Weyl Cosmology and Particle Production

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A little bit of mathematics Differential geometry

Metric tensor $g_{\mu\nu} \Rightarrow \text{interval}$

$$ds^2 = g_{\mu\nu}(x)dx^{\mu}dx^{\nu}$$
 $(g_{\mu\nu}g^{\nu\lambda} = \delta^{\lambda}_{\mu})$

Connections $\Gamma^{\lambda}_{\mu\nu} \Rightarrow$ covariant derivative

$$\nabla_{\mu}I^{\nu} = I^{\nu}_{,\mu} + \Gamma^{\nu}_{\sigma\mu}I^{\sigma}, \ldots$$

Curvature tensor $R^{\mu}_{\nu\lambda\sigma}$

$$R^{\mu}_{\ \nu\lambda\sigma} = \frac{\partial \Gamma^{\mu}_{\nu\sigma}}{\partial x^{\lambda}} - \frac{\partial \Gamma^{\mu}_{\nu\lambda}}{\partial x^{\sigma}} + \Gamma^{\mu}_{\varkappa\lambda} \Gamma^{\varkappa}_{\nu\sigma} - \Gamma^{\mu}_{\varkappa\sigma} \Gamma^{\varkappa}_{\nu\lambda}$$

Ricci tensor $R_{\mu\nu}=R^{\lambda}_{\mu\lambda\nu}$

Curvature scalar $R = g^{\mu\lambda}R_{\mu\lambda}$

$$g_{\mu
u}(x), \quad \Gamma^{\lambda}_{\mu
u}(x)$$

A little bit of mathematics Differential geometry II

Three tensors
$$g_{\mu\nu}(x)$$
, $S^{\lambda}_{\mu\nu}$, $Q_{\lambda\mu\nu}$

Torsion $S^{\lambda}_{\mu\nu} = \Gamma^{\lambda}_{\mu\nu} - \Gamma^{\lambda}_{\nu\mu}$

Nonmetricity $Q_{\lambda\mu\nu} = \nabla_{\lambda}g_{\mu\nu}$

Connections $\Gamma^{\lambda}_{\mu\nu} = C^{\lambda}_{\mu\nu} + K^{\lambda}_{\mu\nu} + L^{\lambda}_{\mu\nu}$
 $C^{\lambda}_{\mu\nu} = \frac{1}{2}g^{\lambda\varkappa}(g_{\varkappa\mu,\nu} + g_{\varkappa\nu,\mu} - g_{\mu\nu,\varkappa})$
 $Q_{\lambda\mu\nu} = \nabla_{\lambda}g_{\mu\nu} \implies Q_{\lambda\mu\nu} = Q_{\lambda\nu\mu}$
 $K^{\lambda}_{\mu\nu} = \frac{1}{2}(S^{\lambda}_{\mu\nu} - S^{\lambda}_{\mu\nu} - S^{\lambda}_{\nu\mu})$
 $L^{\lambda}_{\mu\nu} = \frac{1}{2}(Q^{\lambda}_{\mu\nu} - Q^{\lambda}_{\mu\nu} - Q^{\lambda}_{\nu\mu})$

A little bit of mathematics Differential geometry III

Riemann geometry

$$\begin{split} S_{\mu\nu}^{\lambda} &= 0, \quad Q_{\lambda\mu\nu} = 0 \\ S_{\mu\nu}^{\lambda} &= 0, \quad Q_{\lambda\mu\nu} = 0, \quad \Rightarrow \quad \Gamma_{\mu\nu}^{\lambda} = C_{\mu\nu}^{\lambda} \end{split}$$

Weyl geometry

$$\begin{split} S^{\lambda}_{\mu\nu} &= 0 \quad \Rightarrow \quad \Gamma^{\lambda}_{\mu\nu} = \Gamma^{\lambda}_{\nu\mu} \\ Q_{\lambda\mu\nu} &= A_{\lambda}(x)g_{\mu\nu}(x) \\ \Gamma^{\lambda}_{\mu\nu} &= C^{\lambda}_{\mu\nu} + W^{\lambda}_{\mu\nu} \\ W^{\lambda}_{\mu\nu} &= -\frac{1}{2}(A_{\mu}\delta^{\lambda}_{\nu} + A_{\nu}\delta^{\lambda}_{\mu} - A^{\lambda}g_{\mu\nu}) \end{split}$$

$$A_{\lambda}(x)$$
 — "Weyl vector"

