

# On the Wigner quasiprobability function for N-level quantum system

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# Overview

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# The quasi-probability distributions

Quasi-Probability Distributions (QPD)'s – “quantum analogue” of the statistical distribution on the phase space of a classical system

We are interested in QPD's for finite dimensional quantum systems

THE AIM OF TALK:

To build the  $s$ -parametric family of Wigner functions  $W_\varrho(\Omega; s)$  over the classical phase space ( $\Omega$ ) for a generic  $N$ -level quantum system with density matrix  $\varrho$ .

# The standard form of the Wigner function

For a given state, described by the density operator  $\varrho$ , the Wigner function  $W(\mathbf{q}, \mathbf{p})$  defined on a classical  $2n$ -dimensional phase space spanned by the canonical coordinates  $\mathbf{q}$  and momentum  $\mathbf{p}$  reads

$$W(\mathbf{q}, \mathbf{p}) := \int d^n z e^{\frac{i}{\hbar} \mathbf{z} \cdot \mathbf{p}} \langle \mathbf{q} + \frac{\mathbf{z}}{2} | \varrho | \mathbf{q} - \frac{\mathbf{z}}{2} \rangle,$$

Representation for  $W$  via displacement operator  $D$

$$W = \text{Tr} [\varrho D \Pi D^\dagger]$$

where  $D$  and  $\Pi$  are the displacement and parity operators respectively.

# Quasiprobability distributions

- For any operator  $A$  on the Hilbert space  $\mathcal{H}$  of quantum system one can define a family of functions  $F_A(\Omega; s)$  onto the phase space  $\Omega$ . The index  $s$  labels functions in the family.
- When the operator  $A$  represents the density matrix,  $A = \varrho$ , the corresponding phase-space functions  $F_\varrho(\Omega; s) := P(\Omega; s)$  are named as **Quasiprobability Distributions**.

## The Stratonovich-Weyl correspondence

The physically motivated properties of  $P(\Omega; s)$  were formulated by R.L.Stratonovich more than sixty years ago (1955) and are usually referred to as the **Stratonovich-Weyl correspondence**

# The Stratonovich-Weyl Correspondence

Clauses of SW correspondence:

- **Mapping** • For a density matrix  $\varrho$  the Wigner function  $W_\rho$  on the classical phase-space ( $\Omega$ ) is given by the map:

$$W_\rho(\Omega) = \text{tr}(\varrho \Delta(\Omega))$$

defined by the Hermitian kernel  $\Delta(\Omega) = \Delta(\Omega)^\dagger$ , with a unit norm

$$\int_{\Omega} d\Omega \Delta(\Omega) = 1$$

- **Reconstruction** • The state  $\rho$  can be reconstructed as

$$\rho = \int_{\Omega} d\Omega \Delta(\Omega) W_\rho(\Omega).$$

- **Covariance** • The unitary symmetry  $\varrho' = U(\alpha)\varrho U^\dagger(\alpha)$  induces the kernel transformation:

$$\Delta(\Omega') = U(\alpha)\Delta(\Omega)U^\dagger(\alpha)$$

# The Wigner distribution kernel

The Wigner distribution  $W_\rho(\Omega)$  over a phase space parametrized by the set  $\Omega$  is determined by the kernel  $\Delta(\Omega)$ :

$$W_\rho(\Omega) = \text{tr}(\rho \Delta(\Omega))$$

Since the kernel  $\Delta(\Omega)$  is Hermitian it can be written as

$$\Delta(\theta_1, \dots, \theta_{d_F} | \pi_1, \dots, \pi_N) = U(\theta) P U^\dagger(\theta).$$

Here  $P = \kappa N \Pi$ , with  $\kappa$  - real normalization factor, a diagonal matrix  $\Pi = \text{diag}||\pi_1, \pi_2, \dots, \pi_N||$  with a unit trace,  $\text{tr}(\Pi) = 1$ .

In accordance with the  $SU(n)$ -covariance of kernel we identify:

$d_F$  - parameters of unitary matrix  $U(\theta) \in SU(N)$  with the coordinates of classical phase-space,  $\Omega = (\theta_1, \dots, \theta_{d_F})$ .

