# Energy scan and beam polarization as a tools to study heavy quarkonium production mechanism

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#### Outline

- Short theoretical introduction: NRQCD-factorization and Color-Evaporation model
- Status of NRQCD+LO PRA predictions for low energies (COMPASS, NICA), NA-3 problem
- CEM vs. NRQCD-factorization. Beam polarization as a tool to distinguish production models

### A bit of history

Historically, the first model of heavy-quarkonium production was the **color-singlet model**: The production of state  $X_Q$  $(J/\psi, \chi_{cJ}, ..., \Upsilon(nS), \chi_{bJ}, ...)$  is dominated by production of **color-singlet**  $Q\bar{Q}$ -pair with L and S quantum numbers given by NR potential model for this state. Probability of hadronization is proportional to  $|R^{(k)}(0)|^2$ , (k = 0, 1, ...) from potential model. This model has two problems:

- Leads to a wrong shape of  $p_T$ -spectrum at high energies (Tevatron, LHC) both at LO and NLO of CPM and in  $k_T$ -factorization, which **under-estimates** the cross-section for  $p_T > 10$  GeV by factor of 30 (*Tevatron*  $\psi(2S)$  puzzle).
- ► Is theoretically inconsistent at NLO for production of *P*-wave states: In QCD, non-cancelling IR-divergences arise at NLO.



#### NRQCD and Color-Evaporation model

To solve above-mentioned problems, two approaches have been proposed: **NRQCD-factorization** and **Color-Evaporation Model**.

▶ NRQCD-factorization: Different L, S and color states of  $Q\bar{Q}$ -pair hadronize to X with different "probability" – long-distance matrix element (LDME):  $\langle \mathcal{O}^X \begin{bmatrix} 2S+1L_J^{(\text{color})} \end{bmatrix} \rangle.$ 

► LDME-s of states different from CSM-state are suppressed by powers of  $v^2$  (~ 0.3 for  $J/\psi$ , ~ 0.1 for  $\Upsilon$ ) – *velocity-scaling rules* for LDMEs. E.g. for  $J/\psi$  and  $\psi(2S)$ : CSM= ${}^{3}S_{1}^{(1)} = O(1)$  and  ${}^{3}P_{J}^{(8)} = O(v^2)$  and  ${}^{3}S_{1}^{(8)}$ ,  ${}^{1}S_{0}^{(8)}$ , contribute at  $O(v^4)$ .

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- ▶ In Improved-Color-Evaporation Model: all  $Q\bar{Q}$  states with  $M_X < M_{Q\bar{Q}} < 2M_{(\text{open flav. Q-meson})}$  hadronize to quarkonium X with the same probability  $F_X$
- ▶ Optionally [Ma, Vogt, 2016] ICEM takes into account kinematic (soft-gluon recoil) corrections from the difference of masses  $M_{Q\bar{Q}}$  and  $M_X$  using simple relation  $p_T(X) = p_T(Q\bar{Q}) \times M_X/M_{Q\bar{Q}}$ .
- ▶ ICEM can be viewed as NRQCD-factorization without velocity-scaling rules for probabilities  $F_X$ .

Both models are well-defined to all orders in  $\alpha_s$ , but

NRQCD-factorization is viewed as more "rigorous" approach by the community.

Both models are rather successful in description of  $p_T$ -spectra at high energies. The polarization issue is complicated and requires separate talk...

#### Situation at low energies: NA-3 problem



- ► LO PRA [M.N., V.A.S.] predictions with LDMEs well-describing high-energy data under-estimate the cross-section in pp and  $\pi p$ -collisions at  $\sqrt{S} = 19 23$  GeV.
- The fraction of  $q\bar{q}$ -initiated subprocesses  $(q\bar{q} \rightarrow^3 S_1^{(8)})$  is very low: ~ 1% of total cross-section. 6 / 16

#### Situation at low energies: COMPASS



- LO PRA [M.N., V.A.S.] and NLO CPM [Kniehl, Butenshön] predictions with NRQCD-factorization agree with each-other
- However PRA predictions will probably underestimate
  COMPASS data, since they do not agree with NA-3 data

#### NRQCD prediction is robust

Many options had been tried to improve situation with NA-3 data:

- Various LDME fit strategies: use/not-use HQSS-relations between LDMEs; allow negative color-octet LDMEs; fit color-singlet LDMEs as well.
- LO vs. "NLO" formulas for KMRW UPDF, different input PDF sets (MSTW, CT18, GRV LO/NLO for pion).
- On/off kinematic shift in cascade decays  $X_1 \to X_2$ :  $p_{T2} = p_{T1} \times M_{X_2}/M_{X_1}.$
- ▶ fixed quark mass prescription in the short-distance part vs.  $m_Q = M_X/2$  prescription.

In all cases LDMEs change in such a way that prediction for prompt- $J/\psi$  spectrum stays the essentially same.

