

Multipoint Conformal Blocks and their Vertex Integrable Systems

with I. Burić, S. Lacroix, L. Quintavalle & V. Schomerus

based on:

- [2112.10827],
- [2108.00023], *JHEP* 11 (2021) 182,
- [2009.11882], *JHEP* 10 (2021) 139,
- [1910.10427], *Phys.Rev.Lett.* 126 (2021) 2, 021602.

Motivation

Correlation functions in conformal field theory

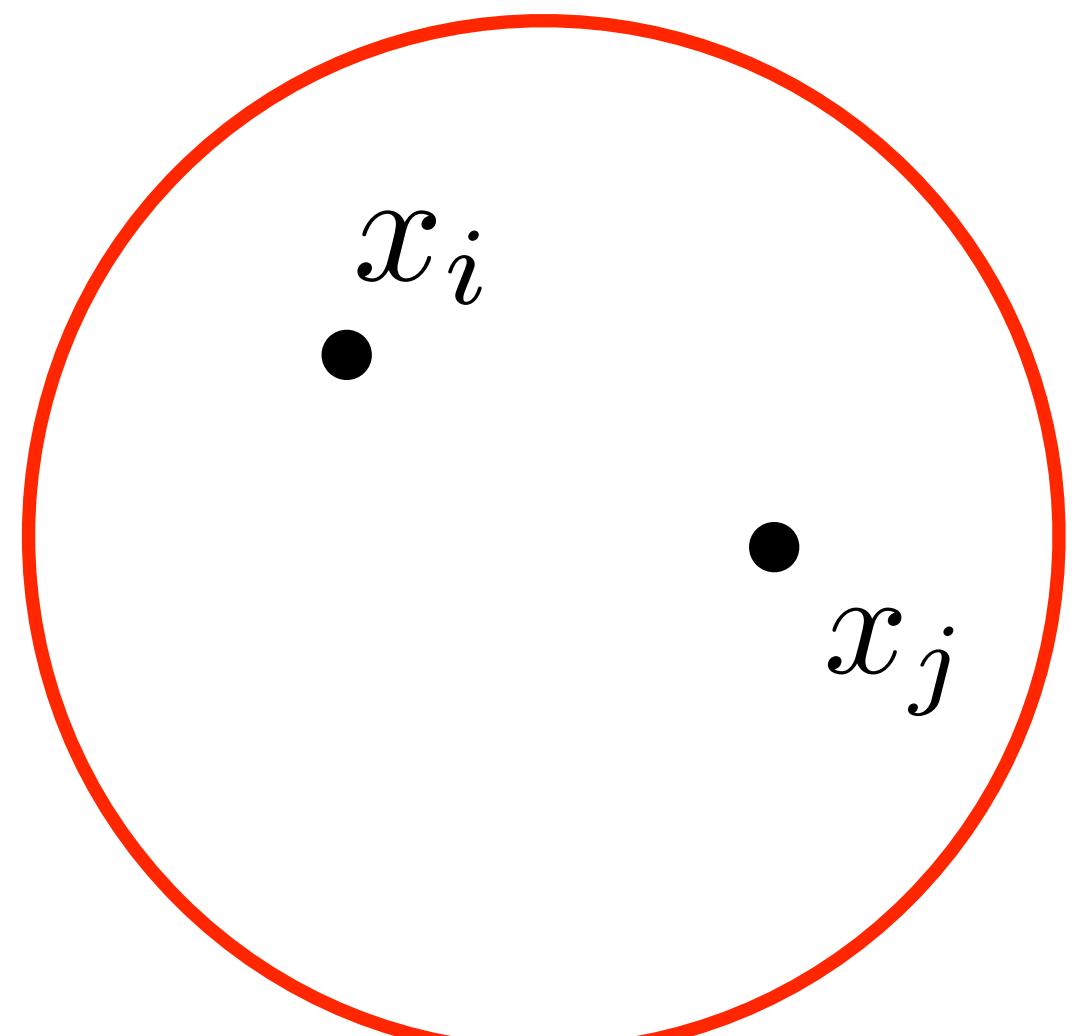
- Ingredients of a CFT in \mathbb{R}^d :

1. **Primary fields:** $x \rightarrow x + \xi(x)$, $ds^2 \rightarrow (1 - \omega_\xi(x))ds^2$, $\xi_A{}^B \in \mathfrak{so}(1, d+1)$

$$\mathcal{O}_i(x_i) \rightarrow \mathcal{O}_i(x_i) + \xi_A{}^B \underbrace{\underline{L_i(x_i, \partial_{x_i})}{}_A{}^B}_{\text{operator product expansion}} \mathcal{O}_i(x_i)$$

2. **Operator product expansion:**

$$\mathcal{O}_i(x_i)\mathcal{O}_j(x_j) = \sum_{[\mathcal{O}_k] \in [\mathcal{O}_i] \otimes [\mathcal{O}_j]} \sum^{\text{multiplicity}}_{n=1} C_{ijk}^{(n)} f_{ijk}^{(n)}(x_i - x_j, \partial_{x_j}) \mathcal{O}_k(x_j)$$



Motivation

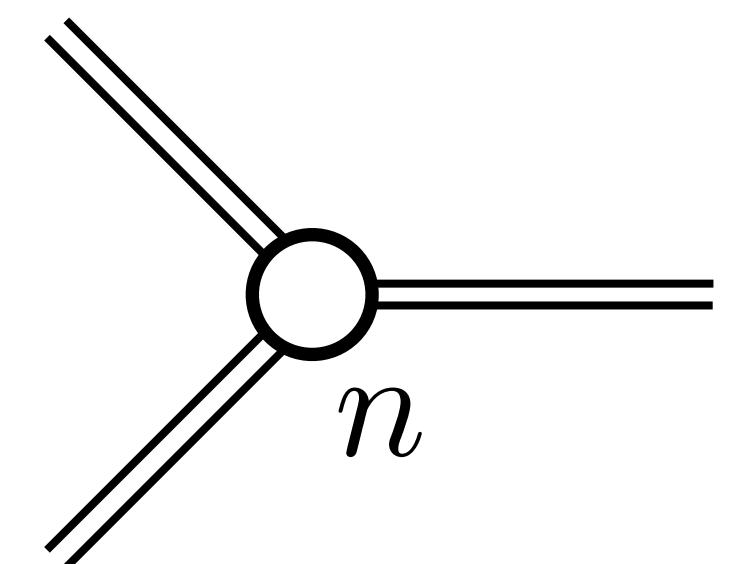
Correlation functions in conformal field theory

- OPE determines **all** correlation functions

$$\langle \mathcal{O}_1 \rangle = \begin{cases} 1 & \text{if } \mathcal{O}_1 = 1 \\ 0 & \text{otherwise} \end{cases}$$

$$\langle \widehat{\mathcal{O}_1 \mathcal{O}_2} \rangle = \begin{cases} C_{\mathcal{O}_1 \mathcal{O}_1 1}^{\mathcal{O}_1 \mathcal{O}_1 1} f_{\mathcal{O}_1 \mathcal{O}_1 1}(x_1 - x_2) & \text{if } \mathcal{O}_1 = \mathcal{O}_2 \\ 0 & \text{otherwise} \end{cases}$$

$$\langle \widehat{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3} \rangle = \sum_n C_{123}^{(n)} \underbrace{f_{123}^{(n)}(x_1 - x_2, \partial_{x_2})}_{\text{---}} f_{331}(x_2 - x_3) = \sum_n C_{123}^{(n)}$$



Motivation

Correlation functions in conformal field theory

- OPE determines **all** correlation functions

$$\langle \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3 \mathcal{O}_4 \rangle = \sum_{\mathcal{O}_k, n, n'} C_{12k}^{(n)} C_{k34}^{(n')} \quad \text{Diagram: } \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} \quad \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array}$$

\mathcal{O}_k

n n'

$$\dots \quad \dots \quad \dots$$
$$\langle \mathcal{O}_1 \dots \mathcal{O}_N \rangle = \sum_{\mathcal{O}_k, n, n'} \underline{C} \dots \underline{C} \quad \times (N-2) \quad \text{Diagram: } \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} \quad \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array}$$

\mathcal{O}_k

n n'

Motivation

Correlation functions in conformal field theory

- OPE determines **all** correlation functions

$$\langle \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3 \mathcal{O}_4 \rangle = \sum_{\mathcal{O}_k, n, n'} C_{12k}^{(n)} C_{k34}^{(n')} \quad \text{Diagram: } \mathcal{O}_k \text{ at } n, \mathcal{O}_k \text{ at } n'$$
$$\dots \quad \dots \quad \dots$$
$$\langle \mathcal{O}_1 \dots \mathcal{O}_N \rangle = \sum_{\mathcal{O}_k, n, n'} \underline{C \dots C} \times (N-2) \quad \text{Diagram: } N \text{ operators } \mathcal{O}_k \text{ connected in a circle}$$

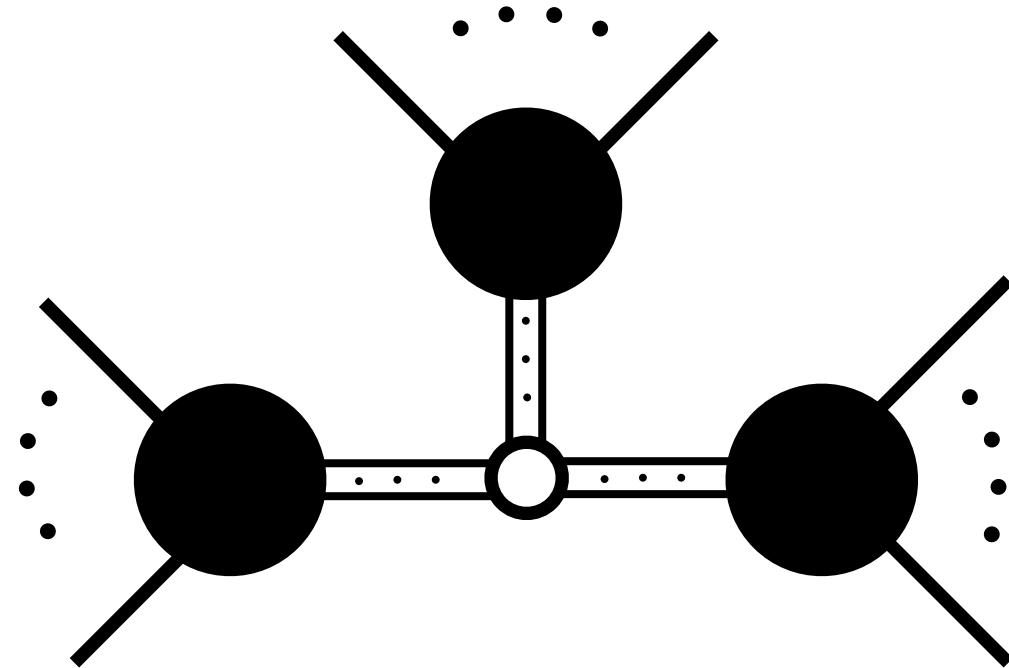
Integrability based approach N-point Gaudin models