# Constructing the kernel $\Delta(\Omega)$

- **Step 1** • The  $SU(N)$  symmetry allows to define the “reconstruction” integral for  $\varrho$  over the  $SU(N)$  group with the Haar measure:

$$\varrho = \int_{SU(N)} d\mu_{SU(N)} \Delta(\Omega) W_\rho(\Omega).$$

- **Step 2** • Substitute decomposition  $\Delta = U(\theta) P U^\dagger(\theta)$  into the identity

$$\varrho = \int_{SU(N)} d\mu_{SU(N)} \Delta(\Omega) \text{tr} [\Delta(\Omega) \varrho].$$

and evaluate the integral using the Weingarten formula:

$$\int d\mu U_{i_1 j_1} U_{i_2 j_2} U_{k_1 l_1}^\dagger U_{k_2 l_2}^\dagger = \frac{1}{N^2 - 1} (\delta_{i_1 k_1} \delta_{i_2 k_2} \delta_{j_1 l_1} \delta_{j_2 l_2} + \delta_{i_1 k_2} \delta_{i_2 k_1} \delta_{j_1 l_2} \delta_{j_2 l_1})$$

$$-\frac{1}{N(N^2 - 1)} (\delta_{i_1 k_1} \delta_{i_2 k_2} \delta_{j_1 l_2} \delta_{j_2 l_1} + \delta_{i_1 k_2} \delta_{i_2 k_1} \delta_{j_1 l_1} \delta_{j_2 l_2}).$$

# Constructing the kernel

- **Step 3** • Derive the algebraic equation for kernel

The equation for  $n \times n$  matrix  $P$  reads:

$$P_{ki}P_{sj} - \frac{1}{N}P_{si}P_{kj} = \frac{N^2 - 1}{2}\delta_{is}\delta_{jk}.$$

- **Step 4** • Solve these equations classifying its solution by the values of adjoint  $SU(n)$ -invariants constructed from  $P$ ,

Solutions are characterized by the Casimir invariants obeying relations:

$$\mathfrak{C}_r + \mathfrak{C}_{r+1} = 0, \quad r = 2, 3, \dots, N-1.$$

Table: Invariant characteristics of kernel  $\Delta$  for low  $N$

Group	$\mathfrak{C}_2$	$\kappa$	Spectrum $\sigma(\Pi)$
SU(3)	2	$\sqrt{6}$	$\{-\frac{1}{3}; \frac{2}{3}, \frac{2}{3}\}$
SU(4)	$\frac{3}{4}$	$2\sqrt{10}$	$\{-\frac{1}{8}; \frac{3}{8}, \frac{3}{8}, \frac{3}{8}\}$
SU(5)	$\frac{4}{9}$	$3\sqrt{15}$	$\{-\frac{1}{15}; \overbrace{\frac{4}{15}, \dots, \frac{4}{15}}^4\}$
	6	$\sqrt{10}$	$-\frac{2}{5}, -\frac{2}{5}; \overbrace{\frac{3}{5}, \dots, \frac{3}{5}}^3$
SU(6)	$\frac{5}{16}$	$4\sqrt{21}$	$\{-\frac{1}{24}; \overbrace{\frac{5}{24}, \dots, \frac{5}{24}}^5\}$
	2	$\frac{\sqrt{210}}{2}$	$-\frac{1}{6}, -\frac{1}{6}; \overbrace{\frac{1}{3}, \dots, \frac{1}{3}}^4\}$

Table II: Spectrum of kernel  $\Delta$  for given  $N$

$$N^* = \left[ \frac{N-1}{2} \right] \text{ types degeneracy of spectrum}$$

	$\mathfrak{C}_2$	$\sigma(\Pi)$
1	$\frac{N-1}{(N-2)^2}$	$\left[ \underbrace{-\frac{1}{N(N-2)}, \dots, -\frac{1}{N(N-2)}}_1 ; \underbrace{\frac{N-1}{N(N-2)}, \dots, -\frac{N-1}{N(N-2)}}_{N-1} \right]$
2	$\frac{2(N-2)}{(N-4)^2}$	$\left[ \underbrace{-\frac{2}{N(N-4)}, \dots, -\frac{2}{N(N-4)}}_2 ; \underbrace{\frac{N-2}{N(N-4)}, \dots, -\frac{N-2}{N(N-4)}}_{N-2} \right]$
$\vdots$	$\vdots$	$\vdots \vdots$
$N^*$	$\frac{N^*(N-N^*)}{(N-2N^*)^2}$	$\left[ \underbrace{\frac{-N^*}{N(N-2N^*)}, \dots, \frac{-N^*}{N(N-2N^*)}}_{N^*} ; \underbrace{\frac{N-N^*}{N(N-2N^*)}, \dots, \frac{N-N^*}{N(N-2N^*)}}_{N-N^*} \right]$

## Few examples of the Wigner distributions

- Qubit – 2-level system;
- Qubit-Qubit– binary system composed from two 2-level. subsystems

# Qubit

## Parametrization of the qubit state by Bloch Sphere

The state of a qubit may be represented as a point on/in a unit sphere of coefficients in decomposition  $\rho = \frac{1}{2} (I + \vec{r} \vec{\sigma})$ , where the points on the sphere correspond to the pure states, while the interior points represent the mixed states.

