

Charmonium polarization in pp collisions

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- 1. The emblematic case of the J/ψ polarization at high- p_T**
- 2. Low- p_T : puzzle or opportunity?**

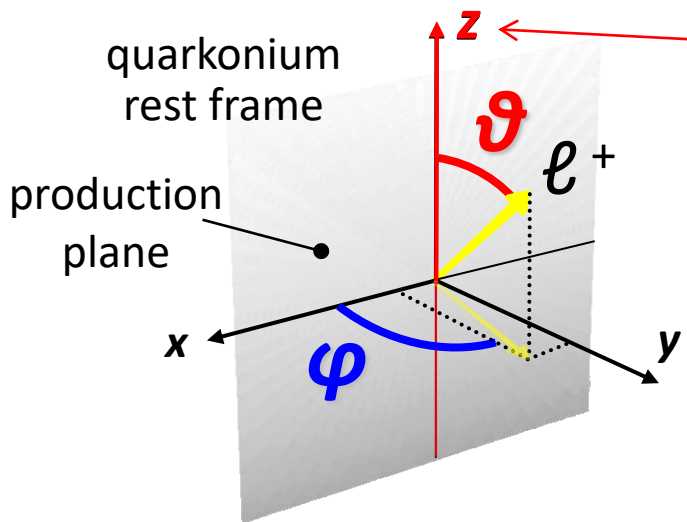
In collaboration with Carlos Lourenço, CERN

Workshop

“Gluon content of proton and deuteron with the Spin Physics Detector at the NICA collider”

JINR, Dubna, October 1st, 2020

Vector particle polarization: frames and parameters



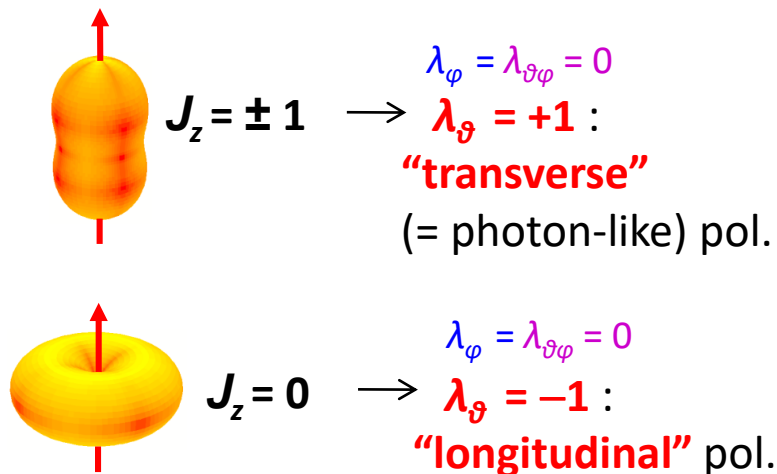
Collins-Soper axis (CS): \approx direction of colliding partons

Gottfried-Jackson (GJ): dir. of one beam (or the target)

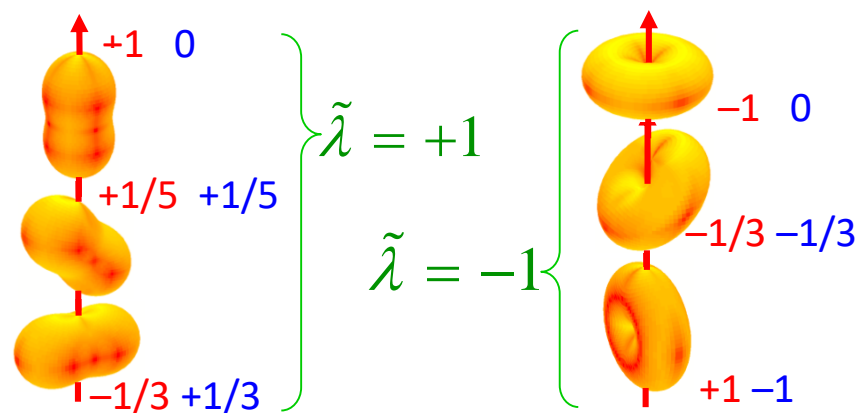
pp-helicity axis (HX): dir. of particle momentum w.r.t. to the pp c.o.m.

$$\frac{dN}{d\Omega} \propto 1 + \lambda_{\theta} \cos^2 \theta + \lambda_{\varphi} \sin^2 \theta \cos 2\varphi + \lambda_{\theta\varphi} \sin 2\theta \cos \varphi$$

λ_{θ} λ_{φ} $\lambda_{\theta\varphi}$ frame-dependent



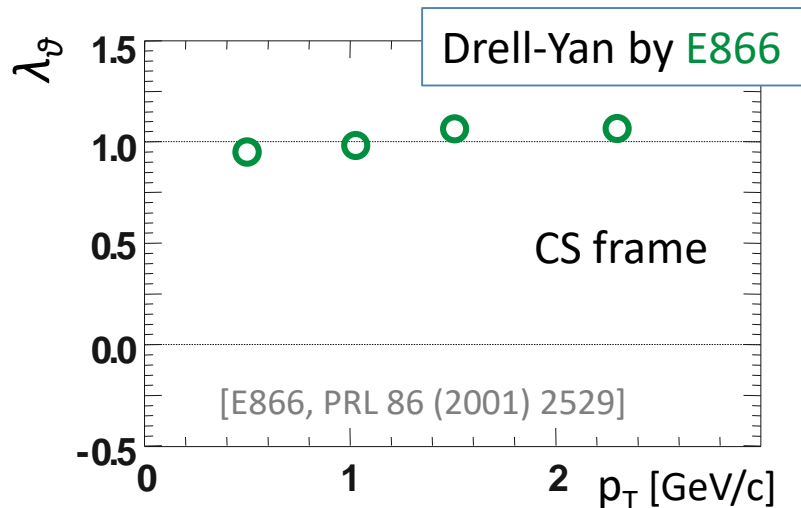
$\tilde{\lambda} = \frac{\lambda_{\theta} + 3\lambda_{\varphi}}{1 - \lambda_{\varphi}}$ frame-independent



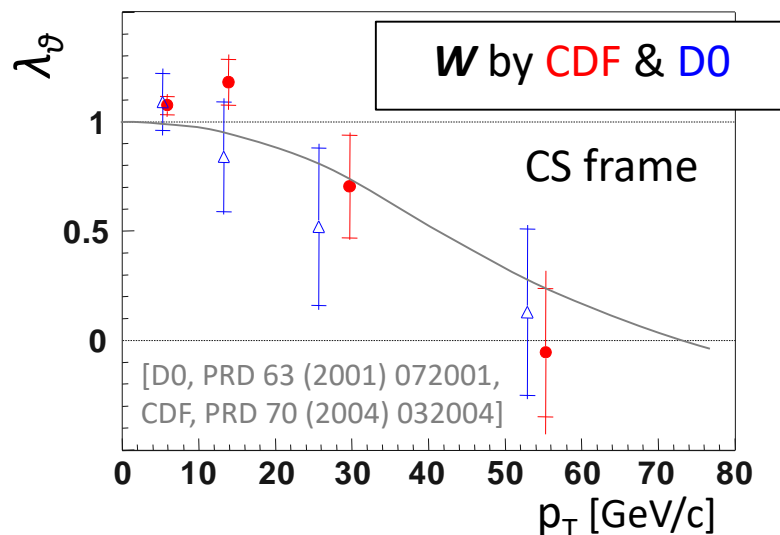
Vector particles are always polarized

The production of Z , W , γ and γ^* (Drell-Yan) is generally well explained by the short-distance coupling of quarks and gluons.

In particular, for **helicity conservation** the polarization is always **transverse** along some natural axis z



At very low p_T and for dominant **2-to-1 processes**, of order $\mathcal{O}(\alpha_s^0)$, a fully transverse polarization is seen in the **Collins-Soper** frame



At high p_T and for dominant **2-to-2 processes**, of order $\mathcal{O}(\alpha_s^1)$, a fully transverse polarization **would be** seen in the **helicity** frame.

The CS frame *smears* λ_γ away from $p_T = 0$. As a recognizable consequence, the polarization becomes **strongly p_T dependent** (from **+1** to **-1/3**)

Is “unpolarized” even possible?

Vector states are intrinsically polarized for any given elementary process

Theorem [P.F. et al., PRL 105, 061601]

For any subprocess producing a **$J = 1$ state**

$$|V; J, J_z\rangle = a_{-1} |1, -1\rangle + a_0 |1, 0\rangle + a_{+1} |1, +1\rangle,$$

there exists a quantization axis

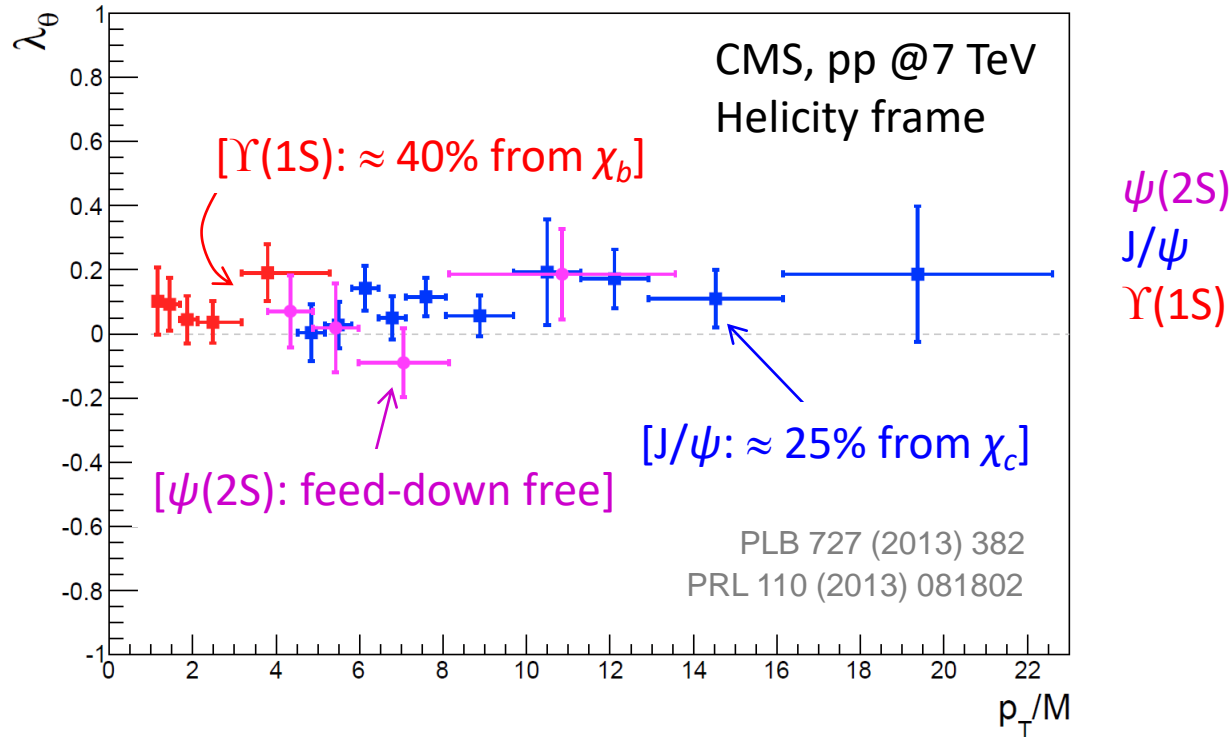
along which the $J_z = 0$ component **a_0 vanishes**

Intuitively consistent with classical expectation:
a vector of modulus 1 has always projection ± 1 along some axis

