

# Functional reduction of one-loop Feynman integrals with arbitrary masses

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A method of functional reduction for the dimensionally regularized one-loop Feynman integrals with massive propagators is described in detail. The method is based on a repeated application of the functional relations proposed by the author. Explicit formulae are given for reducing one-loop scalar integrals to a simpler ones, the arguments of which are the ratios of polynomials in the masses and kinematic invariants. We show that a general scalar  $n$ -point integral, depending on  $n(n + 1)/2$  generic masses and kinematic variables, can be expressed as a linear combination of integrals depending only on  $n$  variables. The latter integrals are given explicitly in terms of hypergeometric functions of  $(n - 1)$  dimensionless variables. Analytic expressions for the 2-, 3- and 4-point integrals, that depend on the minimal number of variables, were also obtained by solving the dimensional recurrence relations. The resulting expressions for these integrals are given in terms of Gauss' hypergeometric function  ${}_2F_1$ , the Appell function  $F_1$  and the hypergeometric Lauricella - Saran function  $F_s$ . A modification of the functional reduction procedure for some special values of kinematical variables is considered.

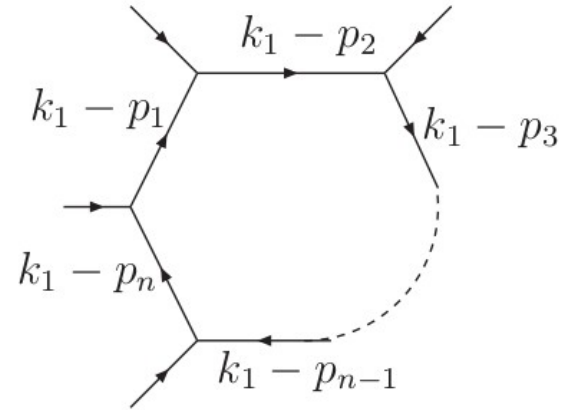
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# Algebraic relation between propagators

$$I_n^{(d)}(\{m_j^2\}; \{s_{ik}\}) = \frac{1}{i\pi^{d/2}} \int \frac{d^d k_1}{D_1 \dots D_n}$$

$$D_j = (k_1 - p_j)^2 - m_j^2$$

$p_j$  – external momenta



$$\prod_{r=1}^n \frac{1}{D_r} = \frac{1}{D_0} \sum_{r=1}^n x_r \prod_{\substack{j=1 \\ j \neq r}}^n \frac{1}{D_j}.$$

new propagator

$p_0, m_0^2$  and  $x_j (j = 1, \dots, n)$  must be chosen so that the ratio above holds  
to do this, they must satisfy the system of equations

# Algebraic relation between propagators

$$\prod_{r=1}^n \frac{1}{D_r} = \frac{1}{D_0} \sum_{r=1}^n x_r \prod_{\substack{j=1 \\ j \neq r}}^n \frac{1}{D_j} \quad \Bigg| \quad \times \prod_{j=0}^n D_j$$

$$D_0 = \sum_{r=1}^n x_r D_r, \quad \longrightarrow \quad k_1^2 - 2k_1 p_0 + p_0^2 - m_0^2 = \sum_{r=1}^n x_r (k_1^2 - 2k_1 p_r + p_r^2 - m_r^2).$$

Differentiating the above  
by  $k_1$  and given that this  
is the independent  
variable of integration  
we find

$$1 = \sum_{r=1}^n x_r,$$

$$p_0 = \sum_{j=1}^n x_j p_j.$$

$$m_0^2 - \sum_{k=1}^n x_k m_k^2 + \sum_{j=2}^n \sum_{l=1}^{j-1} x_j x_l s_{lj} = 0,$$

where

$$s_{ij} = s_{i,j} = (p_i - p_j)^2.$$

the solution depends on  $(n - 2)$  remaining parameters  $x_i$  and one arbitrary mass  $m_0$

$$I_n^{(d)}(\{m_r^2\}; \{s_{ik}\}) = \sum_{j=1}^n x_j I_n^{(d)}(\{m_r^2\}; \{s_{ik}\}) \Big|_{m_j^2 \rightarrow m_0^2, s_{jk} \rightarrow s_{0k}}.$$

# Method of functional reduction

## Sincov's functional equation

$$f(x, y) = f(x, z) - f(y, z).$$

The general solution

$$f(x, y) = g(x) - g(y),$$

where

$$g(x) = f(x, 0)$$

I.e. the function  $f(x, y)$  is a combination of its 'boundary values', which may be completely arbitrary.

