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On the generalized Sundman transformations
and integrable Lienard type equations

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Liénard–type equations

- We consider the following family of equations

$$y_{zz} + f(y)y_z^2 + g(y)y_z + h(y) = 0$$

- Liénard-type equations have a vast range of applications in physics, biology, mechanics e.t.c.
- Recently there has been a great interest in finding classes of integrable Liénard-type equations : the Jacobi last multiplier, the Prelle-Singer method, the classical Lie method, the Chiellini lemma:

W. Nakpim, S.V. Meleshko, SIGMA, (2010); V.K. Chandrasekar et al., Proc. R. Soc. A Math. Phys. Eng. Sci. (2005); S.N. Pandey et al., J. Math. Phys., (2009); S.C. Mancas, H.C. Rosu, Phys. Lett. A., (2013); T. Harko et al., J. Eng. Math. (2014); A.K. Tiwari et al., Nonlin. Dynamics (2015)

Aim and approach

- Our main aim is to find new classes of integrable Liénard-type equations and to explain some of the previous results from a general perspective
- Our approach is based on the generalized Sundman transformations

$$y_{zz} + f(y)y_z^2 + g(y)y_z + h(y) = 0$$

$$w = F(y) \quad \begin{array}{c} \updownarrow \end{array} \quad \begin{array}{l} d\zeta = G(y)dz \\ F_y G \neq 0 \end{array}$$

$$w_{\zeta\zeta} + \tilde{f}(w)w_\zeta^2 + \tilde{g}(w)w_\zeta + \tilde{h}(w) = 0$$

Canonical form

- We consider the following equation

$$y_{zz} + f(y)y_z^2 + g(y)y_z + h(y) = 0$$

- **Proposition 1.** This equation can be transformed into the Liénard equation by means of GST with

$$F(y) = y, \quad G(y) = \exp\{-\int f dy\}.$$

As a result we get

$$y_{\zeta\zeta} + \tilde{g}(y)y_{\zeta} + \tilde{h}(y) = 0$$

$$\tilde{g}(y) = \exp\{\int f dy\}g(y), \quad \tilde{h}(y) = \exp\{2\int f dy\}h(y)$$

- We assume that $g(y) \not\equiv 0$, i.e. we consider a dissipative case.

Linearization conditions

- **Proposition 2.** The mixed Liénard-type equation can be transformed into $w_{\zeta\zeta} + \sigma w_{\zeta} + 2\sigma^2/9w = 0$ via GST if

$$h(y) = \frac{2g(y)}{9} e^{-\int f(y)dy} \left[\int e^{\int f(y)dy} g(y)dy + \kappa \right].$$

- As a consequence we obtain criterion for existence of a non-standard Lagrangian (Z.E. Musielak, J.Phys. A. (2008), CSF (2009), J.L. Cieřlinski, J.L., Nikiciuk, J.Phys.A. (2010), A.K. Tiwari et al., Acta Mech. 2016)

$$L = \frac{1}{w_{\zeta} + \frac{2\sigma}{3}w} \quad \xleftrightarrow{\text{GST}} \quad L = \frac{\sigma}{e^{\int f(y)dy} y_z + \frac{2}{3} \left[\int e^{\int f(y)dy} g(y)dy + \kappa \right]}$$

- Moreover, this class of linerizable mixed Liénard-type equations contains equations, which are completely integrable by the classical Lie approach, i.e. admit at least 2 dimensional Lie algebra and Chiellini integrable equations.

Connections with Painlevé–Gambier equations

- We study connections of the Liénard equations family with its subcases that are of the Painlevé–Gambier type: we consider Painlevé–Gambier equation of types I-III.
- We omit equations that can be linearized via the GST or are equivalent via the GST. Consequently, we need to consider equations II, V, VII, XV, XXIII, XXV, XXVII and XXVIII.
- As a result, we have obtained 9 criteria for the integrability of Liénard equations.

