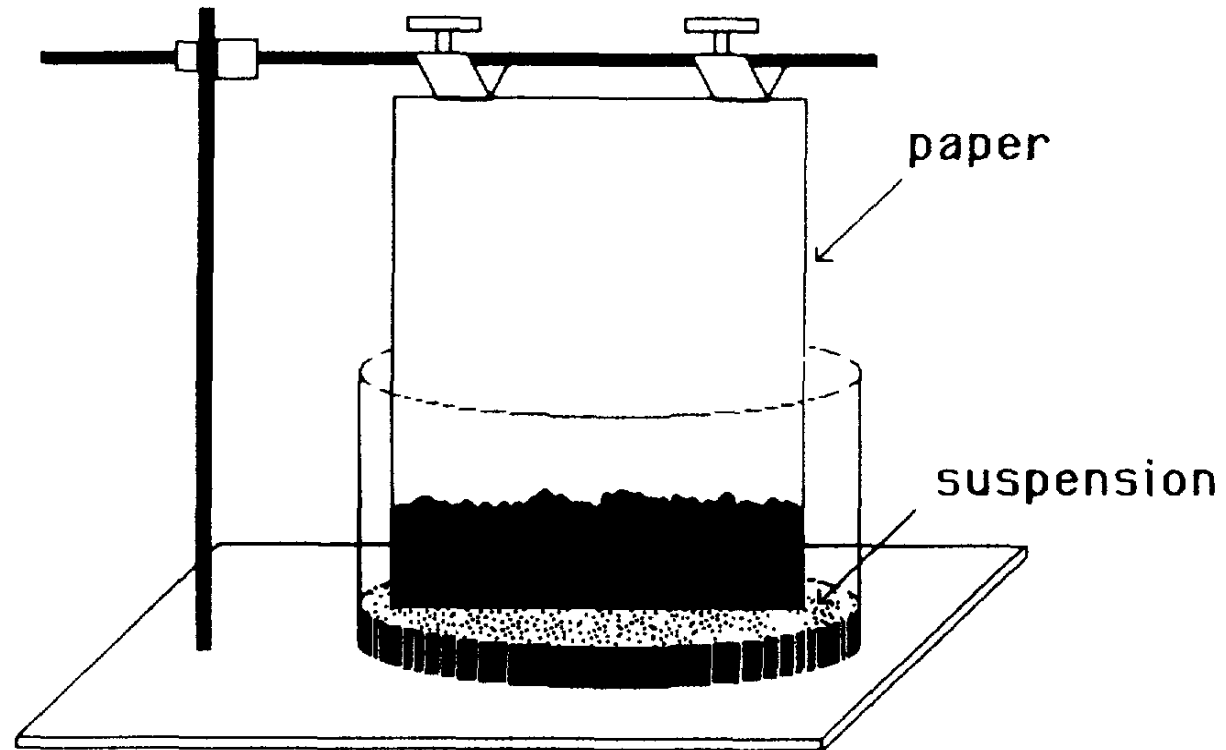


Domain walls ($d = 2$) with long-range elasticity



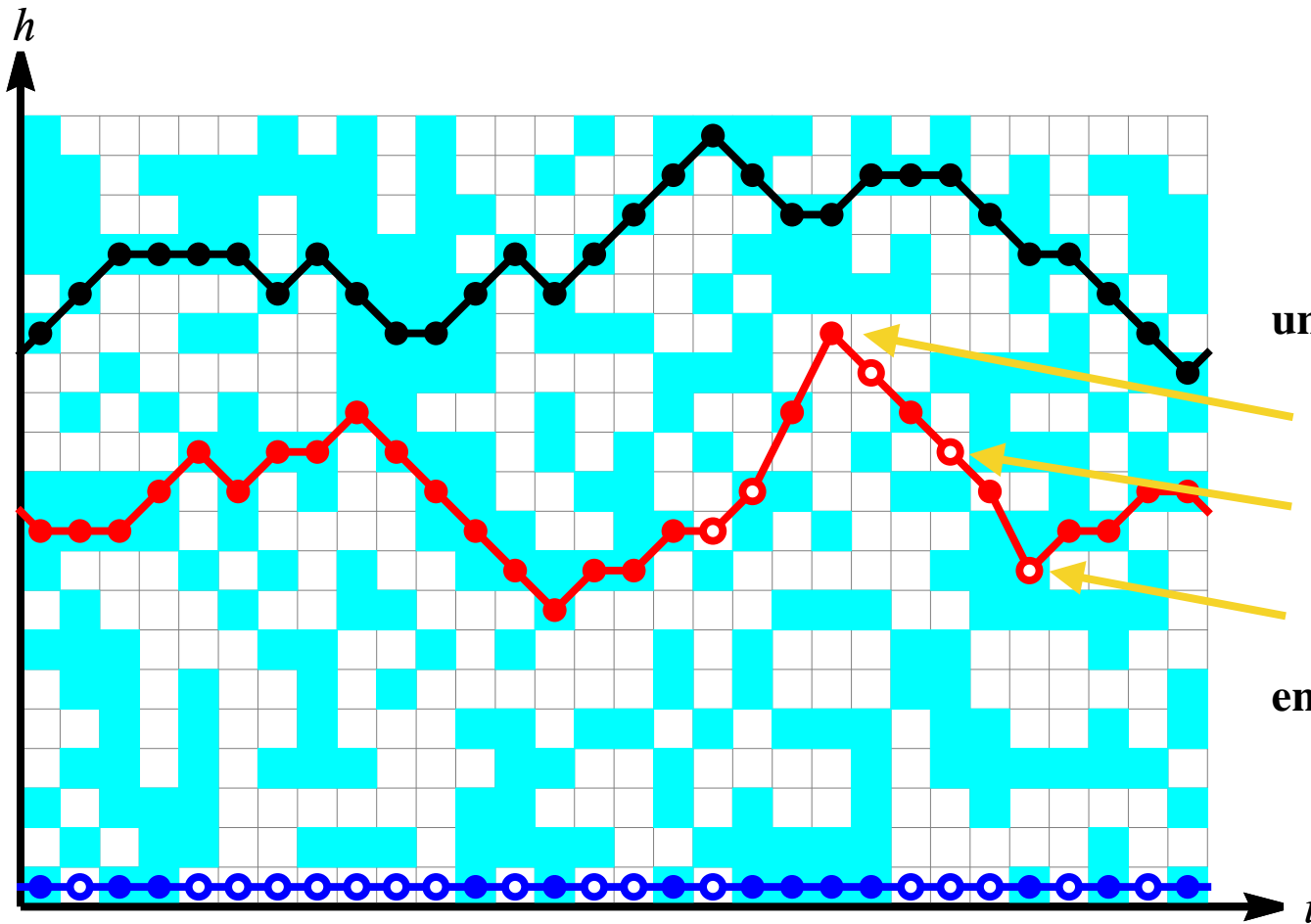
- 1-loop FRG gives fixed point.
- this is not ABBM disorder: $\Delta(0) - \Delta(w) \neq \sigma |w|$
- ABBM only gives short-scale behavior correctly

