

# Holographic Model for Light Quarks in Anisotropic Background

**K.A. Rannu**

Peoples Friendship University of Russia (RUDN)  
Steklov Mathematical Institute (MI RAS)



**RFBR Grants for NICA**

*Dubna*

23.10.2020



RFBR grant №18-02-40069 mega

with I.Ya. Aref'eva and P. Slepov  
*arXiv:2009.05562 [hep-th]*

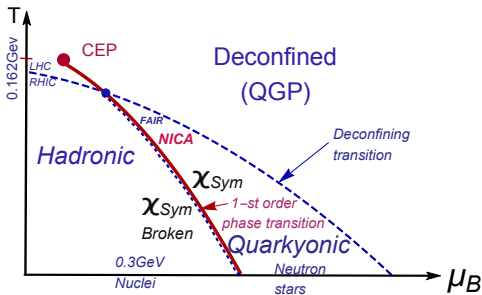
# Motivation

Goal of Holographic QCD — describe QCD phase diagram

## Requirements:

- reproduce the QCD results from perturbative theory at short distances,
- reproduce Lattice QCD results at large distances ( $\sim 1$  fm) and small density.

I.A. talk, Wednesday



## Action and metric

$$\mathcal{S} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left[ R - \frac{f_1(\phi)}{4} F_{(1)}^2 - \frac{f_2(\phi)}{4} F_{(2)}^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

$$F_{\mu\nu}^{(1)} = \partial_\mu A_\nu - \partial_\nu A_\mu, \quad A_\mu^{(1)} = A_t(z) \delta_\mu^0 \quad F_{\mu\nu}^{(2)} = q \, dy^1 \wedge dy^2, \quad F_{23}^{(2)} = q$$

$$ds^2 = \frac{L^2}{z^2} \mathbf{b}(z) \left[ -g(z) dt^2 + dx^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dy_1^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dy_2^2 + \frac{dz^2}{g(z)} \right]$$

## Action and metric

$$\mathcal{S} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left[ R - \frac{f_1(\phi)}{4} F_{(1)}^2 - \frac{f_2(\phi)}{4} F_{(2)}^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

$$F_{\mu\nu}^{(1)} = \partial_\mu A_\nu - \partial_\nu A_\mu, \quad A_\mu^{(1)} = A_t(z) \delta_\mu^0 \quad F_{\mu\nu}^{(2)} = q dy^1 \wedge dy^2, \quad F_{23}^{(2)} = q$$

$$ds^2 = \frac{L^2}{z^2} \mathbf{b}(z) \left[ -g(z) dt^2 + dx^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dy_1^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dy_2^2 + \frac{dz^2}{g(z)} \right]$$

I.A., K.R. JHEP **1805** 206 (2018)

## Action and metric

$$S = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left[ R - \frac{f_1(\phi)}{4} F_{(1)}^2 - \frac{f_2(\phi)}{4} F_{(2)}^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

$$F_{\mu\nu}^{(1)} = \partial_\mu A_\nu - \partial_\nu A_\mu, \quad A_\mu^{(1)} = A_t(z) \delta_\mu^0 \quad F_{\mu\nu}^{(2)} = q dy^1 \wedge dy^2, \quad F_{23}^{(2)} = q$$

$$ds^2 = \frac{L^2}{z^2} \mathbf{b}(z) \left[ -g(z) dt^2 + dx^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dy_1^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dy_2^2 + \frac{dz^2}{g(z)} \right]$$

I.A., K.R. JHEP **1805** 206 (2018)

$\mathbf{b}(z) = e^{2\mathcal{A}(z)}$  — quarks mass

## Action and metric

$$S = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left[ R - \frac{f_1(\phi)}{4} F_{(1)}^2 - \frac{f_2(\phi)}{4} F_{(2)}^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

$$F_{\mu\nu}^{(1)} = \partial_\mu A_\nu - \partial_\nu A_\mu, \quad A_\mu^{(1)} = A_t(z) \delta_\mu^0 \quad F_{\mu\nu}^{(2)} = q dy^1 \wedge dy^2, \quad F_{23}^{(2)} = q$$

$$ds^2 = \frac{L^2}{z^2} \mathbf{b}(z) \left[ -g(z) dt^2 + dx^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dy_1^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dy_2^2 + \frac{dz^2}{g(z)} \right]$$

I.A., K.R. JHEP **1805** 206 (2018)

$\mathbf{b}(z) = e^{2\mathcal{A}(z)}$  — quarks mass

$\mathcal{A}(z) = -cz^2/4 \rightarrow$  heavy quarks (b, t) [Andreev, Zakharov \(2006\)](#)

## Action and metric

$$\mathcal{S} = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left[ R - \frac{f_1(\phi)}{4} F_{(1)}^2 - \frac{f_2(\phi)}{4} F_{(2)}^2 - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

$$F_{\mu\nu}^{(1)} = \partial_\mu A_\nu - \partial_\nu A_\mu, \quad A_\mu^{(1)} = A_t(z) \delta_\mu^0 \quad F_{\mu\nu}^{(2)} = q \, dy^1 \wedge dy^2, \quad F_{23}^{(2)} = q$$

$$ds^2 = \frac{L^2}{z^2} \mathbf{b}(z) \left[ -g(z) dt^2 + dx^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dy_1^2 + \left( \frac{z}{L} \right)^{2-\frac{2}{\nu}} dy_2^2 + \frac{dz^2}{g(z)} \right]$$

I.A., K.R. JHEP **1805** 206 (2018)

$\mathbf{b}(z) = e^{2\mathcal{A}(z)}$  — quarks mass

$\mathcal{A}(z) = -cz^2/4 \rightarrow$  heavy quarks (b, t) [Andreev, Zakharov \(2006\)](#)

$\mathcal{A}(z) = -a \ln(bz^2 + 1) \rightarrow$  light quarks (d, u) [Li, Yang, Yuan \(2017\)](#)

# Solution for anisotropic metric ansatz

$$\mathcal{A}(z) = -a \ln(bz^2 + 1) \quad f_1 = e^{-cz^2 - \mathcal{A}(z)} z^{-2 + \frac{2}{\nu}}$$

*Boundary conditions:*

$$A_t(0) = \mu, \quad A_t(z_h) = 0, \quad g(0) = 1, \quad g(z_h) = 0, \quad \phi(z_0) = 0$$



# Solution for anisotropic metric ansatz

$$\mathcal{A}(z) = -a \ln(bz^2 + 1) \quad f_1 = e^{-cz^2 - \mathcal{A}(z)} z^{-2 + \frac{2}{\nu}}$$

*Boundary conditions:*

$$A_t(0) = \mu, \quad A_t(z_h) = 0, \quad g(0) = 1, \quad g(z_h) = 0, \quad \phi(z_0) = 0$$

next talk by P.S.