[1910.10427]
[2009.11882]

basis of N-point
blocks

=

eigenbasis of
Gaudin Hamiltonians



$$H_I \psi = E_I \psi \quad [H_I, H_J] = 0$$

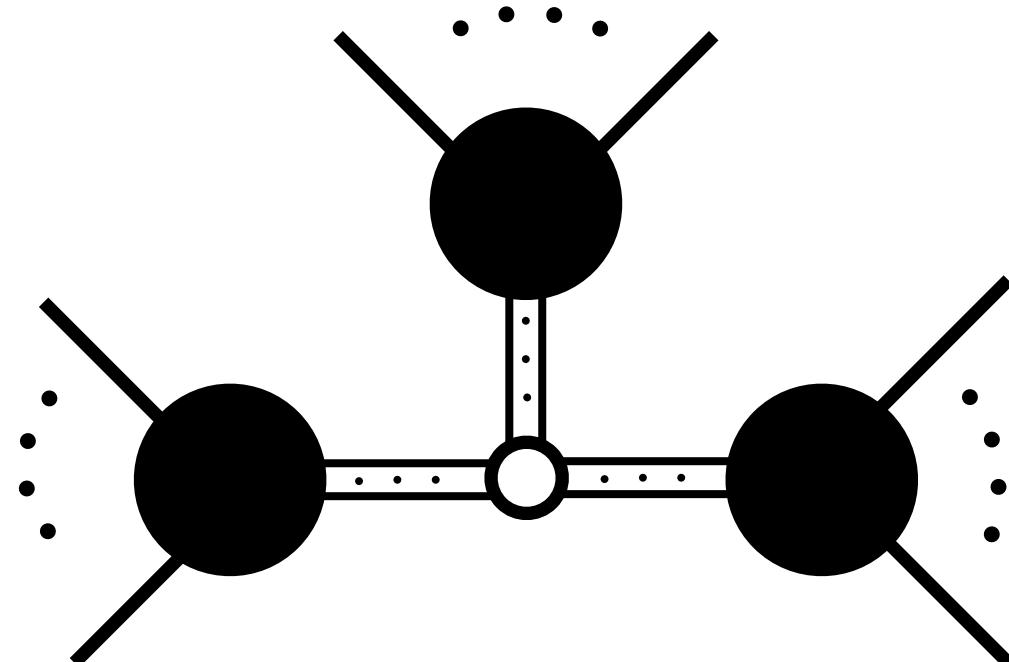
Integrability based approach N-point Gaudin models

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$$H_I \psi = E_I \psi \quad [H_I, H_J] = 0$$
$$\sum_{i_1, \dots, i_p=1}^N c_{i_1 \dots i_p}^{(I)} \text{tr} L_{i_1}(x_{i_1}, \partial_{x_{i_1}}) \dots L_{i_p}(x_{i_p}, \partial_{x_{i_p}})$$

Integrability based approach N-point Gaudin models

[1910.10427]
[2009.11882]

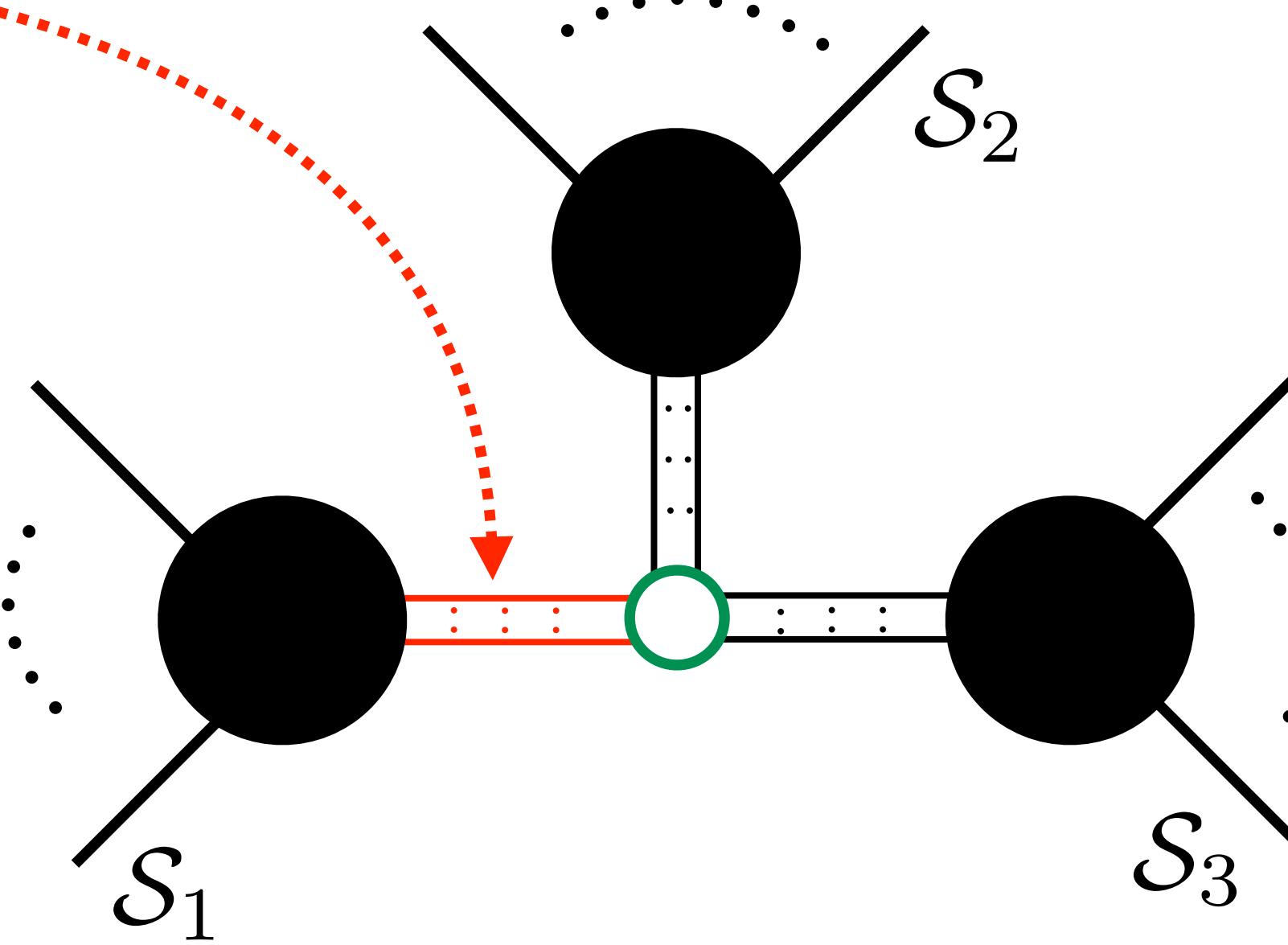
basis of N-point
blocks

=

eigenbasis of
Gaudin Hamiltonians

$$\left\{ \text{tr} \left(\sum_{i \in \mathcal{S}_1} L_i \right)^p \right\}_{p=2,4,\dots}$$

⇒ quantum
numbers
of \mathcal{O}_k



Integrability based approach N-point Gaudin models

[1910.10427]
[2009.11882]

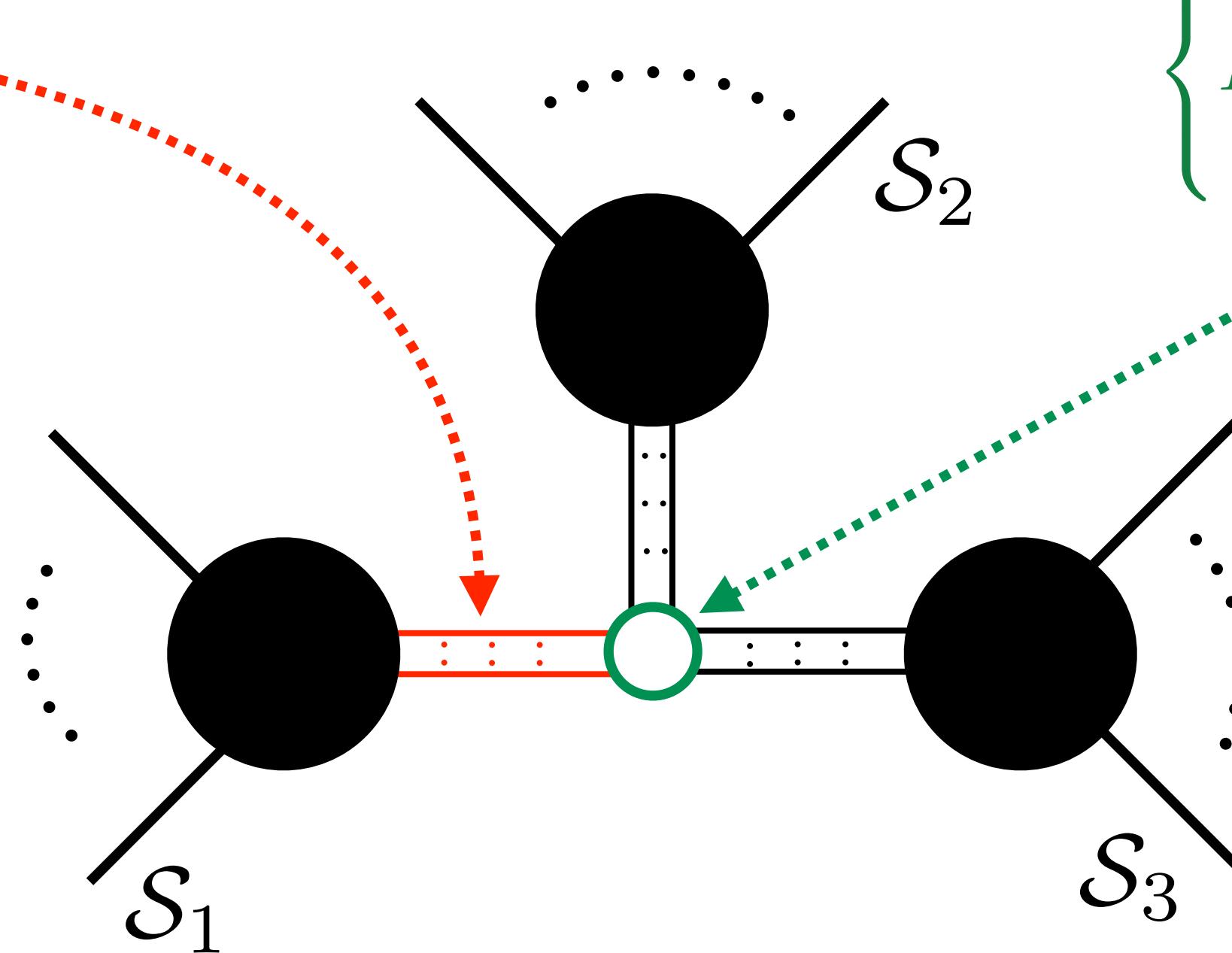
basis of N-point
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$$\left\{ \text{tr} \left(\sum_{i \in \mathcal{S}_1} L_i \right)^p \right\}_{p=2,4,\dots}$$

