

CRITICAL PHENOMENA NEAR THE QCD PHASE TRANSITION AT FINITE DENSITY

BORIS O. KERBIKOV ^{1,2,3}

¹Lebedev Physical Institute

²Moscow Institute for Physics and Technology

³Alikhanov Institute for Theoretical and Experimental Physics

*RFBR grant 18-02-40054 (V.V. Kulikov)

**20/10/20, ed2

Dubna, Russia

20 October – 23 October, 2020

PLAN OF THE TALK

- ① Pre-critical Fluctuations at Finite Density
- ② Fluctuation Driven Anomalies and Instabilities
 - ① Bulk viscosity critical behavior and related instabilities
 - ② Anomalous sound attenuation
 - ③ Electrical conductivity growth and soft photon emission excess
- ③ Negative Bulk Viscosity and Rayleigh sound instability

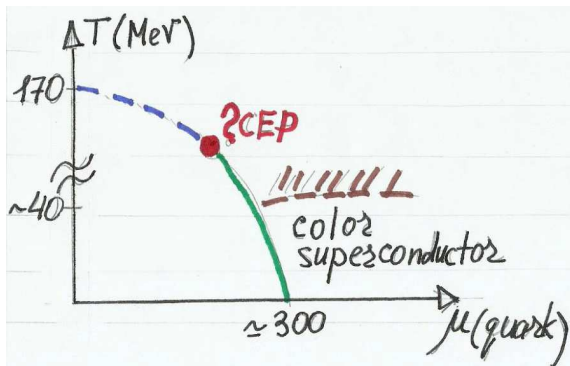


Figure 1

Physics depends on the (T, μ) values – the QCD phase diagram
 (“theorist’s science fiction” – Owe Philipsen)

Fluctuation region: $\mu_q \sim 300\text{--}400$ MeV, $T \rightarrow T_c (\sim 40)$ MeV

Pre-critical fluctuations are inherent to the II-nd or weak I-st order transitions. Wide fluctuation region:

$$Gi = \frac{\delta T}{T_c} \sim 10^{-4} - 10^{-2} \text{ vs } 10^{-12} \text{ for BCS}$$

The field of fluctuating quark pairs. Slow relaxation collective mode with a propagator

$$\hat{L} = -\frac{1}{\nu} \left[t + \frac{\pi}{8T_c} (-i\omega + Dq^2) \right]^{-1} \quad \text{as} \quad t = \frac{T - T_c}{T_c}$$

The basic dynamical quantity is $\Pi_{\mu\nu}^R$ –
the retarded response operator, or Wightman correlator, or polarization
tensor

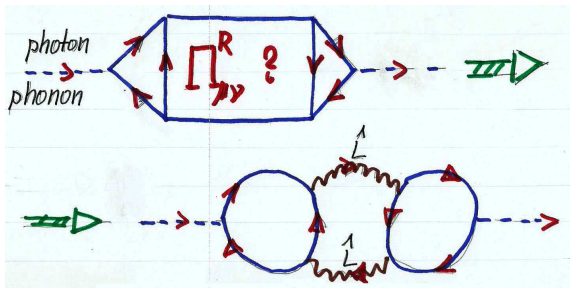


Figure 2

The famous Aslamazov-Larkin diagram

(1.) BULK VISCOSITY: CRITICAL BEHAVIOR

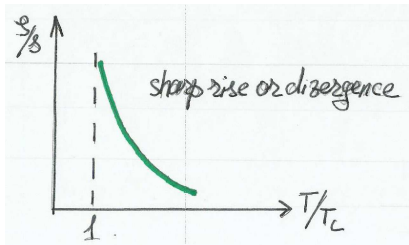


Figure 3

$$T \gg T_c \rightarrow \eta/\zeta \gg 1, \text{ conformal limit: } T \sim T_c \rightarrow \eta/\zeta < 1$$
$$P = P_{equil} - \zeta \vec{\nabla} \mathbf{v}, \text{ where } \mathbf{v} \text{ is flow velocity}$$

Ising-like universality,
(4- ϵ) renormalization,
modes interaction hydro, etc.