Imbibition



The Tang-Leschhorn cellular automaton of 1992

TL92



unstable(i)

links cannot be longer than 2

if $h(i) - h(\text{neighbor}) \geq 2$ **return** false

move forward if open

if $f(i, h(i)) > f_c$ **return** true

move forward if a neighbor is 2 ahead

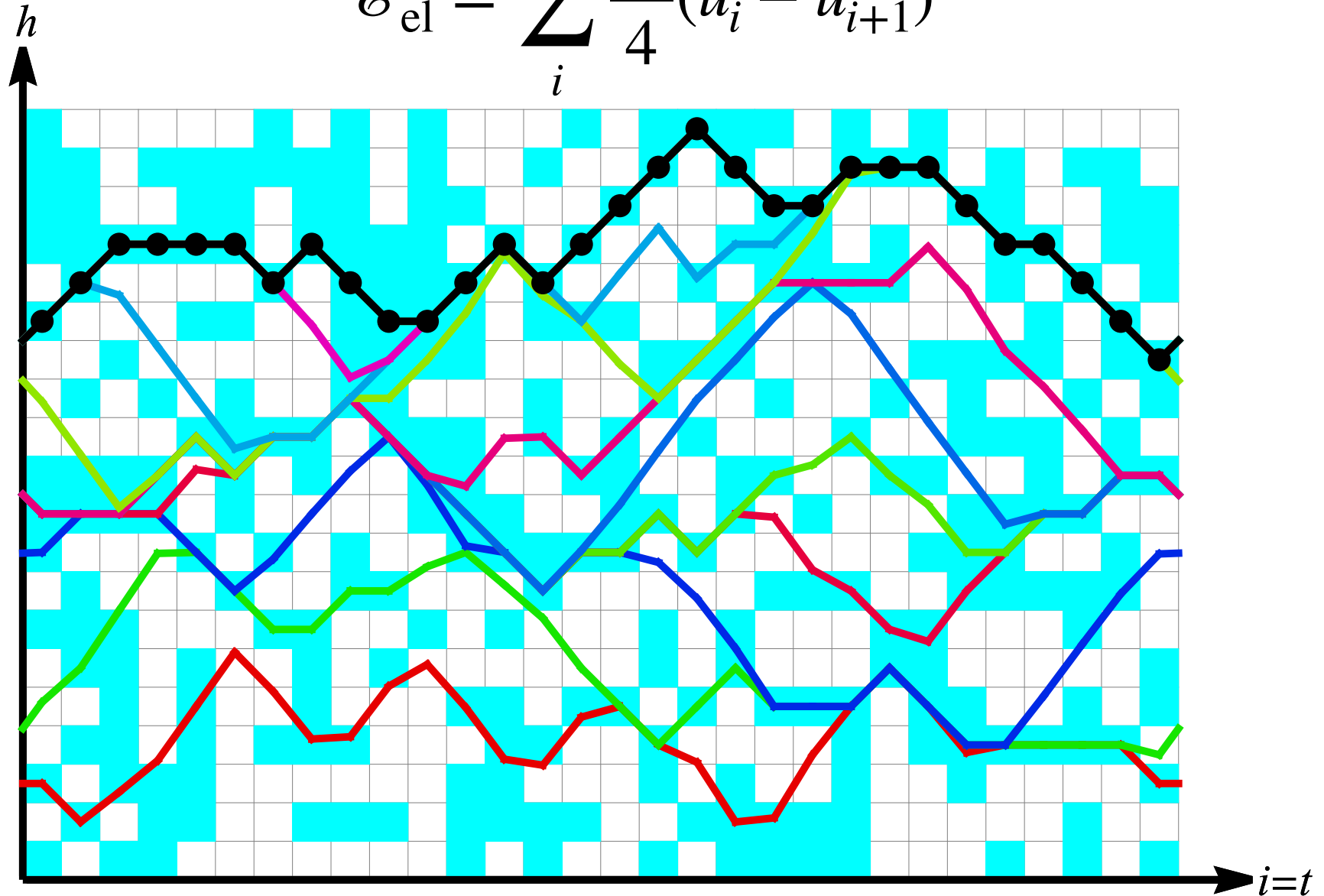
if $h(\text{neighbor}) - h(i) \geq 2$ **return** true

