

Dissipation in the formalism of time-dependent Green's functions at finite temperature

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«Models in Quantum Field Theory»

The known methods

The main questions:

- How to describe dissipation in quantum systems?
- How to calculate the kinetic coefficients from first principles?

The known methods:

- Non-Hamiltonian dissipative formalism (Lindblad equation, ...).
- Projection operators, $\hat{\rho} = \hat{\rho}_{relevant} + \hat{\rho}_{irrelevant}$.
- Linear response theory (Green – Kubo formulas, ...).
- Phenomenological models (stochastic differential equations),

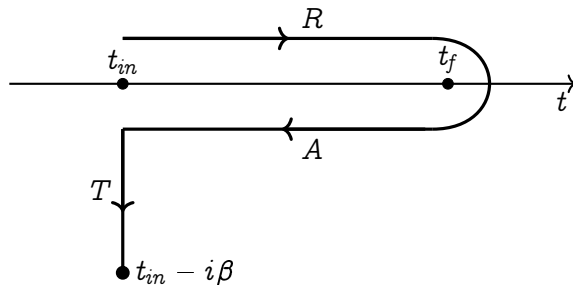
$$\partial_t \varphi_a = (\alpha_{ab} + \beta_{ab}) \frac{\delta S^{st}}{\delta \varphi_b} + \zeta_a. \quad (1)$$

- Motivation: to provide the method of calculation of kinetic coefficients and probe it on a simple model.
- Model systems: fermions or bosons with weak local interactions,

$$\hat{H} = \int d^3\mathbf{x} \hat{\psi}^\dagger(\mathbf{x}, t) \left(-\frac{\Delta}{2m} - \mu \right) \hat{\psi}(\mathbf{x}, t) + g \int d^3\mathbf{x} \hat{\psi}^\dagger(\mathbf{x}, t) \hat{\psi}(\mathbf{x}, t) \hat{\psi}^\dagger(\mathbf{x}, t) \hat{\psi}(\mathbf{x}, t). \quad (2)$$

- We want to describe the dynamics of equilibrium fluctuations on large scales and calculate the kinetic coefficient.

The main object



The 2-point Green's function:

$$\langle \hat{\psi}(x_1, t_1) \hat{\psi}^+(x_2, t_2) \rangle = \text{Sp}(\text{T}\{\hat{\psi}(x_1, t_1) \hat{\psi}^+(x_2, t_2)\} \hat{\rho}). \quad (3)$$

Action:

$$\mathcal{S} = \int_C dt \left[\psi^+(x, t) \left(i\partial_t + \frac{\Delta}{2m} + \mu \right) \psi(x, t) - g \psi^+(x, t) \psi(x, t) \psi^+(x, t) \psi(x, t) \right]. \quad (4)$$

Keldysh variables

$$\xi = \frac{\psi_R + \psi_A}{\sqrt{2}}, \quad \eta = \frac{\psi_R - \psi_A}{\sqrt{2}}, \quad \xi^+ = \frac{\psi_R^+ + \psi_A^+}{\sqrt{2}}, \quad \eta^+ = \frac{\psi_R^+ - \psi_A^+}{\sqrt{2}}. \quad (5)$$

Propagators in Keldysh variables:

$$G_{\eta\xi^+} = -e^{-i\varepsilon(t-t')}\theta(t' - t), \quad (6)$$

$$G_{\xi\eta^+} = e^{-i\varepsilon(t-t')}\theta(t - t'), \quad (7)$$

$$G_{\xi\xi^+} = e^{-i\varepsilon(t-t')}(1 \pm 2n(\varepsilon)), \quad (8)$$

$$G_{\eta\eta^+} = 0. \quad (9)$$

Notations:

$$G_{\eta\xi^+} = \text{---}\blacktriangleleft\times, \quad G_{\xi\eta^+} = \text{---}\blacktriangleright\times, \quad G_{\xi\xi^+} = \text{---}\times. \quad (10)$$

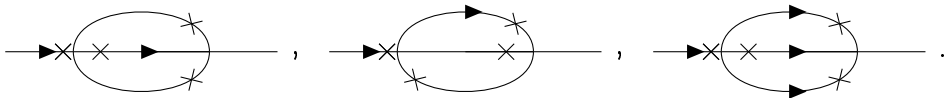
Interactions:

$$S_{int} = g(\xi^+\xi\xi^+\eta + \xi^+\xi\eta^+\xi + \xi^+\eta\eta^+\eta + \eta^+\xi\eta^+\eta). \quad (11)$$

«Pinch» singularities

$$\begin{aligned}
 & \left[\text{Diagram: a horizontal line with an oval loop containing two vertices marked with 'x' and arrows pointing right} \right] \Big|_{p=0, \omega=0} = \int_{-\infty}^{\infty} d(t-t') \left[\text{Diagram: same as above} \right] \Big|_{p=0} = \quad (12) \\
 & = \frac{\mp \zeta g^2}{(2\pi)^6} \int d^3k d^3q \left(\pi \delta(\varepsilon(k+q) - \varepsilon(k) - \varepsilon(q)) + \underbrace{\frac{i}{\varepsilon(k+q) - \varepsilon(k) - \varepsilon(q)}}_{\text{Singular}} \right) \times \\
 & \times (1 \pm 2n(k))(1 \pm 2n(q)).
 \end{aligned}$$