→

$$\zeta \sim t^{z/\nu+\alpha} \simeq t^{-1.7} \text{ vs AL: } \zeta \sim t^{-3/2},$$

where $t = \frac{T-T_c}{T_c}$.

(2.) BULK VISCOSITY: CONTINUATION

Bulk viscosity as the instabilities driving force

- Fragmentation into droplets, spinodal decomposition (I. Mishustin et al.)
- Negative pressure and cavitation
- Anomalous sound attenuation (Mandelstam-Leontovich, Landau VI Hydro, B.K.)

$$A(t) = A_0 e^{-\gamma t} = A_0 e^{\frac{\alpha \lambda}{\chi} x}, \text{ or } T_{\mu\nu}(t) = \exp\left(-\frac{\zeta}{2sT} k^2 t\right)$$

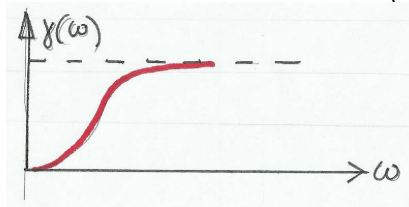


Figure 4

(3.) ELECTRICAL CONDUCTIVITY GROWTH
SOFT PHOTON EMISSION EXCESS

$$\sigma_{el} = -\frac{1}{\omega_k} \text{Im} \Pi_{\mu\mu}^R, \quad \sigma_{AL} \sim t^{-1/2}$$

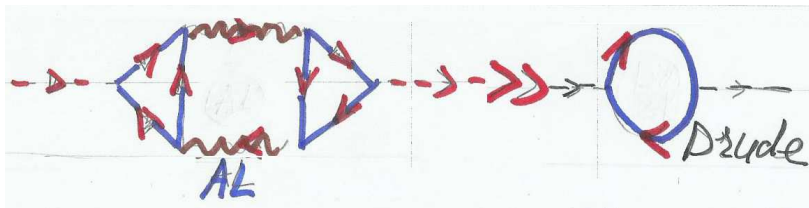


Figure 5: AL $\gg$ Drude

(4.) SOFT PHOTON EMISSION – THE SAME DIAGRAM

$$\omega \frac{dR_\gamma}{d^3p} = -\frac{1}{(2\pi)^3} \frac{1}{e^{\omega/T} - 1} \text{Im} \Pi_{\mu\mu}^R \Rightarrow \text{soft} \Rightarrow \omega \ll T \Rightarrow$$

$$\Rightarrow \omega \frac{dR_\gamma}{d^3p} \simeq \frac{2}{(2\pi)^3} T \sigma_{el} \Rightarrow \text{entrancement at } T \rightarrow T_c$$

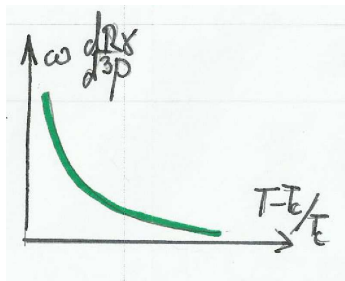


Figure 6

(III.) Negative Bulk Viscosity and Rayleigh Sound Instability (9/9)

$\zeta > 0 \Rightarrow$ sound amplification instead of attenuation

$\zeta < 0$ – how can it be?

$$\text{Landau } \gamma = \frac{\omega^2}{2\zeta C_1^3} \left[\left(\frac{4}{3}\eta + \zeta \right) + \kappa \left(\frac{1}{C_v} - \frac{1}{C_p} \right) \right] \Rightarrow \frac{\omega^2}{2\varepsilon C_s^3} \zeta,$$

$\omega \ll 1/\tau$, τ – relaxation time

$$\zeta(\omega) = \frac{\varepsilon\tau(C_\infty^2 - C_0^2)}{1 + \omega^2\tau^2}$$

$C_0^2 = \left(\frac{\partial p}{\partial \delta} \right)_s$ – the ordinary hydrodynamics speed of sound

$C_\infty^2 = \left(\frac{\partial p}{\partial \delta} \right)_\varphi$ – the speed of sound at a given value of the order parameter

Landau $C_\infty^2 > C_0^2 \Rightarrow \zeta > 0$. **Not always true out of equilibrium.**