Example of integrability criterion (1)

- We consider non-canonical form of Painlevé–Gambier equation VII

$$w_{\zeta\zeta} + 3w_{\zeta} + w^3 + 2w = 0$$

$$w = e^{-(\zeta+\zeta_0)} \operatorname{cn}\{e^{-(\zeta+\zeta_0)} - C_1, 1/\sqrt{2}\}$$

- The Liénard-type equation can be transformed into equation VII of the Painlevé–Gambier type via GST if

$$h(y) = \frac{g(y)}{9} \mathcal{G}(\lambda^2 \mathcal{G}^2 + 2) e^{-\int f(y) dy}$$

- The corresponding GST have the form

$$F = \lambda \mathcal{G}, \quad G = \frac{1}{2}g, \quad \mathcal{G} = \int e^{\int f(y) dy} g(y) dy + \kappa$$

Example of integrability criterion (2)

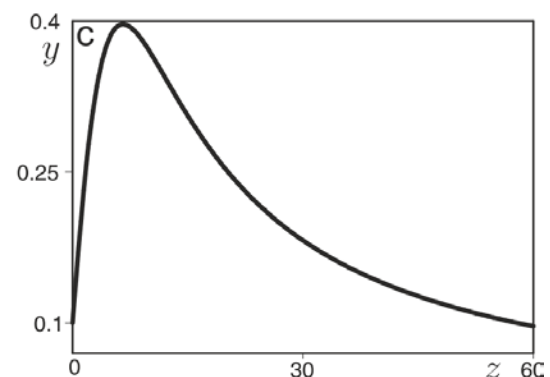
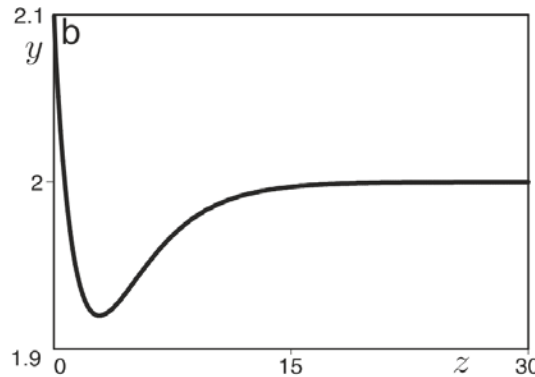
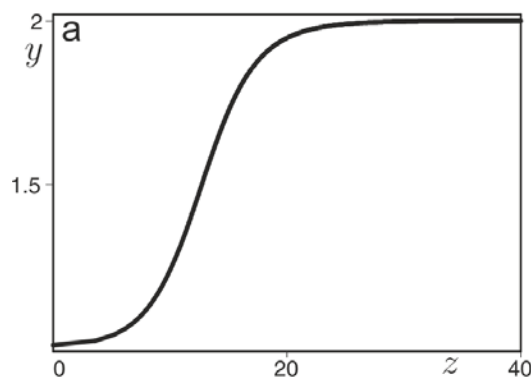
- Traveling wave reduction of the generalized Burgers-Huxley equation

$$y_{zz} + (\alpha y + \beta)y_z + \frac{y}{18}(\alpha y + \beta)(\alpha y + 2\beta)(y^2[\alpha y + 2\beta]^2 + 2) = 0$$

- The general solution has the form

$$y = \frac{1}{\alpha} \left[\pm \left(\beta^2 + \alpha e^{-(\zeta - \zeta_0)} \operatorname{cn} \left\{ e^{-(\zeta - \zeta_0)} - C_1, \frac{1}{\sqrt{2}} \right\} \right)^{1/2} - \beta \right],$$

$$z = \int \frac{3}{\alpha y + \beta} d\zeta.$$



The Liénard equation: Ince XXVII (1)

- We consider equation XXVII of the Painlevé–Gambier type

$$w_{\zeta\zeta} - \frac{4}{3} \frac{w_{\zeta}^2}{w} + \frac{5}{3} \frac{w_{\zeta}}{w} + 3\epsilon w - \frac{1}{3w} = 0$$

$$w = \frac{\sinh\{\sqrt{\epsilon}(\zeta - \zeta_0)\}(\cosh^2\{\sqrt{\epsilon}(\zeta - \zeta_0)\} + 2) + C_1}{3\sqrt{\epsilon} \cosh^3\{\sqrt{\epsilon}(\zeta - \zeta_0)\}}$$

- Lagrangian, first integral and Jacobi multiplier

$$L = -\frac{2}{3} \frac{9\epsilon w^2 - 2(w_{\zeta} - 1)^2}{\sqrt{9\epsilon w^2 - (w_{\zeta} - 1)^2}}$$

$$I = \frac{2(81\epsilon^2 w^4 + 27\epsilon(w_{\zeta} - 1)w^2 - 2(w_{\zeta} - 1)^3)}{3(9\epsilon w^2 - (w_{\zeta} - 1)^2)^{3/2}}$$

$$M = \frac{162\epsilon^2 w^4}{(9\epsilon w^2 - (w_{\zeta} - 1)^2)^{5/2}}$$

The Liénard equation: Ince XXVII (2)