\Rightarrow quantum
numbers
of \mathcal{O}_k

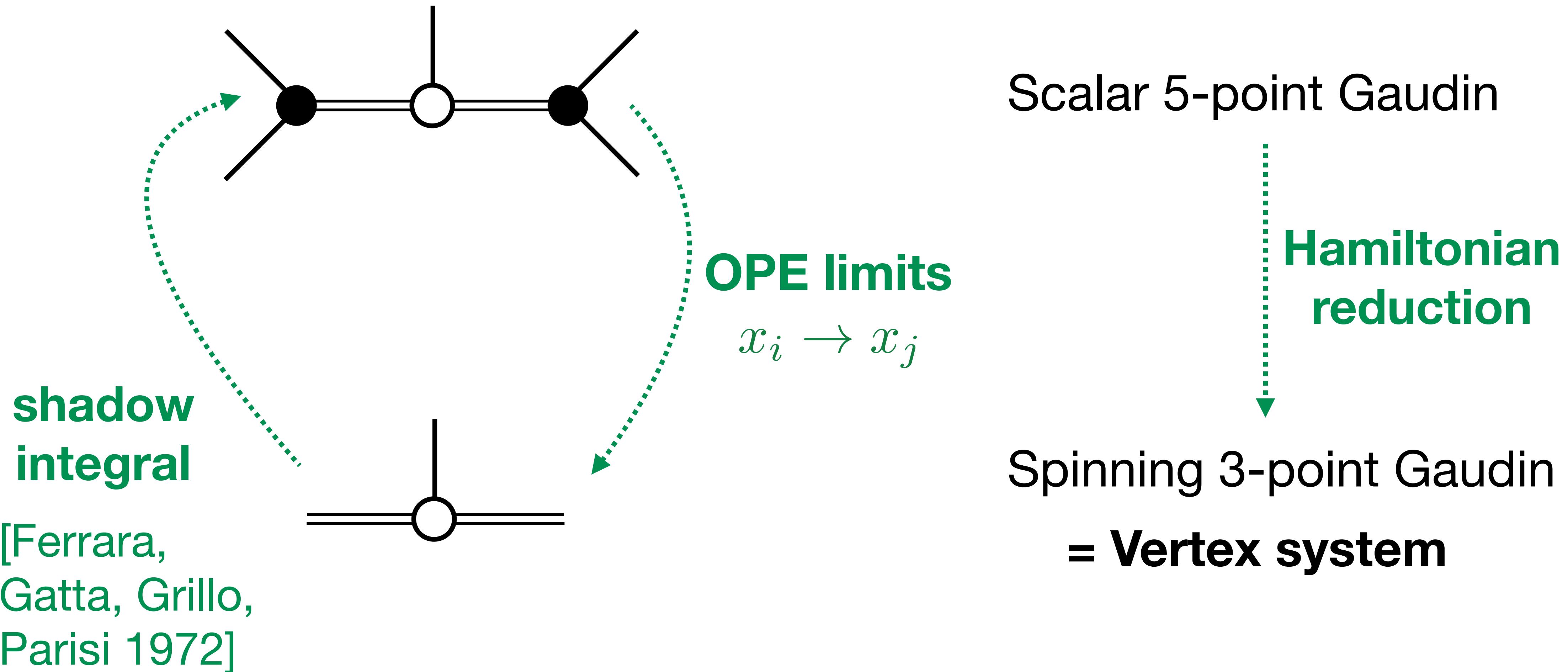


$$\left\{ H_{\text{vertex}}^{(1)} = \text{tr} \left(\sum_{i \in \mathcal{S}_1} L_i \right)^3 \left(\sum_{i \in \mathcal{S}_2} L_i \right), \dots \right\}$$

\Rightarrow basis of
 $\{t_n : [\mathcal{O}_k] \hookrightarrow [\mathcal{O}_i] \otimes [\mathcal{O}_j]\}$

Integrability based approach

From N-point blocks to vertex systems and back

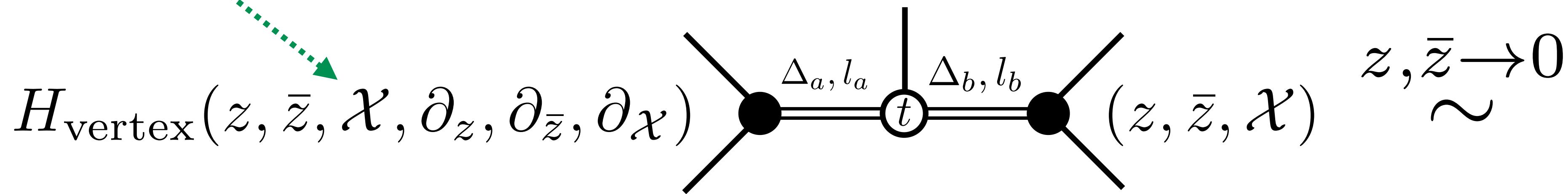
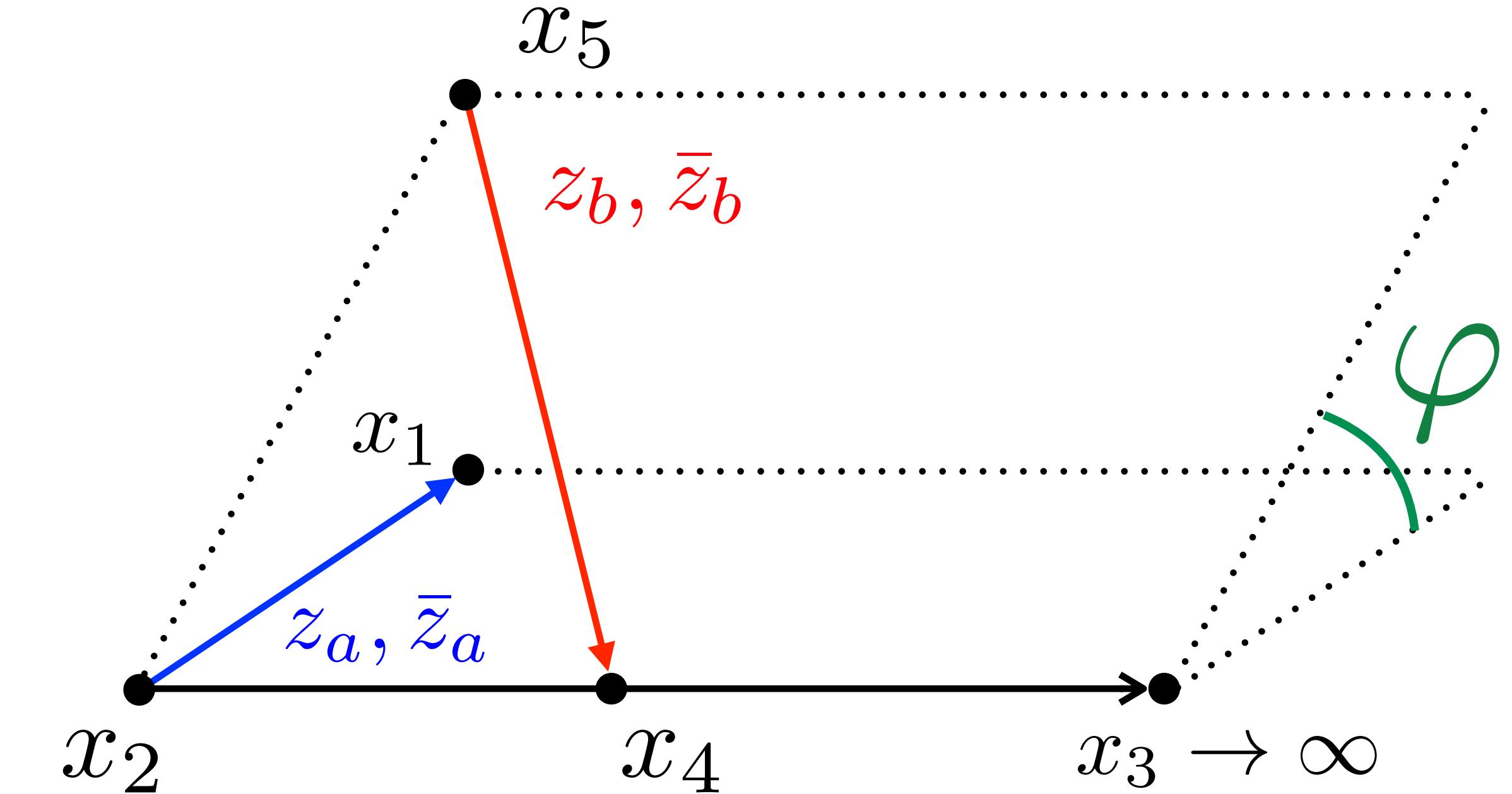


