

Solution of differential equations for multiloop integrals with **Libra**¹ package

Roman N. Lee

Budker Institute of Nuclear Physics, Novosibirsk

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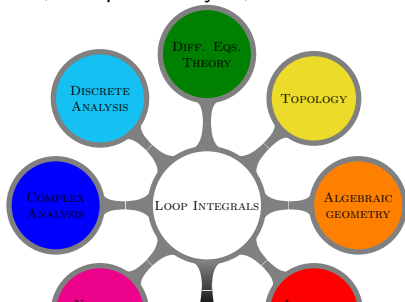
¹RL, CPC **267** (2021) 108058.

Congratulations!

My congratulations to all heros of the event!
I am very grateful to Vladimir Smirnov for pleasant
and fruitful collaboration during many years.

Motivation

- If we want to detect deviations (new physics) from SM, we need to know the predictions of the latter with high precision. In particular, we have to be able to calculate two-loop radiative corrections to various processes [A. Arbuzov's talk yesterday].
- Only quite recently the methods of multiloop calculations have reached the point where they can be really helpful with this goal.
- Besides these practical purposes, multiloop calculations provide a perfect polygon for trying the methods from various mathematical fields: differential equations, complex analysis, number theory



Modern approach to diagrams calculation

- 1 Consider a family of integrals

$$j(n_1, \dots, n_N) = \int \frac{d^d l_1 \dots d^d l_L}{D_1^{n_1} \dots D_N^{n_N}}.$$

Integrals are functions of kinematic variables x_i and $d = 4 - 2\epsilon$.

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- 3 Find differential equations [Kotikov, 1991, Remiddi, 1997] (and/or dimensional recurrences [Tarasov, 1996]) for master integrals

Differential equations

$$\frac{\partial}{\partial x_i} \mathbf{j} = M(\mathbf{x}, \epsilon) \mathbf{j}$$

Dimensional recurrences

$$\mathbf{j}(\epsilon + 1) = R(\mathbf{x}, \epsilon) \mathbf{j}(\epsilon)$$

M and R are $n \times n$ matrices rational in \mathbf{x} and ϵ .

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- 4 Find **general solution**.
- 5 Use other methods for **boundary conditions**.

IBP reduction

Note on phase-space integrals

- Phase-space integrals can be transformed in a standard way into loop integrals with cut propagator:

$$\frac{d^{d-1}p}{(2\pi)^{d-1}2\varepsilon_p} = \frac{d^d p}{(2\pi)^d} 2\pi\delta_+(p^2 - m^2) = \frac{d^{d-1}p}{(2\pi)^d} (-i) \oint_{\varepsilon_p} dp_0 (p^2 - m^2)^{-1}$$

- As IBP identities are insensitive to the integration contour, provided that it does not lead to surface terms, we can treat the cut propagators in the same way as uncut ones.
- There are two things to take care of: first, shift symmetries which mix cut with uncut denominators should be omitted, second, positive (integer) powers of cut denominators are equal to zero.

ϵ -form

Differential equation

$$\partial \mathbf{J}(x) / \partial x = M(x, \epsilon) \mathbf{J}(x)$$

Function change

$$\mathbf{J}(x) = T(x, \epsilon) \tilde{\mathbf{J}}(x)$$

In particular, we can choose master integrals in infinitely many ways.

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Remarkable observation [Henn, 2013]

There often exists a choice of master integrals such that

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This form makes finding a solution in the form of ϵ -expansion very simple.

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How to find canonical basis?

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Algorithmic approach [RL, 2015]: general idea

Perform many “elementary” transformations gradually improving properties of the system.

General structure of reduction algorithm

Algorithm proceeds in three major stages, each involving a sequence of “elementary” transformations.

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1. “Fuchsification”: Eliminating higher-order poles

Input: Rational matrix $M(x, \epsilon)$

Output: Rational matrix with only simple poles on the extended complex plane, $M(x, \epsilon) = \sum_k \frac{M_k(\epsilon)}{x - a_k}$.

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2. Normalization: Normalizing eigenvalues

Input: Matrix from the previous step, $M(x, \epsilon) = \sum_k \frac{M_k(\epsilon)}{x-a_k}$.

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3. Factorization: Factoring out ϵ

Input: Matrix from the previous step.

