

Recent results related to Feynman integrals calculus

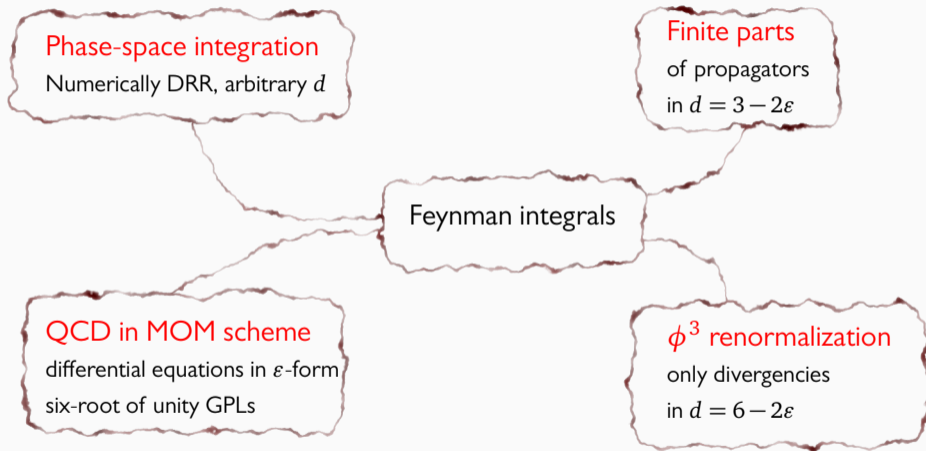
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Setting the stage

Different problems solved with similar techniques

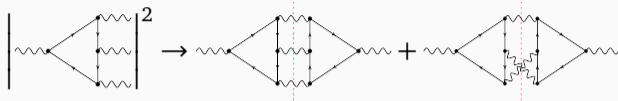


Results to be discussed

1. Fully inclusive phase-space integrals needed for NNLO QCD cross-section calculations for e^+e^- annihilation into hadrons
2. Four-loop fermion propagator in $d = 3 - 2\epsilon$ QED
3. Four-loop QCD beta-function in MOM scheme and operators anomalous dimensions relevant for quark masses at three-loop order
4. Five-loop renormalization and critical exponents in φ^3 theory and its generalizations

Phase space integration

Main idea: treat integration over phase space as loop integrals with cut propagators



Fully inclusive massless phase-space integrals use cases:

- Total cross-section calculation, $R(s)$
- Boundary conditions for semi-inclusive observables, time-like splitting functions
- NNLO subtraction terms for $q\bar{q}$ with jets production in e^+e^- annihilation
- Boundary conditions for heavy quark pair production in e^+e^- annihilation

$1 \rightarrow 4$ @tree, $1 \rightarrow 3$ @1-loop and $1 \rightarrow 2$ @2-loop are known

[Gehrmann et al. '04]

Our goal: $1 \rightarrow 5$ @tree, $1 \rightarrow 4$ @1-loop, $1 \rightarrow 3$ @2-loop and $1 \rightarrow 2$ @3-loop

Details and results

- **Final result:** Solution of DRR in the form of fast convergent series for arbitrary d , analytical results in terms of MZVs for the interesting $d = 4 - 2\epsilon$ case
- Intermediate steps:
 - DRR system for all integrals corresponding to all possible cuts of four-loop propagators
 - Optical theorem relations connecting imaginary part of virtual integral with cuts
 - Two-loop integrals for $1 \rightarrow 3$ to higher weight in $d = 6$ and $d = 8$, w/o IR divergencies
 - Divergencies of massive one-loop $1 \rightarrow 4$ integrals in logarithmic dimension
- Fix arbitrary periodic functions in DRR solution with all available data
- Use PSLQ to reconstruct $d = 4 - 2\epsilon$ result using MZV basis and high precision numbers

Summary:

[Gituliar, Magerya, AP' JHEP18][Magerya, AP' JHEP19]

1. Integrals from all possible cuts available as fast convergent series for arbitrary d
2. For $d = 4 - 2\epsilon$ numerical results successfully reconstructed with MZV upto weight 12

IR divergencies in $3 - 2\epsilon$ dimensional QED

- Long standing question about IR finiteness of 3d QED [Jackiw, Templeton '81]
- Landau-Khalatnikov-Fradkin transformation predicts relations between different loop orders for fermion propagator self-energy [Gusynin, Kotikov, Teber '20]

Our goal: to calculate bare fermion propagator at 3- and 4-loop order in the $n_f \rightarrow 0$ limit

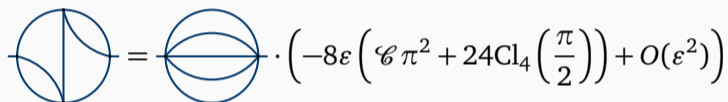
$$S_F(p, \xi) = \frac{iS(\xi)}{\hat{p}}, \quad S(\xi) = 1 + AS_1(\xi) + A^2S_2(\xi) + A^3S_3(\xi) + A^4S_4(\xi) + \mathcal{O}(A^5)$$

LKF predictions:

1. If three-loop term $S_3(\xi) \neq 0$ then four-loop term $S_4(\xi)$ is divergent
2. From known result e.g. in Landau gauge $S_i(0)$ all ξ dependence can be reconstructed from lower loop results