Example

Scalar five-point function

- Use conformal symmetry to fix frame:

$$\chi = \cos^2 \frac{\varphi}{2}$$



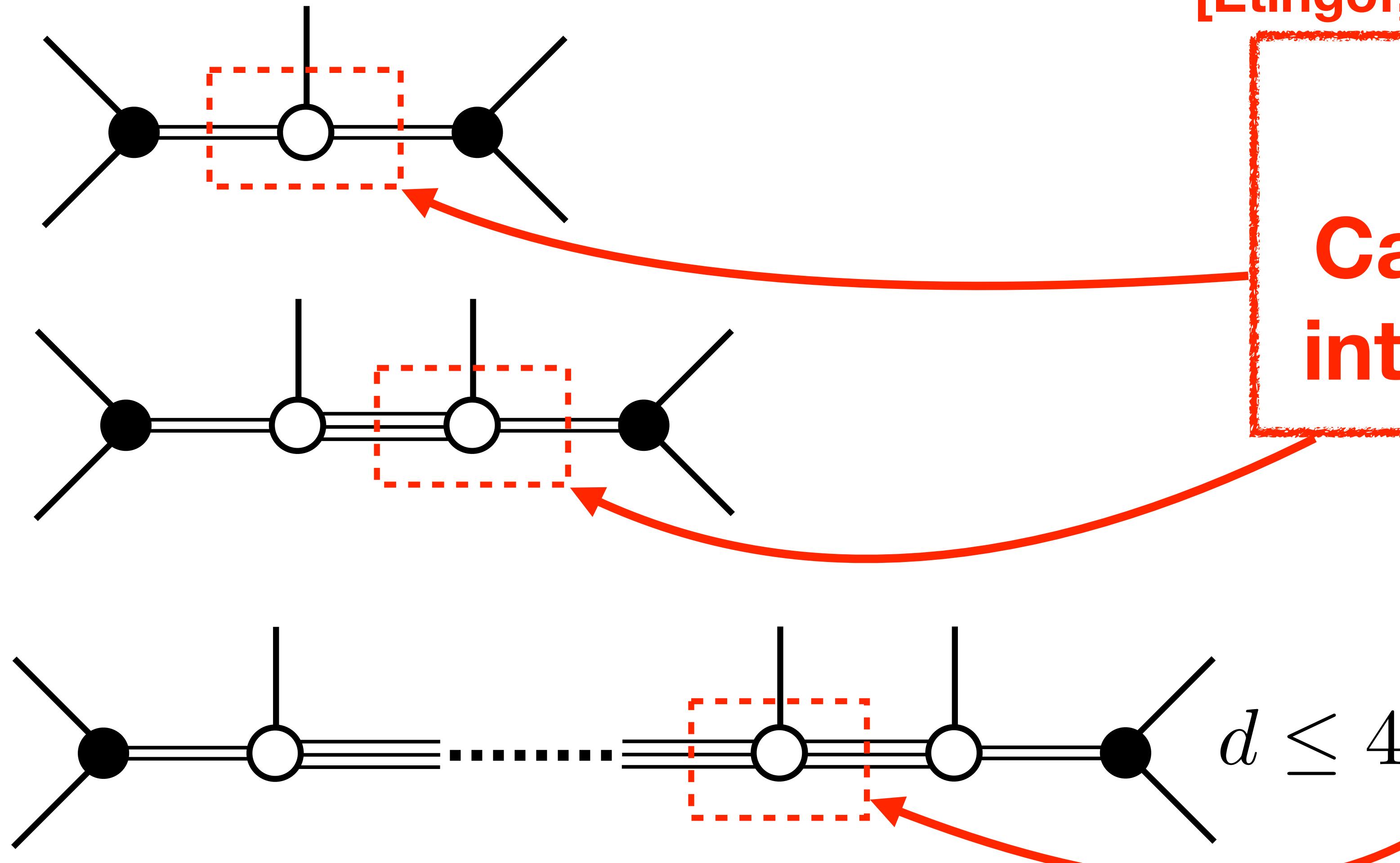
$$z_a^{\frac{\Delta_a + l_a}{2}} \bar{z}_a^{\frac{\Delta_a - l_a}{2}} \left[H_{\text{reduced}}(\chi, \partial \chi) \left. \right|_{\substack{\Delta_a, l_a \\ \Delta_b, l_b}} (\chi) \right] z_b^{\frac{\Delta_b + l_b}{2}} \bar{z}_b^{\frac{\Delta_b - l_b}{2}}$$

Main result

Comb channel vertex systems

[2108.00023]

[Etingof, Felder, Ma, Veselov 2010]



Elliptic $\mathbb{Z}/4\mathbb{Z}$

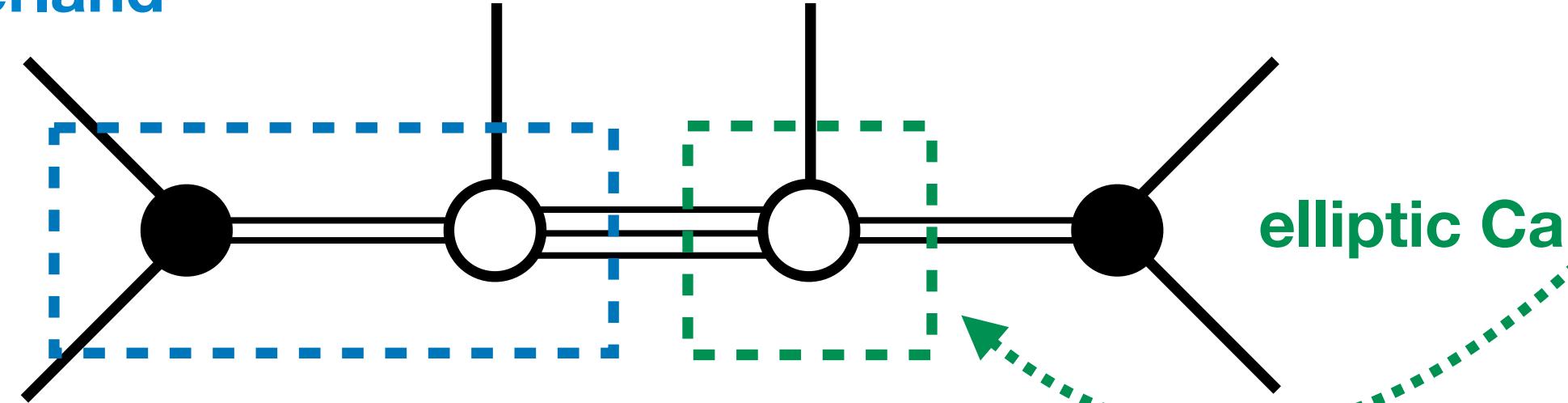
Calogero-Moser
integrable model

$d \leq 4$

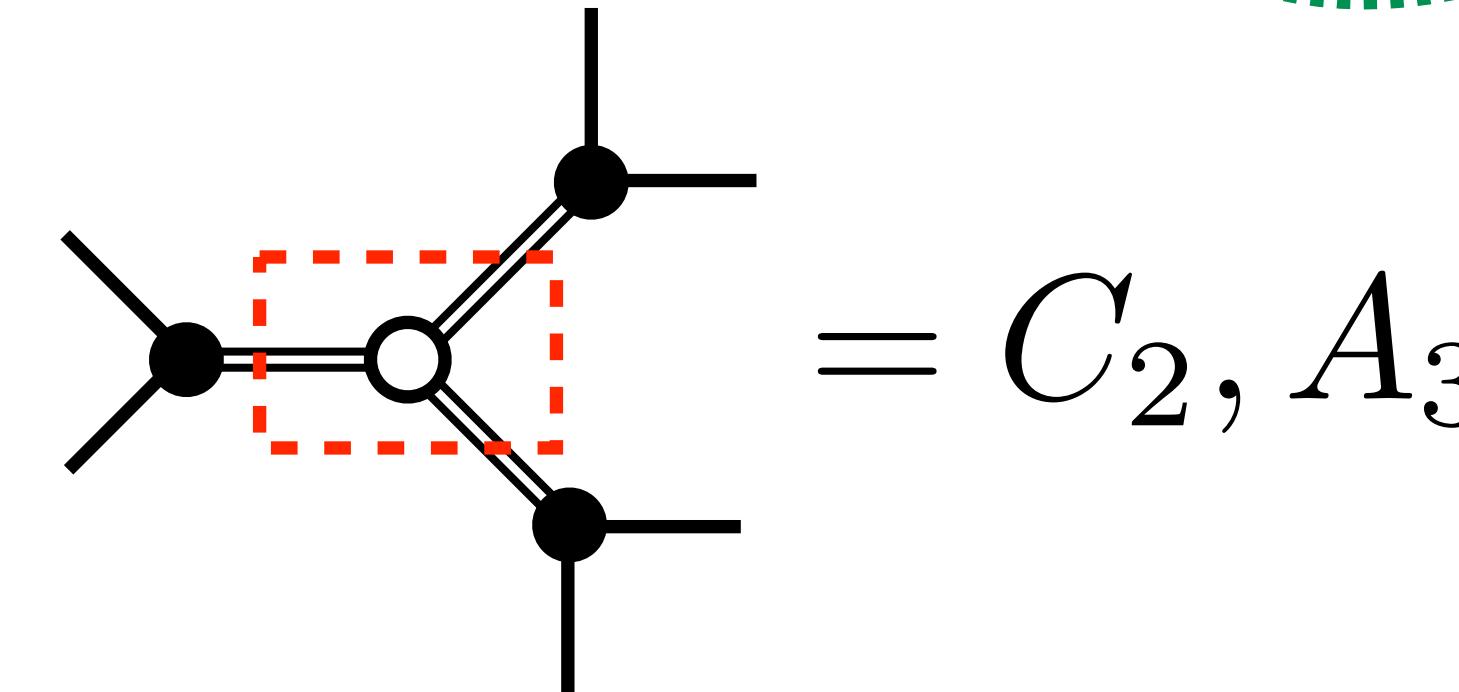
Ongoing research

spinning Calogero-Sutherland
see [2112.10827]

- Solving N-point Gaudin models:



- OPE channels of different topology:
using quivers & cluster algebras



- N-point blocks in lightcone limits:

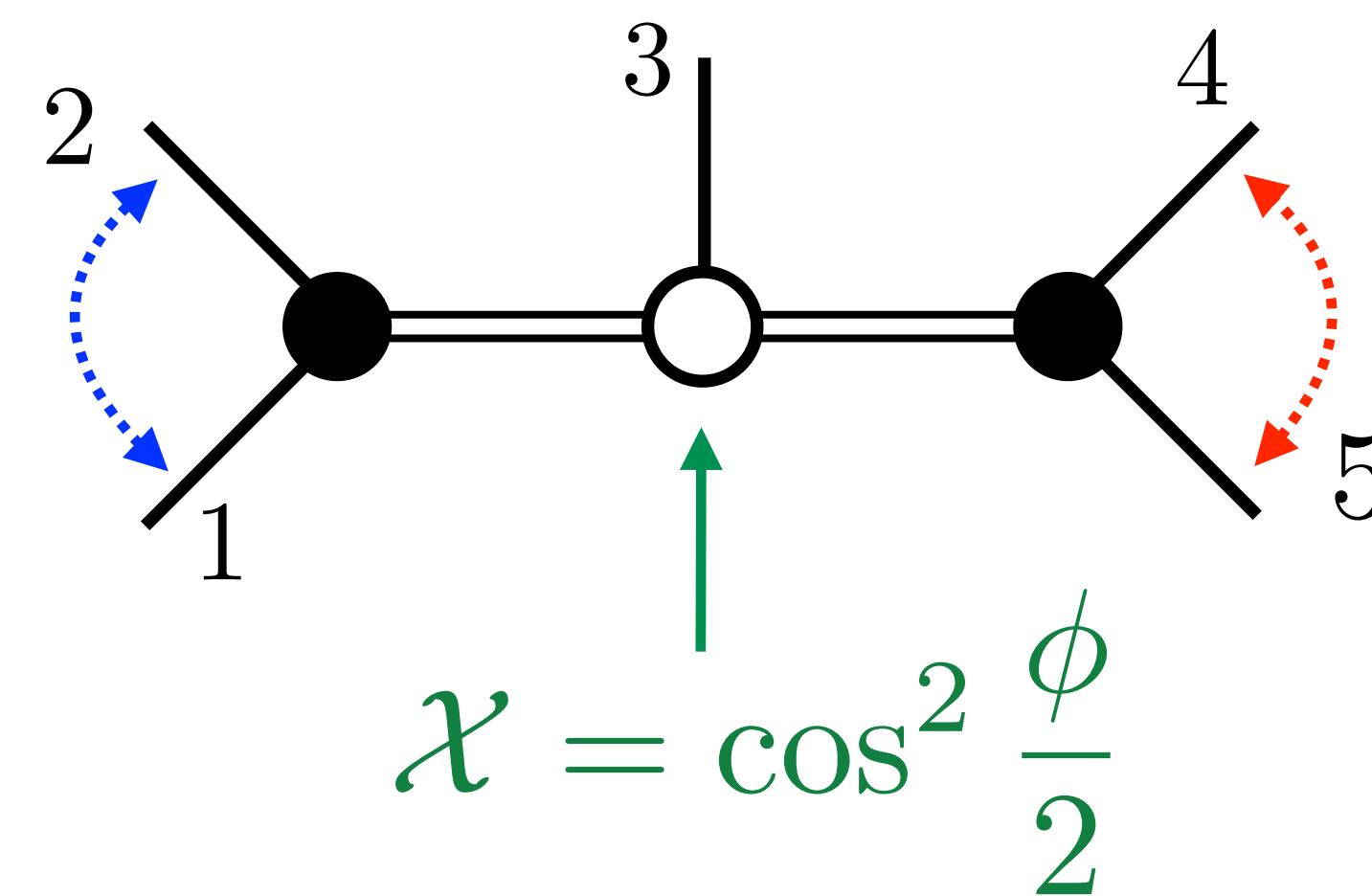
$$\bar{z}_{a,b} \rightarrow 0, \quad z_{a,b} \rightarrow 1, \quad l_{a,b} \rightarrow \infty$$

A diagram showing three nodes connected by solid black lines. One node is black (top-left) and two are white (top-right and bottom-right). Above the top-right white node is the label Δ_a, l_a . Above the bottom-right black node is the label Δ_b, l_b . Below the bottom-right black node is the equation $t(\chi) = \chi^n$. To the right of the diagram is a mathematical expression involving χ^n , $\prod_{s=a,b}$, $\frac{\bar{z}_s^{\frac{\Delta_s - l_s}{2}}}{(1 - z_s)^{\frac{n}{2} + \frac{\Delta_s - l_s - \Delta_3}{4}}} K_{n+\frac{\Delta_s - l_s - \Delta_3}{2}} (2l_s \sqrt{1 - z_s})$.

Appendix 1

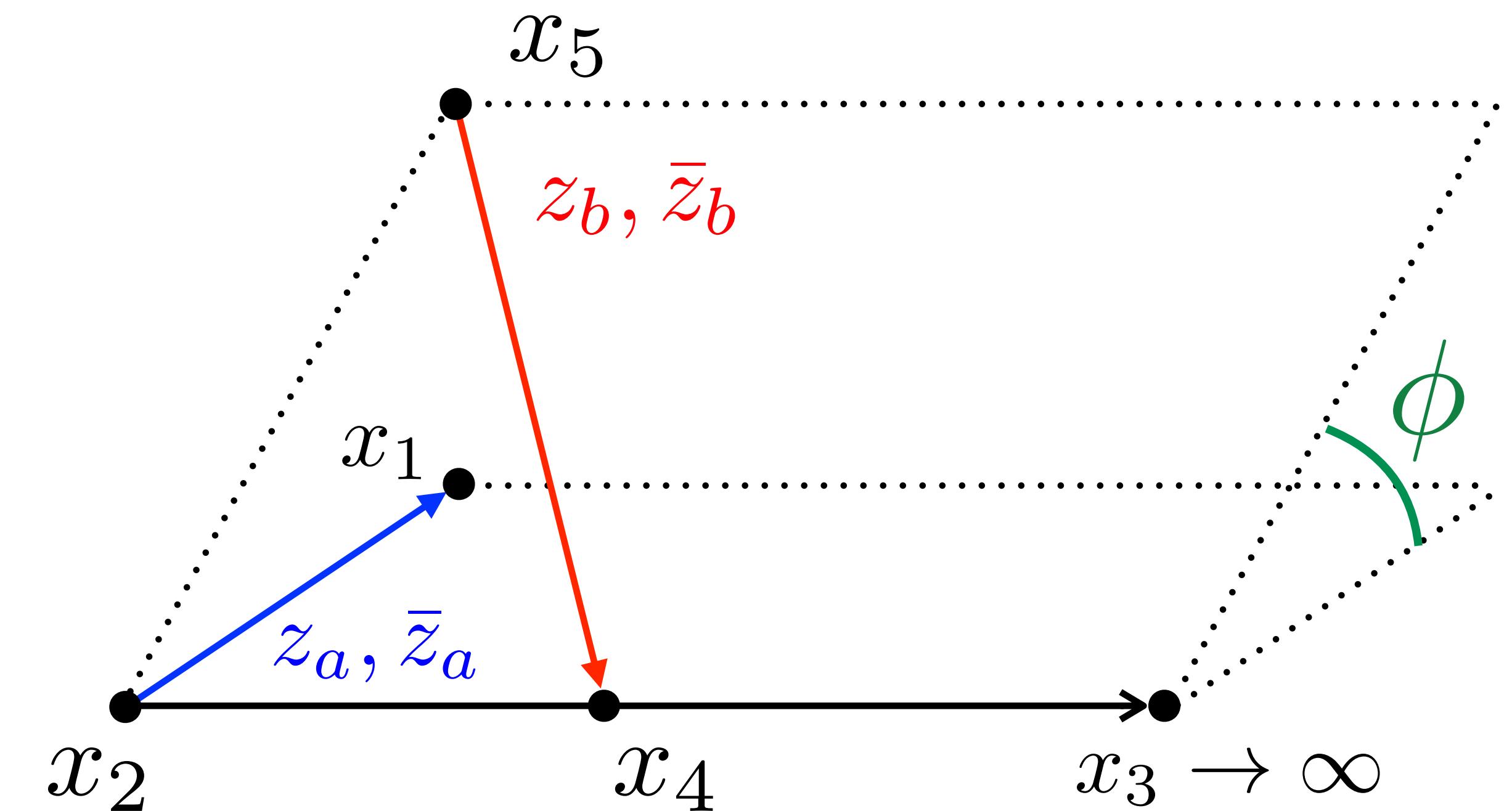
OPE limit

- Most singular contribution to OPE:
- In a N-point function:



$$\phi_1(x_1)\phi_2(x_2) \xrightarrow{x_1 \rightarrow x_2}$$

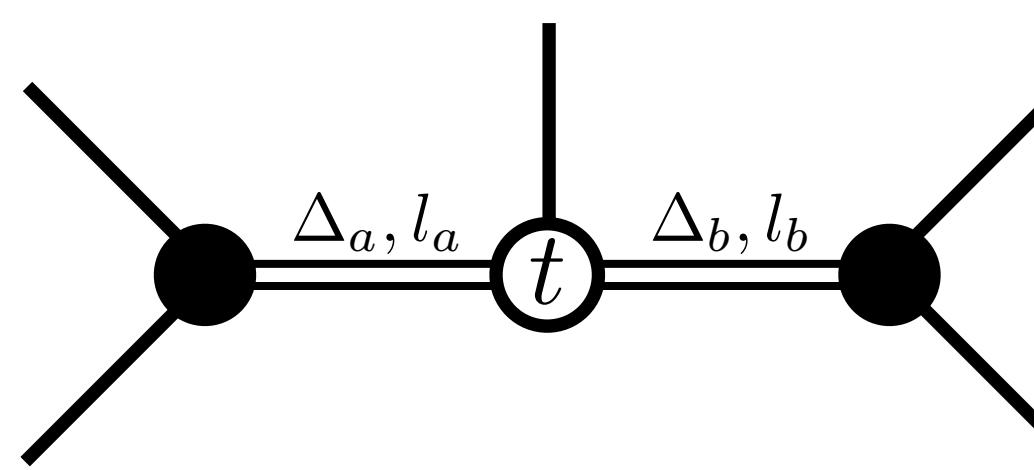
$$(x_1 - x_2)^{2q_a} C_{\phi_1 \phi_2} \mathcal{O} \mathcal{D}(x_1 - x_2, \partial_{x_2}) \mathcal{O}_{\Delta_a, l_a}(x_2) + (\text{less singular})$$



Appendix 1

OPE limit

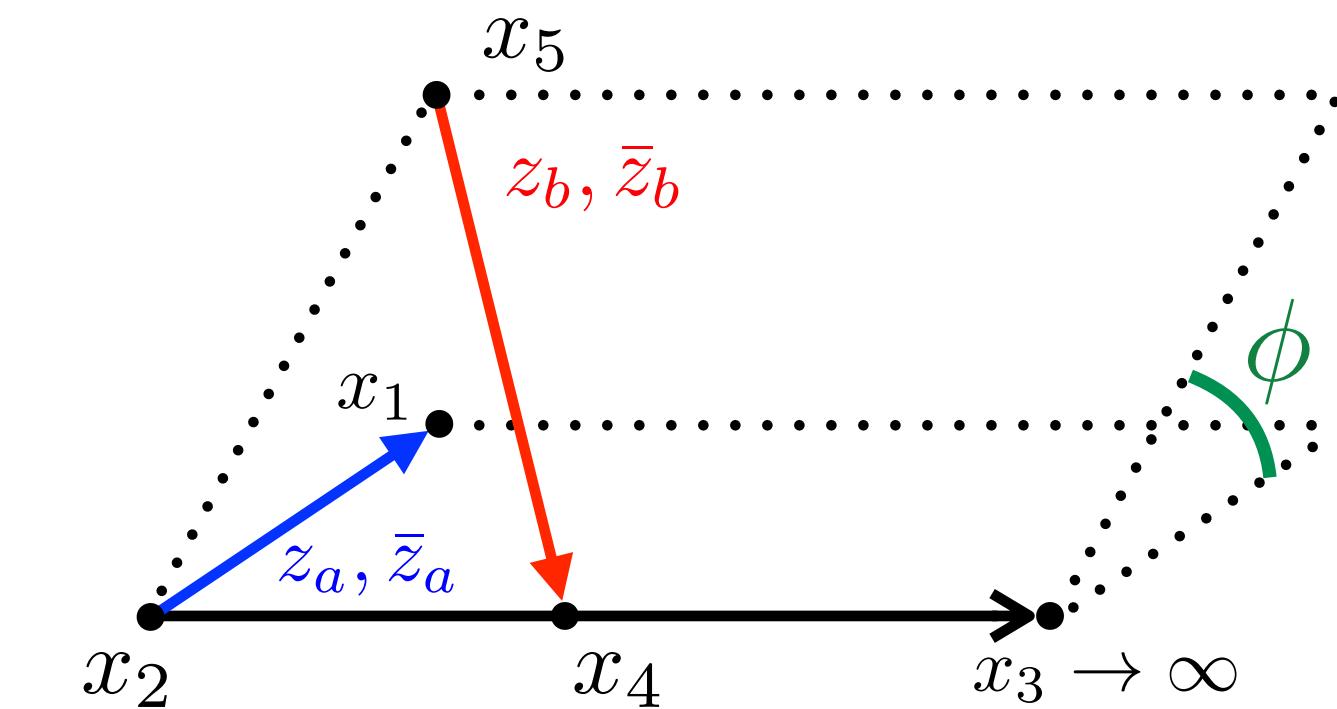
Ansatz:



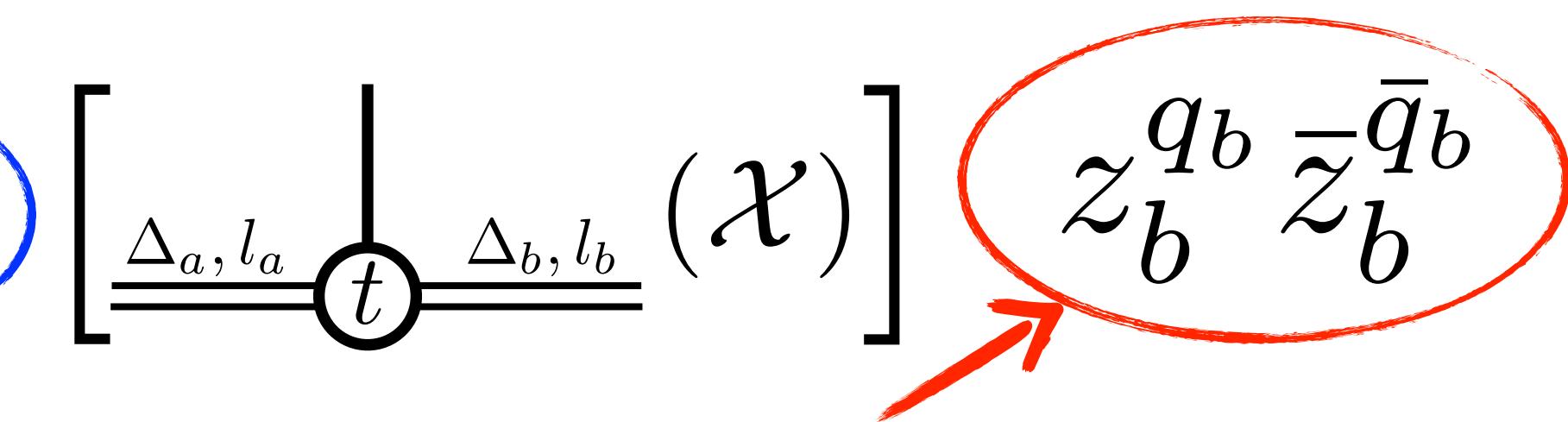
$$(z, \bar{z}, \chi) \quad z, \bar{z} \rightarrow 0 !$$

$$\text{tr } (L_1 + L_2)^{2,4} \sim \text{Cas}^{2,4}(\Delta_a, l_a)$$

$$\Rightarrow q_a, \bar{q}_a = \frac{\Delta_a \pm l_a}{2}$$



$$\chi = \cos^2 \frac{\phi}{2}$$

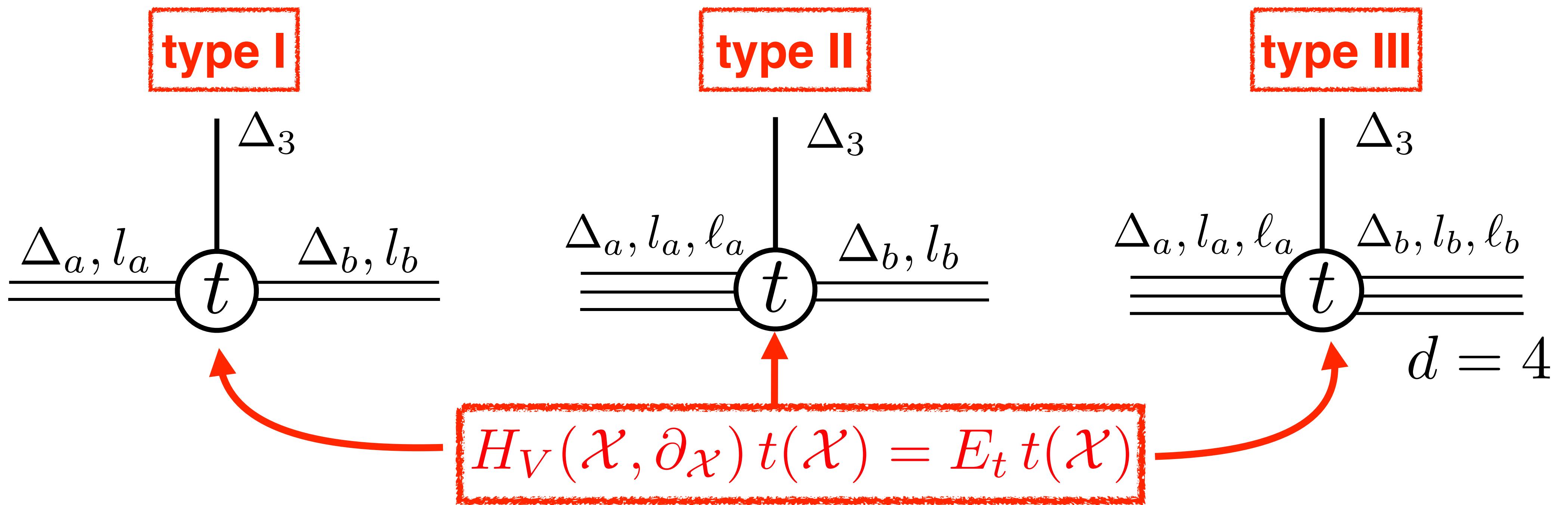


$$\text{tr } (L_4 + L_5)^{2,4} \sim \text{Cas}^{2,4}(\Delta_b, l_b)$$

$$\Rightarrow q_b, \bar{q}_b = \frac{\Delta_b \pm l_b}{2}$$

Appendix 2

From vertex system to Calogero-Moser model



- = OPE limits & shadow integrands of N-point blocks
- = Gaudin/Hitchin systems on the three-punctured sphere

Appendix 2

From vertex system to Calogero-Moser model

- **Conjugation by** $\Theta = \Theta_0 \chi^{\frac{l_1+l_2-2(\ell_1+\ell_2)+\Delta_3+(1-d)/2}{4}} (1-\chi)^{\frac{l_1+l_2-2(\ell_1+\ell_2)-\Delta_3+(1+d)/2}{4}}$

- **Parameters:** $a_\nu := -11 + \sum_{1 \leq i < j \leq 4} m_{i,\nu} m_{j,\nu}$,

$$\sum_{i=1}^4 m_{i,\nu} = 6$$

$$b_\nu := \frac{1}{2} \left(-a_\nu - 6 + \sum_{1 \leq i < j < k \leq 4} m_{i,\nu} m_{j,\nu} m_{k,\nu} \right)$$

$$c_\nu := \prod_{i=1}^4 m_{i,\nu}$$

$$m_{i,2} = m_{i,3}$$

$$m_{1,2} = k + 1$$

$$m_{2,2} = -k$$

$$m_{3,2} = k + 4$$

$$m_{4,2} = -k + 2$$

Appendix 2

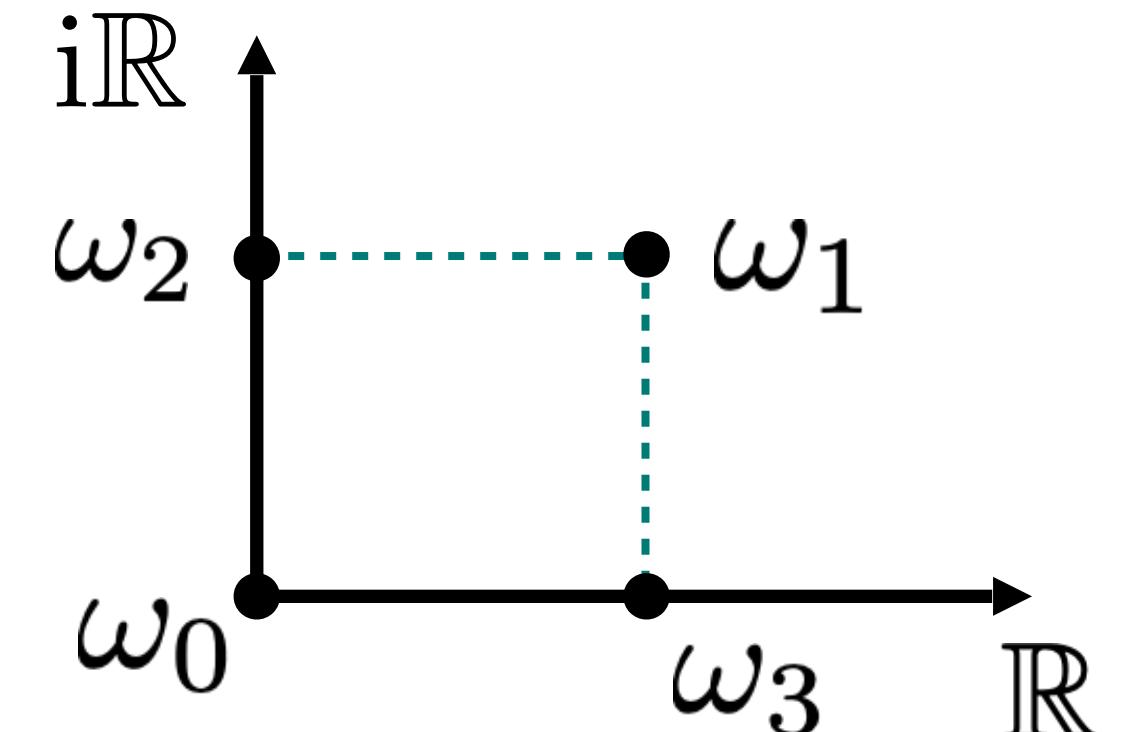
From vertex system to Calogero-Moser model