Output: Matrix in ϵ -form, $M(x, \epsilon) = \epsilon S(x) = \epsilon \sum_k \frac{S_k}{x-a_k}$.

Balance

Balance transformation

$$T(x) = \bar{P} + \frac{x - x_2}{x - x_1} P,$$

where P is some projector and $\bar{P} = I - P$. When $x_1 = \infty$ or $x_2 = \infty$ omit denominator/numerator.

Balance transformation changes properties (pole order and eigenvalues of matrix residue) of the differential system at $x = x_1$ and $x = x_2$ only.



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Programs

Further details of the algorithm deserve a special talk, but the good news is that we now have **programs**! At least, three public ones: **epsilon**, **Fuchsia**, **Libra**.

Libra interface for reduction

Automatic tool (useful for simple cases)

```
In[1]: t=Rookie[M,x,ϵ];
```

Interactive tool (useful for most cases)

```
In[1]: t=VisTransformation[M,x,ϵ];
```

The screenshot shows a window titled "VisTransformation" with a close button (X) in the top left. The main area contains a 3x5 grid of input matrices and their corresponding reduced forms. Each input matrix is followed by a "Fuchsfity" button. The reduced forms are also 3x5 matrices with some entries highlighted in red or green. Below the matrices, there are three buttons: "Apply balance transformation (7b)", "Paste overall transformation", and "0-dimensional u-space and 0-dimensional v-space".

$\begin{bmatrix} 0 & -1+2\epsilon & -1+2\epsilon & -1+2\epsilon & -1-4\epsilon \\ 0 & 0 & 0 & 0 & 0 \\ -1+2\epsilon & -1+2\epsilon & -1-2\epsilon & -2(1+\epsilon) & 1-6\epsilon \end{bmatrix}$	Fuchsfity	$x=-1, pr=0$	$\begin{bmatrix} -1-4\epsilon & -1+2\epsilon & -1+2\epsilon & -1+2\epsilon & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 1-6\epsilon & -2(1+\epsilon) & -1-2\epsilon & -1+2\epsilon & -1+2\epsilon \end{bmatrix}$
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$\begin{bmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{bmatrix}$	Fuchsfity	$x=1, pr=0$	$\begin{bmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{bmatrix}$

Buttons: Apply balance transformation (7b), Paste overall transformation, 0-dimensional u-space and 0-dimensional v-space

Manual tool (useful for really hard cases)

```
In[1]: u=GetSubspaces[M,{x,0},ϵ][[1]];v=GetSubspaces[M...
```

Factoring ϵ

```
In[2]: t=FactorOut[M,x,ϵ,μ];
```

General solution

Differential system in ϵ -form

$$\partial_x \mathbf{J} = \epsilon S(x) \mathbf{J}, \quad S(x) = \sum_k \frac{S_k}{x - a_k}.$$

General solution (evolution operator)

$$U(x, x_0) = \text{Pexp} \left[\epsilon \int_{x_0}^x dx_1 S(x_1) \right]$$

In [3]: `U=PexpExpansion[{S,6},x];`

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Goncharov polylogarithms

So, $U(x, x_0)$ is expressed via Goncharov polylogs [Goncharov, 1998]

$$G(a_n, \dots, a_1 | x) = \int_{x > x_n > \dots > x_1 > 0} dx_1 \dots dx_n \frac{dx_1}{x_1 - a_1} \dots \frac{dx_n}{x_n - a_n}$$

Perfect class of functions: numerical evaluation, analytic continuation, series representation, functional identities, relations to classical polylogs.

Boundary conditions

Suppose we have found a transformation $T(x) = T(x, \epsilon)$ to ϵ -form, $\mathbf{j} = T\mathbf{J}$. Then we can write

$$\begin{aligned}\mathbf{J}(x) &= U(x, x_0)\mathbf{J}(x_0), \\ \mathbf{j}(x) &= T(x)U(x, x_0)[T(x_0)]^{-1}\mathbf{j}(x_0)\end{aligned}$$

But the point x_0 should be somewhat **special** to simplify the evaluation of $\mathbf{j}(x_0)$ as compared to $\mathbf{j}(x)$. With no known exceptions, "**special**" boils down to "**singular**", i.e., we can expect simplifications for x_0 being a singular point of the differential system. Let it be $x_0 = 0$ for simplicity.