# Solution for anisotropic metric ansatz

$$\mathcal{A}(z) = -a \ln(bz^2 + 1) \quad f_1 = e^{-cz^2 - \mathcal{A}(z)} z^{-2 + \frac{2}{\nu}}$$

*Boundary conditions:*

$$A_t(0) = \mu, \quad A_t(z_h) = 0, \quad g(0) = 1, \quad g(z_h) = 0, \quad \phi(z_0) = 0$$

next talk by P.S.

$$A_t = \mu \frac{e^{cz^2} - e^{cz_h^2}}{1 - e^{cz_h^2}}$$

$$g = 1 - \frac{\int_0^z (1 + b\xi^2)^{3a} \xi^{1 + \frac{2}{\nu}} d\xi}{\int_0^{z_h} (1 + b\xi^2)^{3a} \xi^{1 + \frac{2}{\nu}} d\xi} + \frac{2\mu^2 c}{L^2 (1 - e^{cz_h^2})^2} \int_0^z e^{c\xi^2} (1 + b\xi^2)^{3a} \xi^{1 + \frac{2}{\nu}} d\xi \times$$
$$\times \left[ 1 - \frac{\int_0^z (1 + b\xi^2)^{3a} \xi^{1 + \frac{2}{\nu}} d\xi}{\int_0^{z_h} (1 + b\xi^2)^{3a} \xi^{1 + \frac{2}{\nu}} d\xi} \frac{\int_0^{z_h} e^{c\xi^2} (1 + b\xi^2)^{3a} \xi^{1 + \frac{2}{\nu}} d\xi}{\int_0^z e^{c\xi^2} (1 + b\xi^2)^{3a} \xi^{1 + \frac{2}{\nu}} d\xi} \right].$$

## Solution for anisotropic metric ansatz

$$\phi = \int_{z_0}^z \frac{2\sqrt{\nu - 1 + (2(\nu - 1) + 9a\nu^2)b\xi^2 + (\nu - 1 + 3a(1 + 2a)\nu^2)b^2\xi^4}}{(1 + b\xi^2)\nu\xi} d\xi$$

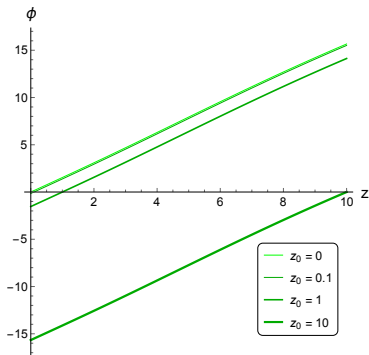
$$\nu > 1 : z_0 \rightarrow 0 \rightarrow \phi(z) \sim \int_0^z dz/z$$

# Solution for anisotropic metric ansatz

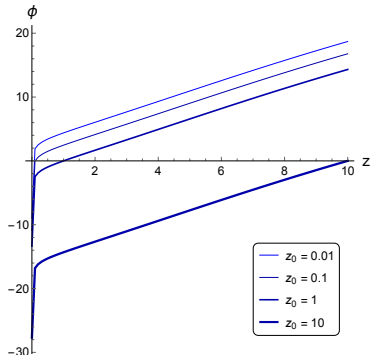
$$\phi = \int_{z_0}^z \frac{2\sqrt{\nu - 1 + (2(\nu - 1) + 9a\nu^2)b\xi^2 + (\nu - 1 + 3a(1 + 2a)\nu^2)b^2\xi^4}}{(1 + b\xi^2)\nu\xi} d\xi$$

$$\nu > 1: z_0 \rightarrow 0 \rightarrow \phi(z) \sim \int_0^z dz/z$$

$\nu = 1$



$\nu = 4.5$



# Thermodynamics

Black hole temperature  $\longleftrightarrow$  QGP temperature (Maldacena)

$$T = \frac{|g'|}{4\pi} \Big|_{z=z_h} = \frac{1}{4\pi} \left[ - \frac{(1+bz_h^2)^{3a} z_h^{1+\frac{2}{\nu}}}{\int_0^{z_h} (1+b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi} \left[ 1 - \frac{2\mu^2 c e^{2cz_h^2}}{L^2 (1-e^{cz_h^2})^2} \times \right. \right. \\ \left. \left. \times \left( 1 - e^{-cz_h^2} \frac{\int_0^{z_h} e^{c\xi^2} (1+b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi}{\int_0^{z_h} (1+b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi} \right) \int_0^{z_h} (1+b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi \right] \right]$$

# Thermodynamics

**Black hole temperature**  $\longleftrightarrow$  **QGP temperature (Maldacena)**

$$T = \left. \frac{|g'|}{4\pi} \right|_{z=z_h} = \frac{1}{4\pi} \left[ - \frac{(1 + bz_h^2)^{3a} z_h^{1+\frac{2}{\nu}}}{\int_0^{z_h} (1 + b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi} \left[ 1 - \frac{2\mu^2 c e^{2cz_h^2}}{L^2 (1 - e^{cz_h^2})^2} \times \right. \right. \\ \left. \left. \times \left( 1 - e^{-cz_h^2} \frac{\int_0^{z_h} e^{c\xi^2} (1 + b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi}{\int_0^{z_h} (1 + b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi} \right) \int_0^{z_h} (1 + b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi \right] \right]$$

**Entropy**  $\longleftrightarrow$  **multiplicity of process (Landau)**

$$s = \left( \frac{L}{z_h} \right)^{1+\frac{2}{\nu}} \frac{(1 + bz_h^2)^{-3a}}{4}$$