[c.f. Etingof, Felder, Ma, Veselov 2010 – eq. (4.3)]

$$\Theta^{-1} H_V \Theta + \text{const} =$$

$$\partial_\zeta^4 + \sum_{\nu=0}^3 \left[a_\nu \wp(\zeta - \omega_\nu) \partial_\zeta^2 + b_\nu \wp'(\zeta - \omega_\nu) \partial_\zeta + c_\nu \wp^2(\zeta - \omega_\nu) \right]$$

$$+ \wp(\omega_3)(a_0 - a_1)k(k+1)[\wp(\zeta - \omega_2) - \wp(\zeta - \omega_3)]$$



$$\chi = \left(1 - \frac{\wp^2(\zeta)}{\wp^2(\omega_3)} \right)^{-1}$$

$$a_\nu, b_\nu, c_\nu = \text{polynomial}(m_{i,\nu})$$

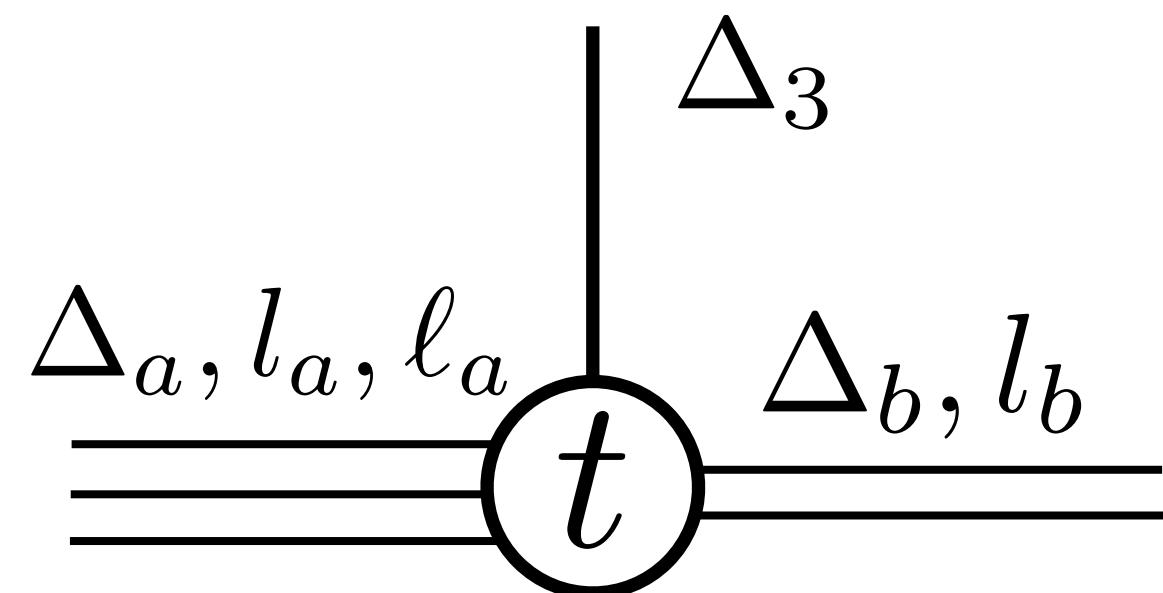
$$\sum_{i=1}^4 m_{i,\nu} = 6$$

$$m_{i,2} = m_{i,3} \propto k$$

$$\mathbb{Z}/4\mathbb{Z} : \zeta \mapsto i^n \zeta$$

Appendix 2

Calogero-Moser multiplicities for type I,II vertices

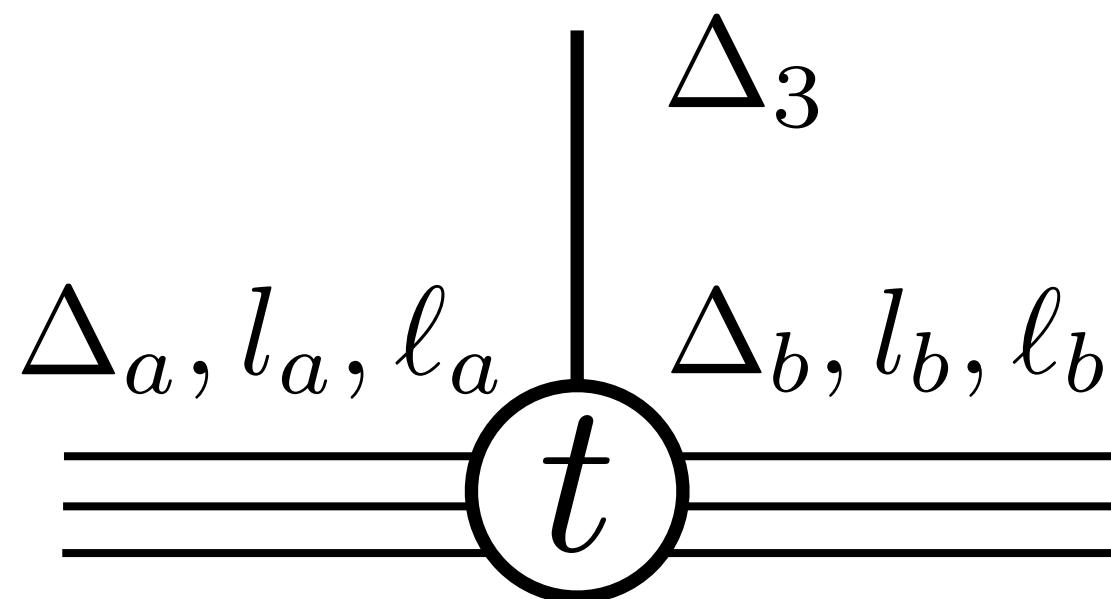


$$\begin{aligned}
 m_{1,0} &= 3 \frac{5-d}{2} - (l_a + l_b) - \Delta_3 - 2\ell_a , & k = |l_a - l_b| - \frac{1}{2} \\
 m_{2,0} &= \frac{d-1}{2} - (l_a + l_b) - \Delta_3 + 2\ell_a , \\
 m_{3,0} &= \frac{d-1}{2} + (l_a + l_b) + \Delta_3 + 2(\Delta_a - \Delta_b) , \\
 m_{4,0} &= \frac{d-1}{2} + (l_a + l_b) + \Delta_3 - 2(\Delta_a - \Delta_b) , \\
 m_{1,1} &= -5 \frac{d-3}{2} - (l_a + l_b) + \Delta_3 - 2\ell_a , \\
 m_{2,1} &= -\frac{d+1}{2} - (l_a + l_b) + \Delta_3 + 2\ell_a , \\
 m_{3,1} &= -\frac{d+1}{2} + (l_a + l_b) - \Delta_3 + 2(\Delta_a + \Delta_b) , \\
 m_{4,1} &= \frac{7d-1}{2} + (l_a + l_b) - \Delta_3 - 2(\Delta_a + \Delta_b) ,
 \end{aligned}$$

Appendix 2

Calogero-Moser multiplicities for type III vertices

$$d = 4$$



$$l = \frac{j + \bar{j}}{2}, \quad \ell = \frac{j - \bar{j}}{2}$$

$$\begin{aligned}
m_{1,0} &= \frac{3}{2} - (l_a + l_b) - \Delta_3 - 2(\ell_a - \ell_b), & k = |l_a - l_b| - \frac{1}{2} \\
m_{2,0} &= \frac{3}{2} - (l_a + l_b) - \Delta_3 + 2(\ell_a - \ell_b), \\
m_{3,0} &= \frac{3}{2} + (l_a + l_b) + \Delta_3 + 2(\Delta_a - \Delta_b), \\
m_{4,0} &= \frac{3}{2} + (l_a + l_b) + \Delta_3 - 2(\Delta_a - \Delta_b), \\
m_{1,1} &= -\frac{5}{2} - (l_a + l_b) + \Delta_3 - 2(\ell_a + \ell_b), \\
m_{2,1} &= -\frac{5}{2} - (l_a + l_b) + \Delta_3 + 2(\ell_a + \ell_b), \\
m_{3,1} &= -\frac{5}{2} + (l_a + l_b) - \Delta_3 + 2(\Delta_a + \Delta_b), \\
m_{4,1} &= \frac{27}{2} + (l_a + l_b) - \Delta_3 - 2(\Delta_a + \Delta_b).
\end{aligned}$$

Appendix 3

Scalar N-point kinematics in embedding space

→ **Embedding space:** $x \in \mathbb{R}^d \rightarrow X \in \mathbb{R}^{1,d+1}, \quad X^2 = 0, \quad X \sim \lambda X$

→ **Scalar primary:** $\phi_i(\lambda_i X_i) = \lambda_i^{-\Delta_i} \phi_i(X_i), \quad (L_i)^A{}_B = X_i^A \frac{\partial}{\partial X_i^B} - [A \leftrightarrow B]$

→ **N-point function:** $\langle \phi_1 \dots \phi_N \rangle = \text{homogeneous}(X_i \cdot X_j) =$

$$\prod_{i < j} (X_i \cdot X_j)^{a_{ij}} F \left(\left\{ u_I = \frac{\prod X_i \cdot X_j}{\prod X_i \cdot X_j} \right\}_{I=1}^{n_{\text{CR}}} \right)$$

→ **Counting cross-ratios:** $n_{\text{CR}} = \begin{cases} N(N-1)/2 - N & \text{if } d+2 \geq N \\ Nd - (d+2)(d+1)/2 & \text{if } d+2 \leq N \end{cases}$

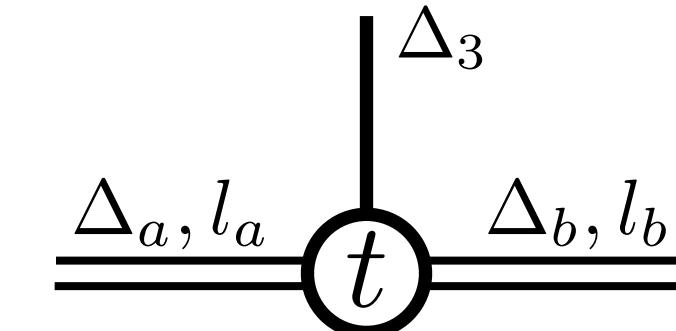
Appendix 3

Spinning three-point kinematics in embedding space

→ **Spinning embedding space:** $X, Z \in \mathbb{R}^{1,d+1}$ $X^2 = X \cdot Z = Z^2 = 0$

→ **Spinning primary:** $\mathcal{O}_{\Delta,l}(\lambda X, \alpha Z + \beta X) = \lambda^{-\Delta} \alpha^l \mathcal{O}_{\Delta,l}(X, Z)$