# Method of functional reduction

$$I_n^{(d)}(\{m_r^2\}; \{s_{ik}\}) = \sum_{j=1}^n x_j I_n^{(d)}(\{m_r^2\}; \{s_{ik}\}) \Big|_{m_j^2 \rightarrow m_0^2, s_{jk} \rightarrow s_{0k}}.$$

Additional conditions designed to reduce the number of variables:

$$s_{0j} = 0, \quad s_{0j} - s_{0i} = 0, \quad s_{0j} \pm s_{ik} = 0, \quad s_{0j} \pm m_0^2 = 0, \quad m_j^2 \pm m_0^2 = 0, \\ m_0^2 = 0, \quad s_{0j} \pm m_0^2 \pm m_k^2 = 0, \quad (i, j, k = 1 \dots n).$$

Solutions of these systems of equations and analysis of these solutions were performed using computer algebra system MAPLE. The number of these systems depends on  $n$  and varied from  $10^3$  to  $10^6$ . CPU execution time ranged from a few minutes to several hours. Many solutions of these equations have been found. Some of them lead to a simultaneous decrease in the number of variables in all integrals on the right-hand side of the functional equation

# Functional reduction of the 2-point integral

$$I_2^{(d)}(m_1^2, m_2^2; s_{12}) = \int \frac{d^d k_1}{i\pi^{d/2}} \frac{1}{[(k_1 - p_1)^2 - m_1^2][(k_1 - p_2)^2 - m_2^2]}.$$

Algebraic relation between propagators:

$$\frac{1}{D_1 D_2} = \frac{x_1}{D_0 D_2} + \frac{x_2}{D_1 D_0}$$

$$x_1 + x_2 = 1, \quad p_0 = x_1 p_1 + x_2 p_2,$$

Algebraic conditions on parameters :

$$m_0^2 - x_1 m_1^2 - x_2 m_2^2 + x_1 x_2 s_{12} = 0.$$

The only arbitrary parameter will be  $m_0$

$$x_1 = \frac{m_2^2 - m_1^2 + s_{12}}{2s_{12}} \pm \frac{\sqrt{4s_{12}(m_0^2 - r_{12})}}{2s_{12}}, \quad x_2 = 1 - x_1$$

$$s_{10} = m_1^2 + m_0^2 - 2r_{12} \pm \frac{m_2^2 - m_1^2 - s_{12}}{2s_{12}} \sqrt{4s_{12}(m_0^2 - r_{12})}, \quad s_{20} = m_2^2 + m_0^2 - 2r_{12} \pm \frac{m_2^2 - m_1^2 + s_{12}}{2s_{12}} \sqrt{4s_{12}(m_0^2 - r_{12})},$$

$$r_{12} = -\frac{\lambda_{12}}{g_{12}} = \frac{2m_1^2 m_2^2 + 2s_{12} m_1^2 + 2s_{12} m_2^2 - m_1^4 - m_2^4 - s_{12}^2}{4s_{12}}.$$

# Notations for determinants

modified Cayley determinant:  $\Delta_n \equiv \Delta_n(\{p_1, m_1\}, \dots, \{p_n, m_n\}) = \begin{vmatrix} Y_{11} & Y_{12} & \dots & Y_{1n} \\ Y_{12} & Y_{22} & \dots & Y_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ Y_{1n} & Y_{2n} & \dots & Y_{nn} \end{vmatrix},$

$$Y_{ij} = m_i^2 + m_j^2 - s_{ij},$$

Gram determinant:  $G_{n-1} \equiv G_{n-1}(p_1, \dots, p_n) = -2 \begin{vmatrix} S_{11} & S_{12} & \dots & S_{1 \ n-1} \\ S_{21} & S_{22} & \dots & S_{2 \ n-1} \\ \vdots & \vdots & \ddots & \vdots \\ S_{n-1 \ 1} & S_{n-1 \ 2} & \dots & S_{n-1 \ n-1} \end{vmatrix}$