Local conformal transformation

$$egin{align} ds^2 &= \Omega^2(x) d\hat{s}^2 = \Omega^2 \hat{g}_{\mu
u}(x) dx^\mu dx^
u \ &g_{\mu
u} &= \Omega^2 \hat{g}_{\mu
u}, \quad g^{\mu
u} &= rac{1}{\Omega^2} \hat{g}^{\mu
u} \ &\sqrt{-g} &= \Omega^4(x) \sqrt{-\hat{g}} \ &C^\lambda_{\mu
u} &= \hat{C}^\lambda_{\mu
u} + \left(rac{\Omega_{,
u}}{\Omega} \delta^\lambda_\mu + rac{\Omega_{,\mu}}{\Omega} \delta^\lambda_
u - g^{\lambda
u} rac{\Omega_{,
u}}{\Omega} \hat{g}_{\mu
u}
ight) \end{aligned}$$

If $A_{\mu} = \text{ gauge field}$

$$A_{\mu} = \hat{A}_{\mu} + 2\frac{\Omega_{,\mu}}{\Omega} \implies$$

$$\Gamma^{\lambda}_{\mu\nu} = \hat{\Gamma}^{\lambda}_{\mu\nu} \implies$$

$$R^{\mu}_{\nu\lambda\sigma} = \hat{R}^{\mu}_{\nu\lambda\sigma}, \quad R_{\mu\nu} = \hat{R}_{\mu\nu}$$

Strength tensor

$$F_{\mu
u} =
abla_{\mu}A_{
u} -
abla_{
u}A_{\mu} = A_{
u,\mu} - A_{\mu,
u}$$



Weyl gravity

$$\mathcal{L}_{W} = \alpha_{1} R_{\mu\nu\lambda\sigma} R^{\mu\nu\lambda\sigma} + \alpha_{2} R_{\mu\nu} R^{\mu\nu} + \alpha_{3} R^{2} + \alpha_{4} F_{\mu\nu} F^{\mu\nu}$$
$$S_{W} = \int \mathcal{L}_{W} \sqrt{-g} d^{4}x, \quad \frac{\delta S_{W}}{\delta \Omega} = 0$$

Dynamical variables: $g_{\mu\nu}(x)$, $A_{\mu}x$

Matter fields $S_{\rm tot} = S_{\rm W} + S_{\rm m}, \quad S_{\rm m} = \int \mathcal{L}_{\rm m} \sqrt{-g} \; d^4x \; S_{\rm m}$ is not necessary conformal invariant $\delta S_{\rm m}$ does!

$$\delta \mathcal{S}_{\mathrm{m}} \stackrel{\mathrm{def}}{=} -rac{1}{2} \int T^{\mu
u} (\delta g_{\mu
u}) \sqrt{-g} \ d^4x - \int G^{\mu} (\delta A_{\mu}) \sqrt{-g} \ d^4x + \int rac{\delta \mathcal{L}_{\mathrm{m}}}{\delta \Psi} (\delta \Psi) \sqrt{-g} \ d^4x = 0$$

 Ψ — collective dynamical variable for matter fields, $\delta \mathcal{L}_{\rm m}/\delta \Psi = 0$ $T^{\mu\nu}$ — energy-momentum tensor, G^{μ} — "Weyl current"