Details and results

- Main difficulty to calculate 3- and 4-loop propagators in $d = 3 - 2\varepsilon$
- Numerical results with high precision available for arbitrary d
- For odd d ε -expansion contains MZVs only, but not the case for $d = 3 - 2\varepsilon$


$$\text{Diagram} = \text{Diagram} \cdot \left(-8\varepsilon \left(\zeta(2) \pi^2 + 24\text{Cl}_4\left(\frac{\pi}{2}\right) \right) + O(\varepsilon^2) \right)$$

- Analytical expressions reconstructed using basis of fourth-root of unity GPLs

Summary:

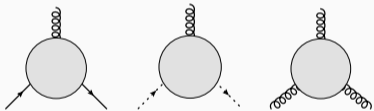
[AP, Gusynin, Kotikov, Teber 'PRD20]

1. Quenched ($n_f = 0$) 3d QED is divergent at 4-loop level and 3-loop part of fermion propagator is not zero $S_3(\xi) \neq 0$
2. All gauge dependent terms are in agreement with LKF predictions
3. We conjecture that all massless propagators in $d = 3 - 2\varepsilon$ are expressible through 4-th root of unity GPLs

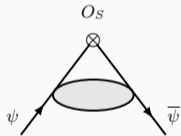
QCD renormalization in MOM scheme

- β_{QCD} known to 5-loop order in \overline{MS} scheme [Baikov,Chetyrkin,Kuhn '17]
- \overline{MS} scheme is simplest to perform calculations, but not unique
- MOM scheme $m = 0$, but off-shell legs provides matching with lattice data

Beta functions from renormalization of three-point vertices



Mass MOM renormalization

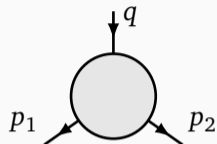


- Most difficult to calculate 3-loop 3-pt integrals with all legs off-shell
- Four-loop beta-functions from conversion relation $a_{MOM} = a_{\overline{MS}} \left(1 + \sum_l X_l \alpha_{\overline{MS}}^l \right)$
- Quark masses from lattice data using three-loop O_S renormalization in MOM scheme

Details and results

DE system reducible to ε -form - solution using GPLs with fixed alphabet

[Henn '13]



- Differential equations in q^2 variable to connect limits
- Symmetric point $p_1^2 = p_2^2 = q^2$, from arbitrary q^2 result
- 2-pt function limit $q^2 \rightarrow 0$, used to fix boundary conditions

All three-loop integrals expressible through 6-th root of unity GPLs upto weight 6 and even more restricted basis of HPLs with $e^{i\pi/3}$ argument

[Kniehl, AP, Veretin' JHEP17]

Summary:

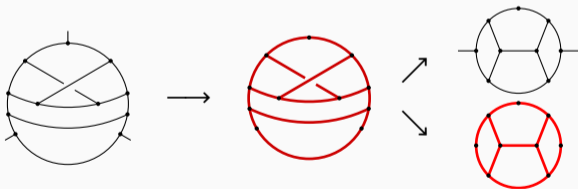
[Bednyakov, AP' PRD20a][Bednyakov, AP' PRD20b]

1. For the first time calculated three-loop 3-pt integrals in symmetric point
2. With 3-loop conversion factor derived 4-loop QCD beta-function in MOM scheme
3. NNNLO relation connecting $\overline{\text{MS}}$ quark mass to MOM/RI mass from lattice data

Five-loop critical exponents in φ^3 theory

- Scalar φ^3 theory in $d = 6$ is a well known playground together with φ^4 in $d = 4$
- Has many generalizations interesting in practice:
Potts model, Lee-Yang edge singularity, model with $O(n)$ symmetry [Fei, et al. '14]
- Existing four-loop calculations are due to known four-loop propagator integrals

TODO: Reduce divergencies of 5-loop vertices and propagators to 4-loop integrals only



Most difficult part: derivatives of propagators in external momenta generating numerators

Details and results

Method of solution:

1. Apply $\mathcal{H}\mathcal{R}'G$ to each logarithmically divergent diagram in massive theory
2. Using IBP rewrite $\mathcal{H}G$ part through the set of master integrals w/o spurious poles and easy to calculate, in our case: 4-pt functions with 16 additional 3-pt functions by hand
3. For 4-pt functions $\mathcal{H}G$ is easy to calculate with 4-loop integrals only
4. For 3-pt functions with IRR reduce problem to calculation of $\mathcal{H}\mathcal{R}'G$ for 4-loop propagator insertion into one-loop diagram

Summary:

[Kompaniets, AP' to appear]

- $\mathcal{H}\mathcal{R}'$ for all φ^3 diagrams, can be used for other more complicated models
- Results for $O(n)$ model checked with existing large n expansion to $\mathcal{O}(1/n^3)$
- φ^3 theory results are in agreement with independent calculation [Schnetz et al.]

Conclusion

With the help of modern methods of Feynman integrals calculation we achieved following improvements in different QFT problems:

1. PS integrals for e^+e^- total crosssection

before: $1 \rightarrow 4$ after: $1 \rightarrow 5$

2. 3d QED fermion propagator

before: 2-loop after: 4-loop

3. Integrals for symmetric point MOM scheme calculations

before: 2-loop after: 3-loop

4. Scalar φ^3 theory in $d = 6$ renormalization

before: 4-loop after: 5-loop

Thank you for attention!