- Criterion of equivalence between the Liénard equation and Painlevé–Gambier XXVII

$$h(y) = -\frac{4}{25}g(y)\mathcal{G} \left[9\epsilon\lambda^2\mathcal{G}^{-3/2} - 1 \right] e^{-\int f(y)dy}$$

$$F = \lambda\mathcal{G}^{-3/4}, G(y) = \frac{3\lambda}{5}g(y)\mathcal{G}^{-3/4}, \mathcal{G} = \int e^{\int f(y)dy}g(y)dy + \kappa$$

- The corresponding Lagrangian, first integral and Jacobi multiplier

$$L = -\frac{2}{3} \frac{9\epsilon\lambda^2\mathcal{G}^{-3/2} - 2\left[\frac{5}{4}\mathcal{G}^{-1}e^{\int f dy}y_z + 1\right]^2}{\left(9\epsilon\lambda^2\mathcal{G}^{-3/2} - \left[\frac{5}{4}\mathcal{G}^{-1}e^{\int f dy}y_z + 1\right]^2\right)^{1/2}} \quad M = \frac{2025\epsilon^2\lambda^4\mathcal{G}^{-5}e^{2\int f dy}}{8\left(9\epsilon\lambda^2\mathcal{G}^{-3/2} - \left[\frac{5}{4}\mathcal{G}^{-1}e^{\int f dy}y_z + 1\right]^2\right)^{5/2}}$$

$$I = \frac{2}{3} \left[81\epsilon^2\lambda^4\mathcal{G}^{-3} - 27\epsilon\lambda^2\mathcal{G}^{-3/2} \left(\frac{5}{4}\mathcal{G}^{-1}e^{\int f dy}y_z + 1 \right) + 2 \left(\frac{5}{4}\mathcal{G}^{-1}e^{\int f dy}y_z + 1 \right)^3 \right] \left(9\epsilon\lambda^2\mathcal{G}^{-3/2} - \left[\frac{5}{4}\mathcal{G}^{-1}e^{\int f dy}y_z + 1 \right]^2 \right)^{-3/2}$$

The Liénard equation: Ince XXVIII (1)

- We consider equation XXVIII of the Painlevé–Gambier type

$$ww_{\zeta\zeta} - \frac{1}{2}w_{\zeta}^2 + w^2w_{\zeta} - \frac{1}{2}w^4 + 72H = 0$$

$$w = \frac{6(\wp^2\{\zeta - \zeta_0, 12H, g_3\} - H)}{\wp_z(\zeta - \zeta_0, 12H, g_3)}$$

- Lagrangian, first integral and Jacobi multiplier

$$L = -\frac{(w_{\zeta} + w^2)^2}{3w^2} \left[(w_{\zeta} + w^2)^2 - 864H \right] + \frac{20736H^2}{w^2}$$

$$I = pw_{\zeta} - L = -\frac{(3w_{\zeta} - w^2)(w_{\zeta} + w^2)^3}{3w^2} + \frac{288H(w_{\zeta}^2 - w^4 - 72H)}{w^2}$$

$$M = 4[144H - (w_{\zeta} + w^2)^2]w^{-2}$$

The Liénard equation: Ince XXVIII (2)

- Criterion of equivalence between the Liénard equation and Painlevé–Gambier XXVIII

$$h(y) = (3/4)g(y)\lambda^{-4}\mathcal{G}^{-5/3} [144H - \lambda^4\mathcal{G}^{8/3}]$$

$$F(y) = \lambda\mathcal{G}^{2/3}, \quad G = g(y)\lambda^{-1}\mathcal{G}^{-2/3}, \quad \mathcal{G} = \int e^{\int f(y)dy} g(y)dy + \kappa$$