$$S_{ij} = s_{in} + s_{jn} - s_{ij},$$

$$\lambda_{i_1 i_2 \dots i_n} = \Delta_n(\{p_{i_1}, m_{i_1}\}, \{p_{i_2}, m_{i_2}\}, \dots, \{p_{i_n}, m_{i_n}\})$$

$$r_{ij\dots k} = -\frac{\lambda_{ij\dots k}}{g_{ij\dots k}}.$$

$$g_{i_1 i_2 \dots i_n} = G_{n-1}(p_{i_1}, p_{i_2}, \dots, p_{i_n}).$$

and

$$K_{j_r j_1 \dots j_{r-1} j_{r+1} \dots j_n} = \frac{\partial r_{j_1 \dots j_r \dots j_n}}{\partial m_{j_r}^2}.$$

# Functional reduction of the 2-point integral

$$I_2^{(d)}(m_1^2, m_2^2; s_{12}) = \int \frac{d^d k_1}{i\pi^{d/2}} \frac{1}{[(k_1 - p_1)^2 - m_1^2][(k_1 - p_2)^2 - m_2^2]}.$$

Algebraic relation between propagators:

$$\frac{1}{D_1 D_2} = \frac{x_1}{D_0 D_2} + \frac{x_2}{D_1 D_0}$$

Integrating algebraic relation over momentum  $k_1$  we get the functional relation

$$I_2^{(d)}(m_1^2, m_2^2; s_{12}) = x_1 I_2^{(d)}(m_2^2, m_0^2; s_{20}) + x_2 I_2^{(d)}(m_1^2, m_0^2; s_{10}).$$

The only arbitrary parameter will be  $m_0$

- Case 1.  $m_0^2 = 0$
- Case 2.  $m_0^2 = r_{12}$
- Case 3. Combination of two equations



# Case 1. $m_0^2 = 0$

$$I_2^{(d)}(m_1^2, m_2^2; s_{12}) = \bar{x}_1 I_2^{(d)}(m_2^2, 0; \bar{s}_{20}) + \bar{x}_2 I_2^{(d)}(m_1^2, 0; \bar{s}_{10}),$$

$$\bar{x}_{1,2} = x_{1,2}|_{m_0^2=0}, \quad \bar{s}_{01} = s_{01}|_{m_0^2=0}, \quad \bar{s}_{02} = s_{02}|_{m_0^2=0}$$

$$I_2^{(d)}(m^2, 0; p^2) = -\Gamma\left(1 - \frac{d}{2}\right) m^{d-4} {}_2F_1\left[1, 2 - \frac{d}{2}; \frac{d}{2}; \frac{p^2}{m^2}\right]$$

## Case 2. $m_0^2 = r_{12}$

$$I_2^{(d)}(m_1^2, m_2^2; s_{12}) = \kappa_{12} I_2^{(d)}(r_{12}, r_2; r_2 - r_{12}) + \kappa_{21} I_2^{(d)}(r_{12}, r_1; r_1 - r_{12})$$

where

$$\kappa_{12} = \frac{\partial r_{12}}{\partial m_1^2}, \quad \kappa_{21} = \frac{\partial r_{12}}{\partial m_2^2}, \quad r_i = m_i^2.$$

$$(d-1)I_2^{(d+2)}(r_{12}, r_j; r_j - r_{12}) = -2r_{12}I_2^{(d)}(r_{12}, r_j; r_j - r_{12}) - I_1^{(d)}(r_j).$$

$$I_2^{(d)}(r_{12}, r_j; r_j - r_{12}) = \frac{-\pi^{\frac{3}{2}} r_{12}^{\frac{d}{2}-2}}{2 \sin \frac{\pi d}{2} \Gamma\left(\frac{d-1}{2}\right)} \sqrt{\frac{r_{12}}{r_{12} - r_j}} + \frac{\pi}{2r_{12}} \frac{r_j^{\frac{d}{2}-1}}{\sin \frac{\pi d}{2} \Gamma\left(\frac{d}{2}\right)} {}_2F_1\left[1, \frac{d-1}{2}; \frac{r_j}{r_{12}}\right]$$