...which implies that **$\lambda_g = +1$ along that axis**

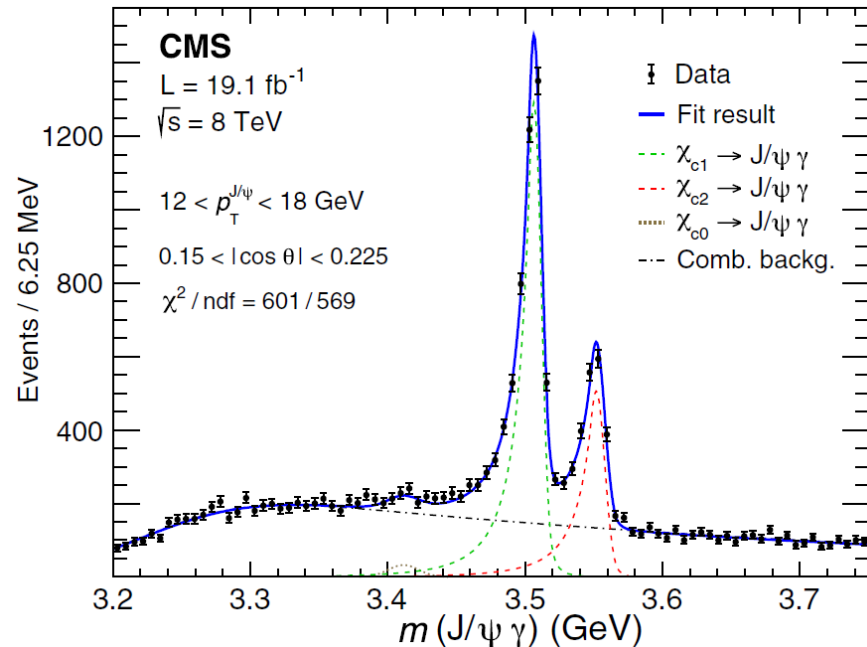
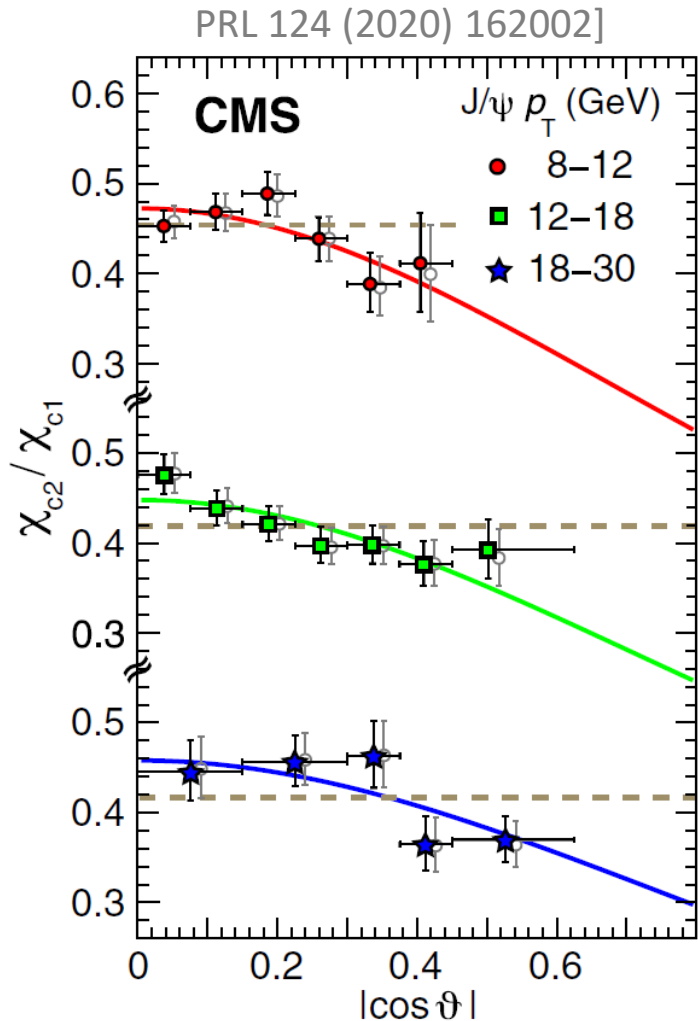
Vector quarkonia: a paradigmatic exception

Mid-rapidity LHC data show unpolarized production of vector quarkonia



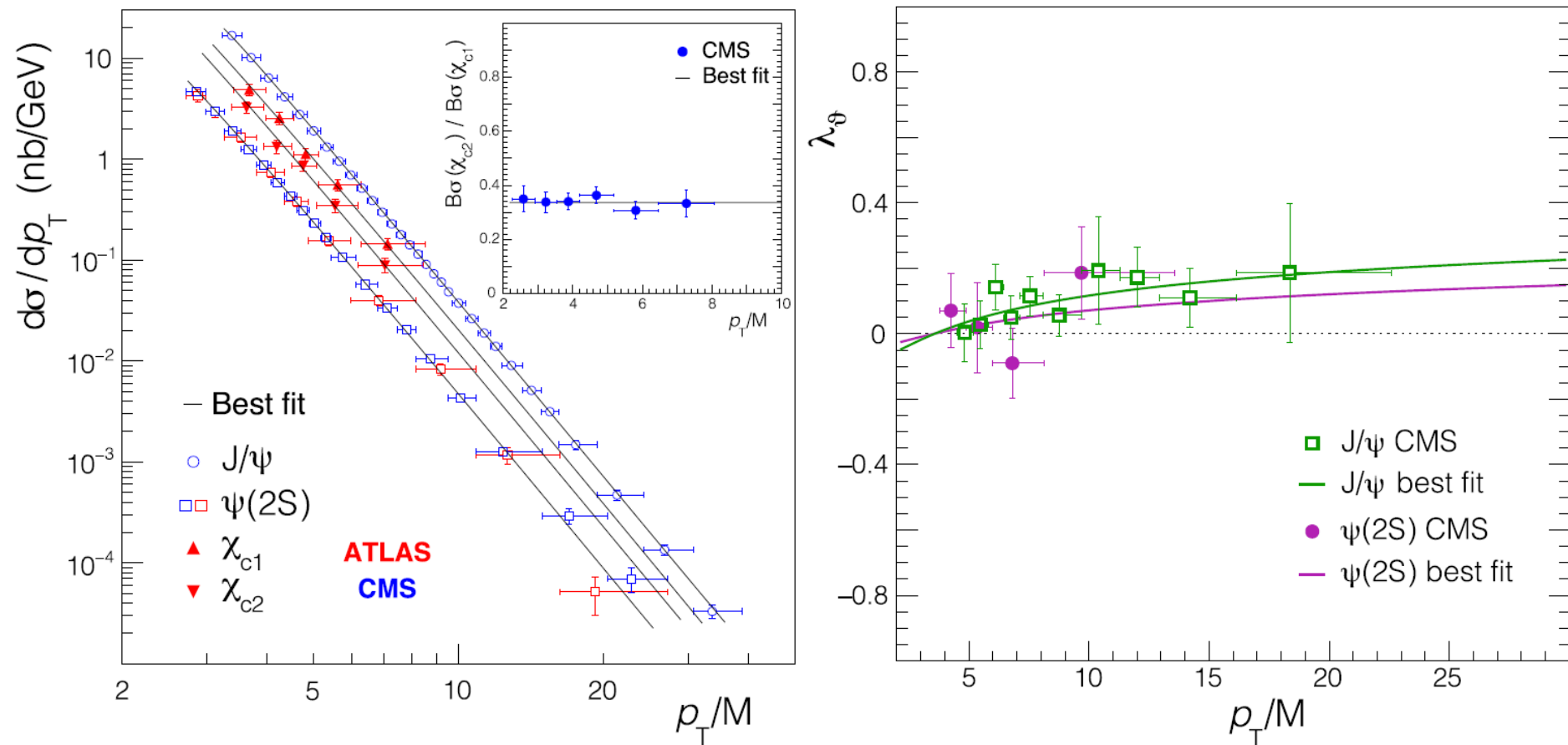
- None of the parameters λ_θ , λ_ϕ , $\lambda_{\theta\phi}$, $\tilde{\lambda}$ is significantly $\neq 0$
- There is **no visible dependence on p_T** : seemingly not a transition domain
- No visible difference between states despite different χ feed-downs \longrightarrow

The role of χ_c decays: finally from data



CMS measured the ratio between the (J/ψ from) χ_{c2} and $\chi_{c1} \cos\vartheta$ distributions. This provides a constraint on the *difference* between the two polarizations

Indirect experimental constraints



ATLAS and CMS measurements of J/ψ , $\psi(2S)$, χ_{c1} and χ_{c2} cross sections, together with the J/ψ and $\psi(2S)$ polarizations, constrain the *sum* of the χ_{c1} and χ_{c2} polarizations

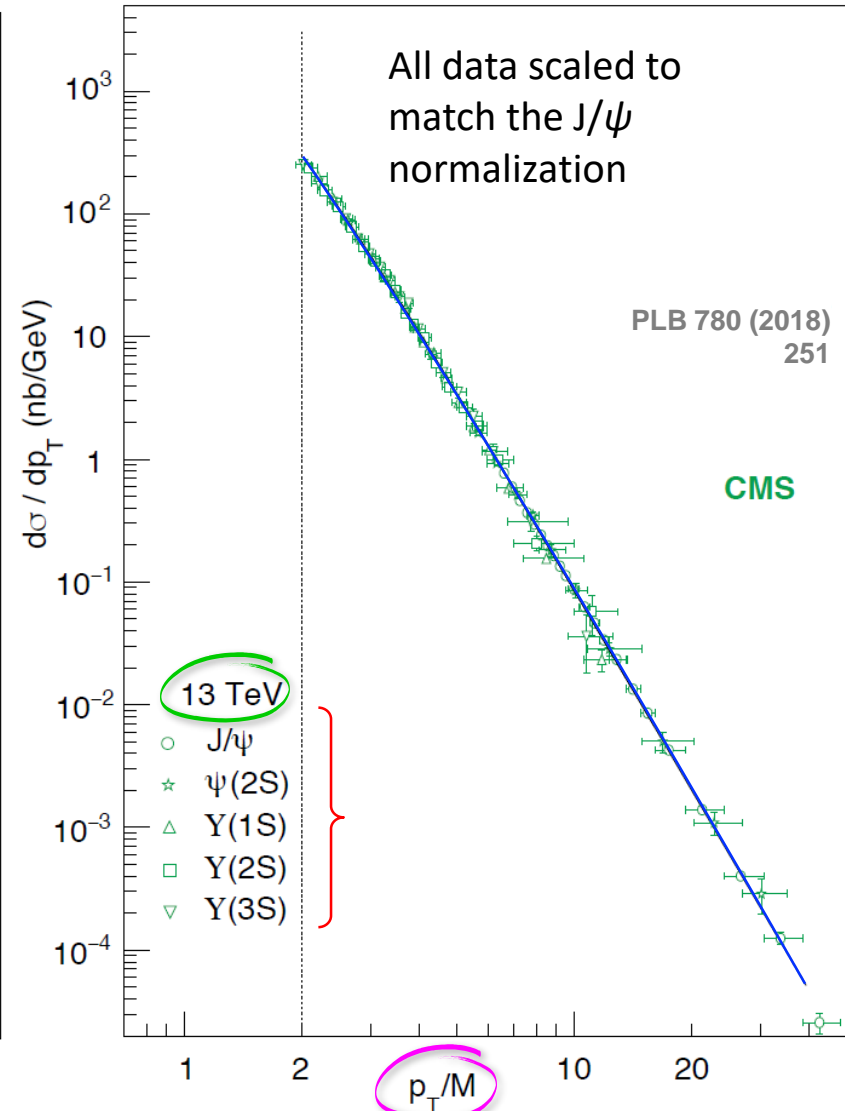
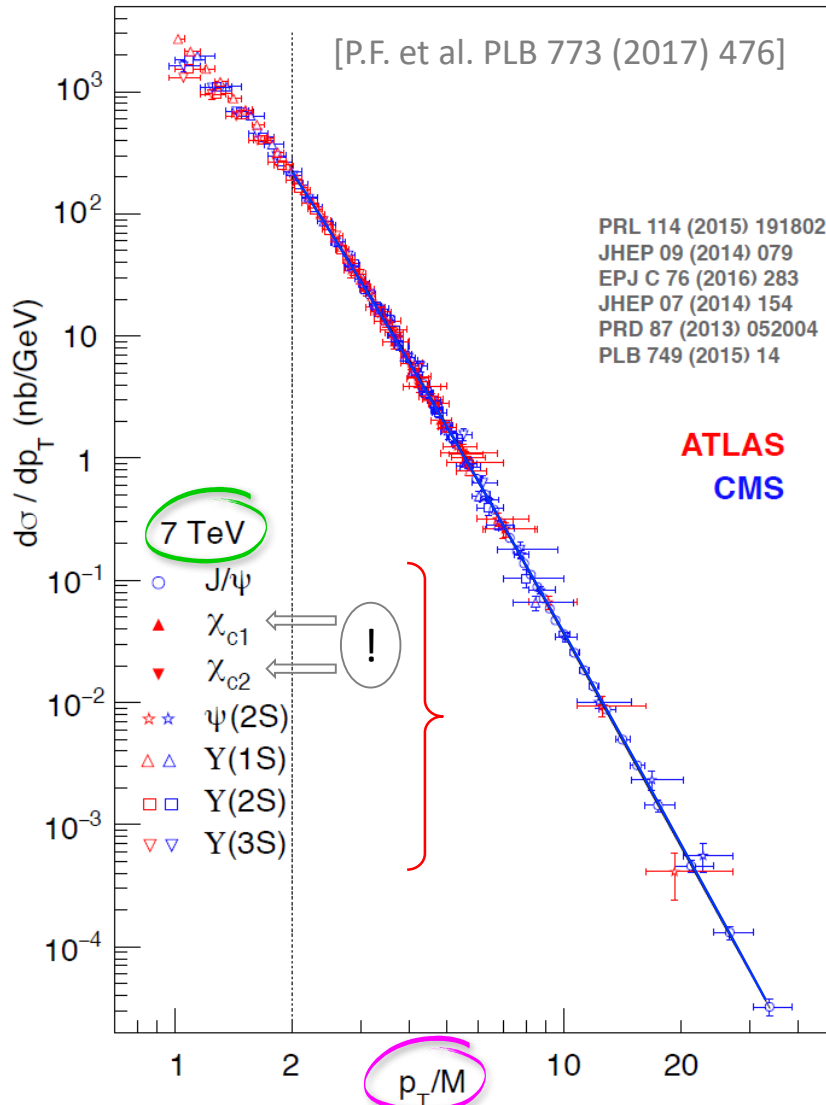
(*) Only assumption: *directly* produced J/ψ and $\psi(2S)$ have the same polarization vs p_T/M