- The corresponding Lagrangian, first integral and Jacobi Multiplier

$$L = -(\lambda^2/3) \left[\frac{2}{3}y_z + \mathcal{G} \right]^2 \mathcal{G}^{-2/3} \left[\mathcal{G}^{2/3}\lambda^4 \left[\frac{2}{3}y_z + \mathcal{G} \right]^2 - 864H \right] + \frac{20736H^2}{\lambda^2\mathcal{G}^{4/3}}.$$

$$I = -\frac{\lambda^6(2y_z - \mathcal{G})}{3} \left(\frac{2}{3}y_z + \mathcal{G} \right)^3 + \frac{288H\lambda^2}{\mathcal{G}^{2/3}} \left(\frac{4}{9}y_z^2 - \mathcal{G}^2 \right) - \frac{20736H^2}{\lambda^2\mathcal{G}^{4/3}}.$$

$$M = \frac{576H}{\lambda^2\mathcal{G}^{4/3}} - 4\lambda^2\mathcal{G}^{-2/3} \left(\frac{2}{3}y_z + \mathcal{G} \right)^2$$

Example: the modified Emden equation

- We assume that $f(y) = 0$, $g(y) = \alpha y$, $\kappa = 0$, $\lambda^2 = (25/36)(2\alpha)^{-1/2}$

$$y_{zz} + \alpha y y_z + \frac{2\alpha^2}{25} y^3 - \epsilon = 0$$

- Lagrangian, first integral and Jacobi multiplier

$$L = -\frac{2(25\epsilon y - (5y_z + 2\alpha y^2)^2)}{3\alpha y^2 \sqrt{50\epsilon y - (5y_z + 2\alpha y^2)^2}}$$

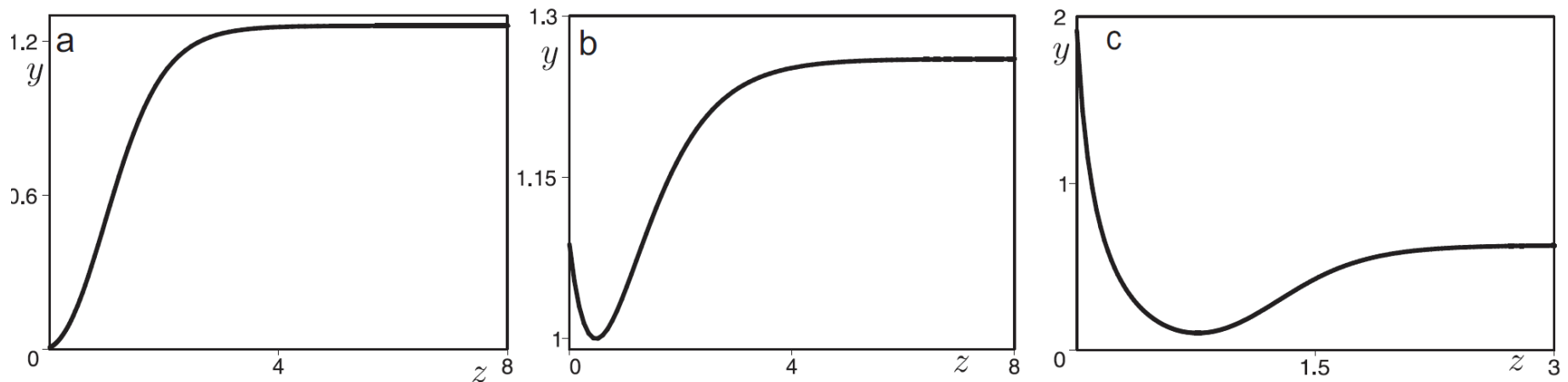
$$I = \frac{2500\epsilon^2 - 300\epsilon\alpha y(5y_z + 2\alpha y^2) + 4\alpha(5y_z + 2\alpha y^2)^3}{3\alpha(50\epsilon y - (5y_z + 2\alpha y^2)^2)^{3/2}}$$

$$M = \frac{62500\epsilon^2}{(50\epsilon y - (5y_z + 2\alpha y^2)^2)^{5/2}}$$

Example: the modified Emden equation

- The general solution has the form

$$y = \left[\frac{5\sqrt{2\epsilon}}{2\alpha} \frac{\cosh^3\{\sqrt{\epsilon}(\zeta - \zeta_0)\}}{\sinh\{\sqrt{\epsilon}(\zeta - \zeta_0)\}(\cosh^2\{\sqrt{\epsilon}(\zeta - \zeta_0)\} + 2) + C_1} \right]^{2/3}, \quad z = \int \sqrt{2y} d\zeta.$$