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Problem

$U(x, x_0)$ diverges when x_0 tends to zero. Therefore, we have to consider not the values, but the asymptotics of $\mathbf{j}(x_0)$ at $x = 0$.

Boundary conditions

Regularized evolution operator

$$U(x, \underline{0}) = \lim_{x_0 \rightarrow 0} U(x, x_0) x_0^{\epsilon S_0},$$

where $S_0 = \text{Res}_{x=0} S(x)$.

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- ① $U(x, \underline{0})$ has no divergences.
- ② $U(x, \underline{0})$ is a general solution.
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Specific solution reads $\mathbf{j}(x) = T(x)U(x, \underline{0})\mathbf{C}$. The column of boundary constants \mathbf{C} can be fixed by evaluating some coefficients in the asymptotics of $\mathbf{j}(x)$ when $x \rightarrow 0$.

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Good news

Libra can determine which asymptotic coefficients, \mathbf{c} , are sufficient to calculate and find the “adapter” matrix L in $\mathbf{C} = L\mathbf{c}$.

NB: for regular point, of course, $\mathbf{c} = \mathbf{j}(0)$ and $L = T^{-1}(0)$.

```
In[4]: {L, cs}=GetLcs[S, T, {x, 0}];
```

Example: boundary conditions for $\sigma_{e^- \gamma \rightarrow e^- \gamma}$ @NLO

The following threshold ($s \rightarrow 1$, $x \rightarrow 0$) asymptotic coefficients are to be calculated:

$$\begin{aligned}
 & \left[\text{diagram} \right]_{y^{2-4\epsilon}}, \left[\text{diagram} \right]_{y^{-2}}, \left[\text{diagram} \right]_{y^{4-8\epsilon}, y^{2-4\epsilon}, y^{4\epsilon-2}}, \left[\text{diagram} \right]_{y^{-4\epsilon}, y^{4\epsilon-2}}, \\
 & \left[\text{diagram} \right]_{y^{-2}, y^{-4\epsilon}}, \left[\text{diagram} \right]_{y^{-3}}, \left[\text{diagram} \right]_{y^{-2}, y^{-4\epsilon}}, \left[\text{diagram} \right]_{y^{-4\epsilon-2}}, \\
 & \left[\text{diagram} \right]_{y^{-2\epsilon-3}, y^{4\epsilon-2}}, \left[\text{diagram} \right]_{y^{-4}}, \left[\text{diagram} \right]_{y^{4-8\epsilon}}, \left[\text{diagram} \right]_{y^{-4\epsilon}}, \left[\text{diagram} \right]_{y^{-4\epsilon}}, \\
 & \left[\text{diagram} \right]_{y^{-4\epsilon}}, \left[\text{diagram} \right]_{y^{-4}}, \left[\text{diagram} \right]_{y^{-4}}, \left[\text{diagram} \right]_{y^{-4}}, \left[\text{diagram} \right]_{y^{-2\epsilon-5}}.
 \end{aligned}$$

Here $[\text{integral}]_{x^\alpha}$ denotes the coefficient in front of x^α in the small- x asymptotics of integral ($x \approx \frac{1}{2}\sqrt{s-1}$).

Example: Boundary conditions for $\sigma_{e^- \gamma \rightarrow e^- \gamma}$ @NLO

[RL, Lyubyakin, Stocky (2020); RL, Schwartz, Zhang (2021)]

- Selection rule by the fractional power $n\epsilon$ leaves us with 11 possibly nonzero constants. The fractional power -4ϵ corresponds to the hard momentum flowing over the black lines, while -8ϵ — to soft momentum.

$$\begin{aligned}
 & \left[\text{Diagram 1} \right]_{x^{2-4\epsilon}}, \left[\text{Diagram 2} \right]_{x^{4-8\epsilon}, x^{2-4\epsilon}}, \left[\text{Diagram 3} \right]_{x^{-4\epsilon}}, \\
 & \left[\text{Diagram 4} \right]_{x^{-4\epsilon}}, \left[\text{Diagram 5} \right]_{x^{-4\epsilon}}, \left[\text{Diagram 6} \right]_{x^{-4\epsilon-2}}, \\
 & \left[\text{Diagram 7} \right]_{x^{4-8\epsilon}}, \left[\text{Diagram 8} \right]_{x^{-4\epsilon}}, \left[\text{Diagram 9} \right]_{x^{-4\epsilon}}, \left[\text{Diagram 10} \right]_{x^{-4\epsilon}}.
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- Selection rule by the fractional power $n\epsilon$ leaves us with 11 possibly nonzero constants. The fractional power -4ϵ corresponds to the hard momentum flowing over the black lines, while -8ϵ — to soft momentum.
- Selection rule by integer power reduces the number to 5.