Weyl gravity II

$$\begin{split} \delta g_{\mu\nu} &= \frac{2}{\Omega} g_{\mu\nu} (\delta\Omega) + \Omega^2 (\delta \hat{g}_{\mu\nu}) \\ \delta A_{\mu} &= (\delta A_{\mu}) + 2 (\delta (\log \Omega)_{,\mu}) \\ \delta \hat{g}_{\mu\nu} : & \frac{\delta S_{\rm m}}{\delta \hat{g}_{\mu\nu}} \stackrel{\text{def}}{=} -\frac{1}{2} \, \hat{T}_{\mu\nu} \quad \Rightarrow \quad \hat{T}_{\mu\nu} = \Omega^2 \, T_{\mu\nu} \\ (\hat{T}^{\nu}_{\mu} &= \Omega^4 \, T^{\nu}_{\mu}, \ \hat{T}^{\mu\nu} = \Omega^6 \, T_{\mu\nu}) \\ \delta \hat{A}_{\mu} : & \frac{\delta S_{\rm m}}{\delta \hat{A}_{\mu}} \stackrel{\text{def}}{=} -\hat{G}^{\mu} \quad \Rightarrow \quad \hat{G}^{\mu} = \Omega^4 \, G^{\mu} \end{split}$$

Self-consistency condition

$$\delta\Omega$$
: $2G^{\mu}_{;\mu} = Trace[T^{\mu\nu}]$

";" — covariant derivative with $C_{\mu\nu}^{\lambda}$



Cosmology

Cosmology = homogeneity and isotropy Robertson-Walker metric

$$ds^{2} = dt^{2} - a^{2}(t) \left(\frac{dr^{2}}{1 - kr^{2}} + r^{2}(d\theta^{2} + \sin^{2}\theta d\varphi^{2}) \right), \quad \mathbf{k} = 0, \pm 1$$

$$A_{\mu} = (A_{0}(t), 0, 0, 0) \quad \Rightarrow$$

$$F_{\mu\nu} \equiv 0$$

$$T^{\mu}_{\nu} = (T^0_0(t), T^1_1(t) = T^2_2 = T^3_3)$$

Special gauge

$$A_0(t) = 0$$



Field equations

$$\delta A_{\mu}: -6\gamma \dot{R} = G^{0}, \quad \gamma = \frac{1}{3}(\alpha_{1} + \alpha_{2} + 3\alpha_{3})$$

$$R = -6\left(\frac{\ddot{a}}{a} + \frac{\dot{a}^{2} + k}{a^{2}}\right)$$

$$\delta g_{\mu\nu}: -12\gamma\left\{\frac{\dot{a}}{a}\dot{R} + R\left(\frac{R}{12} + \frac{\dot{a}^{2} + k}{a^{2}}\right)\right\} = T_{0}^{0}$$

$$-4\gamma\left\{\ddot{R} + 2\frac{\dot{a}}{a}\dot{R} - R\left(\frac{R}{12} + \frac{\dot{a}^{2} + k}{a^{2}}\right)\right\} = T_{1}^{1}$$

Self-consistency condition:

$$2\frac{(G^0a^3)^{\dot{}}}{a^3}=T_0^0+3T_1^1$$

Perfect fluid, Riemann geometry I

J.R. Ray (1972):

$$S_{\mathrm{m}} = -\int arepsilon(X,n)\sqrt{-g}\ d^4x + \int \lambda_0(u_\mu u^\mu - 1)\sqrt{-g}\ d^4x \ + \int \lambda_1(nu^\mu)_{;\mu}\sqrt{-g}\ d^4x + \int \lambda_2X_{,\mu}u^\mu\sqrt{-g}\ d^4x$$

 $\varepsilon(X, n)$ — energy density

 $u^{\mu}(x)$ — four-velocity

n(x) — particle number density,

X(x) — auxiliary variable,

 $\lambda_i(x)$ — Lagrange multipliers

Constraints

 $u^{\mu}u_{\mu}=1$ — normalization $(nu^{\mu})_{;\mu}=0$ — particle number conservation $X_{,\mu}u^{\mu}=0$ — numbering of the trajectories

Perfect fluid, Riemann geometry II

Energy-momentum tensor

$$T^{\mu\nu} = \varepsilon u^{\mu} u^{\nu} - p g^{\mu\nu}$$

Hydrodynamical pressure

$$p = n \frac{\partial \varepsilon}{\partial n} - \varepsilon$$

V.A.Berezin (1987):

Phenomenological description of particle creation

$$(nu^{\mu})_{;\mu} = 0 \implies (nu^{\mu})_{;\mu} - \Phi(inv) = 0$$

Perfect fluid, Weyl geometry: G^{μ} — ? How to incorporate A_{μ} ?