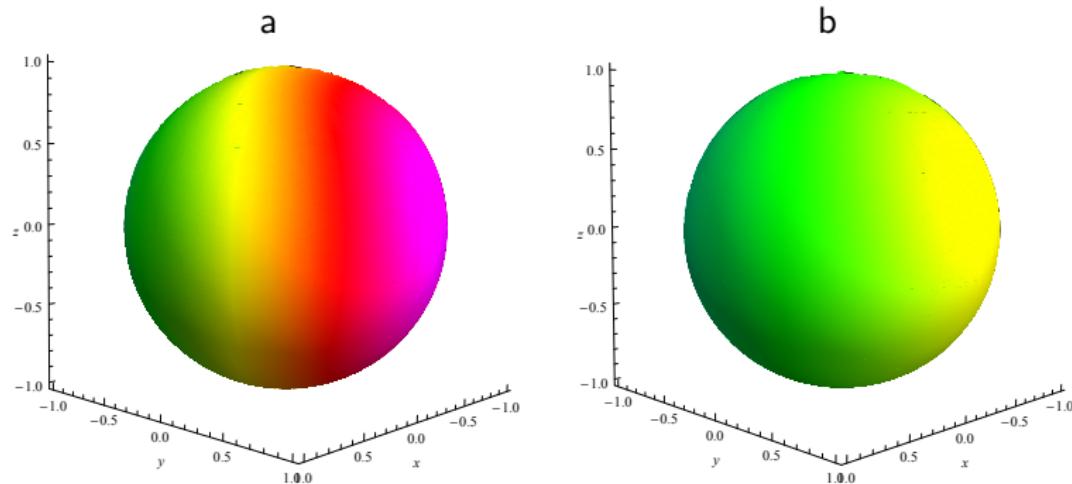
## Qubit kernel

$$\Delta(\alpha_1, \alpha_2) = U(\alpha_1, \alpha_2) \frac{1}{2} (I - \sqrt{3} \sigma_z) U^\dagger(\alpha_1, \alpha_2)$$

## Qubit Wigner function

$$W(\alpha_1, \alpha_2 ; r, \theta, \phi) = \text{tr} [\rho \Delta(\alpha_1, \alpha_2)] = \\ \frac{1}{2} \left( 1 + \sqrt{3} r (\sin(\alpha_2) \sin(\theta) \cos(\alpha_1 + \phi) - \cos(\alpha_2) \cos(\theta)) \right).$$

# Qubit



The Wigner function for the qubit on the unit sphere of "phase angles" (a) for a singlet (b) half way to the singlet along **z** axes.

# Qubit-qubit

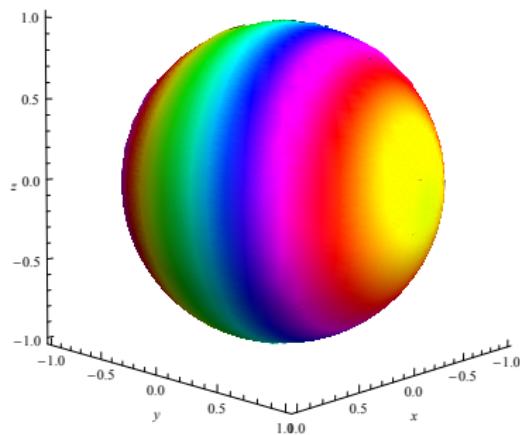
Wigner function for the qubit-qubit pair in the Werner state

For the  $U(2) \otimes U(2)$  invariant Werner state  $\rho_{\text{Werner}} = \frac{1}{4}(I + p\vec{\sigma} \cdot \vec{\sigma})$ , choosing an appropriate parametrization we get the following form for the Wigner function