## Case 3. Combination of two equations

We take  $m_0^2 = m_2^2$

$$I_2(m_1^2, m_2^2; s_{12}) = \frac{m_1^2 - m_2^2}{s_{12}} I_2^{(d)} \left( m_1^2, m_2^2; \frac{(m_1^2 - m_2^2)^2}{s_{12}} \right) + \frac{s_{12} - m_1^2 + m_2^2}{s_{12}} I_2^{(d)} \left( m_2^2, m_2^2; \frac{(s_{12} - m_1^2 + m_2^2)^2}{s_{12}} \right)$$

invariant under interchange  $m_1^2 \longleftrightarrow m_2^2$

$$I_2(m_1^2, m_2^2; s_{12}) = \frac{s_{12} + m_1^2 - m_2^2}{2s_{12}} I_2^{(d)} \left( m_1^2, m_1^2; \frac{(s_{12} + m_1^2 - m_2^2)^2}{s_{12}} \right) + \frac{s_{12} - m_1^2 + m_2^2}{2s_{12}} I_2^{(d)} \left( m_2^2, m_2^2; \frac{(s_{12} - m_1^2 + m_2^2)^2}{s_{12}} \right)$$

$$I_2^{(d)}(m^2, m^2; p^2) = m^{d-4} \Gamma \left( 2 - \frac{d}{2} \right) {}_2F_1 \left[ \begin{matrix} 1, 2 - \frac{d}{2} \\ \frac{3}{2} \end{matrix}; \frac{p^2}{4m^2} \right]$$

# Functional reduction of the 3-point integral

$$I_3^{(d)}(m_1^2, m_2^2, m_3^2; s_{23}, s_{13}, s_{12}) = \frac{1}{i\pi^{d/2}} \int \frac{d^d k_1}{D_1 D_2 D_3}.$$

Algebraic relation between propagators: 
$$\frac{1}{D_1 D_2 D_3} = \frac{x_1}{D_0 D_2 D_3} + \frac{x_2}{D_1 D_0 D_3} + \frac{x_3}{D_1 D_2 D_0}.$$

Algebraic conditions on parameters : 
$$\left\{ \begin{array}{l} p_0 = x_1 p_1 + x_2 p_2 + x_3 p_3, \\ x_1 + x_2 + x_3 = 1, \\ x_1 x_2 s_{12} + x_1 x_3 s_{13} + x_2 x_3 s_{23} - x_1 m_1^2 - x_2 m_2^2 - x_3 m_3^2 + m_0^2 = 0. \end{array} \right.$$


The functional equation depends on two arbitrary parameters

$$I_3^{(d)}(m_1^2, m_2^2, m_3^2; s_{23}, s_{13}, s_{12}) = x_1 I_3^{(d)}(m_0^2, m_2^2, m_3^2; s_{23}, s_{03}, s_{02}) + x_2 I_3^{(d)}(m_1^2, m_0^2, m_3^2; s_{03}, s_{13}, s_{01}) + x_3 I_3^{(d)}(m_1^2, m_2^2, m_0^2; s_{02}, s_{01}, s_{12})$$

Functional reduction goes in two steps

# Step 1

$$\begin{aligned}
 & I_3 (m_1^2, m_2^2, m_3^2; s_{23}, s_{13}, s_{12}) \\
 &= \kappa_{123} I_3 (r_{123}, r_2, r_3; s_{23}, r_3 - r_{123}, r_2 - r_{123}) \\
 &+ \kappa_{213} I_3 (r_{123}, r_1, r_3; s_{13}, r_3 - r_{123}, r_1 - r_{123}) \\
 &+ \kappa_{312} I_3 (r_{123}, r_2, r_1; s_{12}, r_1 - r_{123}, r_2 - r_{123}),
 \end{aligned}$$


  
 four independent variables

where

$$\begin{aligned}
 r_{123} &= -\frac{\lambda_{123}}{g_{123}}, & r_i &= m_i^2, \\
 \kappa_{123} &= \frac{\partial r_{123}}{\partial m_1^2}, & \kappa_{213} &= \frac{\partial r_{123}}{\partial m_2^2}, & \kappa_{312} &= \frac{\partial r_{123}}{\partial m_3^2},
 \end{aligned}$$