# Thermodynamics

**Black hole temperature  $\longleftrightarrow$  QGP temperature (Maldacena)**

$$T = \frac{|g'|}{4\pi} \Big|_{z=z_h} = \frac{1}{4\pi} \left[ - \frac{(1+bz_h^2)^{3a} z_h^{1+\frac{2}{\nu}}}{\int_0^{z_h} (1+b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi} \left[ 1 - \frac{2\mu^2 c e^{2cz_h^2}}{L^2 (1-e^{cz_h^2})^2} \times \right. \right. \\ \left. \left. \times \left( 1 - e^{-cz_h^2} \frac{\int_0^{z_h} e^{c\xi^2} (1+b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi}{\int_0^{z_h} (1+b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi} \right) \int_0^{z_h} (1+b\xi^2)^{3a} \xi^{1+\frac{2}{\nu}} d\xi \right] \right]$$

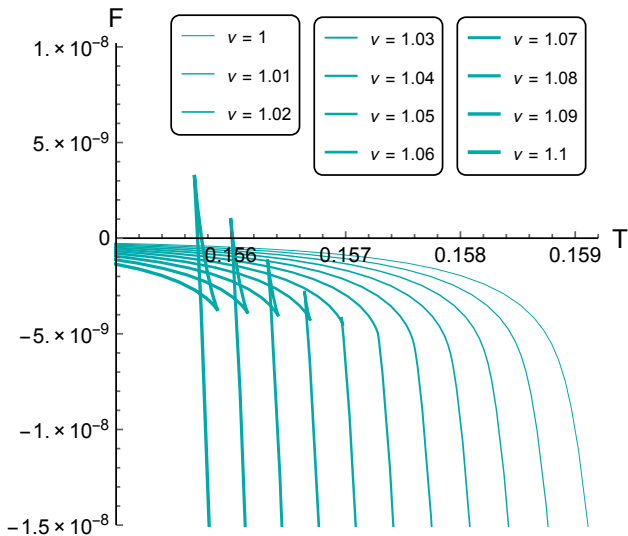
**Entropy  $\longleftrightarrow$  multiplicity of process (Landau)**

**Free energy behavior  $\longrightarrow$  phase transitions**

$$s = \left( \frac{L}{z_h} \right)^{1+\frac{2}{\nu}} \frac{(1+bz_h^2)^{-3a}}{4} \qquad F = \int_{z_h}^{z_{h2}} s dT = \int_{z_h}^{z_{h2}} s T' dz$$

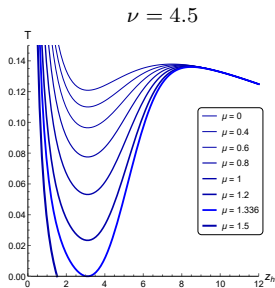
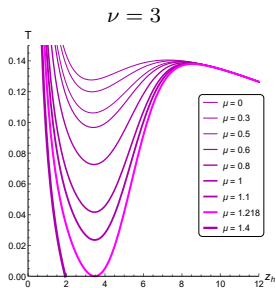
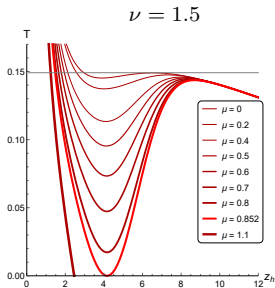
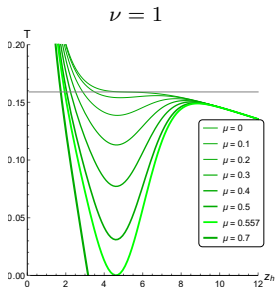
Peculiar features of model  $\longleftarrow$  in  $T(z_h)$  (*I.A. talk, Wednesday*)

# Free energy: $\mu = 0$





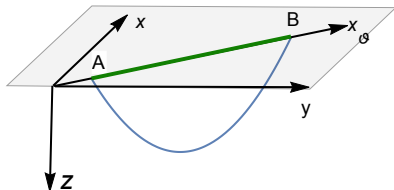
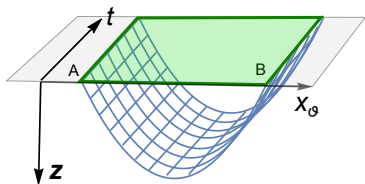
# Temperature



# Temporal Wilson loops

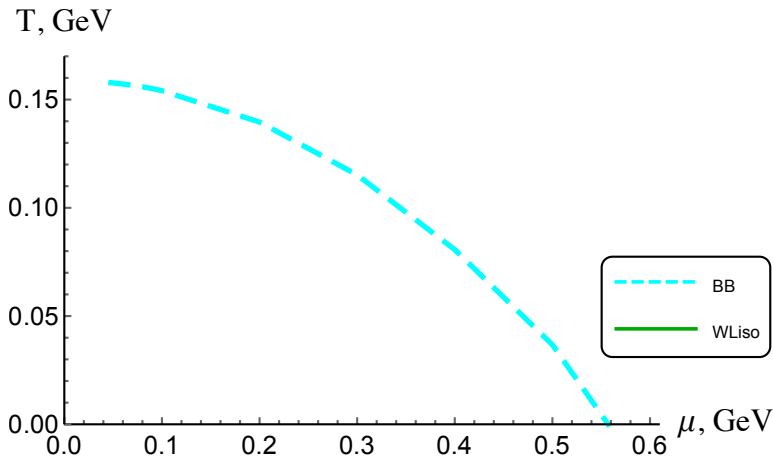
Nambu-Goto action for strings

$$S \sim \sigma_{DW} \ell$$



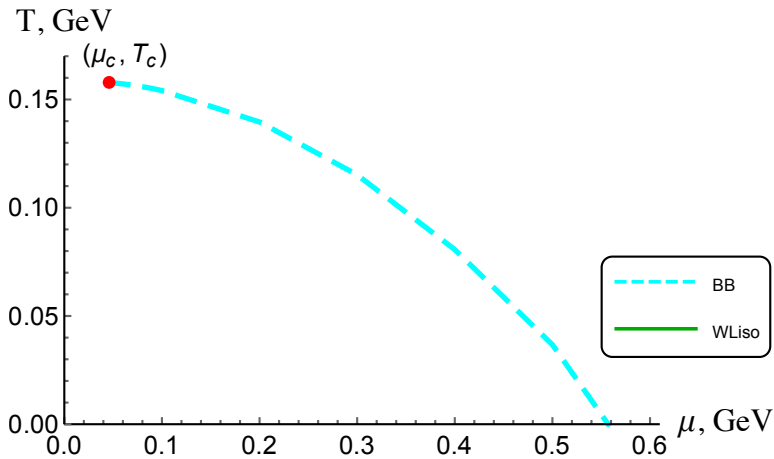
$$X^0 \equiv t, \quad X^1 \equiv x = \xi \cos \theta, \quad X^2 \equiv y_1 = \xi \sin \theta, \quad X^3 \equiv y_2 = \text{const}, \quad X^4 \equiv z = z(\xi)$$