→ **Spinning three-point function:**



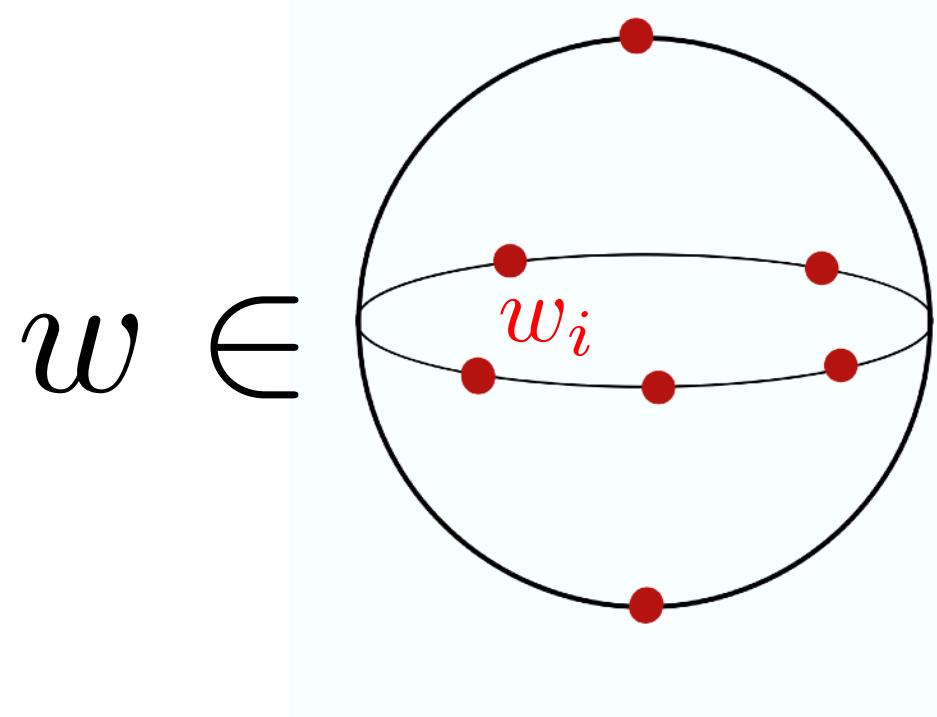
$$\begin{aligned} \langle \mathcal{O}_{\Delta_a, l_a}(X_a, Z_a) \phi_3(X_3) \mathcal{O}_{\Delta_b, l_b}(X_b, Z_b) \rangle = \\ (X_1 \cdot X_2)^{a_{12}} (X_2 \cdot X_3)^{a_{23}} (X_3 \cdot X_1)^{a_{31}} [X_3(X_a \wedge Z_a) X_b]^{b_1} [X_3(X_b \wedge Z_b) X_a]^{b_2} t(\mathcal{X}) \end{aligned}$$

$$\mathcal{X} = \frac{[X_3(X_a \wedge Z_a) X_b] [X_3(X_b \wedge Z_b) X_a]}{[(X_a \wedge Z_a) \cdot (X_b \wedge Z_b)] (X_3 \cdot X_a) (X_3 \cdot X_b)}$$

Appendix 4: N-point Gaudin Models

Lax matrix and Hamiltonians

- Lax matrix: $\mathcal{L}^A{}_B(w) = \sum_{i=1}^N \frac{L_i(x_i, \partial_{x_i})^A{}_B}{w - w_i} \in \bigotimes_{i=1}^N [\mathcal{O}_i] \otimes \mathfrak{so}(1, d+1)$
 - Invariants: $\mathcal{H}_p(w) = \text{tr } \mathcal{L}(w)^p + O(\mathcal{L}^{p-1}) \quad [\mathcal{H}_p(w), \mathcal{H}_q(w')] = 0$
 - Obtain Gaudin Hamiltonians by **partial fraction decomposition**, e.g.
- $$\mathcal{H}_2(w) = \sum_{i=1}^N \left\{ \frac{\text{tr } L_i^2}{(w - w_i)^2} + \frac{\sum_{j \neq i} \frac{\text{tr } L_i L_j}{w_i - w_j}}{w - w_i} \right\}$$

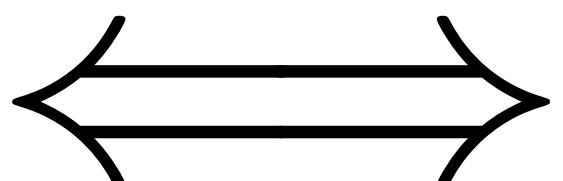
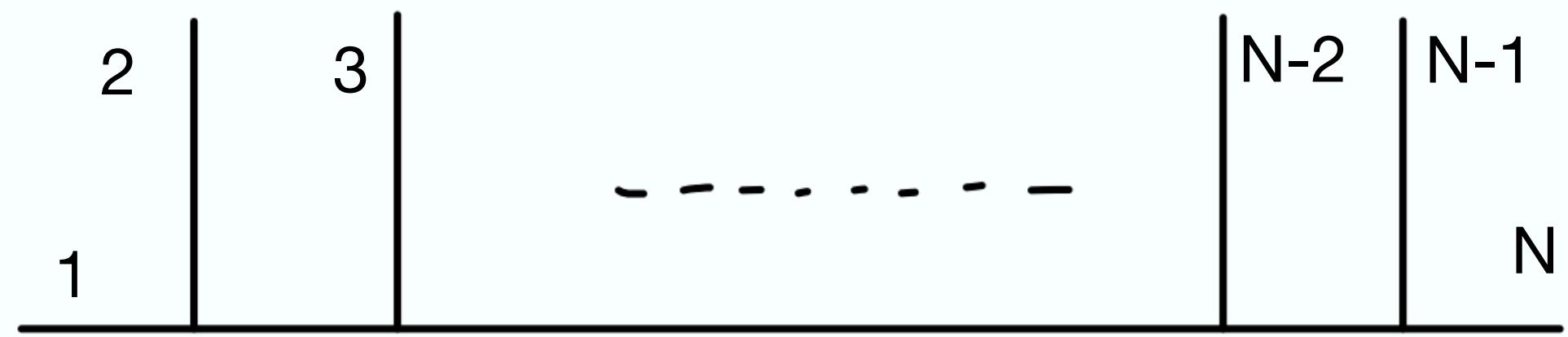


Appendix 4: N-point Gaudin Models

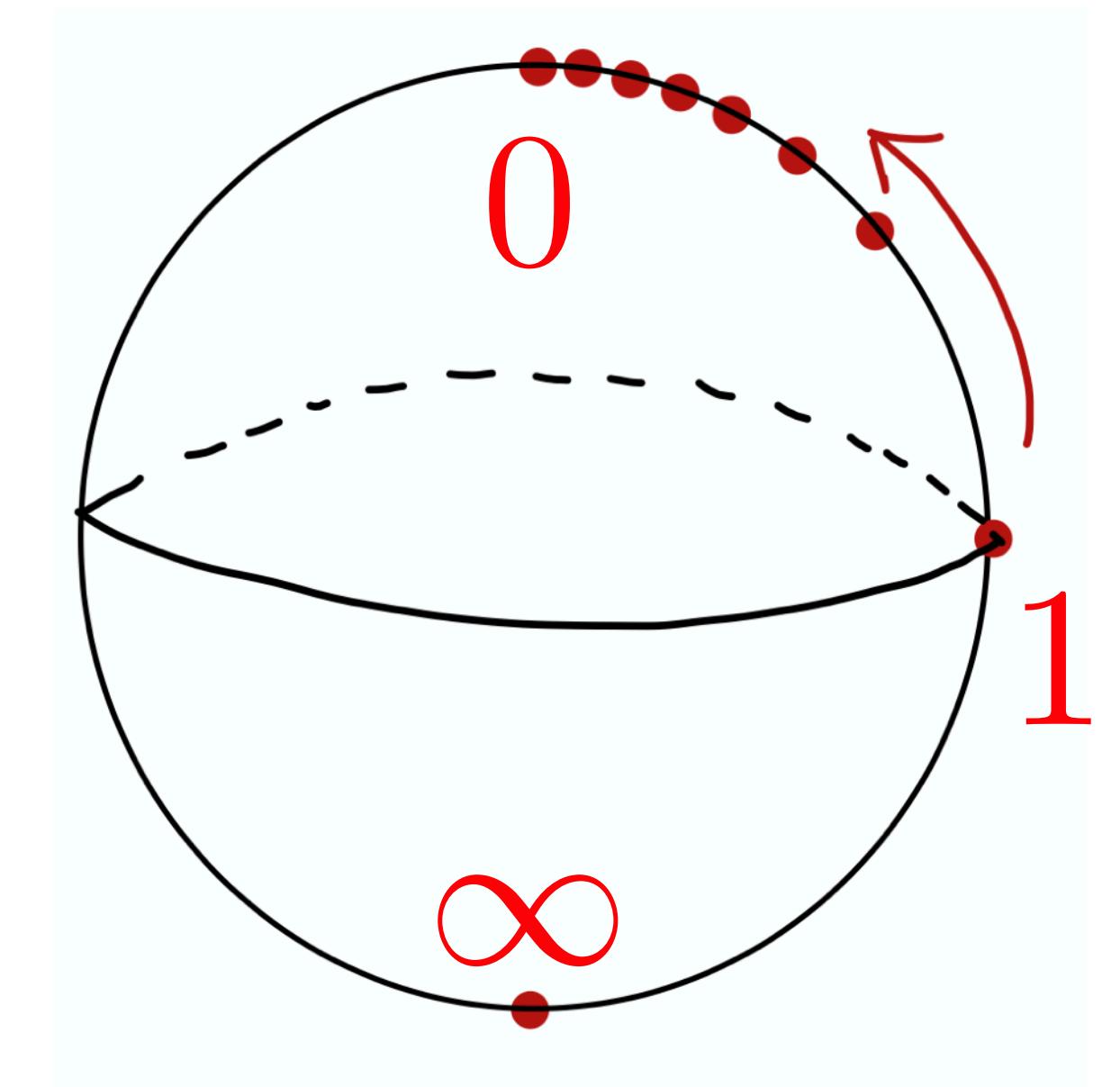
OPE channel limits

- take limits of **colliding punctures**: $w_i - w_j = \epsilon \rightarrow 0$ $\epsilon^{p\lambda} \mathcal{H}_p(\epsilon^\lambda w) \rightarrow \text{finite}$

- Example: **comb channel**:



$$w_1 = 0, w_2 = \epsilon^{N-3}, \dots, w_{N-2} = \epsilon, w_{N-1} = 1, w_N = \infty$$



$$r = 2, 3, \dots, N-2 \quad \mathcal{H}_p^{[r]}(w) = \lim_{\epsilon \rightarrow 0} \epsilon^{p(N-r-2)} \mathcal{H}_p(\epsilon^{N-r-2} w)$$