

# COMPACT STARS

(Cosmic matter in heavy-ion collision laboratories?)





The High-Density Symmetry Energy in Heavy ion Collisions and Compact stars Hermann Wolter, University of Munich



Subtitle of CSQCD VI: Cosmic matter in heavy ion collision laboratories?

Aim and outline of this talk:

- Discuss investigation of EoS, relevant for astrophysics, in heavy ion collisions
- non-equilibrium processes  $\rightarrow$  transport theory (hydrodynamics?)
- issues in transport theory
- Application at very low and supersaturation densities
- status of knowledge of symmetry energy (in the hadronic sector)

Collaborators:

Heavy ion collisions:		Clustering in dilute r
Maria Colonna, Massimo Di Toro, Enzo Greco		Stefan Typel (TU Dar
(Lab. Naz. del Sud, INFN, Catania),		Gerd Röpke (Univ. of
Malgorzata Zielinska-Pfabe (Smith College, USA)		David Blaschke, Tho
Theo Gaitanos (Univ. Thesaloniki)	1	
Tatiana Mikhailova, JINR, Dubna		

Clustering in dilute matter, SN and NS observables: Stefan Typel (TU Darmstadt, GSI) Gerd Röpke (Univ. of Rostock) David Blaschke, Thomas Klähn (Univ. of Wroclaw)

#### Code Comparison:

Jun Xu (SINAP,Shanghai), Yingxun Zhang (CIAE, Beijing), Lie-Wen Chen (Jiao Tong Univ., Shanghai), Betty Tsang (MSU), Yong-Jia Wang (Huzhou Univ.), Pawel Danielewicz (MSU)







# Equation-of-State and Symmetry Energy: Microscopic Results



## The Search for the Nuclear Symmetry Energy

$$E(\rho_B, \delta)/A = E_{nm}(\rho_B) + E_{sym}(\rho_B)\delta^2 + O(\delta^4) + ...$$
  $\delta = \frac{\rho_n - \rho_p}{\rho_n + \rho_p}$ 



#### HIC one way to obtain information on the EoS - but complex processes

Fermi energies: (multi)-fragmentation in central collisions

Intermediate energies: several 100 MeV/A to several GeV/A Vaporization, production of new particles, like pions, strangeness (kaons, hyperons), etc,



#### Transport theory based on a chain of approximations

Martin-Schwinger hierarchy in many-body densities, real time formalism truncation, introduction of self energies (1-body quantities), irreversability

Quantum transport theory: Kadanoff-Baym theory

Semiclassical approximation :

Wigner transform, treat as phase space probabilities Gradient approximation (separation of short and long scales)

**Quasi-particle approximation** 

Spectral function  $\rightarrow$  delta function with effective momenta and masses neglect off-shell effects (or treat approximately)

 $\rightarrow$  kinetic equation

$$\begin{aligned} \frac{\partial f}{\partial t} + \frac{\vec{p}}{m} \vec{\nabla}^{(r)} f - (\vec{\nabla}^{(r)} U(r, p) \vec{\nabla}^{(p)} + \vec{\nabla}^{(p)} U(r, p) \vec{\nabla}^{(r)}) f(\vec{r}, \vec{p}; t) = \\ \int d\vec{p}_2 d\vec{p}_1 d\vec{p}_{2'} v_{21} \sigma_{12}^{in-med}(\Omega) (2\pi)^3 \delta(p_1 + p_2 - p_{1'} - p_{2'}) \Big[ f_{1'} f_2(f_1 f_2) - f_1 f_2(f_{1'} f_{2'}) \\ & \text{Pauli blocking factors,} \\ & \bar{f}_i := (1 - f_i) \end{aligned}$$

Mean field evolution (Vlasov) + uncorr. 2-body collisions (Boltzmann) + Pauli-blocking of final states (Uehling-Uhlenbeck)



## Two families of transport approaches, fluctuations

Boltzmann-Vlasov-like (BUU/BL/BLOB)  $\left(\frac{\partial}{\partial t} + \frac{\vec{p}}{m}\vec{\nabla}^{(r)} - \vec{\nabla}U(r)\vec{\nabla}^{(p)}\right)f(\vec{r},\vec{p};t)$   $= I_{coll}\left[\sigma^{in-med}\right] + \delta I_{fluct}$ 