(*) A “universal” p_T/M scaling

No hint of mass-dependence in mid-rapidity p_T distributions (nor for λ_g)

from J/ψ to $\Upsilon(3S)$ after dimensional scaling, $p_T \rightarrow p_T/M$, at least for $p_T/M > 2$

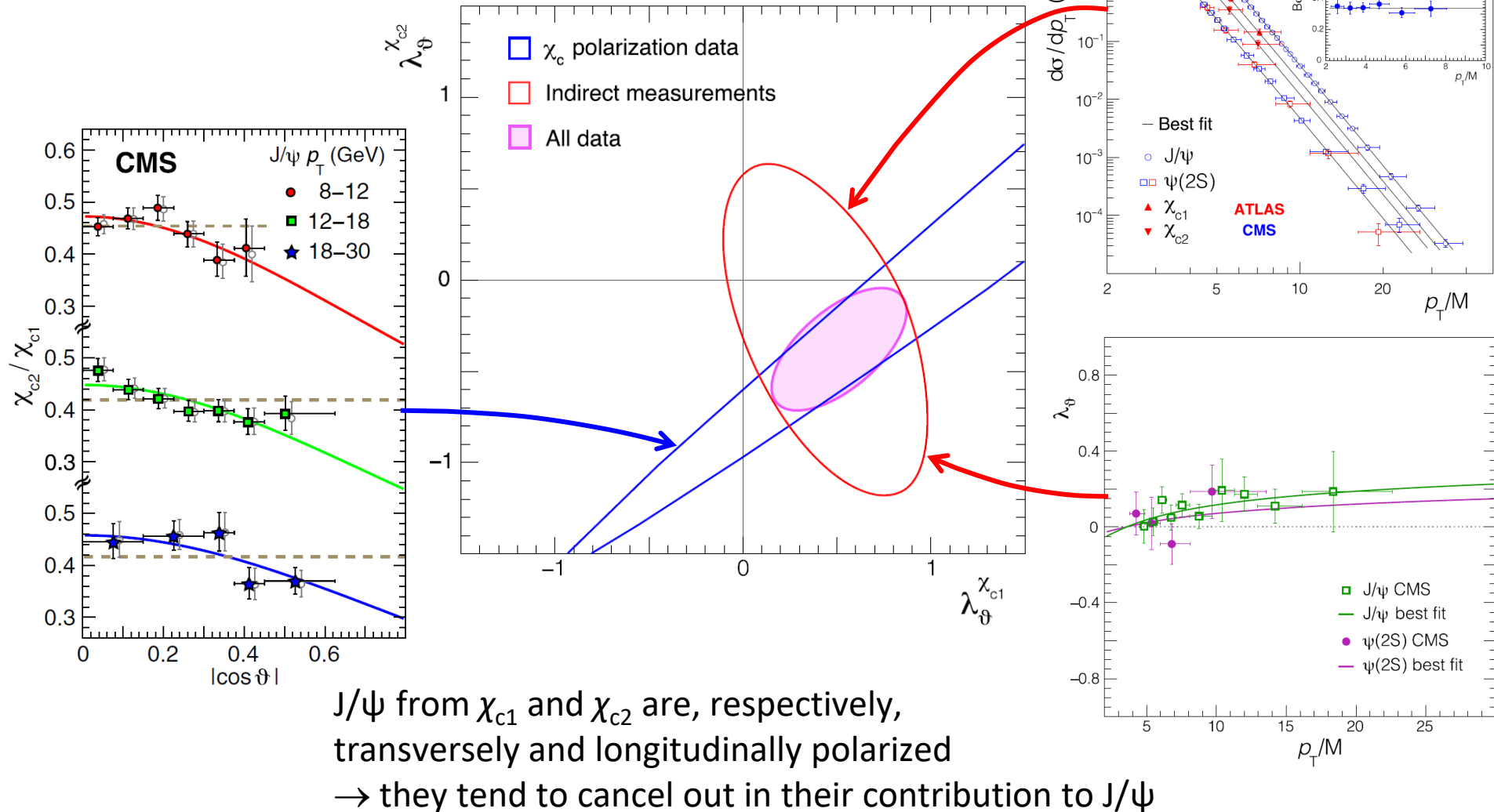
→ no reason to question similarity of production dynamics between direct J/ψ and $\psi(2S)$



The χ_c states are strongly polarized!

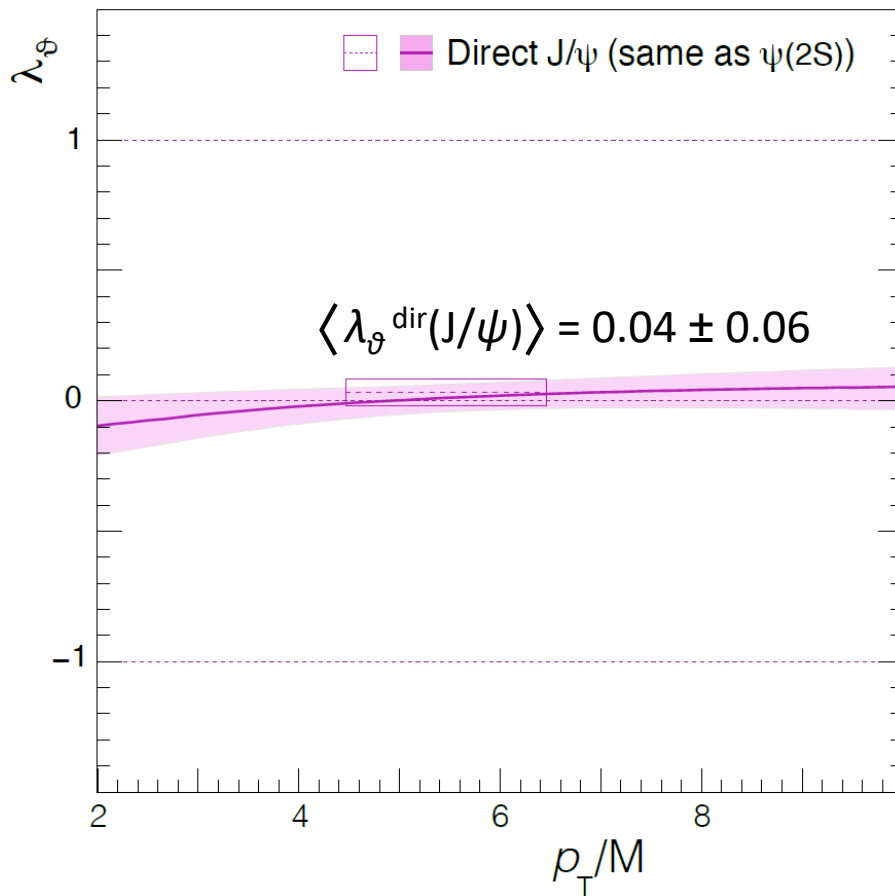
The combination of these two “orthogonal” experimental constraints determine the two individual χ_{c1} and χ_{c2} polarizations

[P.F. et al. EPJC 80 (2020) 623]



...and the J/ψ polarization is even more “zero”!

The global data fit also allows us to extract a measurement of the polarization of the *directly produced* J/ψ



A stronger evidence of unpolarized production!

Zero and **constant** polarization is a big challenge to production models

Only a “fortunate” **mixture of subprocesses** or **randomization effects** can lead to zero polarization

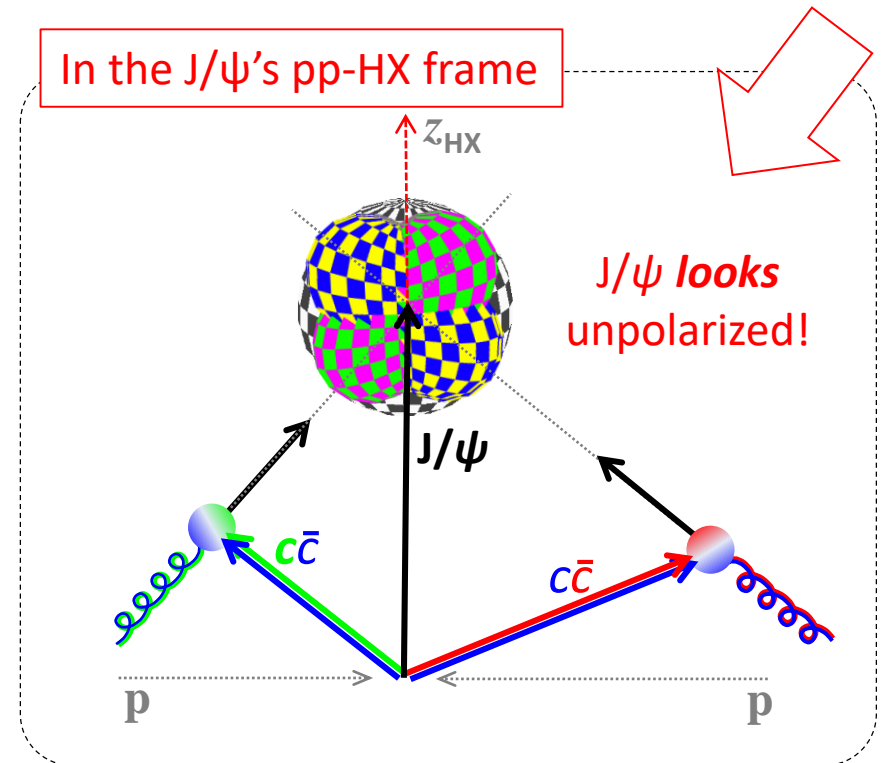
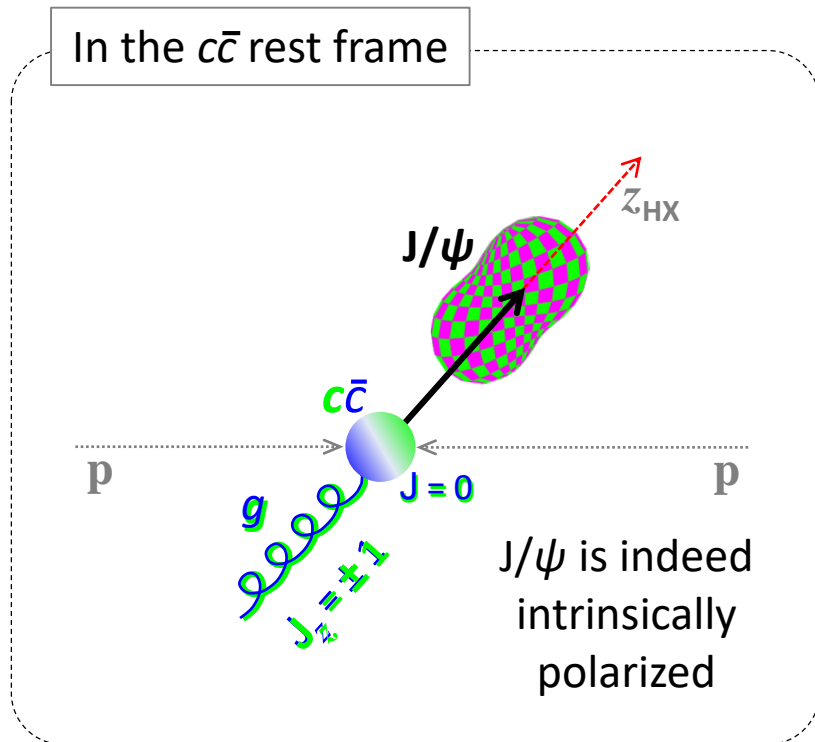
→ a clear sign of the **unique nature and production mechanism** of heavy quarkonia

Are we seeing a cascade mechanism?