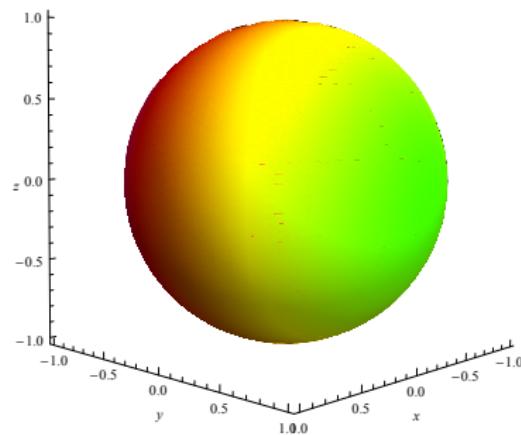
$$\begin{aligned} W(\alpha_0, \alpha_2, \alpha_4, \alpha_6; p) &= \frac{1}{2} \sqrt{\frac{5}{2}} (-p + 8 + p(\cos(\alpha_6)\cos(\alpha_4) - \cos(\alpha_4) \\ &\quad - 3\cos(\alpha_6) - 2(1 - \cos(\alpha_6))(\frac{1}{2}\cos(\alpha_2)(\cos(\alpha_4) + 1) \\ &\quad + 2\sin\left(\frac{\alpha_2}{2}\right)\sin(\alpha_4)\cos\left(\frac{\alpha_0}{2}\right))). \end{aligned}$$

# Qubit-qubit

a



b



The Wigner function for the qubit-qubit pair an Werner maximally ( $p = -1$ ) entangled an separable ( $p = -\frac{1}{4}$ ), (a) and (b) respectively. The slice  $\alpha_4 = \frac{3\pi}{4}$ ,  $\alpha_6 = \frac{5\pi}{6}$  has been taken.

# Concluding remarks

- ① The exact form of the kernel of Wigner-Weyl transformation yielding a quasi-probability distribution function for general  $N$ -level quantum system has been constructed.
- ② For an  $N$ -level quantum system a  $N^* = \lfloor \frac{N-1}{2} \rfloor$  fold family of transformation kernels fulfilling the Stratonovich-Weyl restrictions has been found.
- ③ The family members are classified by the corresponding values of the Casimir invariants.

Thank you!

# Technical details of calculations

$$\boxed{\Pi - \frac{1}{N}I = \frac{\alpha}{N}\Pi^{-1}} \quad \boxed{\text{tr}(\Pi^{-1}) = 0}$$
$$\alpha = (N^2 - 1)/2\kappa^2.$$

Expressing the traces of powers as polynomials in  $\alpha$

$$\text{tr}(\Pi^2) = \frac{1}{N} + \alpha$$

$$\text{tr}(\Pi^3) = \frac{1}{N^2} + \frac{2}{N}\alpha$$

.....

$$\text{tr}(\Pi^k) = C_{Q(k)}^{\frac{k+Q(k)}{2}} \left(\frac{1}{n}\right)^{Q(k)-1} \left(\frac{\alpha}{n}\right)^{\frac{k}{2} - \frac{Q(k)}{2}} \times \\ {}_3F_2 \left(1, \frac{Q(k)-k}{2}, \frac{Q(k)+k+2}{2}; \frac{Q(k)+1}{2}, \frac{Q(k)}{2} + 1; -\frac{1}{4n\alpha}\right)$$

# Classification via Casimir invariants

$$\det[\lambda - \Pi] = \lambda^N - S_1 \lambda^{N-1} + S_2 \lambda^{N-2} + \cdots + (-1)^N S_N = 0$$

$S_{N-1} = 0$  polynomial equation provides  $N^* = \left[ \frac{N-1}{2} \right]$  values for  $\alpha$ .

Cartan subalgebra decomposition of  $P$

$$P = \kappa \left( I + \sqrt{\frac{N(N-1)}{2}} \sum_{\lambda_a \in h} \xi_a \cdot \lambda_a \right).$$

Casimir invariants

$$\mathfrak{C}_2 = (N-1) \xi \cdot \xi \quad \mathfrak{C}_3 = (N-1) (\xi \vee \xi) \cdot \xi$$

$$\mathfrak{C}_4 = (N-1) (\xi \vee \xi) \cdot (\xi \vee \xi) \quad \mathfrak{C}_5 = (N-1) ((\xi \vee \xi) \vee (\xi \vee \xi)) \cdot \xi$$

where  $(U \vee V)_a := c d_{abc} U_a V_b$ .  $d_{abc}$  are symmetric structure constants for  $\mathfrak{su}(N)$  and  $c = \sqrt{N(N-1)/2}$

# Realization of the construction

$$S_k = \frac{1}{k!} \det \begin{pmatrix} t_1 & 1 & 0 & \cdots & 0 \\ t_2 & t_1 & 2 & \cdots & 1 \\ t_3 & t_2 & t_1 & \cdots & \\ \vdots & \vdots & \vdots & \ddots & k-1 \\ t_k & t_{k-1} & t_{k-2} & \cdots & t_1 \end{pmatrix}.$$
$$\det[\lambda - \Pi] = (\lambda - \lambda_1)^{N_1} (\lambda - \lambda_2)^{N_2}$$