end

variants: Buldyrev, S. Havlin and H.E. Stanley 1992

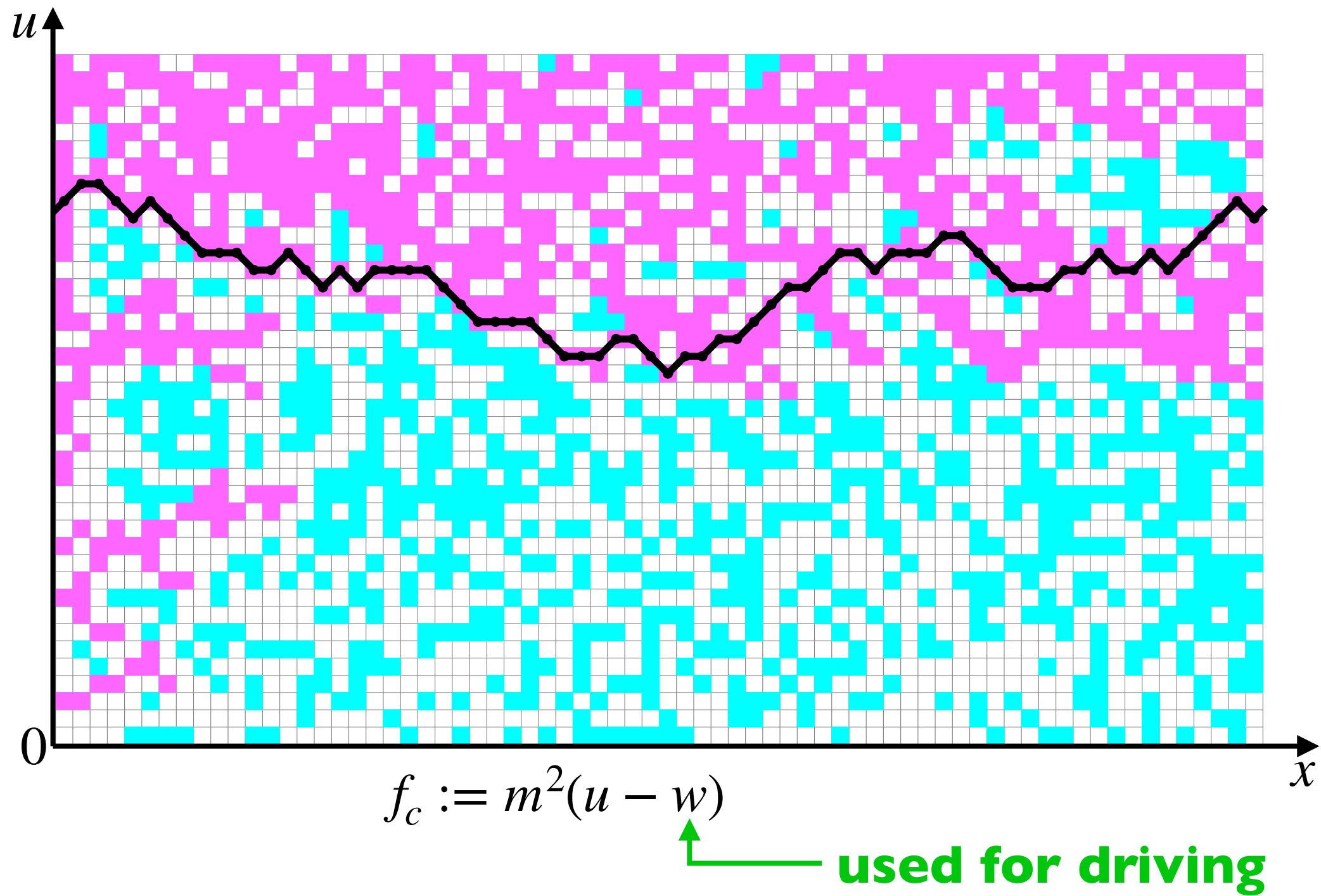
Anharmonic depinning = TL92

$$\mathcal{E}_{\text{el}} = \sum_i \frac{c_4}{4} (u_i - u_{i+1})^4$$



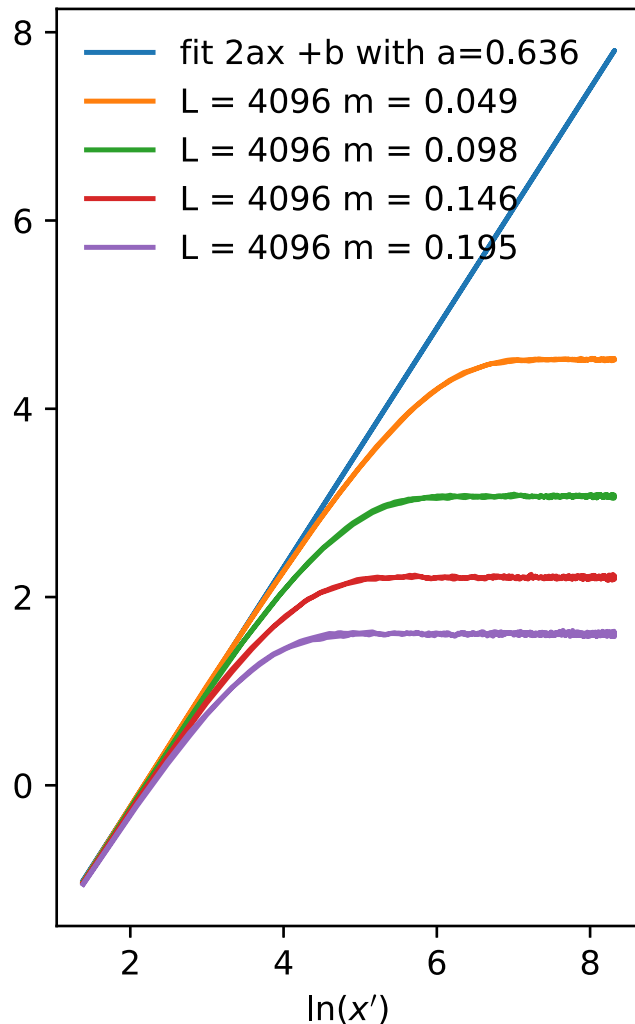
anharmonic depinning respects the Middleton theorem
= return point memory (not guaranteed for qKPZ)

TL92 and directed percolation ($d = 1$)



2-point function

$$\frac{1}{2} \overline{[u(x) - u(y)]^2} \sim \begin{cases} A |x - y|^{2\zeta}, & |x - y| < \xi \\ B m^{-2\zeta_m}, & |x - y| > \xi \end{cases}$$



from directed percolation

$$\zeta^{d=1} = \frac{\nu_{\perp}}{\nu_{\parallel}} = 0.632613(3)$$

$$\zeta_m^{d=1} = \frac{2\nu_{\perp}}{1 + \nu_{\perp}} = 1.046190(4)$$

two distinct exponents in all d

$$\zeta_m > \zeta$$

What is the appropriate long-distance theory?

Can we measure it?

standard elasticity
 $c \rightarrow 0$

non-linear elasticity

$$\eta \partial_t u(x, t) = \boxed{c \nabla^2 u(x, t)} + c_4 \nabla [\nabla u(x, t)]^3 - m^2 [u(x, t) - w]$$

$+ F(x, u(x, t))$

disorder force

confining potential

background field

What is the appropriate long-distance theory?

Can we measure it?

standard elasticity
 $c > 0$

non-linear elasticity

$$\eta \partial_t u(x, t) = c \nabla^2 u(x, t) + c_4 \nabla [\nabla u(x, t)]^3 - m^2 [u(x, t) - w(x, t)] + \lambda [\nabla u(x, t)]^2 + F(x, u(x, t))$$

KPZ term

disorder force

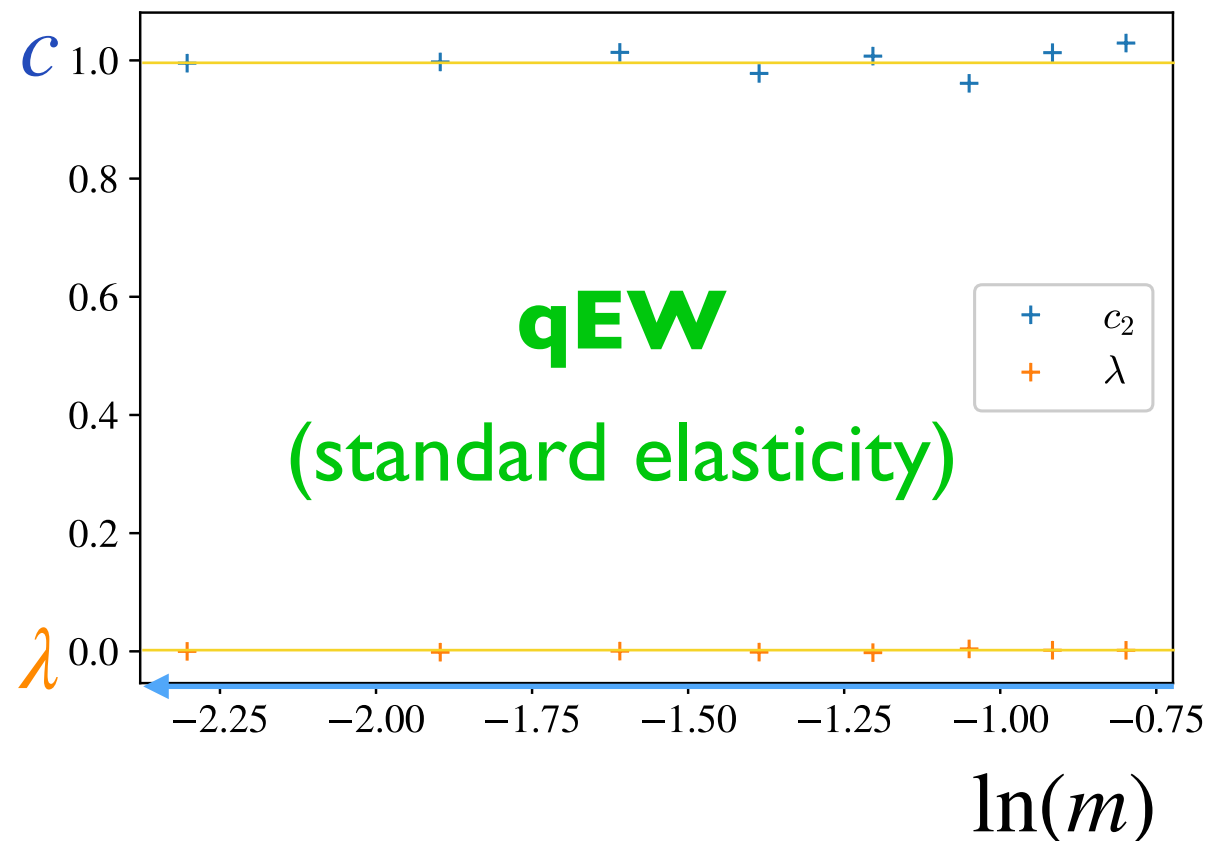
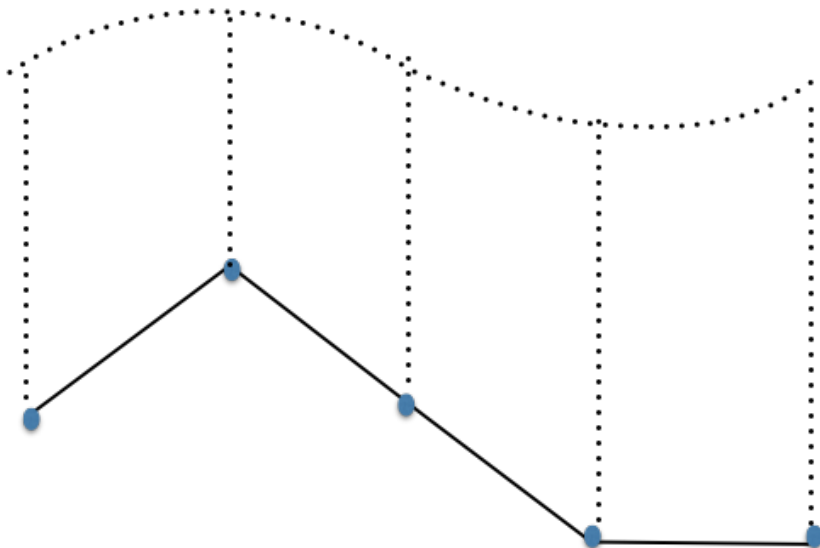
confining potential
(unrenormalized)