Without invoking any theory framework, the most natural way to explain a zero polarization observation is a two-step mechanism with an **unobserved intermediate $J = 0$ state**

$$\text{E.g.: } pp \rightarrow c\bar{c}[J=0] \rightarrow J/\psi g g g$$

In the transition from the $J = 0$ “pre-resonance” to the vector bound state, the polarization is fully **randomized** because we lose connection to its natural reference

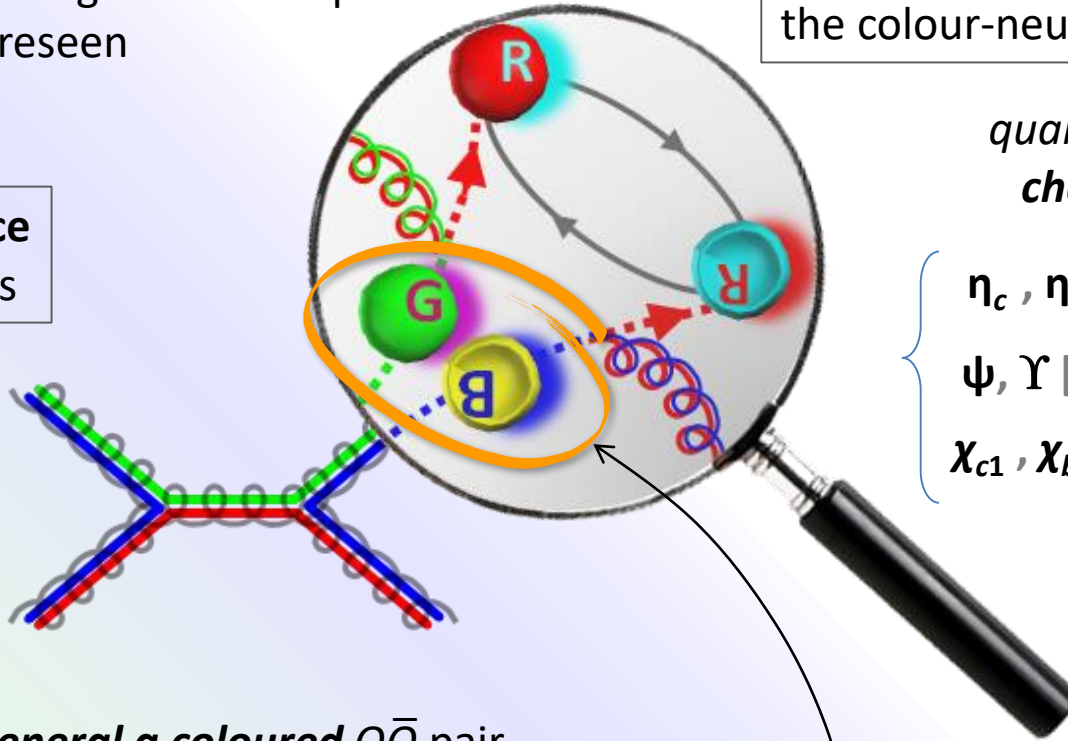


The “cascade” (*factorization*) approach of NRQCD

Non-Relativistic

For **heavy** quarkonia
two distinguishable steps
are foreseen

1) **short-distance**
partonic process



2) **long-distance** evolution to
the colour-neutral bound state

quantum numbers
change to final

$$\left\{ \begin{array}{l} \eta_c, \eta_b [^1S_0] \\ \psi, \Upsilon [^3S_1] \quad \chi_{c0}, \chi_{b0} [^3P_0] \\ \chi_{c1}, \chi_{b1} [^3P_1] \quad \chi_{c2}, \chi_{b2} [^3P_2] \end{array} \right.$$

produces **in general a coloured** $Q\bar{Q}$ pair
of any $^{2S+1}L_J$ quantum numbers

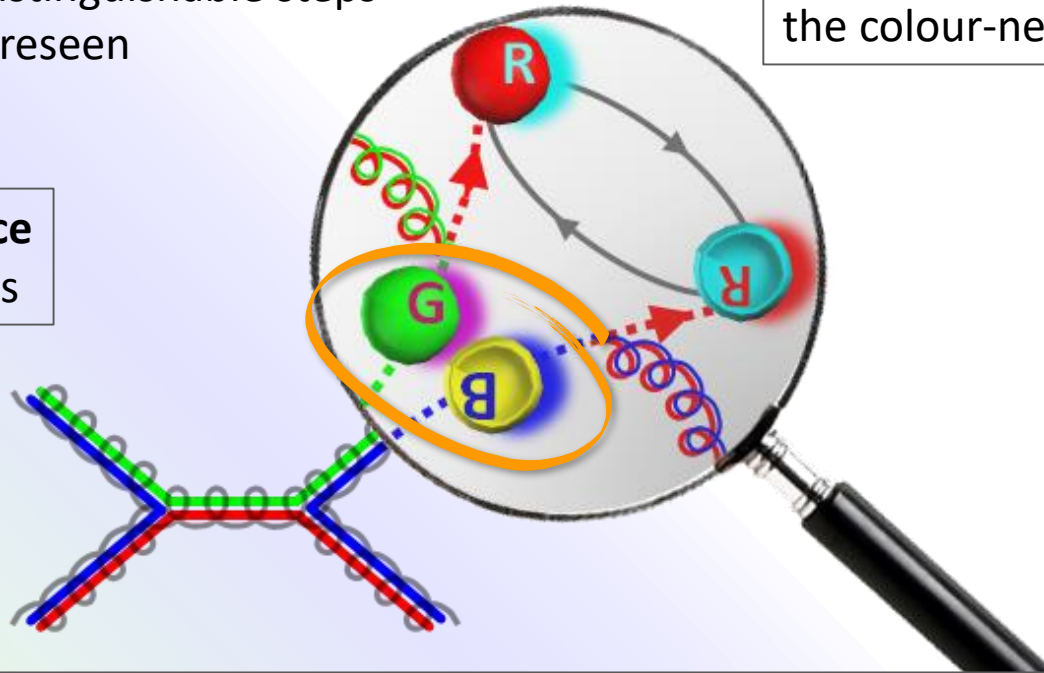
$$\begin{array}{ccccccc} {}^1S_0 & {}^1S_0 & {}^3S_1 & {}^3P_0 & {}^3P_2 & & \\ & {}^1D_2 & {}^3P_2 & {}^3D_3 & {}^1P_1 & {}^3S_1 & \\ {}^3P_1 & & {}^3P_2 & {}^3D_3 & {}^1P_1 & & \\ & {}^3D_2 & {}^3D_1 & {}^3P_1 & & & \end{array}$$

Even if the **pre-resonance** $Q\bar{Q}$ state
is not observed, it determines,
with its own quantum properties,
the observable kinematics and *polarization*

The “cascade” (*factorization*) approach of NRQCD

For **heavy** quarkonia
two distinguishable steps
are foreseen

1) **short-distance**
partonic process



2) **long-distance** evolution to
the colour-neutral bound state

1) *short-distance coefficients (SDCs)*:
 p_T -dependent partonic cross sections

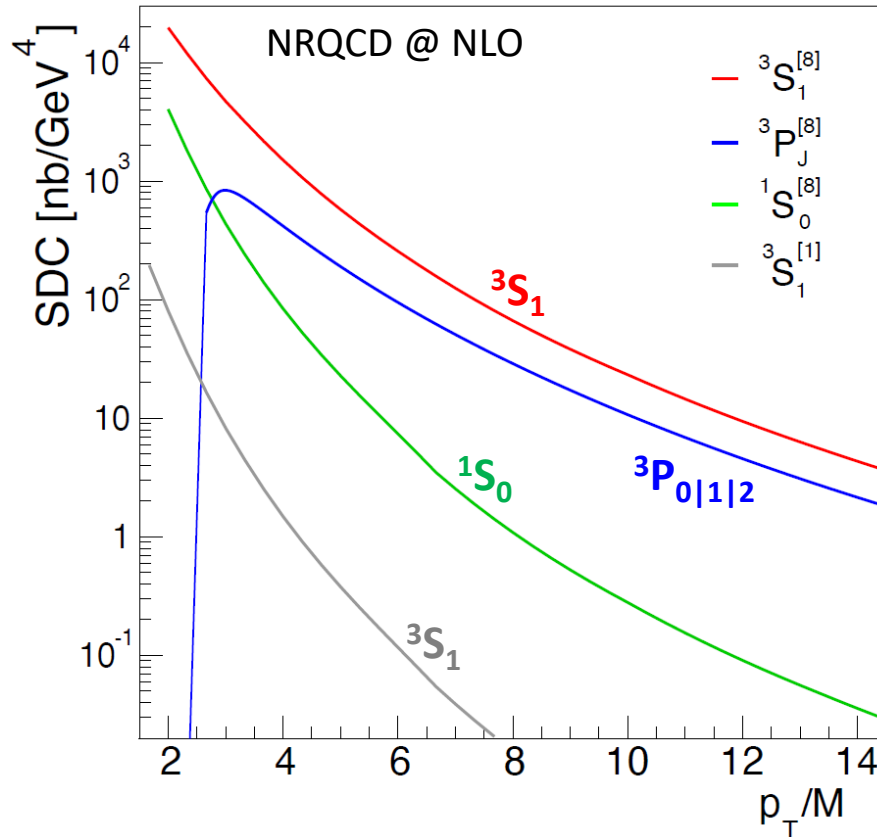
2) *long-distance matrix elements (LDMEs)*:
constant, **fitted from data**

$$\sigma(A + B \rightarrow Q + X) = \sum_{S, L, C} \mathcal{S}\{A + B \rightarrow (Q\bar{Q})_C [{}^{2S+1}L_J] + X\} \cdot \mathcal{L}\{(Q\bar{Q})_C [{}^{2S+1}L_J] \rightarrow Q\}$$

$Q\bar{Q}$ **angular momentum**
and **colour** configurations

Direct J/ψ in NRQCD: the “bricks” of the p_T distribution

A hierarchy in the expansion over the “small” Q - Q bar relative velocity (“ v -scaling”) foresees the dominance of a few of the ${}^{2S+1}L_J$ cascade channels:



1S_0 octet

3S_1 octet

${}^3P_{0|1|2}$ octets

3S_1 singlet

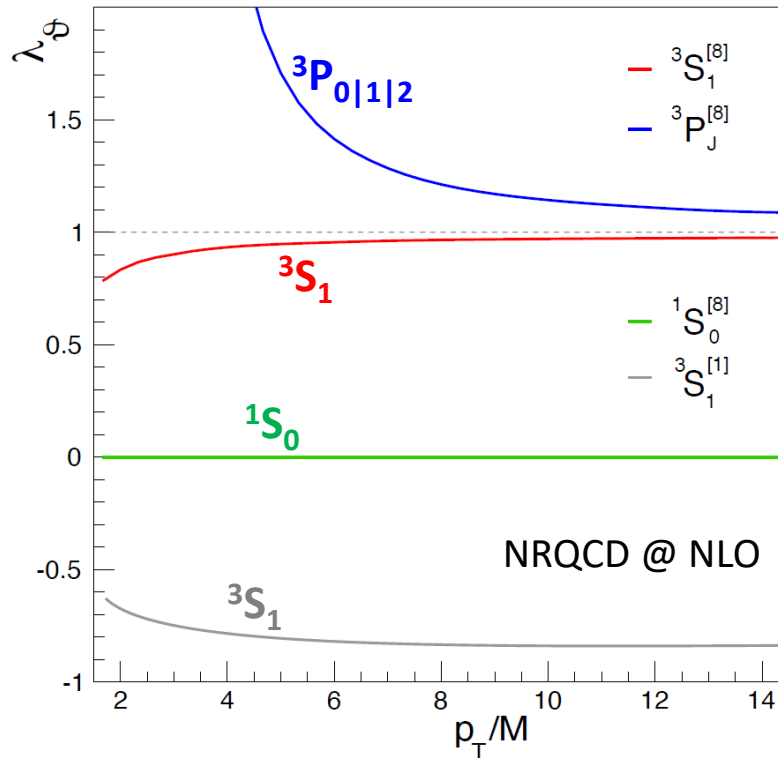
← P-wave term actually negative:
proper cancellation needed
to recover the physical cross section

Mixture of different pre-resonance contributions,
with characteristic p_T spectra (and polarizations: see next slide)

→ by *fitting* the experimental p_T distributions it is possible to determine the coefficients of all terms (LDMEs) and consequently *predict* the polarizations

The polarization terms: pieces of a puzzle?

Of the four contributing terms, only the 1S_0 leads “naturally” to zero polarization:



P-wave term actually unphysical ($> +1$)
proper cancellation needed
to recover the physical polarization

1S_0 octet

3S_1 octet

$^3P_{0|1|2}$ octets

3S_1 singlet

To reproduce the data, the remaining terms must

- either be individually suppressed
→ violation of NRQCD's v^2 hierarchy!
- or sum to \sim zero → redundant expansion basis!

Zero J/ψ polarization
is a **conceptual**
puzzle for NRQCD!

What about χ_{c1} and χ_{c2} ?

In NRQCD, $\chi_{c1,2}$ production has two terms: 3S_1 octet and ${}^3P_{1,2}$ singlet.