## Step 2

$$I_3^{(d)}(r_{123}, r_2, r_3; s_{23}, r_3 - r_{123}, r_2 - r_{123}) = \kappa_{23} I_3^{(d)}(r_{123}, r_{23}, r_3; r_3 - r_{23}, r_3 - r_{123}, r_{23} - r_{123}) \\ + \kappa_{32} I_3^{(d)}(r_{123}, r_{23}, r_2; r_2 - r_{23}, r_2 - r_{123}, r_{23} - r_{123})$$

$$I_3^{(d)}(m_1^2, m_2^2, m_3^2; s_{23}, s_{13}, s_{12}) \\ = \kappa_{123} \kappa_{23} I_3^{(d)}(r_{123}, r_{23}, r_3; r_3 - r_{23}, r_3 - r_{123}, r_{23} - r_{123}) \\ + \kappa_{123} \kappa_{32} I_3^{(d)}(r_{123}, r_{23}, r_2; r_2 - r_{23}, r_2 - r_{123}, r_{23} - r_{123}) \\ + \kappa_{213} \kappa_{31} I_3^{(d)}(r_{123}, r_{13}, r_1; r_1 - r_{13}, r_1 - r_{123}, r_{13} - r_{123}) \\ + \kappa_{213} \kappa_{13} I_3^{(d)}(r_{123}, r_{13}, r_3; r_3 - r_{13}, r_3 - r_{123}, r_{13} - r_{123}) \\ + \kappa_{312} \kappa_{12} I_3^{(d)}(r_{123}, r_{12}, r_2; r_2 - r_{12}, r_2 - r_{123}, r_{12} - r_{123}) \\ + \kappa_{312} \kappa_{21} I_3^{(d)}(r_{123}, r_{12}, r_1; r_1 - r_{12}, r_1 - r_{123}, r_{12} - r_{123}).$$

three independent variables

# Analytic results for integrals depending on the MNV

$$(d-2)I_3^{(d+2)}(r_{123}, r_{23}, r_3; r_3 - r_{23}, r_3 - r_{123}, r_{23} - r_{123}) = \\ -2r_{123}I_3^{(d)}(r_{123}, r_{23}, r_3; r_3 - r_{23}, r_3 - r_{123}, r_{23} - r_{123}) - I_2^{(d)}(r_{23}, r_3; r_3 - r_{23})$$

## Solution

$$I_3^{(d)}(r_{123}, r_{23}, r_3; r_3 - r_{23}, r_3 - r_{123}, r_{23} - r_{123}) = \\ \frac{1}{\sin \frac{\pi d}{2}} \left\{ \frac{r_{123}^{\frac{d-6}{2}}}{\Gamma\left(\frac{d-2}{2}\right)} C_3(x, y) + \frac{\pi^{\frac{3}{2}} r_{23}^{\frac{d-4}{2}}}{4r_{123}\Gamma\left(\frac{d-1}{2}\right)} \sqrt{\frac{r_{23}}{r_{23} - r_3}} {}_2F_1\left[1, \frac{d-2}{2}; \frac{r_{23}}{r_{123}}\right] \right. \\ \left. - \frac{\pi r_3^{\frac{d-2}{2}}}{4\Gamma\left(\frac{d}{2}\right)(r_{23} - r_3)r_{123}} \sqrt{1 - \frac{r_3}{r_{23}}} F_1\left(\frac{d-2}{2}, 1, \frac{1}{2}, \frac{d}{2}; \frac{r_3}{r_{123}}, \frac{r_3}{r_{23}}\right) \right\}$$

where

$$C_3(x, y) = \frac{\pi x y^2}{4(x^2 - y^2)^{\frac{1}{2}}} \ln \left( \frac{x - (x^2 - y^2)^{\frac{1}{2}}}{x + (x^2 - y^2)^{\frac{1}{2}}} \right) \quad x = \sqrt{\frac{r_{123}}{r_{123} - r_3}}, \quad y = \sqrt{\frac{r_{123}}{r_{123} - r_{23}}}.$$