# Phase diagram: $\nu = 1$



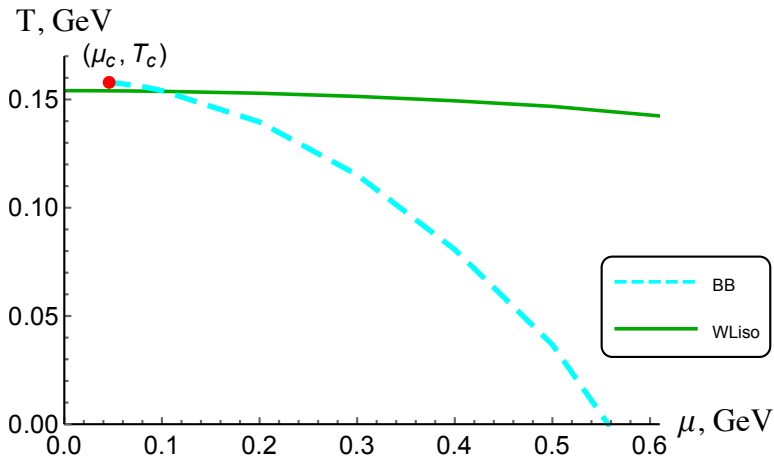
# Phase diagram: $\nu = 1$

CEP:  $\mu_c = 0.04779$ ,  $T_c = 0.1578$



# Phase diagram: $\nu = 1$

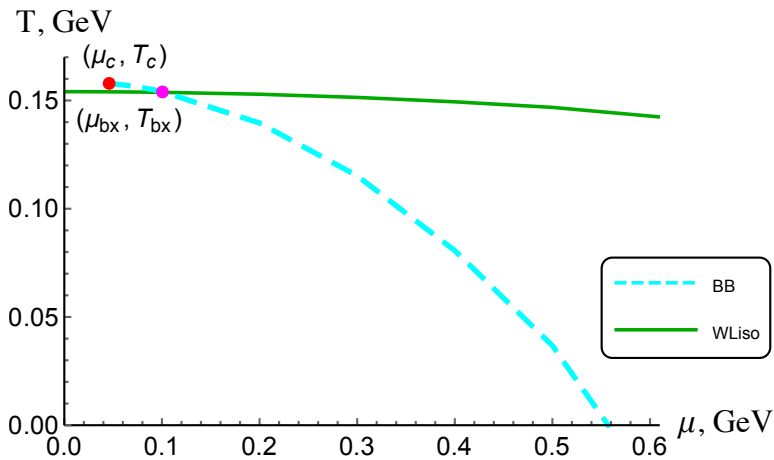
CEP:  $\mu_c = 0.04779$ ,  $T_c = 0.1578$



# Phase diagram: $\nu = 1$

CEP:  $\mu_c = 0.04779$ ,  $T_c = 0.1578$

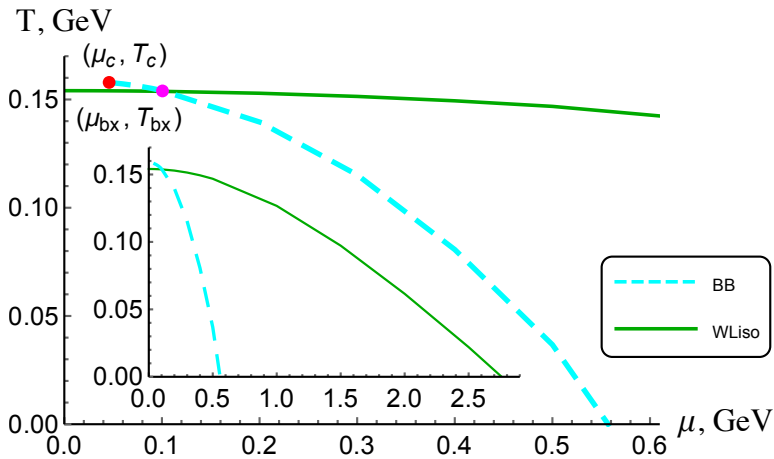
$\mu_{bx} = 0.1027$ ,  $T_{bx} = 0.1538$



# Phase diagram: $\nu = 1$

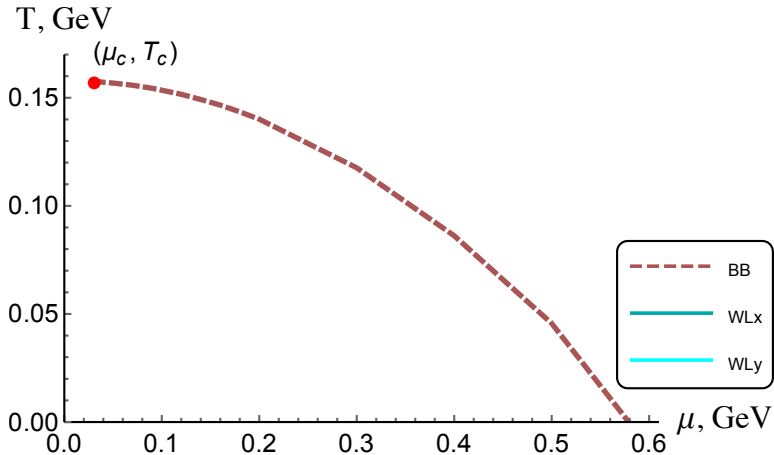
CEP:  $\mu_c = 0.04779$ ,  $T_c = 0.1578$