One parameter r determines

1) the χ_{c2} / χ_{c1} yield ratio

2) $\lambda_{\vartheta}(\chi_{c1})$

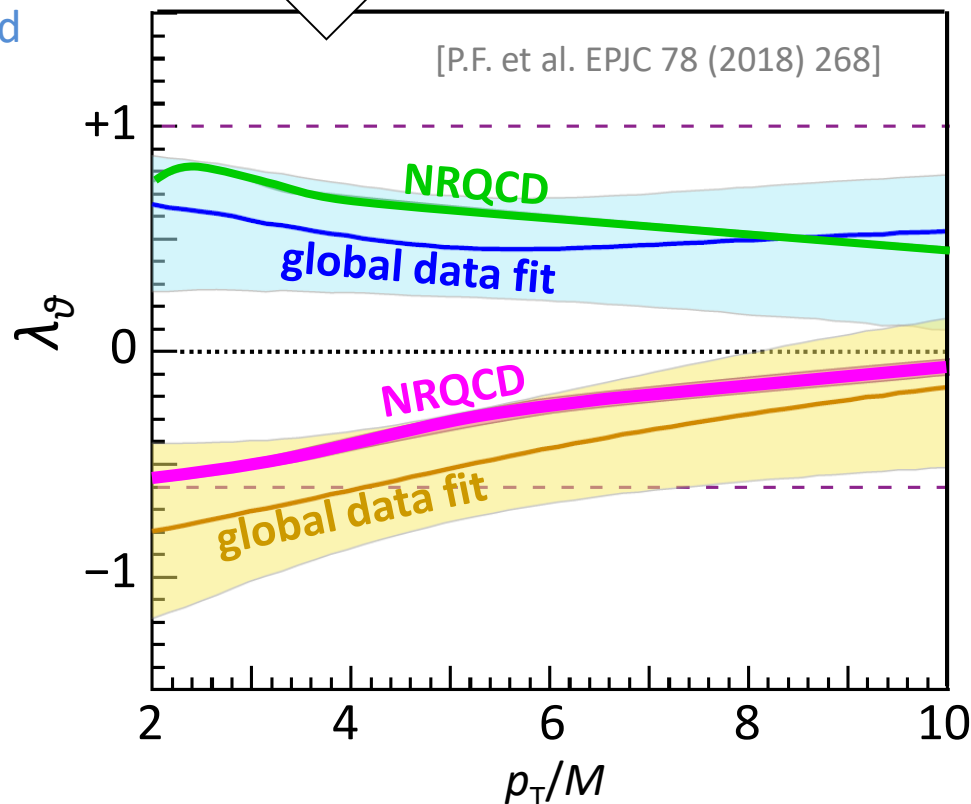
3) $\lambda_{\vartheta}(\chi_{c2})$

$$r \equiv m_c^2 \left\langle \mathcal{O}^{\chi_{c0}}({}^3S_1^{[8]}) \right\rangle / \left\langle \mathcal{O}^{\chi_{c0}}({}^3P_0^{[1]}) \right\rangle$$

= 0.217 ± 0.003 from the CMS + ATLAS

χ_{c2} / χ_{c1} yield ratio (averaged)

A strongly
constrained and
unambiguous
prediction, not
requiring any
“fine-tuning” ...



$\Leftarrow \chi_{c1}$... and
perfectly
agreeing
with data

$\Leftarrow \chi_{c2}$

An out-of-the-box
success of NRQCD!

Part 2

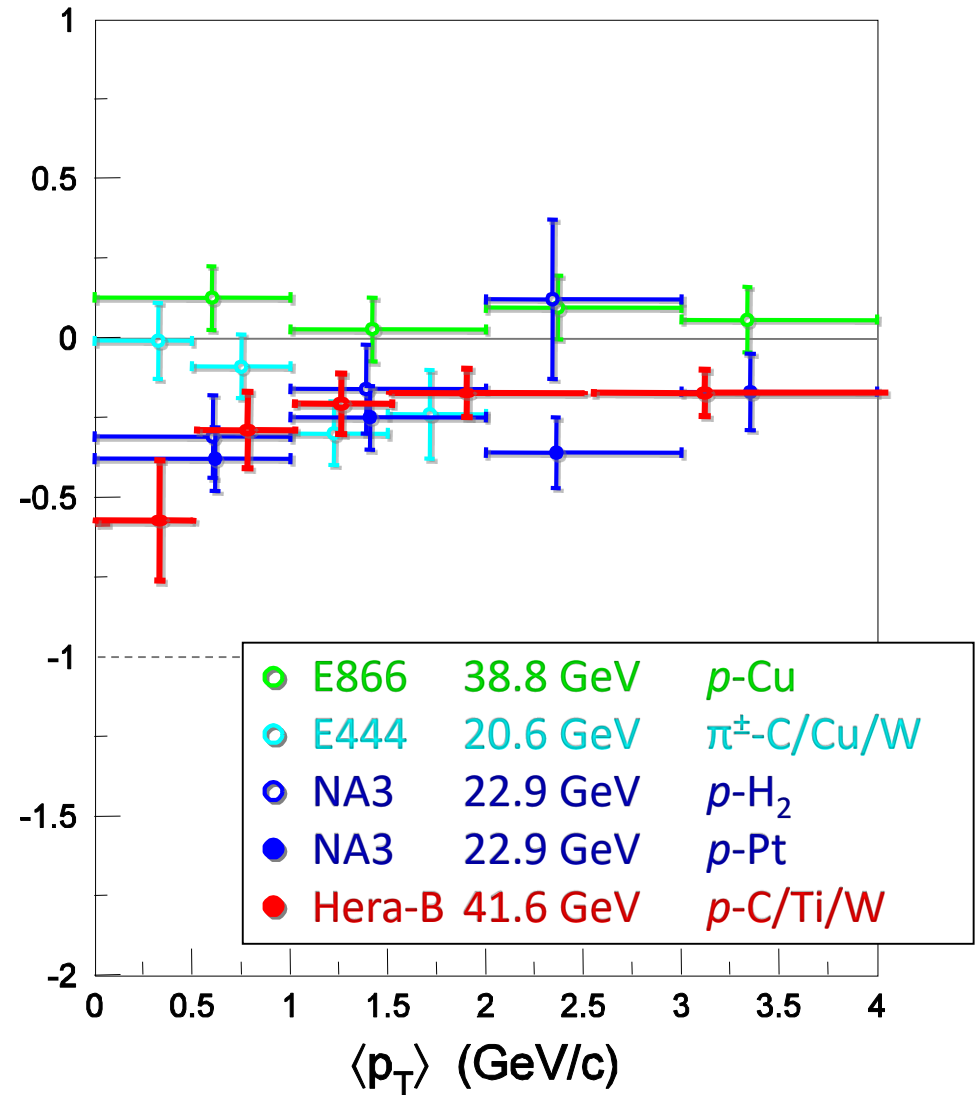
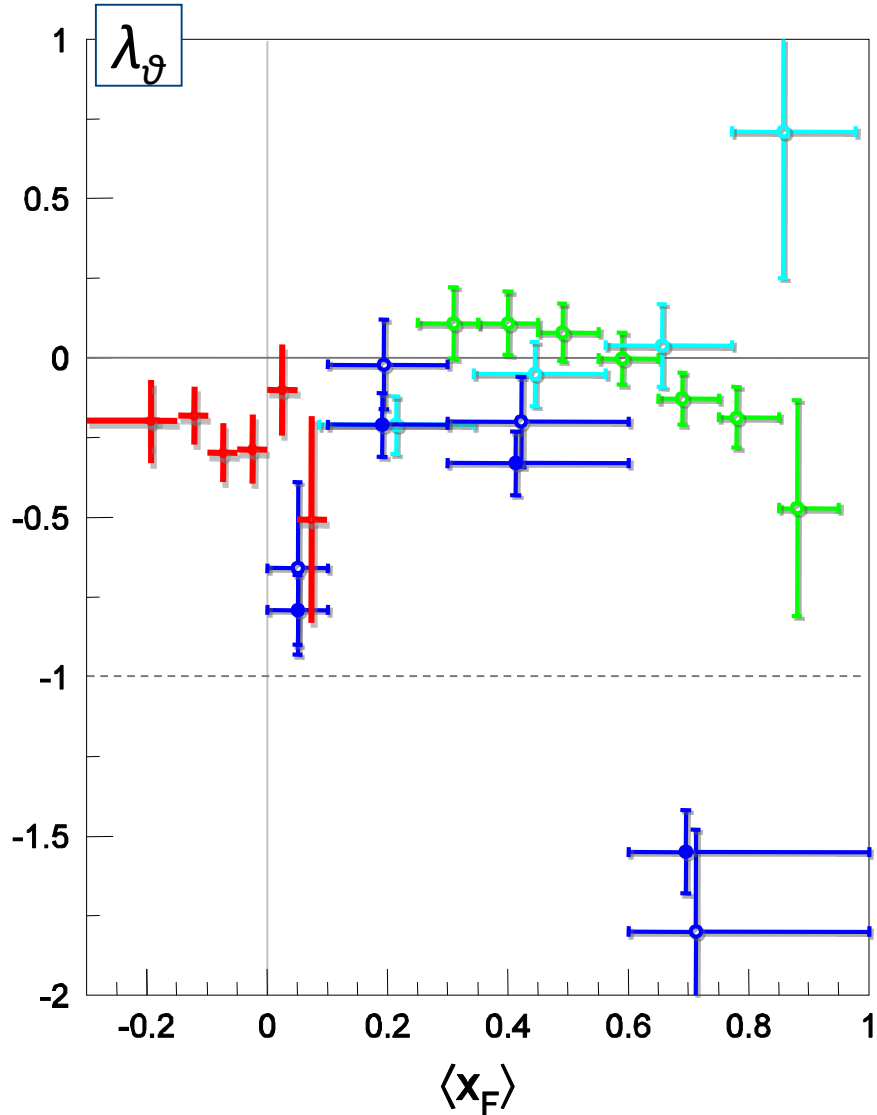
The puzzle of low- p_T (fixed target) data

Past fixed-target experiments provided J/ψ and Υ polarization measurements with different beams and targets, different energies, and in three different reference frames, as functions of p_T and x_F .

They form a very perplexing picture...