Example: trigonometric nonlinearity

- We suppose that $f(y) = -\tanh(y)$, $g(y) = y^3 \cosh(y)$, $\kappa = 0$

$$y_{zz} - \tanh(y)y_z^2 + y^3 \cosh(y)y_z + \frac{1}{25}(y^6 - 72\epsilon\lambda^2)y \cosh^2 y = 0$$

- Lagrangian, first integral and Jacobi multiplier

$$L = \frac{4}{3} \frac{(5y_z + y^4 \cosh(y))^2 - 36\epsilon\lambda^2 y^2 \cosh^2(y)}{\cosh(y)y^4 \sqrt{72\epsilon\lambda^2 y^2 \cosh^2(y) - (5y_z + y^4 \cosh(y))^2}}$$

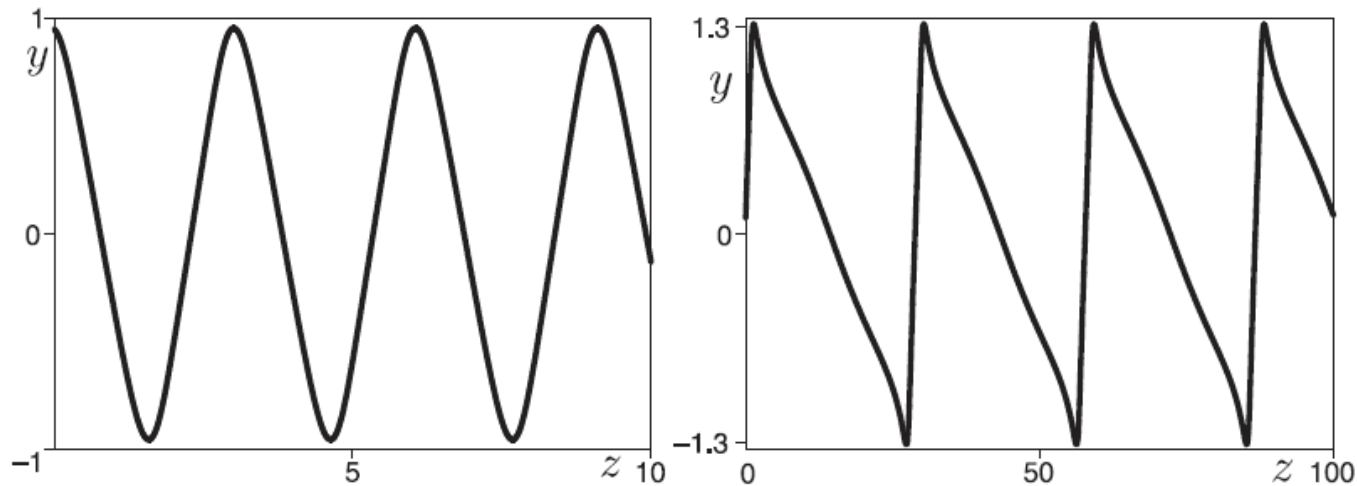
$$M = \frac{259200\epsilon^2 \lambda^4 \cosh^3(y)}{(72\epsilon\lambda^2 y^2 \cosh^2(y) - (5y_z + y^4 \cosh(y))^2)^{5/2}}$$

$$I = \frac{4}{3} [2592\epsilon^2 \lambda^4 \cosh^3(y) - 108\epsilon\lambda^2 y^2 \cosh^2(y)(5y_z + y^4 \cosh(y)) + (5y_z + y^4 \cosh(y))^3] (72\epsilon\lambda^2 y^2 \cosh^2(y) - (5y_z + y^4 \cosh(y))^2)^{-3/2}$$

Example: trigonometric nonlinearity

- The general solution can be expressed as follows

$$y = \left[\frac{6\lambda\sqrt{2\epsilon} \cosh^3\{\sqrt{\epsilon}(\zeta - \zeta_0)\}}{\sinh\{\sqrt{\epsilon}(\zeta - \zeta_0)\}(\cosh^2\{\sqrt{\epsilon}(\zeta - \zeta_0)\} + 2) + C_1} \right]^{1/3}, \quad z = \frac{5}{6\sqrt{2}\lambda} \int \cosh^{-1}(y) d\zeta.$$



Example: rational nonlinearity

- We assume that $f(y) = 1/(2y)$, $g(y) = -a^3/y^3$, $\kappa = 0$, $\lambda = 2^{-2/3}3^{-1/3}$

$$y_{zz} + \frac{1}{2y}y_z^2 - \frac{a^3}{y^3}y_z + \frac{a^6}{2y^5} - \frac{5832H}{a^2y} = 0$$