$\mu_{bx} = 0.1027$ ,  $T_{bx} = 0.1538$



# Phase diagram: $\nu = 1.03$

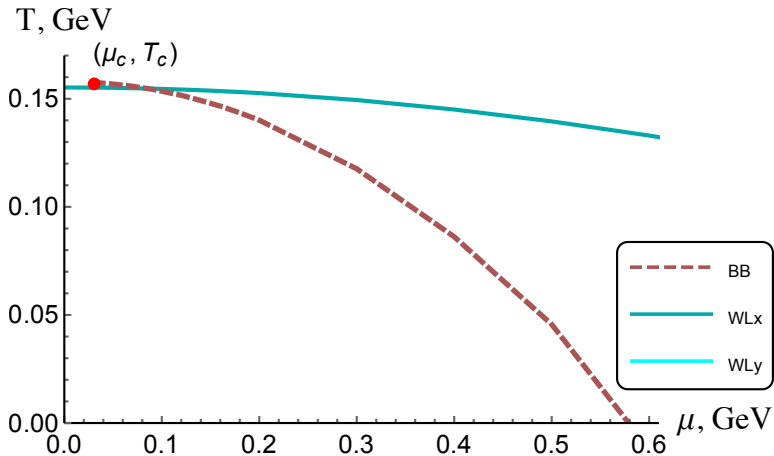
CEP:  $\mu_c = 0.031$ ,  $T_c = 0.1575$





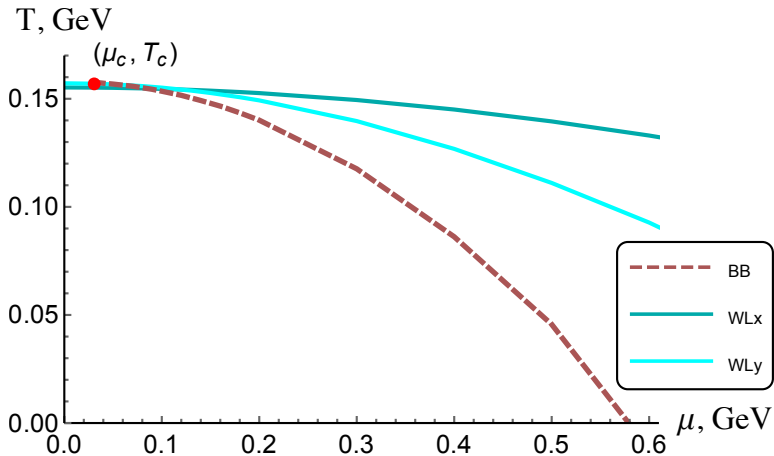
# Phase diagram: $\nu = 1.03$

CEP:  $\mu_c = 0.031$ ,  $T_c = 0.1575$



# Phase diagram: $\nu = 1.03$

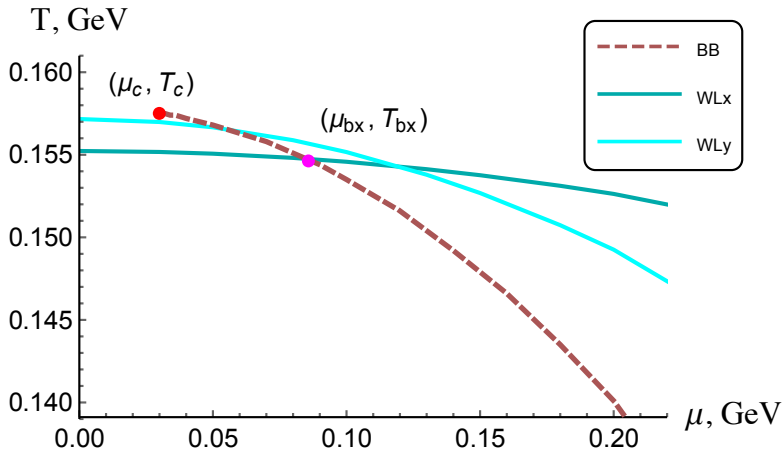
CEP:  $\mu_c = 0.031$ ,  $T_c = 0.1575$



# Phase diagram: $\nu = 1.03$

CEP:  $\mu_c = 0.031$ ,  $T_c = 0.1575$

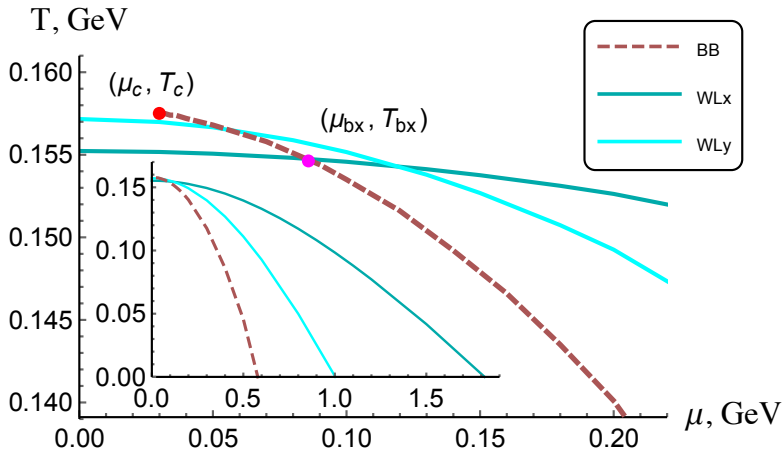
$\mu_{bx} = 0.08497$ ,  $T_{bx} = 0.1548$