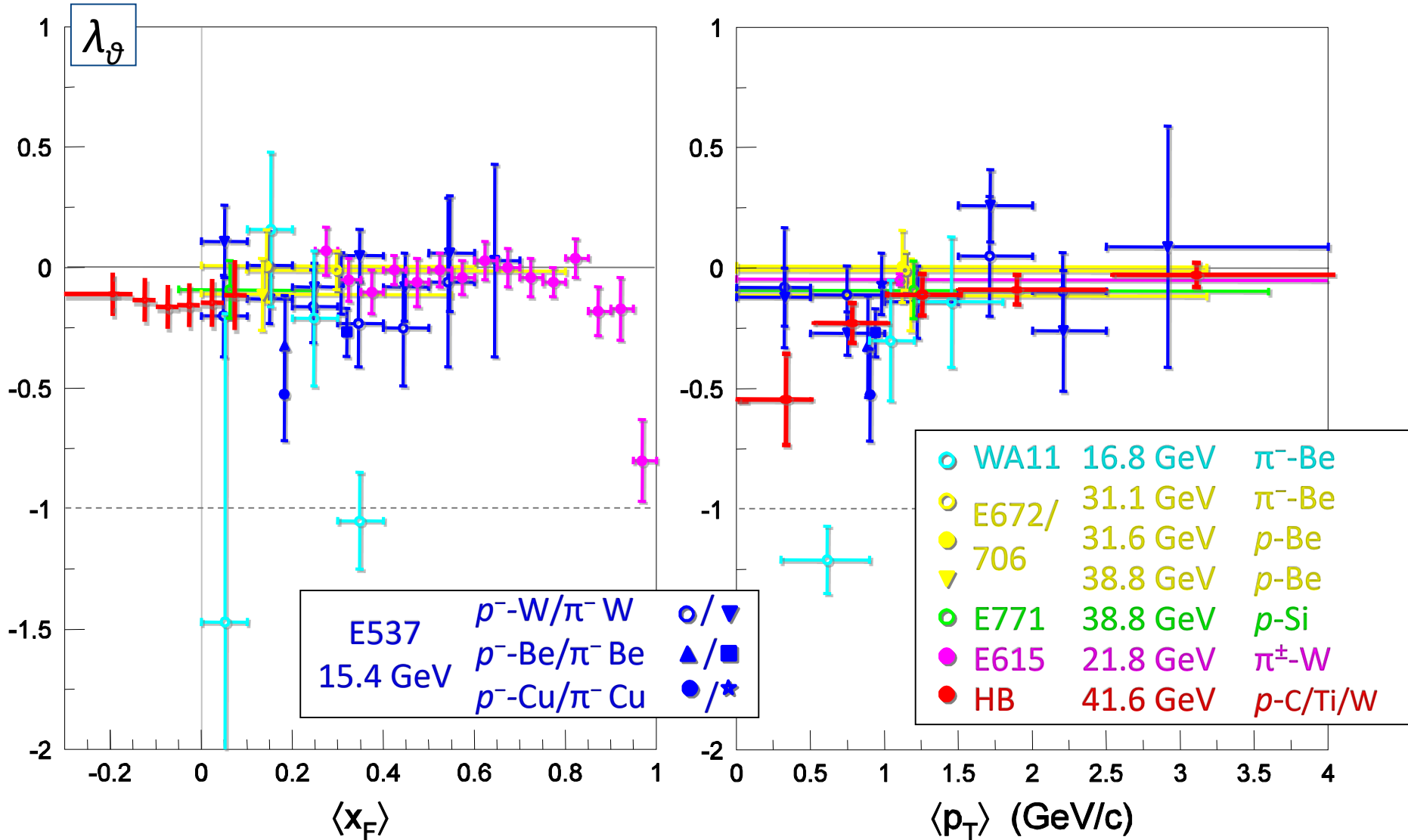
J/ ψ polarization in the CS frame

Collins-Soper



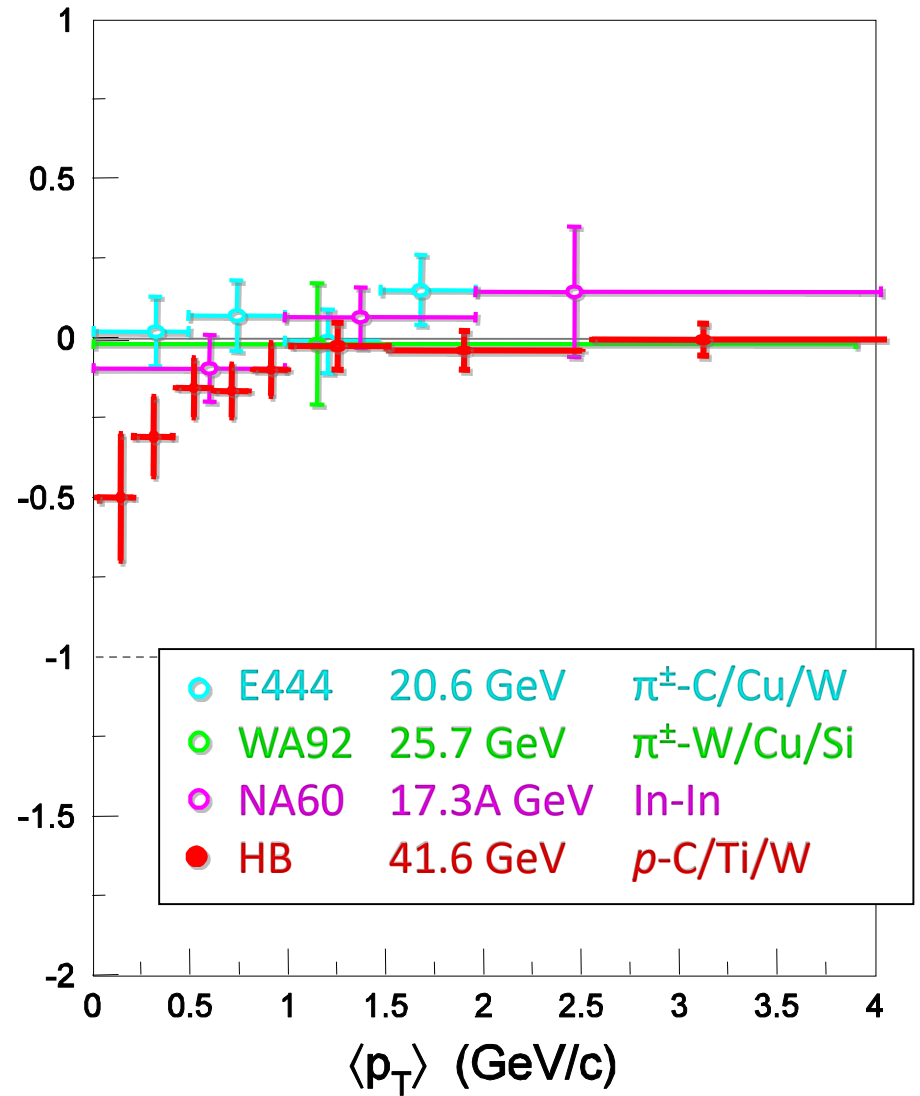
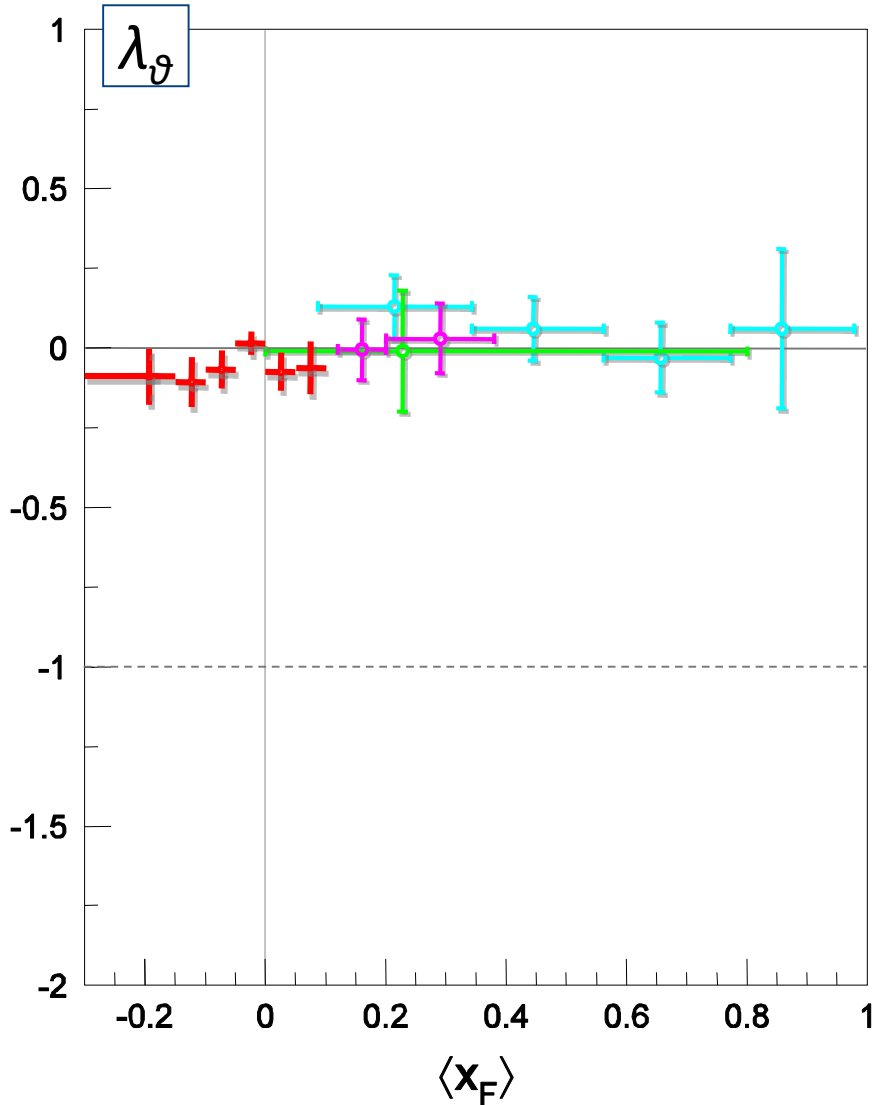
J/ψ polarization in the GJ frame

Gottfried-Jackson

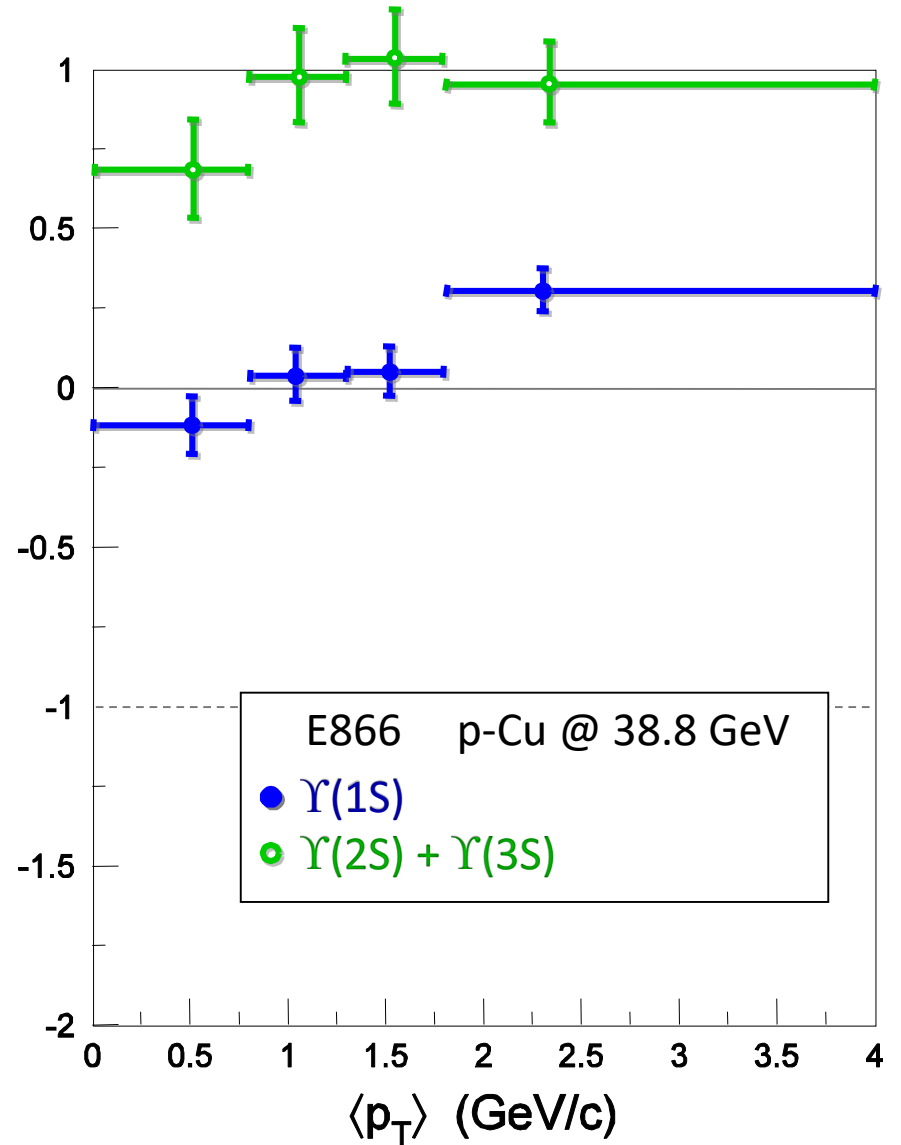
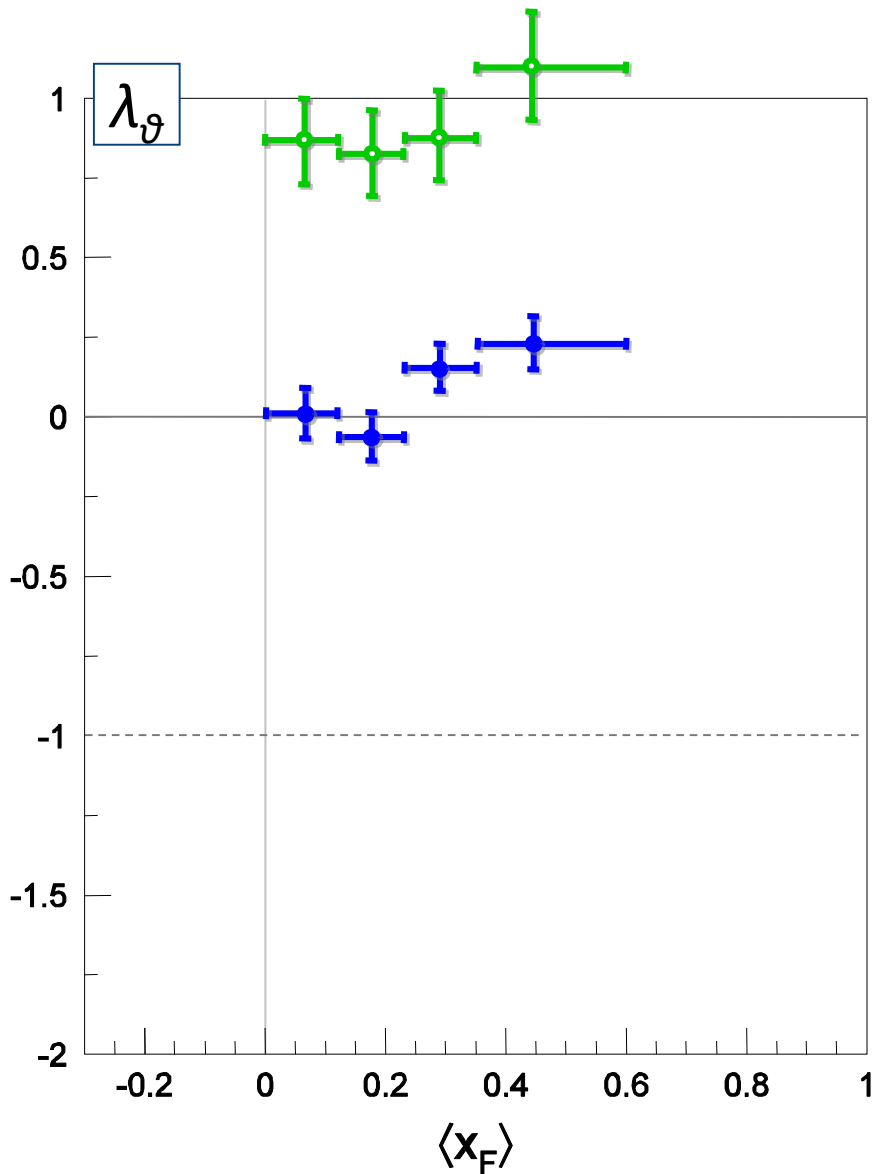