background field
(modulated)

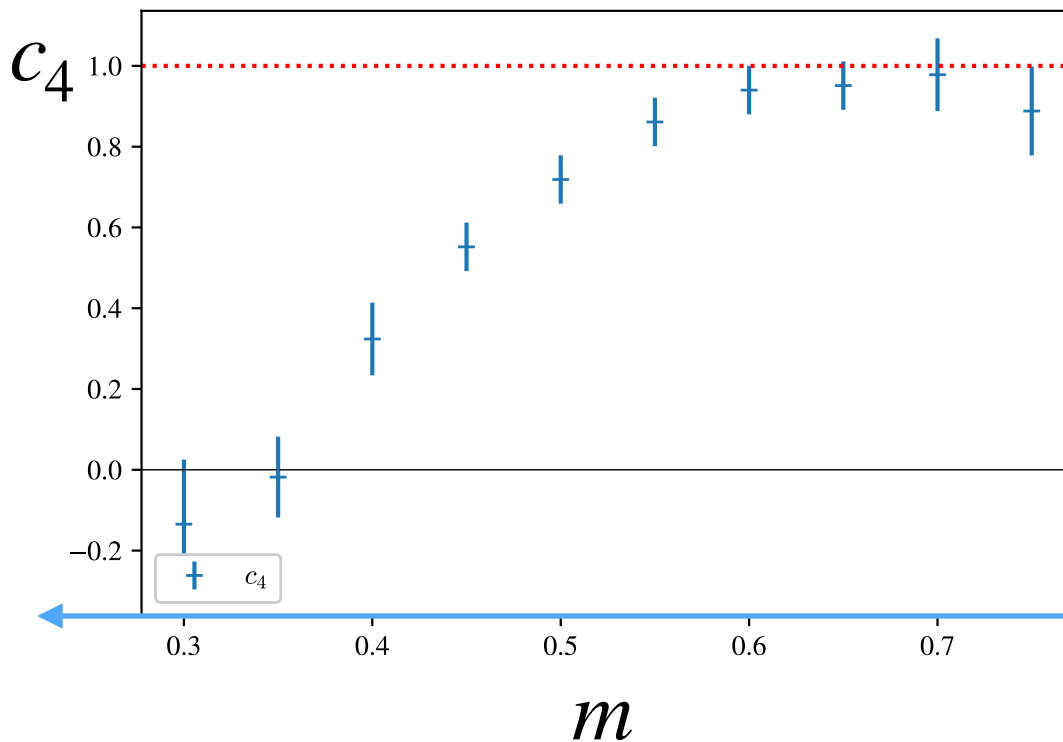
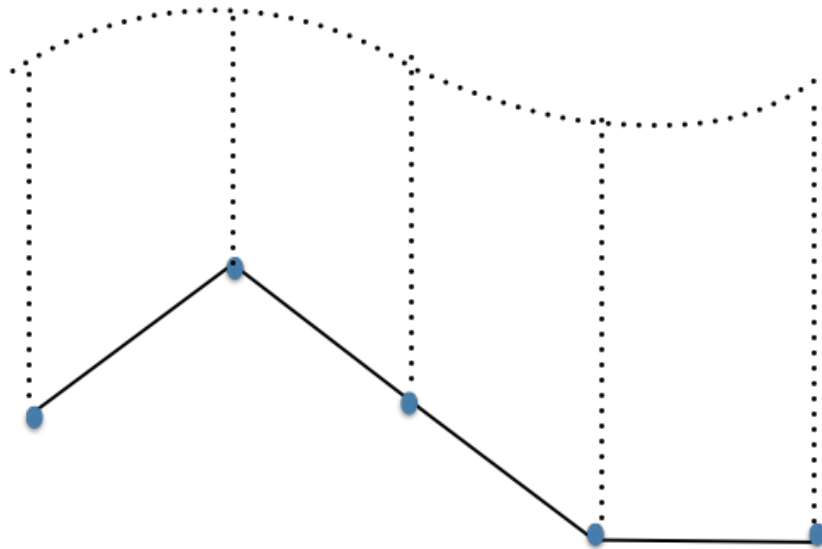
Measuring the elastic constants for harmonic depinning (qEW)

$$\eta \partial_t u(x, t) = c \nabla^2 u(x, t) - m^2 [u(x, t) - w(x, t)] + F(x, u(x, t))$$