#### Example: $m_Q$ -dependence Fit of high-energy data (ATLAS, CMS, CDF, PHENIX):





LO. mO=1.5 GeV:

oJ#1S08=0±0.00153448

oJ#3P08=0.00714443±0.000206544

o#2\$1\$08=0±0.00136388

oxc23S18=0±0.000497784

LO, mO=M/2: oJ#3S18=0.00170806±0.0000185009 oJ#3S18=0.00157612±0.0000169744 oJ#1S08=0±0.00118435 oJ#3P08=0.0121518±0.000245038 #2\$3\$18=0.000927456±0.00002685 #2S3S18=0.00109943±0.000049092 o#2S1S08=0±0.0009566 #2S3P08=0.000371047±0.00016073 od/2S3P08=0.0100972±0.000434641 yc13S18=0.000595678±0.000078324 oyc13S18=0.0021449±0.000122563 yc23S18=0.000843343±0.00019379

#### $\chi^2/d.o.f.=1.44934$

NT. 20 GeV. Librard. L.B. (199-rough-textpTdc/darpell) tab/C

 $\chi^2/d.o.f.=2.65076$ 

#### Example: $m_Q$ -dependence

Predictions of both fits for NA-3 data are essentially the same:



#### Example: $m_Q$ -dependence

What changes is the fraction of prompt  $J/\psi$  produced directly:



In  $m_Q = M_X/2$ -model this fraction is energy-dependent:



while in fixed  $m_Q$ -model it is approximately constant ( $\simeq 60\%$ ).

#### Suggestion I: Energy scan as a tool

- ► NRQCD-factorization predictions in LO of PRA and NLO of CPM agree well. But LO PRA under-estimates NA-3 data. Large NNLO mcorrections? Threshold logarithms (ln(1 − z))?
- It is important to study observables beyond prompt  $J/\psi$  to discriminate production models.
- ▶ Fixed- $m_Q$  model = "factorization picture" between  $Q\bar{Q}$ -production and hadronization, while  $m_Q = M/2$ -model ⇒ more complicated interplay between hard subprocess and hadronization. Which of them is correct?
- ► Two models predict different  $\sigma_{(\text{direct}-J/\psi)}/\sigma_{(J/\psi)}$  or  $\sigma_{(J/\psi \text{ from decays})}/\sigma_{(J/\psi)}$ . But these observables are experimentally challenging.

▶ Energy-dependence of the ratios:

$$\frac{d\sigma_{\psi(2S)}/dp_T}{d\sigma_{J/\psi}/dp_T}, \quad \frac{d\sigma_{\chi_{cJ}}/dp_T}{d\sigma_{J/\psi}/dp_T}$$

probes the same physics.

#### ICEM calculation

Results of [R.Vogt, et al.] where reproduced with our UPDF (KMRW UPDF derived from MSTW-08 LO PDF set instead of CCFM UPDF by [Jung, Hauptmann, 2016]). The same  $F_X$ -probabilities as in the paper where used:



#### ICEM calculation

## Calculation for COMPASS kinematics basically coincides with PRA-NRQCD result:



So there could be no sharp distinction between NRQCD and ICEM predictions for un-polarized cross-section.

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#### Suggestion II: Polarized beam as a tool

**Beam polarization** can be used as a handle to distinguish two models.

Idea is to measure the polarization-asymmetry for the  $p_T$ -spectra of  $J/\psi$  and  $D\bar{D}$ -pairs:

$$A_{J/\psi \text{ or } D\bar{D}}(p_T) = \frac{d\sigma_+/dp_T(p_T) - d\sigma_-/dp_T(p_T)}{d\sigma_+/dp_T(p_T) + d\sigma_-/dp_T(p_T)}$$

Prediction:

► In NRQCD the weights of various  $Q\bar{Q}$  states is very different  $\Rightarrow$  $A_{J/\psi}^{\text{NRQCD}}(p_T) \neq A_{D\bar{D}}(p_T)$ .

▶ In CEM all L, S states contribute equally  $\Rightarrow$  $A_{J/\psi}^{\text{CEM}}(p_T) \simeq A_{D\bar{D}}(p_T)$ .

- ▶ Transverse spin asymmetry is useful as well.
- ▶ Unfortunately, radiative corrections will dilute the difference and make  $A_{J/\psi}(p_T) \rightarrow A_{D\bar{D}}(p_T)$  ("decoherence"), the effect of this should be estimated.

#### Conclusions

- ▶ Heavy quarkonium production at low enegries can become (yet another) challenge for NRQCD-factorization
- ▶ LO PRA and NLO CPM calculations agree but probably will under-estimate the data
- It is important to study not only  $J/\psi$  but also  $\chi_{cJ}$  and  $\psi(2S)$ -cross-sections as function of energy to discriminate various models
- ▶ Beam-polarization asymmetry for  $J/\psi$  and  $D\bar{D}$ -production can be used to discriminate between NRQCD and ICEM. Calculations are in progress...

Thank you for your attention!