J/ ψ polarization in the HX frame

centre-of-mass helicity



Υ Polarization in the CS frame



Picture to be observed “with a grain of salt”:

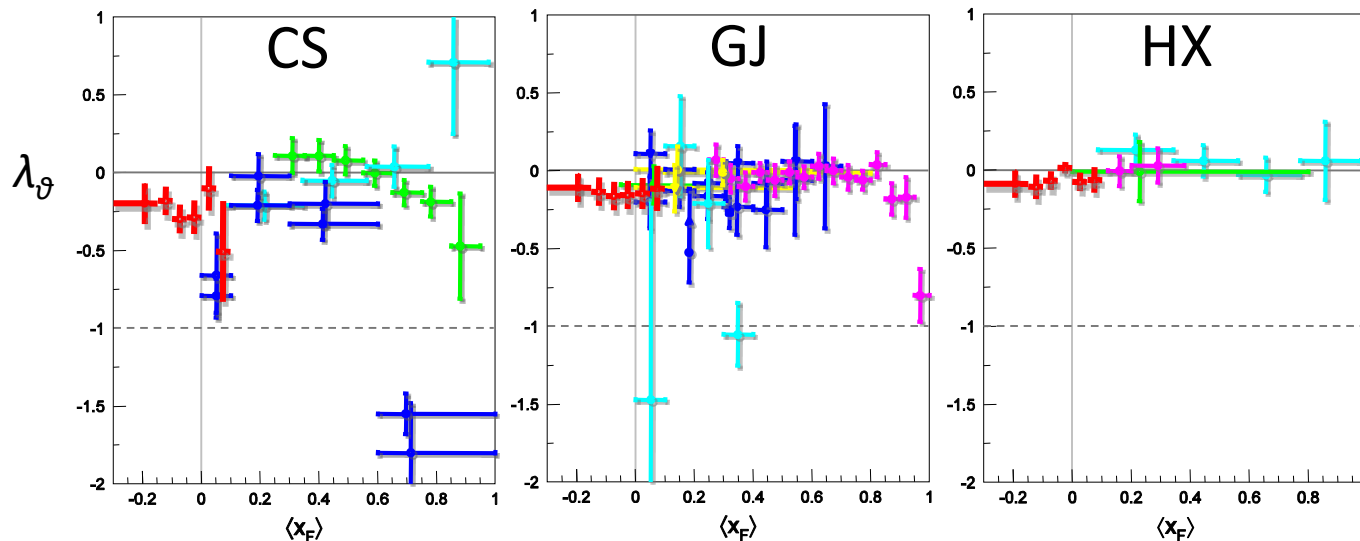
- most of these measurements were obtained from 1D analyses (with risks discussed in [P. Faccioli, Mod. Phys. Lett. A Vol. 27, 1230022 (2012)])
- for some of them systematic uncertainties were never evaluated
- some of them exhibit suspicious fluctuations, even reaching unphysical values
- we are mixing different energies and target nuclei (nuclear effects may exist)

Nevertheless, we can see some indications...

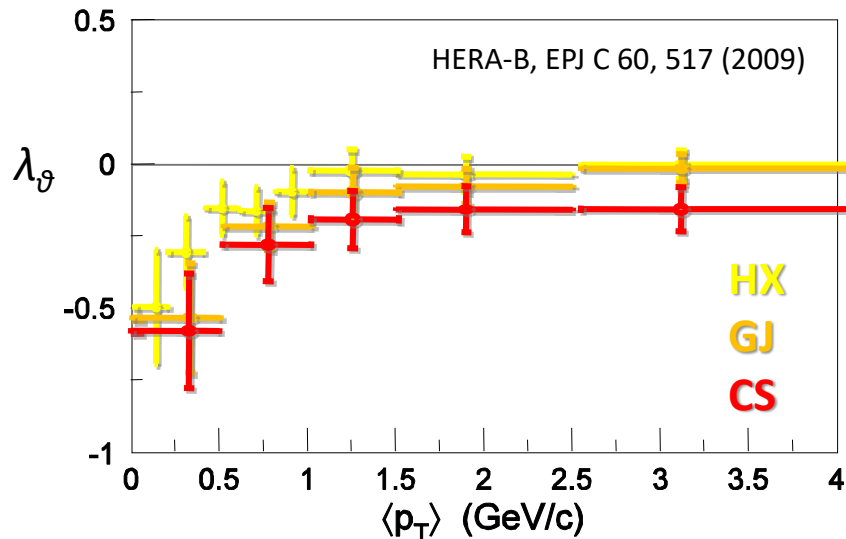
0) Obvious: new, better measurements are welcome

- Looking at polar and azimuthal components
- Using invariant polarization as check
- Possibly disentangling feed-down components: $\psi(2S)$ and/or χ_c

1) CS/GJ > HX hierarchy: dominance of 2 → 1



λ_g looks “flat”
in the HX frame



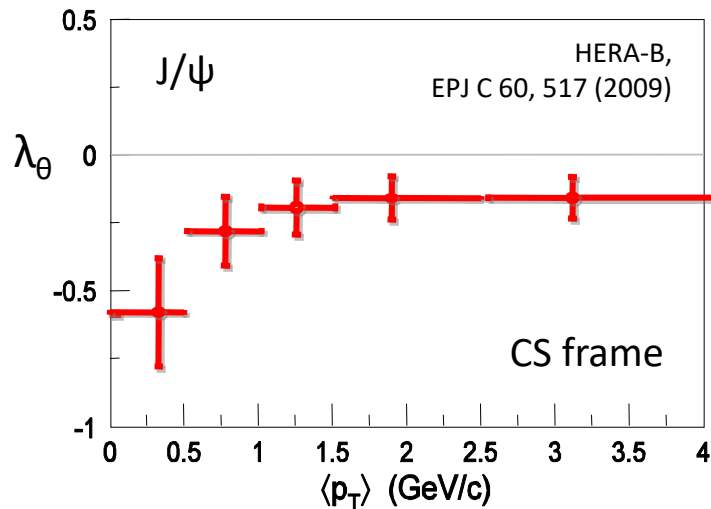
Hierarchy clearly seen by HERA-B
in the three frames (uncertainties $\sim 100\%$ corr.)

→ CS (direction of colliding partons) gives the
“optimal” observation

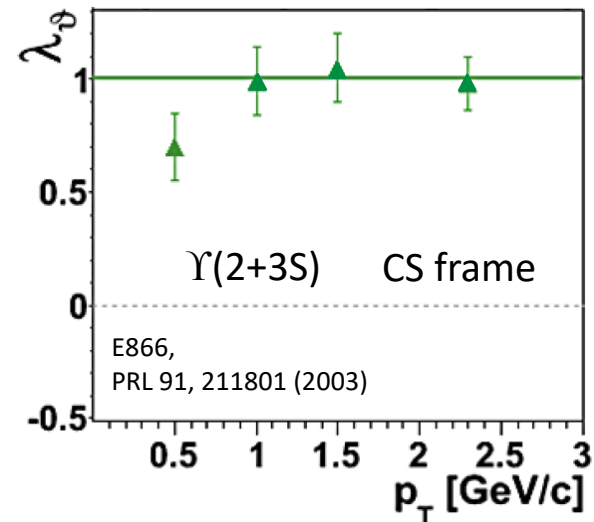
→ probable dominance of
2→1 q - q bar / g - g → QQ bar processes,
where the QQ bar is strongly polarized,
directly inheriting the angular momentum
state of the system of colliding partons:
we see the **partons’ natural polarizations**

2) Smearing with increasing p_T : importance of k_T effects

The J/ψ polarization *magnitude*,



but not the Υ one,



seems to *decrease quickly with increasing* p_T .

In fact the J/ψ measurement reaches higher p_T/M values than the Υ one. Does this mean that in J/ψ production, but not in Υ production, we start seeing $2 \rightarrow 2$ processes “smearing” the polarization?

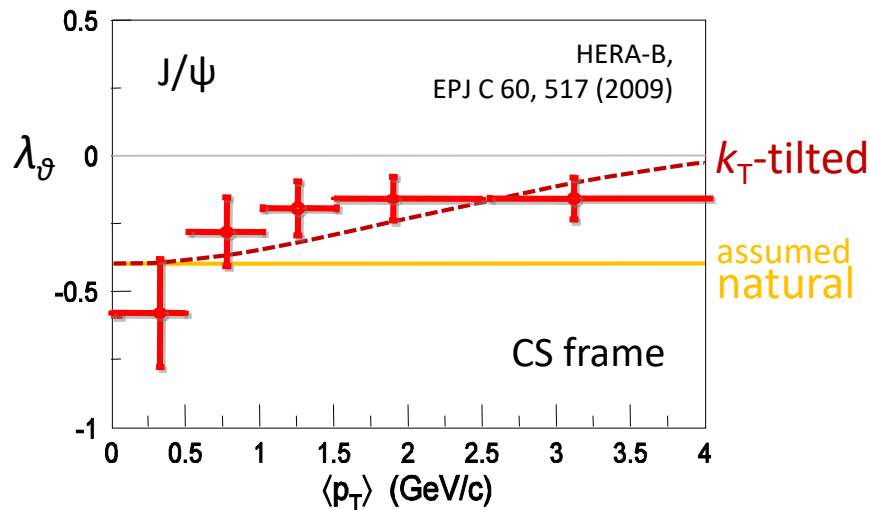
Or, what about the parton k_T ?

2) Smearing with increasing p_T : importance of k_T effects

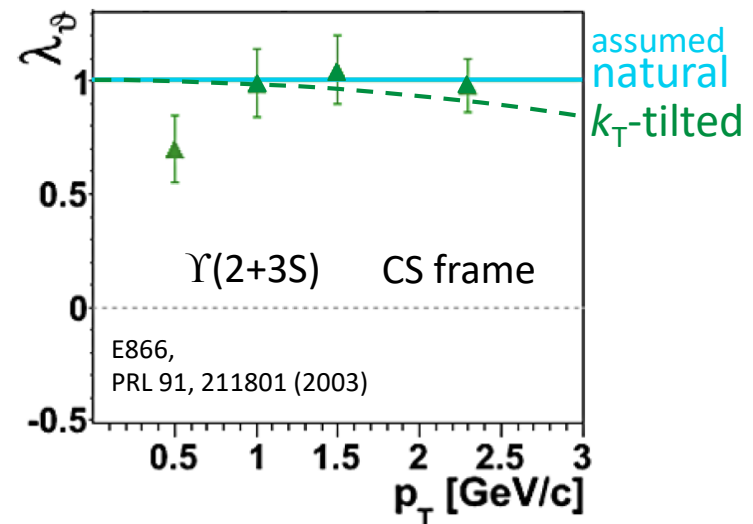
The intrinsic transverse momenta of the partons cause an *event-by-event tilt* between the “natural” polarization axis (relative direction of the colliding partons), and the polarization axis used in the experimental analysis (CS).

The tilt angle δ satisfies $\sin^2 \delta \approx \frac{2k_T^2}{m_T^2} \approx \frac{p_T^2}{M^2 + p_T^2}$ [P. F. et al., EPJC 69, 657 (2010)]

This description approximately accounts for the p_T dependence observed for the $c\bar{c}$...