- The corresponding Lagrangian, first integral and Jacobi last multiplier

$$L = -\frac{(a^3 + y^2y_z)^4}{2187y^6} + \frac{32H(a^3 + y^2y_z)^2}{a^2y^2} + \frac{186624H^2y^2}{a^4}$$

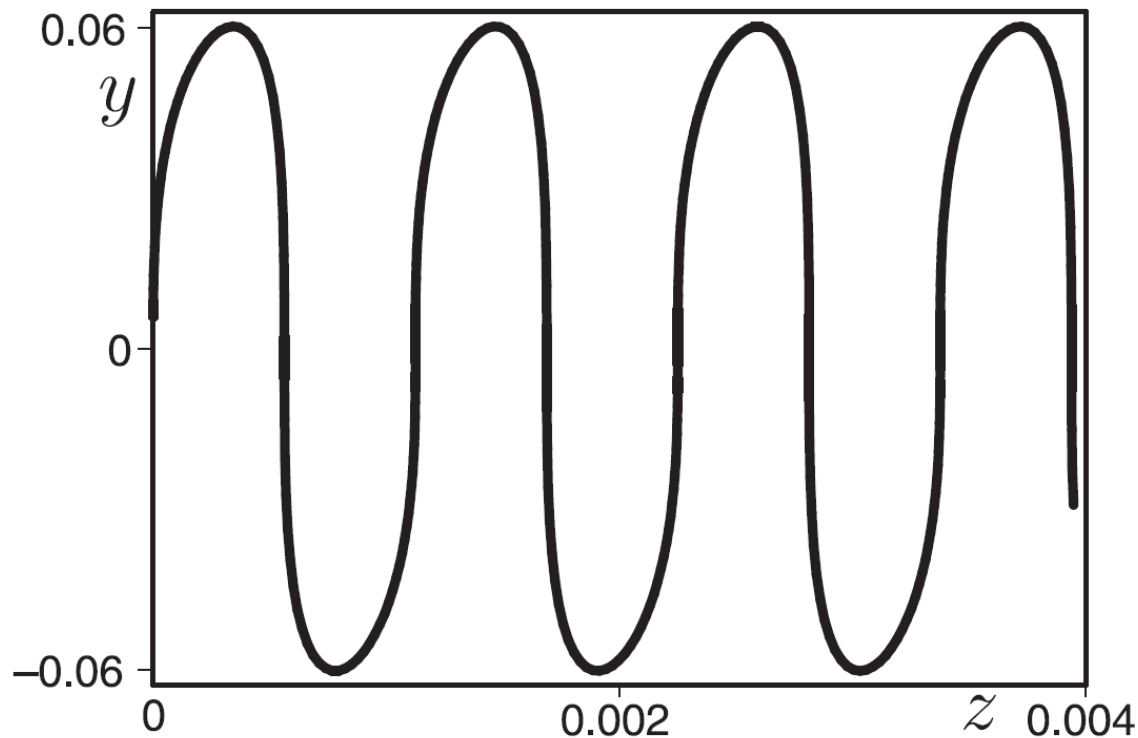
$$I = -\frac{(3y^2y_z - a^3)(y^2y_z + a^3)^3}{2187y^6} + \frac{32H(y^2y_z - a^3)(y^2y_z + a^3)}{y^2a^2} - \frac{186624H^2y^2}{a^4}$$

$$M = \frac{64Hy^2}{a^2} - \frac{4(y^2y_z + a^3)^2}{729y^2}$$

Example: rational nonlinearity

- The general solution has the form

$$y = \frac{a^2 \wp_z(\zeta - \zeta_0, 12\beta, g_3)}{18(\wp^2\{\zeta - \zeta_0, 12\beta, g_3\} - \beta)}, \quad z = -\frac{1}{3a} \int y^2 d\zeta$$



Conclusion

- We have demonstrated that some of previously known analytical results, e.g. integrability criteria and Lagrangian structures, for the Liénard equation follows from linearizability by GST
- We have found new criteria for the integrability of the Liénard-type equations and constructed new examples of integrable Liénard-type equations
- We have demonstrated that our approach can be applicable for finding Lagrangians, first integrals and Jacobi multipliers for the Liénard-type equations



Thank you for your attention!