Dynamics of the 1-body phase space distribution function f with 2-body dissipation

fluctuations around diss. solution  $f(r,p,t) = \overline{f}(r,p,t) + \delta f(r,p,t)$ 



Molecular-Dynamics-like (QMD/AMD)

$$|\Phi\rangle = \bigwedge_{i=1}^{A} \varphi(\mathbf{r}; \mathbf{r}_{i}, \mathbf{p}_{i}) |0\rangle$$
  
$$\dot{\mathbf{r}}_{i} = \{\mathbf{r}_{i}, \mathbf{H}\}; \quad \dot{\mathbf{p}}_{i} = \{\mathbf{p}_{i}, \mathbf{H}\}; \quad \mathbf{H} = \sum_{i} t_{i} + \sum_{i,j} V(\mathbf{r}_{i} - \mathbf{r}_{j})$$

TD-Hartree(-Fock) (or classical molecular dynamics with extended particles) plus stochastic NN collisions No quantum fluctuations, but classical N-body fluctuations, damped by the smoothing.

However, more fluctuations than BUU, since dof are nucleons and not test particles:

 $\rightarrow$  amount controlled by width of single particle packet  $\Delta L$ 





LC's are not stabilized by the mean field but by many body correlations. Introduce as explicit degrees of freedom, generated by the collision term



(P. Danielewicz and Q. Pan, PRC 46 (1992)) (d,t,3He, but no  $\alpha$ !)

Maybe the way to deal with the hadron-quark phase transition in transport approaches

Fluctuation in the instable region are Amplified and stabilized by the mean field



BUU calculation in a box with initial conditions inside the instability region:  $\rho = \rho_0/3$ , T=5 MeV,  $\delta = 0$ 

(V.Baran, et al., Phys.Rep.410,335(05))

## **Particle Production**

Inelastic collisions: Production of particles and resonances



#### What can one learn from different species?

- $\bullet$  pions: production at all stages of the evolution via the  $\Delta\text{-resonace}$
- kaons (strange mesons with high mass): subthreshold production, probe of high density phase
- ratios of  $\pi^+/\pi^-$  and  $K^0/K^+$ :
- $\rightarrow$  probe for symmetry energy

e.g. pion and kaon production;

coupling of  $\Delta$  and strange-ness channels.



Coupled transport equations

Many new potentials, elastic and inelastic cross sections needed, D dynamics in medium

#### symmetry energy effects on $\pi$ , $\Delta$ production

1. Mean field effect:  $\rm U_{\rm sym}$  more repulsive for neutrons, and more for asystiff

$$\frac{n}{p} \downarrow \Rightarrow \frac{Y(\varDelta^{0,-})}{Y(\varDelta^{+,++})} \downarrow \Rightarrow \frac{\pi^{-}}{\pi^{+}} \downarrow$$
  
decrease with asy – stiffness

2. Threshold effect, in medium effective masses:



 $\rightarrow$  can be competing effects!

## Code Comparison Project

Boltzmann-Vlasov-like (BUU/BL/BLOB)

$$\left( \frac{\partial}{\partial t} + \frac{\vec{p}}{m} \vec{\nabla}^{(r)} - \vec{\nabla} U(r) \vec{\nabla}^{(p)} \right) f(\vec{r}, \vec{p}; t)$$
  
=  $I_{coll}[\sigma^{in-med}, f_i]$ 

Dynamics of the 1-body phase space distribution function f with 2-body dissipation Molecular-Dynamics-like (QMD/AMD)

$$|\Phi\rangle = \bigwedge_{i=1}^{A} \varphi(r; r_i, p_i) |0\rangle$$
  
$$\dot{r}_i = \{r_i, H\}; \quad \dot{p}_i = \{p_i, H\}; \quad H = \sum_i t_i + \sum_{i,j} V(r_i - r_j)$$

TD-Hartree(-Fock (AMD)) (or classical molecular dynamics with extended particles) plus stochastic NN collisions

6-dim integro-differential, non-linear eq.