$$\left[\text{Diagram 1} \right]_{x^{2-4\epsilon}}, \left[\text{Diagram 2} \right]_{x^{4-8\epsilon}, x^{2-4\epsilon}},$$

$$\left[\text{Diagram 3} \right]_{x^{4-8\epsilon}}, \left[\text{Diagram 4} \right]_{x^{-4\epsilon}}$$

General solution, Frobenius method

In many applications there is a natural small parameter. E.g., for $e^+e^- \rightarrow X$ the electron mass is small, but can not be put to zero. Instead one should expand the regularized evolution operator

$$U(x, \underline{0}) = \lim_{x_0 \rightarrow 0} U(x, x_0) x_0^{S_0}$$

in generalized power series. **Libra** has tools for it. It closely follows the approach described in Ref. [RL, Smirnov, and Smirnov, 2018].

$U(x, \underline{0})$ as generalized power series

Recursion data for (matrix) coefficients of U :

```
In[1]: sdata=SeriesSolutionData[S,x,x];
```

Using data for expanding to fixed order:

```
In[2]: Uexp=ConstructSeriesSolution[sdata,x,6];
```

Algebraic extensions

- Sometimes, in order to find the transformation to ϵ -form, one has to extend the class of transformations by passing from x to y , such that $x = x(y)$ is some rational function. **Libra** has tool for it:

```
In[1]: ChangeVar[ds, x → (4 - y*y)/(1 - y*y), y];
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- Moreover, in many cases there is no common rationalizing variable. Thus, **Libra** implements a more powerful way to treat such algebraic extensions, with

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In[1]: AddNotation[ds,y → x(1-y*y) - 4 y*y];
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One may add as many notations as needed, and **Libra** will take care of them (minimizing their appearance, correctly treating their differentiation).

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- There is a bunch of functions related to treating the algebraic extensions: **QuolyMod**, **DiffMod**, **SeriesCoefficientMod**, **EValuesMod**, etc. These functions can be used also to treat irreducible denominators, like $x^2 + x + 1$ in a way which do not introduce radicals.

Irreducible cases

- As we know, even with algebraic extensions it is not always possible to reduce the system to ϵ -form. Sometimes the integrals just can not be expressed via polylogs, [see D. Broadhurst's talk yesterday].
- In Ref. [RL and Pomeransky, 2017] the criterion of irreducibility has been derived. The reducibility has been shown to correspond to triviality of some holomorphic vector bundle on the Riemann sphere. Thanks to Birkhoff-Grothendieck theorem, it is possible to constructively decide this. **Libra** has the corresponding tool: the command

```
In[3]: {T1,T2,T3}= BirkhoffGrothendieck[T,x];
```

decomposes Laurent-polynomial matrix T into the product $T_1 T_2 T_3$, where T_1, T_1^{-1} are polynomial in x , T_3, T_3^{-1} are polynomial in x^{-1} , and $T_2 = \text{diag}(x^{n_1}, \dots, x^{n_k})$. The bundle is trivial iff $T_2 = 1$.

Summary

- Modern multiloop calculation techniques can really help in NNLO calculation useful for the experiments.
- **Libra** can really help in applying the differential equations method. It has tools
 - for the reduction of the differential system to ϵ -form,
 - for the construction of general solution in terms of Goncharov's polylogs,
 - for determining the minimal set of asymptotic coefficients to be evaluated to fix the boundary conditions,
 - for constructing Frobenius expansion,
 - for treating the algebraic extensions,
 - for detecting the irreducible cases.
 - It can work with univariate and multivariate systems.

Outlook

- Libra improvements:
 - Improve automatic tool Rookie $[M, x, \epsilon]$.
 - Better treatment of algebraic extensions, especially, for multivariate case.
- Differential equations method:
 - construct a systematic approach to irreducible cases. In particular, Birkhoff-Grothendieck factorizations seems to be carry a lot of information yet to be properly understood and used.

Thank you!

References

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