Singular diagrams:



Regularization:

$$e^{-i\varepsilon(t-t')} \rightarrow e^{-i\varepsilon(t-t') - \gamma|t-t'|}, \quad \gamma = \alpha \frac{p^2}{2m}. \quad (13)$$

The origin of dissipation

Dyson equation:

$$D^{-1} = K - \Sigma \Rightarrow D_{\eta+\xi}^{-1} = i\omega + i\varepsilon + \underline{\gamma - \Sigma_{\eta+\xi}}. \quad (14)$$

The sum of sunset diagrams in hydrodynamic limit:

$$I(\mathbf{p}, \omega) \approx I(\mathbf{p}, \omega) \Big|_{\substack{\mathbf{p}=0 \\ \omega=0}} + \omega \frac{\partial I(\mathbf{p}=0, \omega)}{\partial \omega} \Big|_{\omega=0} + \frac{p_i p_j}{2!} \frac{\partial^2 I(\mathbf{p}, \omega=0)}{\partial p_i \partial p_j} \Big|_{\mathbf{p}=0}. \quad (15)$$

- Propagators and one loop:

$$\langle \hat{\psi}(t_1) \hat{\psi}^+(t_2) \rangle \sim \exp \left(-i \left(\frac{p^2}{2m} - \mu \right) (t_1 - t_2) \right). \quad (16)$$

- Two loops:

$$\langle \hat{\psi}(t_1) \hat{\psi}^+(t_2) \rangle \sim \exp \left(-i \left(\frac{p^2}{2m} - \mu \right) (t_1 - t_2) - \underbrace{\tilde{\alpha} \frac{p^2}{2m} |t_1 - t_2|}_{\text{dissipation}} \right). \quad (17)$$

- By analogy with phenomenological models, $\tilde{\alpha}$ is the kinetic coefficient.

Asymptotic of some integrals

The contributions to $\tilde{\alpha}$ are given by:

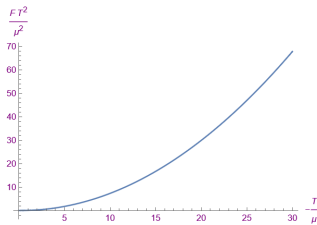
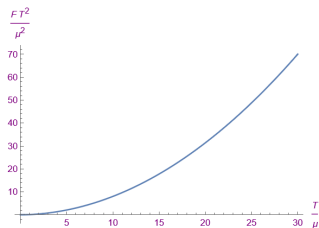
$$\lim_{\alpha \rightarrow 0} \frac{\mp 2\zeta g^2 m^2}{(2\pi)^6} \cdot \text{Re} \int d^3k d^3q \frac{f(k, q)}{[i(k \cdot q) - \alpha(k^2 + k \cdot q + q^2)]^2}, \quad (18)$$

$$\begin{aligned} \lim_{\alpha \rightarrow 0} \frac{\mp 2\zeta g^2 m}{(2\pi)^6} \cdot \text{Re} \int d^3k d^3q & \left(- \frac{if(k, q)}{[i(k \cdot q) - \alpha(k^2 + k \cdot q + q^2)]^2} + \right. \\ & + \frac{2}{3} \frac{f(k, q) \cdot (k + q)^2}{[i(k \cdot q) - \alpha(k^2 + k \cdot q + q^2)]^3} + \frac{f_1(k, q)}{i(k \cdot q) - \alpha(k^2 + k \cdot q + q^2)} + \\ & \left. + \frac{2}{3} \frac{if_2(k, q) \cdot (k + q)^2}{[i(k \cdot q) - \alpha(k^2 + k \cdot q + q^2)]^2} \right), \end{aligned} \quad (19)$$

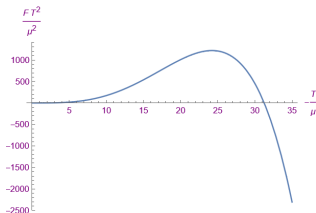
$$f(p, k, q) = 4n(k)n(q) \mp 4n(p - k - q) - 8n(k)n(p - k - q). \quad (20)$$

After calculating the integrals

$$\tilde{\alpha} = g^2 m^2 T^2 F(T/\mu).$$



Fermions with positive (left) or negative (right) chemical potentials.



Bosons with negative chemical potentials.

- Existing methods for calculating the kinetic coefficients are difficult to carry out real calculations.
- Based on the formalism of the time-dependent Green's time functions at finite temperature the new method for calculating the kinetic coefficients in quantum systems has been developed.
- To calculate the kinetic coefficient, it is necessary to carry out a rather complicated asymptotic analysis of a certain class of integrals.