6A-dim many body problem

→ very complex, simulate solutions introduces many technical details



→ Transport Code Evaluation (Comparison) Project

## Code Comparison Project (1st stage):

check consistency of transport codes in calculations with same system (Au+Au), E=100,400 AMeV, identical physical input (mean field (EOS) and cross sections, BUU and QMD codes

idea: establish sort of theoretical systematic error of transport calculations (and hopefully to reduce it ) published: J. Xu, et al. (31 authors); PRC 93, 044609 (2016)

HIC at b=7m (midcentral)
selected contour plots;
different evolution apparent
→ compare collision numbers,
blocking, and observables

quantify spread of simulations by<br/>value of "flow"=slope of transverse<br/>momentum at midrapidityBUU and QMD approx. consistent<br/>uncertainity100 AMeV: ~30%

400 AMeV: ~13%

Understand differences better ("entangled")



#### 2nd stage: Box calculation comparison

simulation of the static system of infinite nuclear matter,  $\rightarrow$  solve transport equation in a periodic box



Useful for many reasons:

- check consistency of calculation e.g. EoS energy dens  $\epsilon$  vs. pressure P
- check consistency of simulation:
   collision numbers, blocking
   (exact limits from kinetic theory)
- check aspects of simulation separately Cascade: only collisions without/with blocking
   Vlasov: only mean field propagation
- check ingredients of particle production e.g. pion production



--> characteristically different fluctuation in BUU and QMD models

#### Check of $\pi$ , $\Delta$ production in box cascade calculation

no pions

- Constant  $\Delta$  mass  $M_{\Delta}$  = 1.232 GeV.
- Constant  $\sigma(NN \rightarrow N\Delta) = 40$  mb for  $\sqrt{s} > M_N + M_{\Delta}$ .



pions and  $\Delta$  $\Delta$  mass distribution  $NN \leftrightarrow N\Delta$  $\Gamma[\Delta \rightarrow N\pi] = 0.115 \text{ GeV}$ 

(very preliminary figure removed)

## Transport theory $\leftarrow \rightarrow$ Experiments (Observables)

Transport theory calculates f(r,p;t), or  $\{r_i(t), p_i(t)\}$  i.e. full information about the complete evolution,

Experiment measures  $f(r \rightarrow \infty, p; t \rightarrow \infty)$ , i.e. asymptotic momentum distribution, of nucleons, and also of newly produced particles (p, K, ...) (inel. cross sections) and of clusters (in principle many-body observables)



## Constraints on EOS of symmetric nuclear matter

(P.Danielewicz, et al., Science 298(02)1592)



soft EOS, pot ChP' ard EOS, pot ChPT

t EOS, IQMD, pot RMF ard EOS, IQMD, pot RMF

soft EOS. IQMD. Giessen cs.

ard EOS, IQMD, Giessen cs

- analysis of elliptic and directed flow for E=.15 to 10 AGeV

- using a particular model (purely hadronic), shaded area constraint in this model

- densities in the range from  $2 - 4.5(?) \rho_0$ 

- eliminates some extreme models

heavy (Au+Au,

compres-sion)

compression) relative to

light system (C+C, small



0.8 1.0 1.2 1.4 1.6 Elah [GeV] Both favor a rather soft EoS,

 $(\mathbf{M}_{\mathbf{K}^+}/\mathbf{A})_{\mathbf{M}_{\mathbf{M}^+},\mathbf{M}_{\mathbf{M}}} / (\mathbf{M}_{\mathbf{K}^+}/\mathbf{M}_{\mathbf{M}_{\mathbf{M}}})$ 

5

4

3

2

1

where the momentum dependence of the potential is important



## Clustering of very dilute nuclear matter



Cluster production can also be important at high energies/densities, e.g. NICA White paper, EPJA 52:

#### #34 Light cluster production at NICA

N.-U. Bastian<sup>1</sup>, P. Batyuk<sup>2</sup>, D. Blaschke<sup>1,2,3</sup>, P. Danielewicz<sup>4</sup>, Yu. B. Ivanov<sup>3,5</sup>, Iu. Karpenko<sup>6,7</sup>, G. Röpke<sup>3,8,a</sup>, O. Rogachevsky<sup>2</sup>, H. H. Wolter<sup>9</sup>



## Constraints on the Symmetry Energy as correlation S<sub>0</sub> vs. L

Limits on the value around saturation density

$$\left| \boldsymbol{E}_{sym}(\rho) = \boldsymbol{S}_{0} + \boldsymbol{L} \boldsymbol{S}_{0} + \boldsymbol{L} \boldsymbol{S}_{0} \boldsymbol{\rho}_{0} \right| + \frac{\boldsymbol{K}_{sym}}{\boldsymbol{18}} \left( \frac{\rho - \rho_{0}}{\rho_{0}} \right)^{2} \right|$$

A consensus towards a rather narrow region in the parameter space is seen.