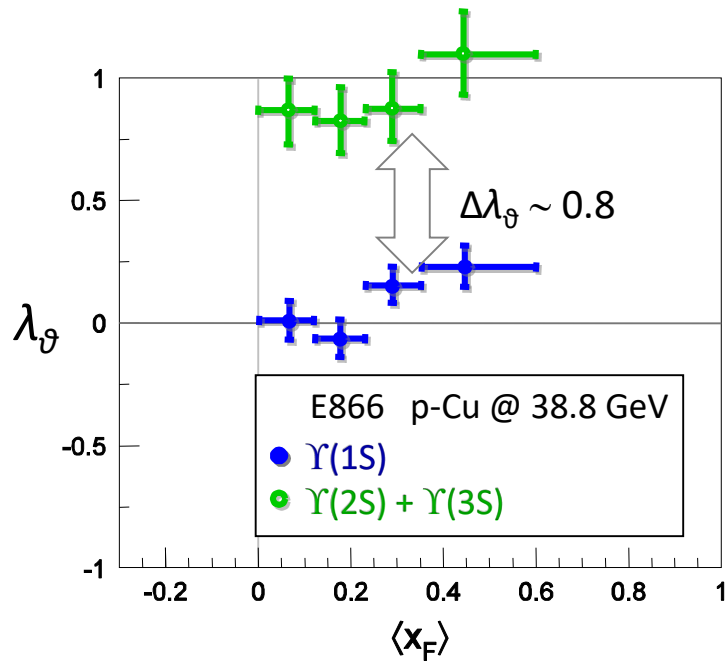


...and for the lack of a corresponding observation in the $b\bar{b}$ case



The $p_T \rightarrow 0$ limit gives the most interesting (unsmearred) polarization measurement!

3) E866's Υ puzzle: importance of the χ feed-down

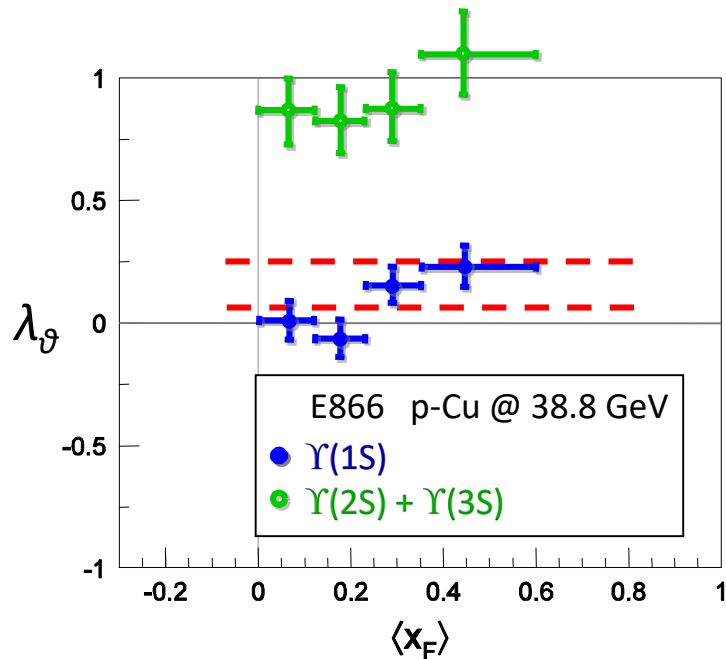


1S, 2S and 3S states should have about the same polarization when **directly produced** (or when coming from heavier Υ) $\rightarrow \lambda_\theta \approx +1$

To justify the large difference between 2-3S and 1S, we must assume that $\chi_1 + \chi_2$ feed-down:

- is **negligible for 2-3S states** and **large for 1S**
- tends to be **longitudinal**

3) E866's Υ puzzle: importance of the χ feed-down



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If...

- **50–60%** of the $\Upsilon(1S)$ come from χ_b (only an example: they have in general different polarizations, as seen for the χ_c at high p_T)
- χ_b states are produced, e.g., with $J_z = \pm 1 \rightarrow \lambda_\theta(\chi_{b1}) = \lambda_\theta(\chi_{b2}) = -1/3$
- \Rightarrow the observed $\Upsilon(1S)$ would have λ_θ in the range **$1/13 - 1/4 = 0.08 - 0.25$**

χ production is **not a second-order correction** for J/ ψ and Υ yields and polarizations!

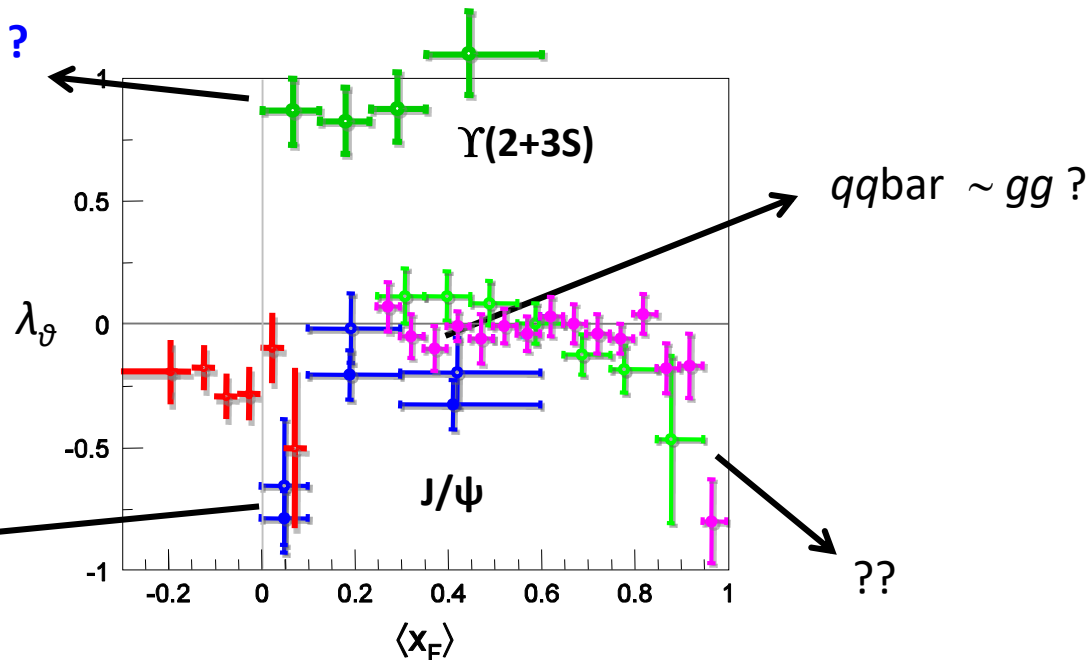
4) Strong x_F (and mass) dependence: $q\text{-}q\text{bar}$ vs gg ?

A trend may be recognized in the perplexing scenario of J/ψ and Υ polarizations vs $x_F = x_1 - x_2$ when we correlate the observed **longitudinal** polarizations with the dominance of $gg \rightarrow QQ\text{bar}$ processes and **transverse** polarizations with the dominance of $q\text{qbar} \rightarrow QQ\text{bar}$ processes

$q\text{qbar}_{\Upsilon(ns)} \gg q\text{qbar}_{J/\psi}$?

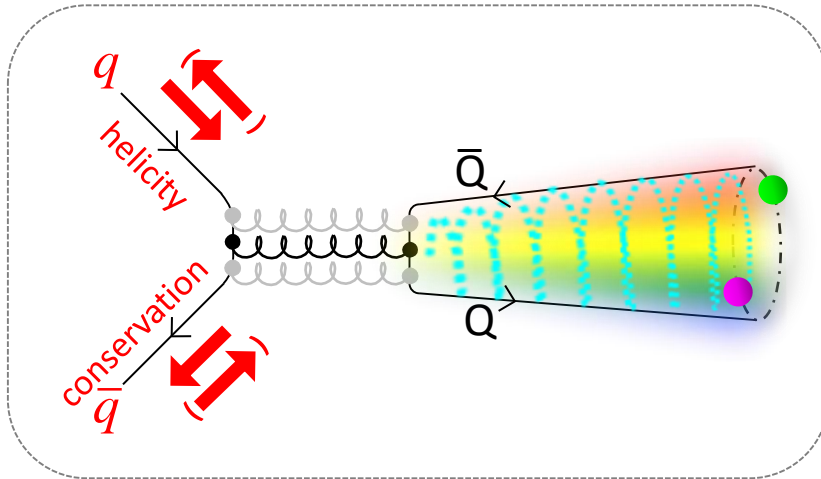
$$\frac{\langle X_1 \cdot X_2 \rangle_{\Upsilon(ns)}}{\langle X_1 \cdot X_2 \rangle_{J/\psi}} \sim 10$$

$x_1 \cong x_2 = \mathcal{O}(0.1)$
gg dominance?

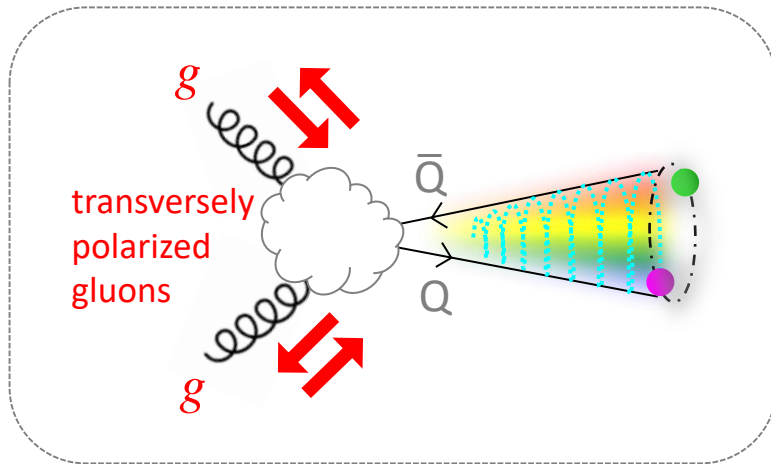
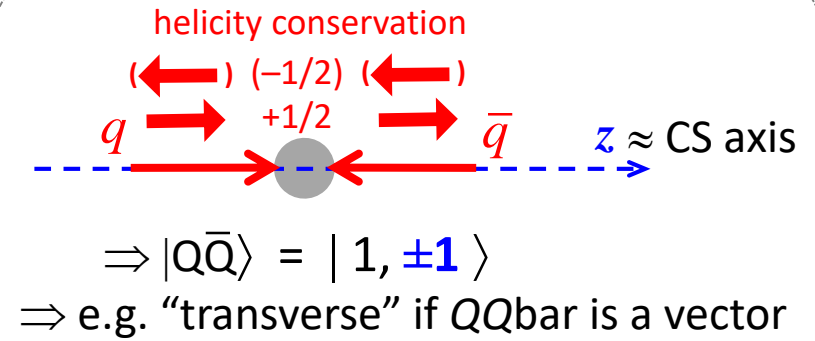


Comparing measurements with predictions for the gg and $q\text{qbar}$ cases can probe the identity of the colliding partons
 \rightarrow use polarization vs x_F as further constraint on gluon distribution!

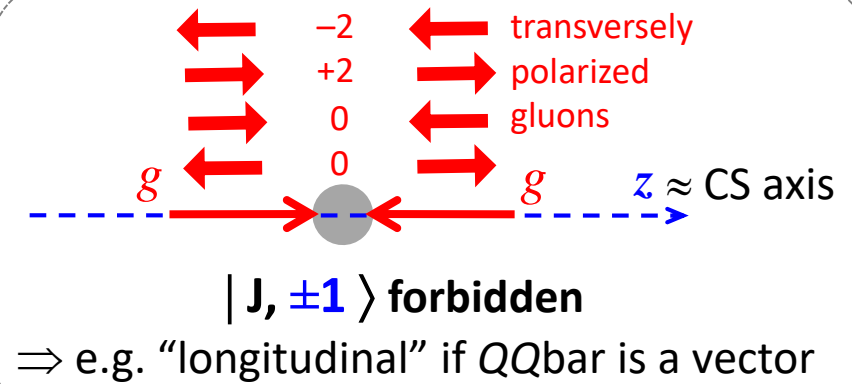
$(gg \text{ vs } qq\bar{q}) \leftrightarrow (\text{longitudinal vs transverse}) ?$



$q\bar{q}$ rest frame = $Q\bar{Q}$ rest frame



g - g rest frame = $Q\bar{Q}$ rest frame



(observed quarkonium polarization will depend on the J^P of the intermediate $Q\bar{Q}$)

Summary

High p_T

Zero polarization for the J/ψ , given that it is a vector (=intrinsically polarized) particle, is an emblematic manifestation of its peculiar production mechanism.

The agreement with NRQCD requires a specific parameter tuning, possibly pointing to the existence of a simpler (more natural) hierarchy of processes.

More precise measurements are needed to assess whether the polarization always remains zero and flat vs p_T .

Low p_T

The puzzling scenario of existing fixed-target data contains interesting indications:

- dominance of $2 \rightarrow 1$ processes
- importance of parton- k_T effects
- necessity to discriminate direct production and χ -state feed-down contributions
- maximal difference between polarizations in q - q bar and gg production:
an opportunity to improve gluon PDF determination?