$$w(x) = w_0 + A \sin\left(\frac{\pi x}{L}\right)$$

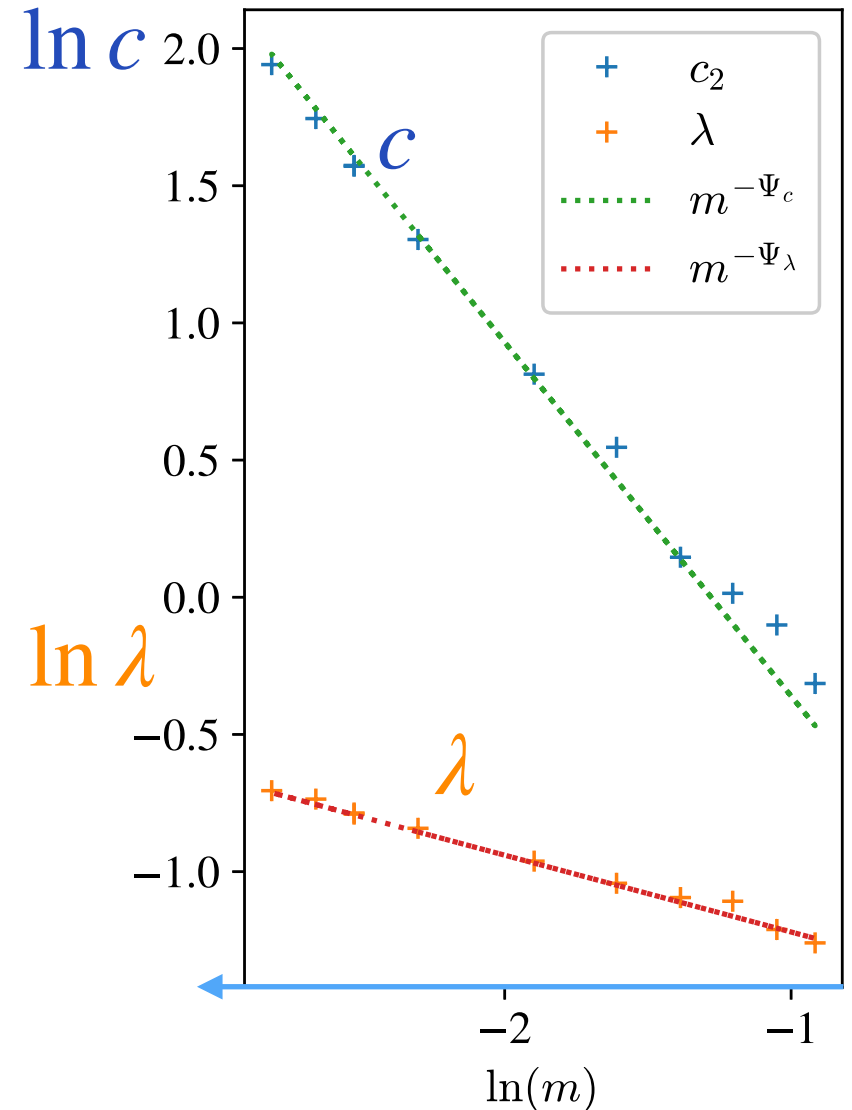


$$w(x) = w_0 + A \sin\left(\frac{\pi x}{L}\right)$$



Measuring the elastic constants

anharmonic depinning ($c_4 > 0$)



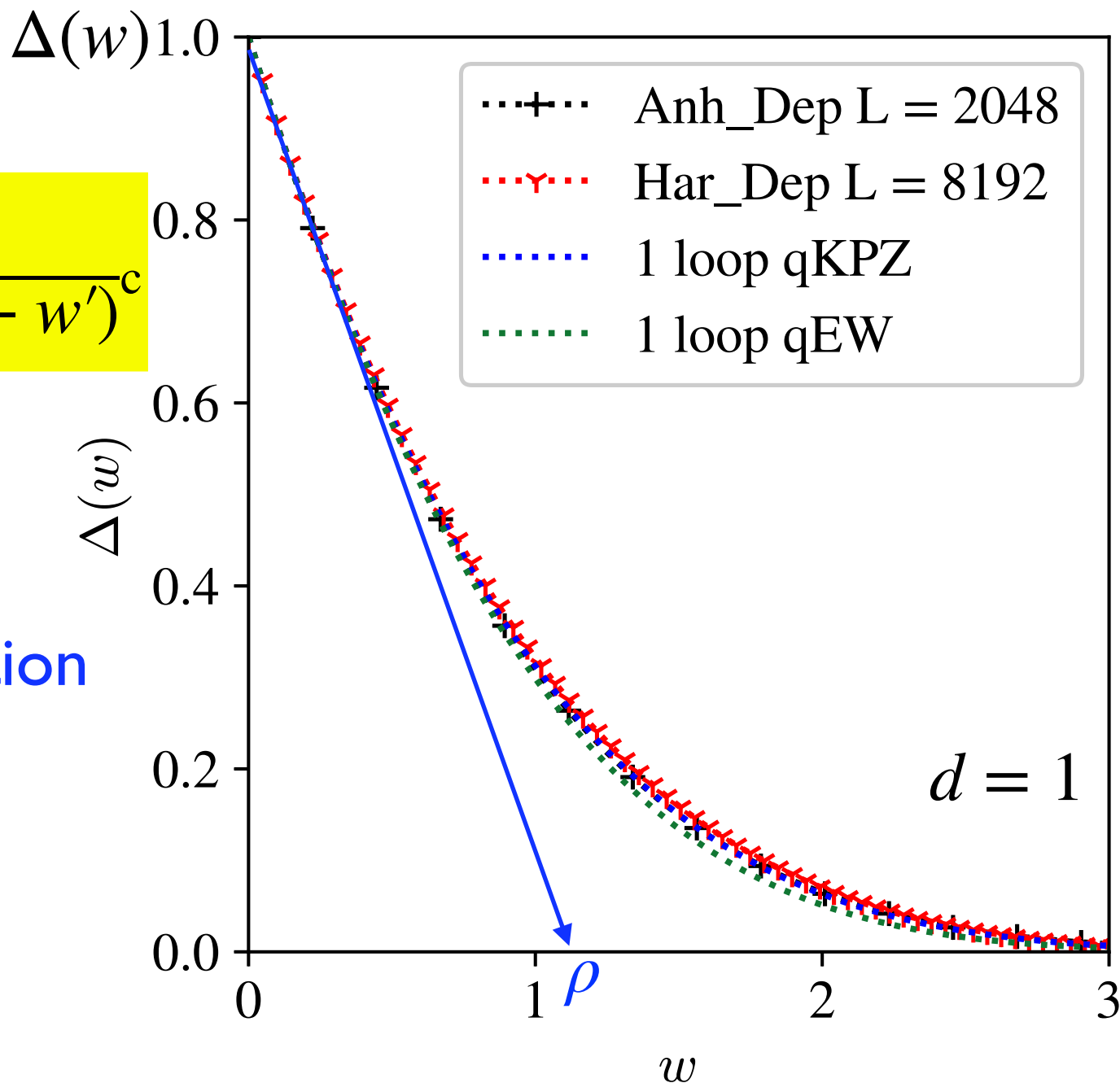
Measuring the effective force correlator