# Functional reduction of the 4-point integral

The integral depending on ten variables is rewritten through a combination of integrals depending on four independent variables

$$\begin{aligned}
 & I_4(m_1^2, m_2^2, m_3^2, m_4^2; s_{12}, s_{23}, s_{34}, s_{14}, s_{24}, s_{13}) \\
 &= \kappa_{1234}\kappa_{234}\kappa_{34}I_4^{(d)}(r_{1234}, r_{234}, r_{34}, r_4; \\
 &\quad r_{234} - r_{1234}, r_{34} - r_{234}, r_4 - r_{34}, r_4 - r_{1234}, r_4 - r_{234}, r_{34} - r_{1234}) \\
 &+ \kappa_{1234}\kappa_{234}\kappa_{43}I_4^{(d)}(r_{1234}, r_{234}, r_{34}, r_3; \\
 &\quad r_{234} - r_{1234}, r_{34} - r_{234}, r_3 - r_{34}, r_3 - r_{1234}, r_3 - r_{234}, r_{34} - r_{1234}) \\
 &+ \kappa_{1234}\kappa_{324}\kappa_{24}I_4^{(d)}(r_{1234}, r_{234}, r_{24}, r_4; \\
 &\quad r_{234} - r_{1234}, r_{24} - r_{234}, r_4 - r_{24}, r_4 - r_{1234}, r_4 - r_{234}, r_{24} - r_{1234}) \\
 &+ \kappa_{1234}\kappa_{324}\kappa_{42}I_4^{(d)}(r_{1234}, r_{234}, r_{24}, r_2; \\
 &\quad r_{234} - r_{1234}, r_{24} - r_{234}, r_2 - r_{24}, r_2 - r_{1234}, r_2 - r_{234}, r_{24} - r_{1234}) + \dots
 \end{aligned}$$

The full answer is almost 2 pages long



# Functional reduction of the 5 and 6-point integrals

The  $I_5$  integral depending on 15 variables can be represented as a linear combination of 120 integrals, each of which depends on only 5 variables

$$I_5^{(d)}(m_i^2, m_j^2, m_k^2, m_l^2, m_r^2; m_j^2 - m_i^2, m_k^2 - m_j^2, m_l^2 - m_k^2, m_r^2 - m_l^2, m_r^2 - m_i^2, m_k^2 - m_i^2, m_l^2 - m_i^2, m_l^2 - m_j^2, m_r^2 - m_j^2, m_r^2 - m_k^2),$$

where  $m_i^2, m_j^2, m_k^2, m_l^2, m_r^2$  are ratios of polynomials in masses and kinematic invariants.

The  $I_6$  integral depending on 21 variables can be represented as a linear combination of 720 integrals, each of which depends on only 6 variables

# General algorithm of the functional reduction

Final functional reduction formulae for the integrals  $I_2^{(d)}, \dots, I_6^{(d)}$  can be obtained by exploiting the following algorithm:

- write down the term

$$\kappa_{1\dots n} \kappa_{2\dots n} \dots \kappa_{n-1 n} I_n^{(d)}(m_1^2, m_2^2, \dots, m_n^2; s_{12}, s_{23}, \dots) \quad (10.1)$$

- replace in the integral  $s_{ij} \rightarrow m_j^2 - m_i^2$  ( $j > i$ )
- replace in the integral  $m_1^2 \rightarrow r_{1\dots n}, m_2^2 \rightarrow r_{2\dots n}, \dots, m_n^2 \rightarrow r_n$
- replace  $\kappa_{ij\dots} \rightarrow \frac{\partial r_{ij\dots}}{\partial m_i^2}$
- generate  $n! - 1$  terms by symmetrizing the term (10.1) with respect to the indices  $1, 2, \dots, n$  and add all these terms to (10.1).

All steps are very straightforward and easily achieved with a computer program. This algorithm works perfectly for integrals  $I_2^{(d)}, \dots, I_6^{(d)}$ . We verified numerically that it is also valid for integrals  $I_7^{(d)}, I_8^{(d)}$ . Notice that the number of terms in the final reduction formula for massless integrals is  $n!/2$ .

# Conclusions

- Author provided a systematic approach for reducing a generic n-point one-loop integral with arbitrary masses and kinematic invariants to a linear combination of integrals that depend on n variables.
- The integrals depending on the MNV encountered at the last stage of the reduction were expressed in terms of multiple hypergeometric series depending on n-1 dimensionless variables.
- Analytic results for integrals with the MNV can be derived by solving dimensional recurrence relations.
- Obtained representations of one-loop integrals can be helpful for deriving  $\epsilon = (4 - d)/2$  expansion of these integrals.

# Thank you for your attention!