# Phase diagram: $\nu = 1.03$

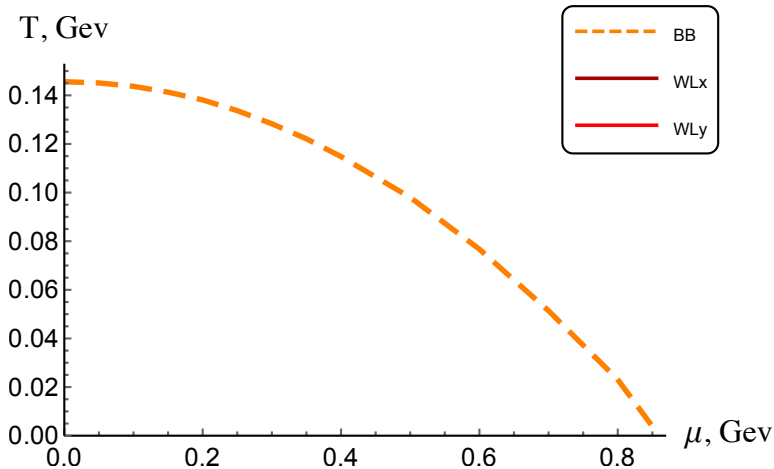
CEP:  $\mu_c = 0.031$ ,  $T_c = 0.1575$

$\mu_{bx} = 0.08497$ ,  $T_{bx} = 0.1548$



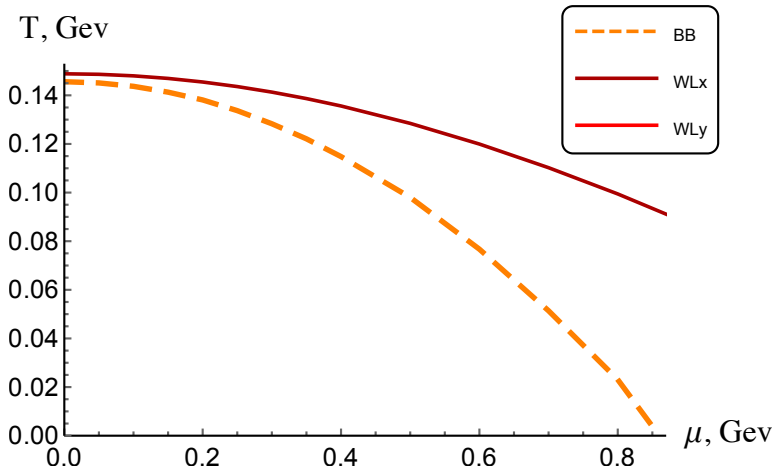
# Phase diagram: $\nu = 1.5$

No CEP



# Phase diagram: $\nu = 1.5$

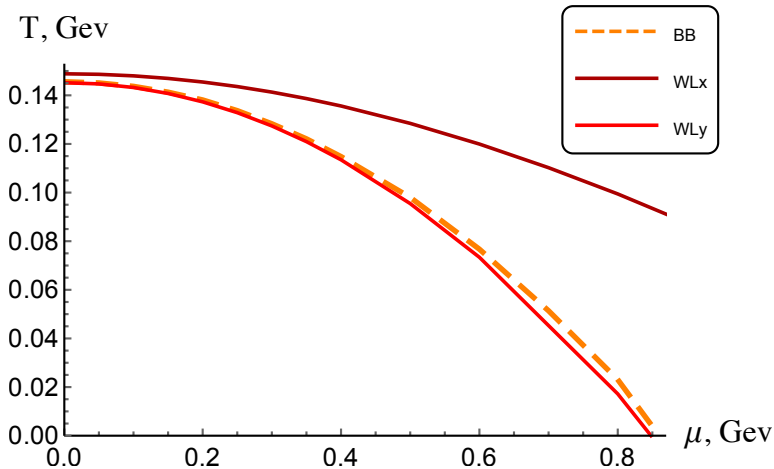
No CEP



# Phase diagram: $\nu = 1.5$

No CEP

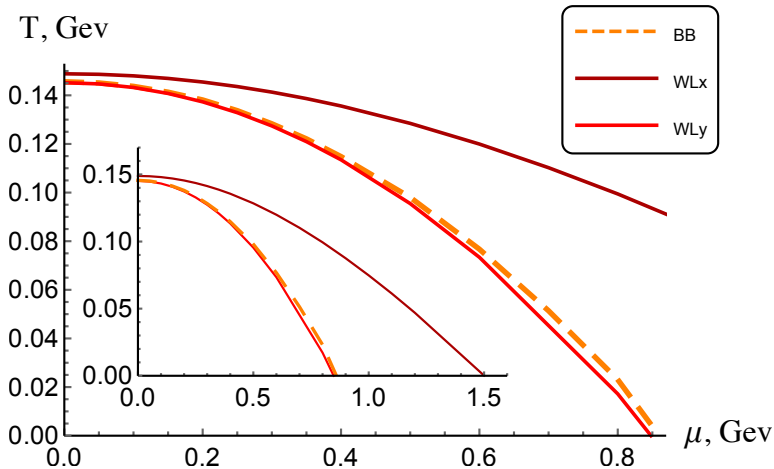
No BY-crosssection



# Phase diagram: $\nu = 1.5$

No CEP

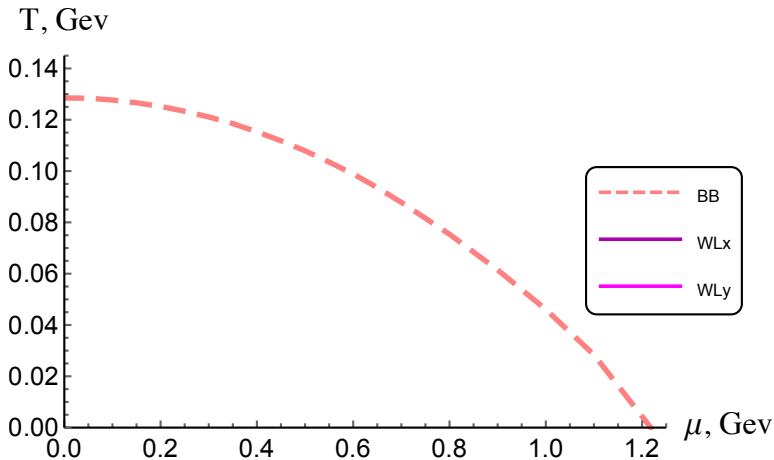
No BY-crossection





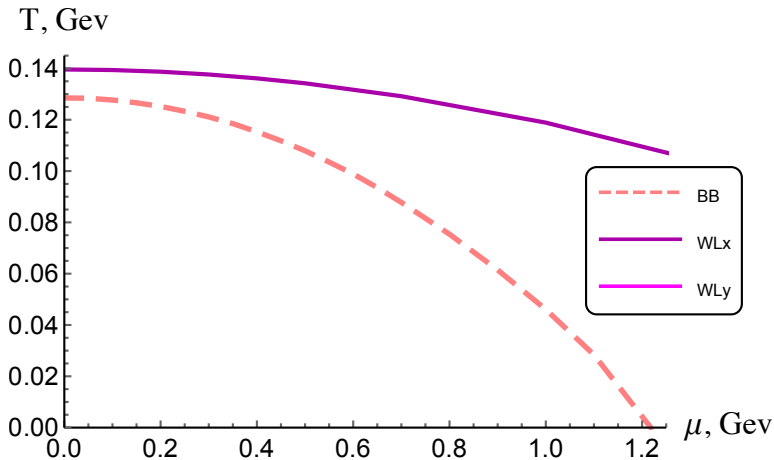
# Phase diagram: $\nu = 3$

No CEP



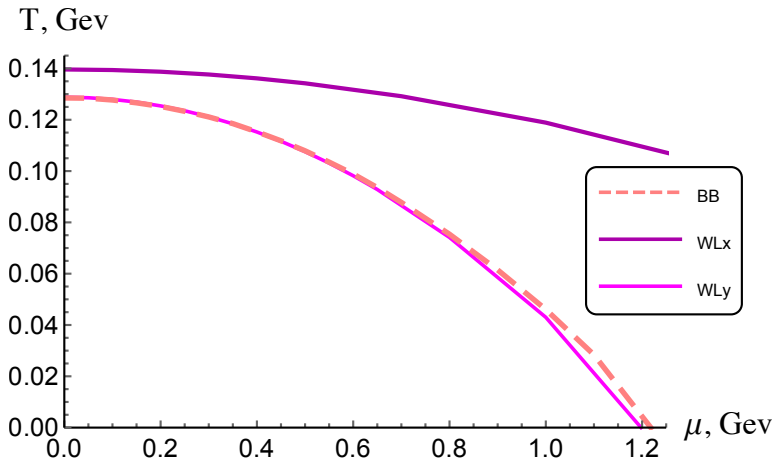
# Phase diagram: $\nu = 3$

No CEP



# Phase diagram: $\nu = 3$

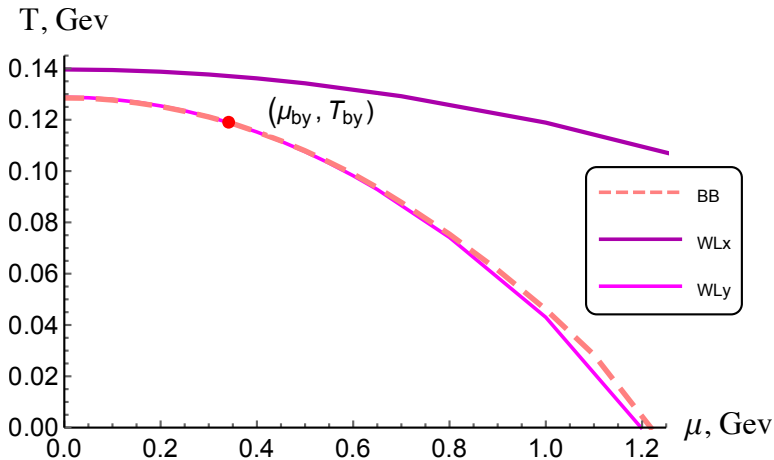
No CEP



# Phase diagram: $\nu = 3$

No CEP

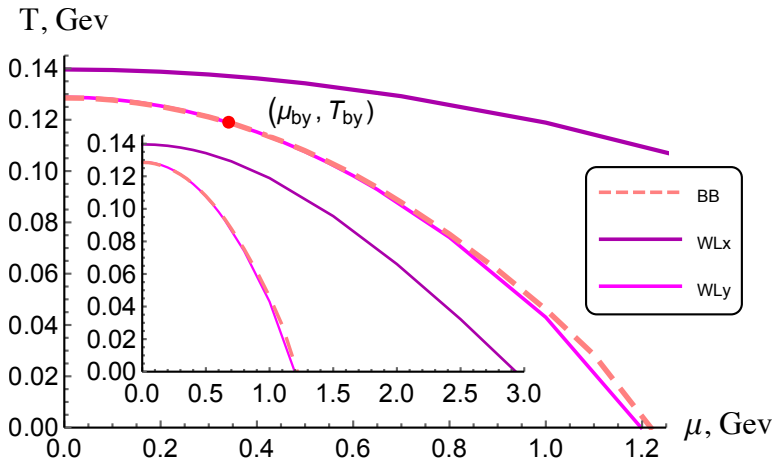
$$\mu_{by} = 0.3244, T_{by} = 0.1198$$



# Phase diagram: $\nu = 3$

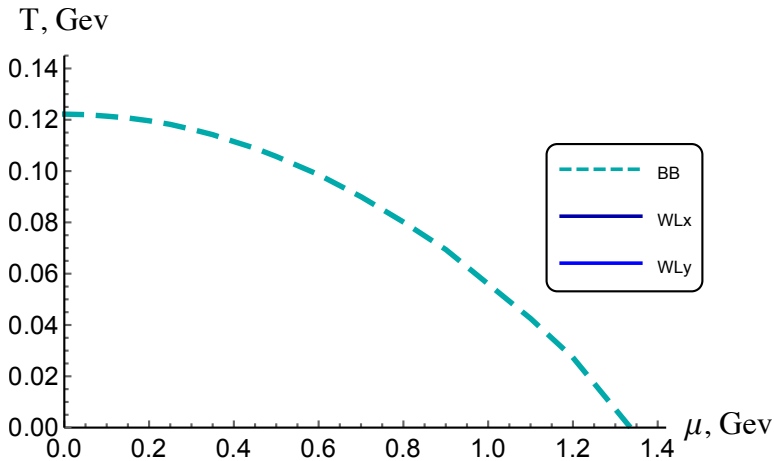
No CEP

$$\mu_{by} = 0.3244, T_{by} = 0.1198$$



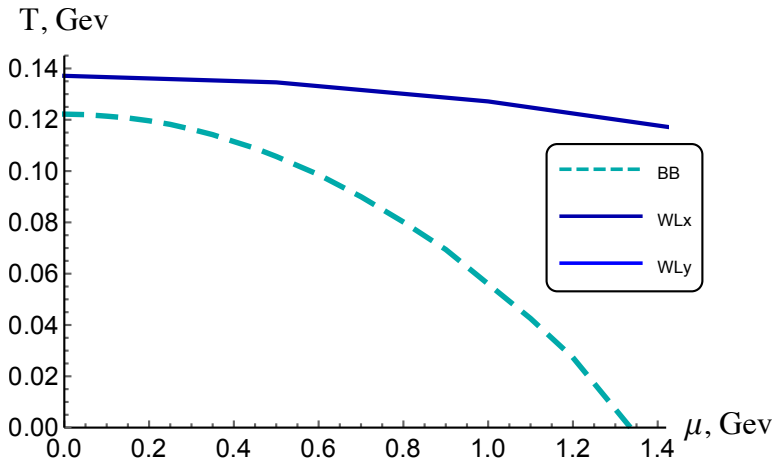
# Phase diagram: $\nu = 4.5$

No CEP



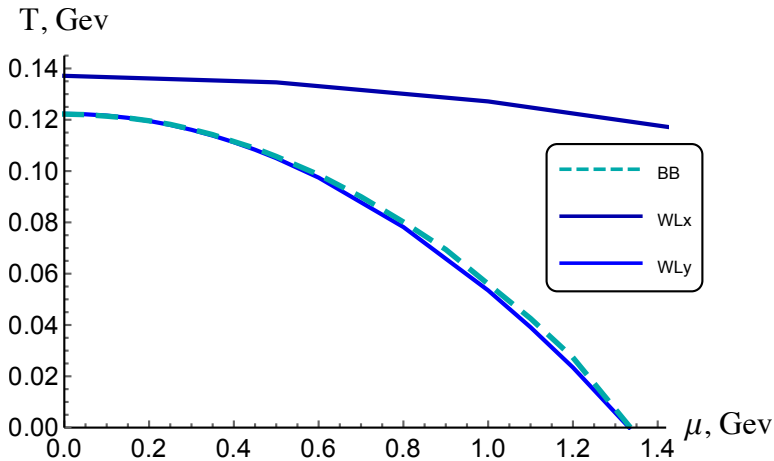
# Phase diagram: $\nu = 4.5$

No CEP



# Phase diagram: $\nu = 4.5$

No CEP