$$\Delta(w - w')$$

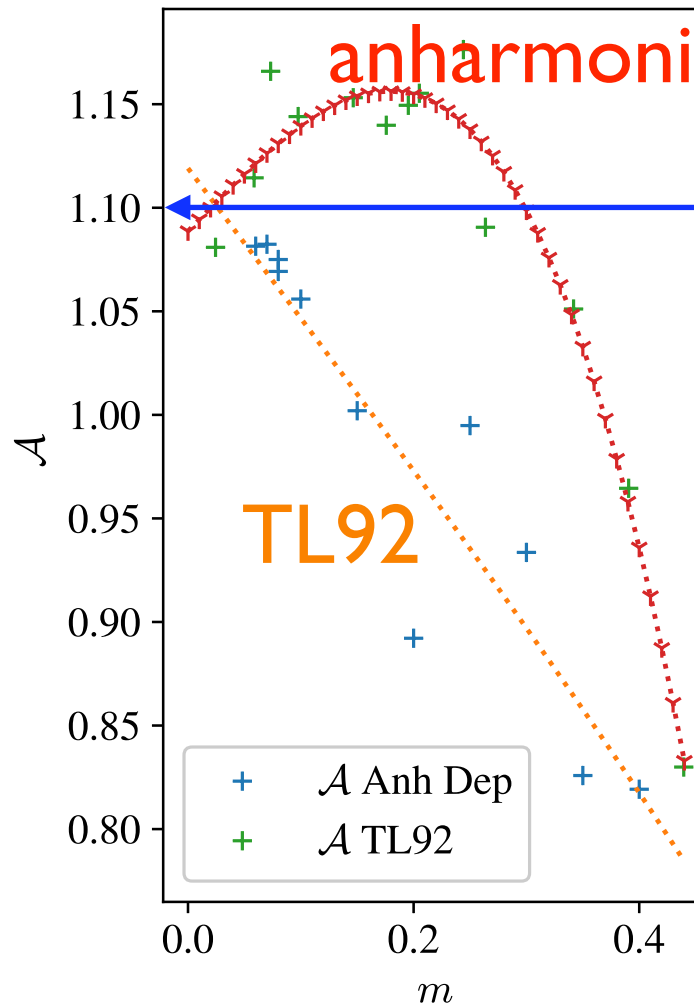
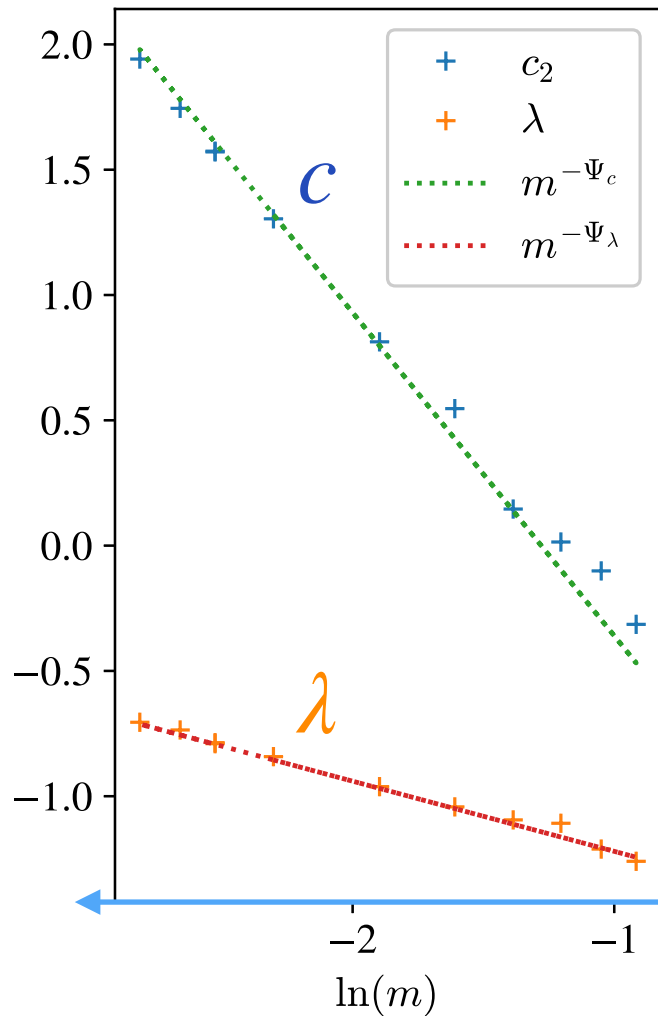
$$= m^4 L^d \overline{(u_w - w)(u_{w'} - w')}^c$$

$$u_w = \frac{1}{L^d} \int_x u_w(x)$$

↑
centre-of-mass position
given w



Coupling constant for qKPZ



anharmonic depinning

$$\mathcal{A}_{d=1} = 1.10(2)$$

scale-free universal
KPZ amplitude

$$\mathcal{A} := \rho \frac{\lambda}{c} \equiv \frac{\Delta(0)}{|\Delta'(0^+)|} \frac{\lambda}{c}$$

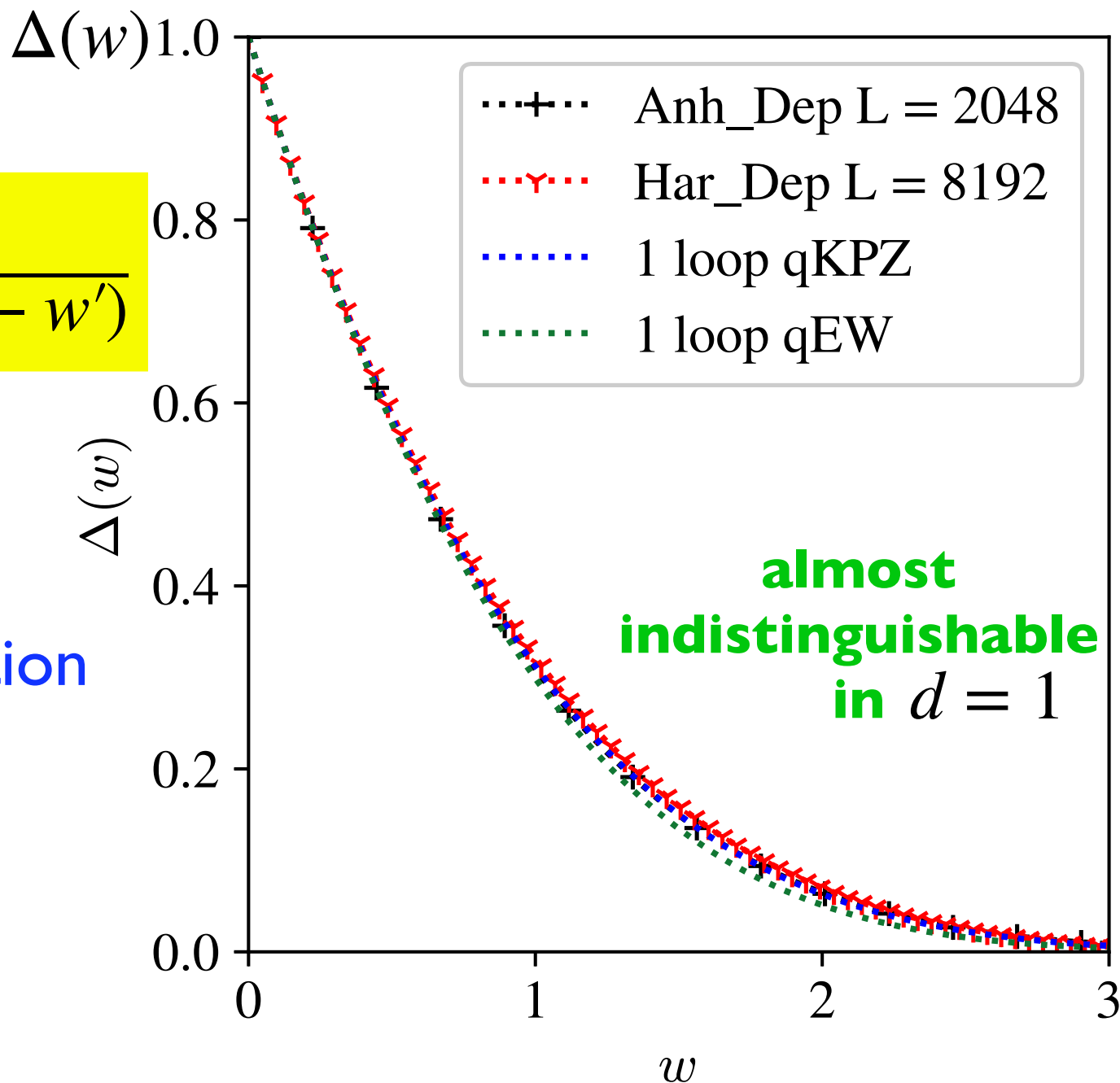
Measuring the effective force correlator

$$\Delta(w - w')$$

$$= m^4 L^d (u_w - w)(u_{w'} - w')$$

$$u_w = \frac{1}{L^d} \int_x u_w(x)$$

↑
centre-of-mass position
given w



Universality classes for depinning

qKPZ
SR-elasticity

$$d = 4$$



$$d = 3$$



$$d = 2$$

magnetic domain wall



$$d = 1$$

imbibition



qEW
SR-elasticity

$$d = 4$$



$$d = 3$$

vortex lattice/CDW



$$d = 2$$

magnetic domain wall



$$d = 1$$

magnetic domain wall



qEW
LR-elasticity

$$d = 2$$

**magnetic domain walls,
earthquakes, knitting**



$$d = 1$$

**contact line,
fracture**



$$d = 0$$

analytically solvable: dragged particle (RNA/DNA peeling)