But includes correlation between parameters due to extrapolation to  $\rho_0$ , e..g. masses





 $S(\rho = \frac{2}{3}\rho_0) = S - \frac{L}{9} + \frac{K_{sym}}{162} + \dots \approx S - \frac{L}{9}$ 

-- e.g. SE that fit nuclear masses cross below saturation density, (some average densitiy of a finite nucleus)
 -- induces a correlation between value and slope at ρ<sub>0</sub>, within the model., eg. in lin. approx.

#### The Symmetry Energy at High Density

Au+Au @ 400 AMeV new experiment ASY-EOS P. Russotto et al., Phys. Rev. C 94, 034608 (2016)

$$N(\Theta; y, p_t) = N_0 (1 + v_1 \cos \Theta + v_2 \cos 2\Theta + ...)$$

ratio of neutron to hydrogen flow - Elliptic flow  $v_2$  in this energy region good probe of high density

analysis of density dependence of SE in terms of power law exponent  $\boldsymbol{\gamma}$ 

$$\boldsymbol{E}_{sym}(\rho) = \frac{1}{3} \varepsilon_{F} \left(\frac{\rho}{\rho_{0}}\right)^{2/3} + \boldsymbol{C} \left(\frac{\rho}{\rho_{0}}\right)^{2}$$

Big step forward in constraining the high-density symmetry energy



# **Particle Production**

## Inelastic collisions: Dynamics of production of particles and resonances



G. Ferini et al., Nucl. Phys. A 762, 147 (2005)

## Results for pion ratios



- → potentials  $U_p$ ,  $U_p$ ,  $U_K$  → effective masses in medium, threshold effects
- → inelastic cross sections, e.g.  $NN \rightarrow N\Delta$ , in medium
- →  $\Delta$  Resonances with decay widths, mass distributions, spectral fcts, more general: off-shell transport



#### $\rightarrow$ present situation of analysis unsatisfactory

- $\rightarrow$  check of  $(\pi, \Delta)$  physics in box calculations
- $\rightarrow$  more sensitivity in spectral distributions, Sprit experiment (MSU-RIKEN)
- → reconsider Kaons, more direct probe of high density region



#### Testing models of the high-density symmetry energy in NS

1 (Tolmann-Oppenheimer-Volkov TOV):

$$\frac{\partial P}{\partial r} = -\frac{m\rho}{r^2} \left(1 + \frac{4\pi \, r^3 \, P}{m}\right) \left(1 + e + \frac{P}{\rho}\right) \left(1 - \frac{2m}{r}\right)^{-1}$$



#### Testing of EoS's constrained by HIC in NS not often performed systematically. One example:

T. Klähn,<sup>1,2,\*</sup> D. Blaschke,<sup>3,4,†</sup> S. Typel,<sup>3</sup> E. N. E. van Dalen,<sup>2</sup> A. Faessler,<sup>2</sup> C. Fuchs,<sup>2</sup> T. Gaitanos,<sup>5</sup> H. Grigorian,<sup>1,6</sup> A. Ho,<sup>7</sup> E. E. Kolomeitsev,<sup>8</sup> M. C. Miller,<sup>9</sup> G. Röpke,<sup>1</sup> J. Trümper,<sup>10</sup> D. N. Voskresensky,<sup>3,11</sup> F. Weber,<sup>7</sup> and H. H. Wolter<sup>5</sup> PHYSICAL REVIEW C 74, 035802 (2006)



Synopsis of constraints from neutron stars, HIC and microscopic calculations (for neutron star matter, i.e.  $\beta$ -equilibrium)



#### Conclusions:

- Heavy ion collisions important to investigate the EoS in a wide range of densities
- Non-equilibrium processes, therefore transport theories are necessary to extract EoS. EoS does not enter directly, rather potentials and cross sections
- Transport theories are well developed, but also open problems fluctuations, finite size effects, change of degrees of freedom, homogeneous <->clustered matter, hadron<->quark mat
- Symmetry energy in hadronic sector starting to converge, but more work needed in the high density region
- HIC important for investigation of NICA, FAIR , J-PARK energy regime