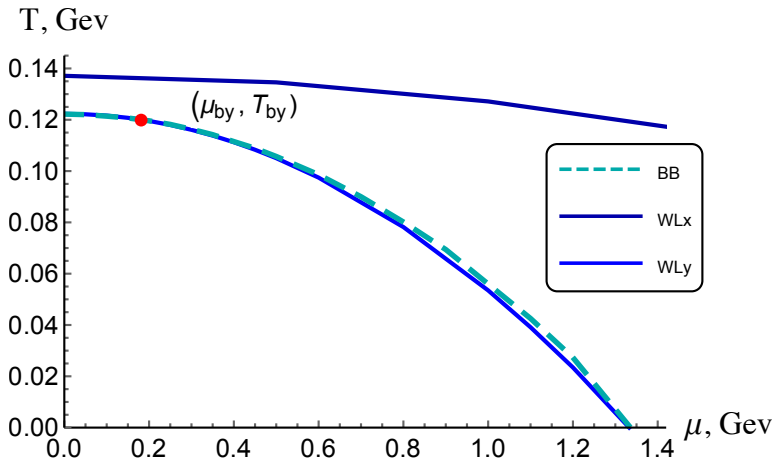




# Phase diagram: $\nu = 4.5$

No CEP

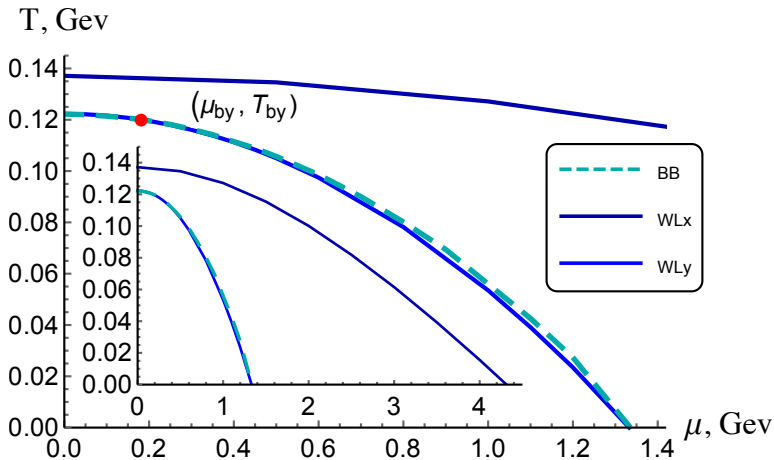
$$\mu_{by} \approx 0.1748, T_{by} \approx 0.1201$$



# Phase diagram: $\nu = 4.5$

No CEP

$\mu_{by} \approx 0.1748, T_{by} \approx 0.1201$



# Conclusions

## For light quarks

- 1 1-st order (Hawking-Page-like) phase transition line
  - starts from critical point for  $\nu < 1.05$  and from  $\mu = 0$  for  $\nu \geq 1.05$ ,
  - does not break at a relatively high temperature, but lasts till  $T = 0$ .
- 2 Longitudinal orientation of quarks pairs does not contribute to confinement/deconfinement phase transition.
- 3 Transfer of the main role in the phase transition smooth, without jumps (as it was in heavy quarks model).

# Conclusions

## For light quarks

- 1 1-st order (Hawking-Page-like) phase transition line
  - starts from critical point for  $\nu < 1.05$  and from  $\mu = 0$  for  $\nu \geq 1.05$ ,
  - does not break at a relatively high temperature, but lasts till  $T = 0$ .
- 2 Longitudinal orientation of quarks pairs does not contribute to confinement/deconfinement phase transition.
- 3 Transfer of the main role in the phase transition smooth, without jumps (as it was in heavy quarks model).

# Conclusions

## For light quarks

- 1 1-st order (Hawking-Page-like) phase transition line
  - starts from critical point for  $\nu < 1.05$  and from  $\mu = 0$  for  $\nu \geq 1.05$ ,
  - does not break at a relatively high temperature, but lasts till  $T = 0$ .
- 2 Longitudinal orientation of quarks pairs does not contribute to confinement/deconfinement phase transition for large anisotropy.
- 3 Transfer of the main role in the phase transition smooth, without jumps (as it was in heavy quarks model).

# Conclusions

## For light quarks

- 1 1-st order (Hawking-Page-like) phase transition line
  - starts from critical point for  $\nu < 1.05$  and from  $\mu = 0$  for  $\nu \geq 1.05$ ,
  - does not break at a relatively high temperature, but lasts till  $T = 0$ .
- 2 Longitudinal orientation of quarks pairs does not contribute to confinement/deconfinement phase transition.
- 3 Transfer of the main role in the phase transition smooth, without jumps (as it was in heavy quarks model).

## What's next?

Hybrid model: heavy&light quarks — all together now

Thank you  
for your attention

The work is supported by RFBR grant №18-02-40069 mega