FRG flow equations

Flow of the disorder for qKPZ

shooting parameter



$$\partial_\ell \tilde{\Delta}(u) = \left(4 - d \frac{\zeta_m}{\zeta} - 2\zeta_m \right) \tilde{\Delta}(u) + u\zeta_m \tilde{\Delta}'(u) + \frac{d(d+2)}{12} \tilde{\lambda}^2 \tilde{\Delta}(u)^2 - \tilde{\Delta}'(u)^2 - \tilde{\Delta}''(u) [\tilde{\Delta}(u) - \tilde{\Delta}(0)]$$

replace ζ_m/ζ

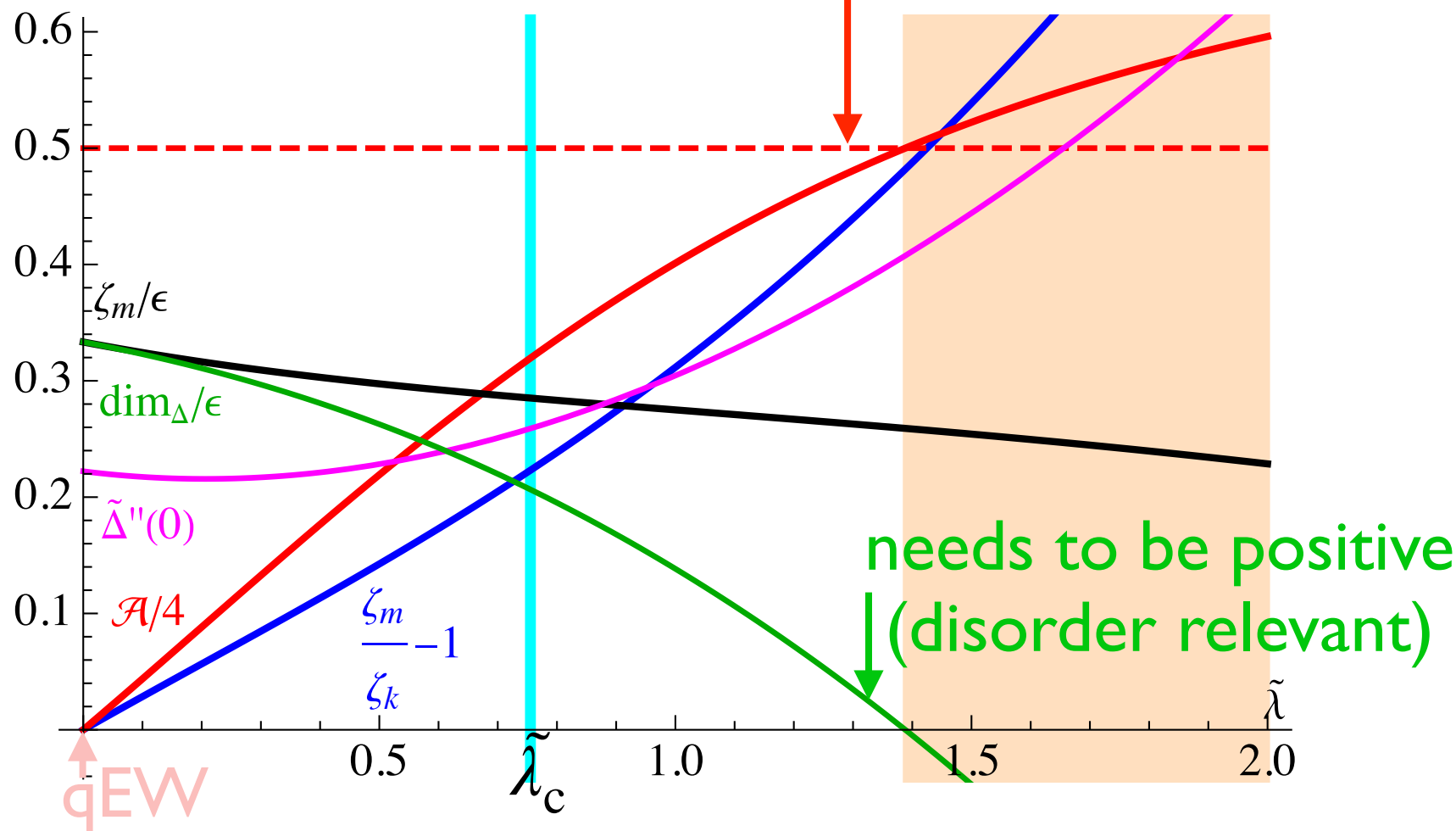
$$\frac{\zeta_m}{\zeta} = 1 + \frac{1}{2} \left[-\tilde{\lambda} \tilde{\Delta}'(0^+) - \frac{d-1}{3} \tilde{\lambda}^2 \tilde{\Delta}(0) \right].$$

flow for $\tilde{\lambda}$ (with confining potential, i.e. massive theory)

$$-m \partial_m \tilde{\lambda} = \zeta_m \tilde{\lambda} - \frac{4-d}{6} \tilde{\lambda}^3 \tilde{\Delta}(0) \implies \tilde{\lambda}_c = \sqrt{\frac{6\zeta_m}{(4-d)\tilde{\Delta}(0)}}$$

Solution in $d = 1$

$\mathcal{A} < 2$ (critical force positive)



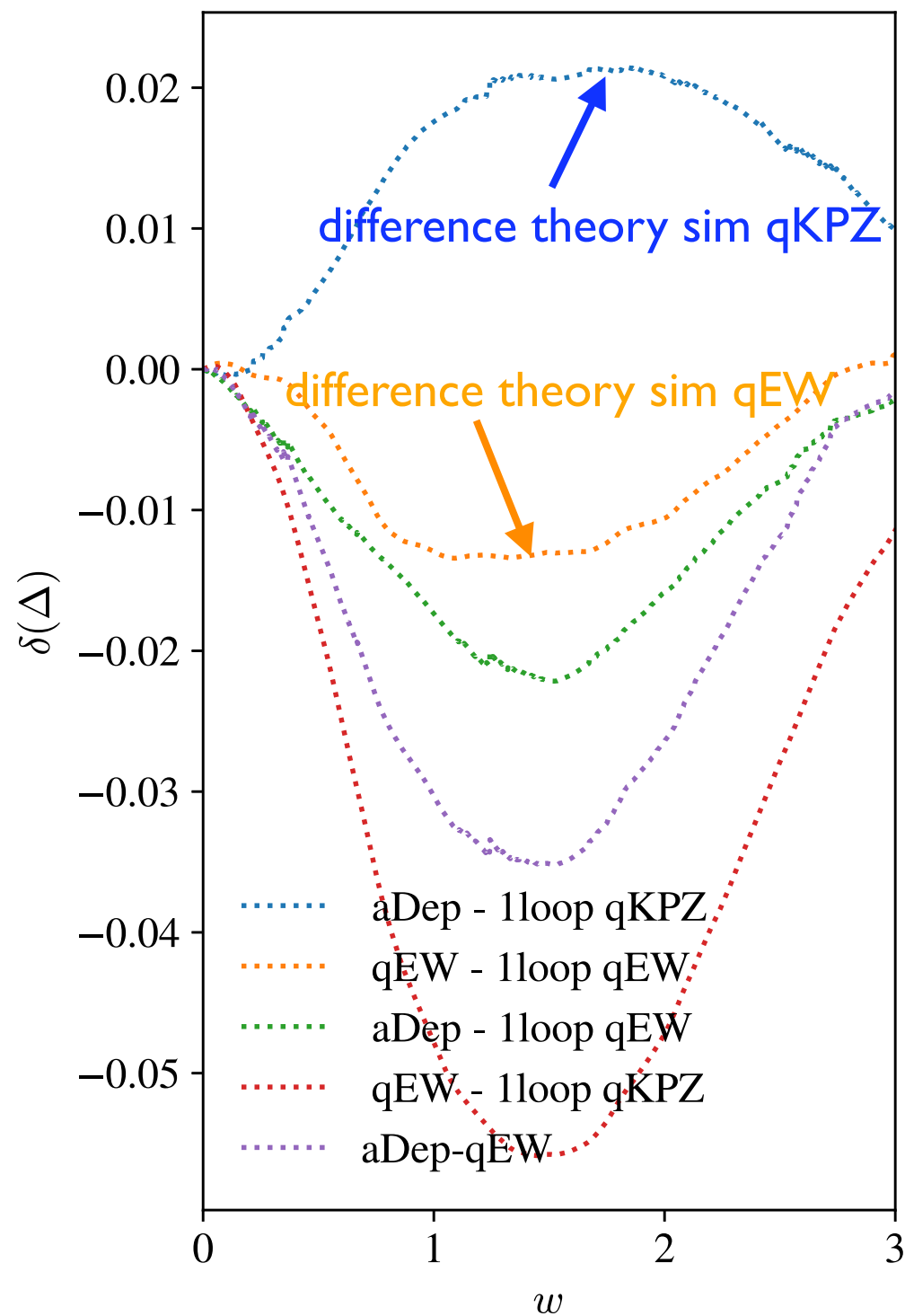
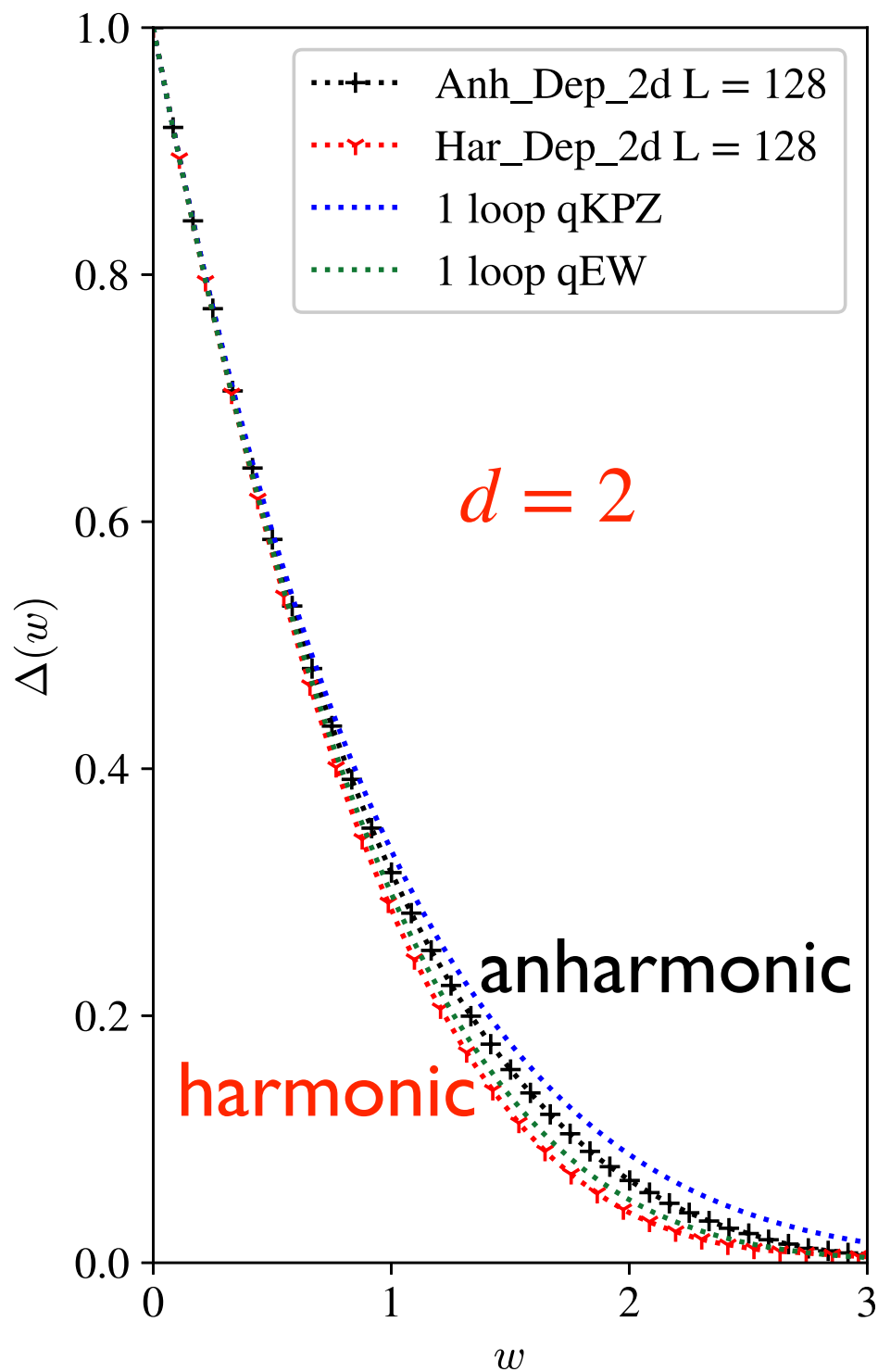
RG:

$$\begin{aligned}\zeta_m^{d=1} &= 0.86 \\ \zeta^{d=1} &= 0.69 \\ z^{d=1} &= 1.27 \\ \mathcal{A}^{d=1} &= 1.27\end{aligned}$$

numerics:

$$\begin{aligned}\zeta_m^{d=1} &= 1.05 \\ \zeta^{d=1} &= 0.63 \\ z^{d=1} &= 1.10(2) \\ \mathcal{A}^{d=1} &= 1.10(2)\end{aligned}$$

Shape of $\Delta(w)$ different in $d = 2$



Conclusions

- when in doubt: measure effective long-distance action (= theory/description)
- standard elastic depinning (**qEW**) has non-trivial disorder correlator given by FRG
- imbibition (e.g. TL92), anharmonic depinning and qKPZ all belong to the same universality class: the effective long-wavelength theory is **qKPZ**
- you need to introduce a confining potential $m^2[w - u(x, t)]$ to measure disorder correlations
 - ⇒ give up the Cole-Hopf transform
 - ⇒ yields an RG fixed point
- a field theory can be build