New possibilities

Single particle u^{μ}

Riemann geometry

$$S_{\text{part}} = -m \int ds$$

Weyl geometry — new invariant $B = A_{\mu}u^{\mu}$

$$S_{\mathrm{part}} = \int f_1(B)ds + \int f_2(B)d\tau = \int \{f_1(B)\sqrt{g_{\mu\nu}u^{\mu}u^{\nu}} + f_2(B)\}d\tau$$

Equations of motion

$$f_{1}(B)u_{\lambda;\mu}u^{\mu} = \left((f_{1}^{'}(B) + f_{2}^{''}(B))A_{\lambda} - f_{1}^{'}(B)u_{\lambda} \right)B_{,\mu}u^{\mu} + (f_{1}^{'}(B) + f_{2}^{'}(B))F_{\lambda\mu}u^{\mu}$$

Since
$$F_{\lambda\mu}u^{\lambda}u^{\mu}\equiv 0$$
 and $u_{\lambda;\sigma}u^{\lambda}\equiv 0$

Either
$$(f_1' + f_2'')B = f_1'$$
 or $B_{,\mu}u^{\mu} = 0$



The invariant $B = A_{\mu}u^{\mu}$ is closely tied to the number density n. Hence,

$$\varepsilon(X, n) \Rightarrow \varepsilon(X, \varphi(B)n)$$

And how about the constraint $(nu^{\mu})_{;\mu} = \Phi(inv)$? Conformal transformation

$$(nu^{\mu})_{;\mu}\sqrt{-g} = (nu^{\mu}\sqrt{-g})_{,\mu}$$

$$n = \frac{\hat{n}}{\Omega^{3}}, \quad u^{\mu} = \frac{\hat{u}^{\mu}}{\Omega}, \quad \sqrt{-g} = \Omega^{4}\sqrt{-\hat{g}} \implies (nu^{\mu}\sqrt{-g})_{,\mu} = (\hat{n}\hat{u}^{\mu}\sqrt{-\hat{g}})_{,\mu} \implies$$

 $\Phi\sqrt{-g}$ — conformal invariant \Rightarrow In the absence of the classical fields, i.e., in the case of the particle creation by the vacuum fluctuations

$$\Phi = \alpha_1' R_{\mu\nu\lambda\sigma} R^{\mu\nu\lambda\sigma} + \alpha_2' R_{\mu\nu} R^{\mu\nu} + \alpha_3' R^2 + \alpha_4' F_{\mu\nu} F^{\mu\nu}$$

Riemann geometry — the only possibility

$$\Phi \propto C^2$$

 $C^2 = C_{\mu\nu\lambda\sigma}C^{\mu\nu\lambda\sigma}$ — square of the Weyl tensor.

Y.B.Zel'dovich and A.A.Starobinsky (1977):

Particle creation by the vacuum fluctuation of the massless scalar field on the homogeneous slightly anisotropic cosmological background.

Now it becomes fundamental!

Gravitational cosmological equations

Vector:

$$-6\gamma\dot{R}=G^0$$

Tensor:

$$-\gamma \left(12\frac{\dot{a}}{a}\dot{R} + R(4R_0^0 - R)\right) = T_0^0$$
$$-12\gamma \left(\ddot{R} + 3\frac{\dot{a}}{a}R\right) = T$$

Self-consistency condition

$$2\frac{(G^0a^3)^2}{a^3} = T_0^0 + 3T_1^1 = T$$

It is quite clear that the self-consistency condition is just the consequence of the vector and trace equations. At last, let us write down the expressions for the $\begin{pmatrix} 0 \\ 0 \end{pmatrix}$ -component of Ricci tensor and the scalar curvature in terms of the scale factor a(t)

$$R_0^0 = -3\frac{\ddot{a}}{a}$$

$$R = -6\left(\frac{\ddot{a}}{a} + \frac{\dot{a}^2 + k}{a^2}\right), \quad k = 0, \pm 1$$

Note that in cosmology we need to know only $T^{00}=T_0^0$ and $T=\operatorname{Trace} T^{\mu\nu}$, since $T^{0i}=0$ and $T^{ij}=T_1^1g^{ij}$, $T_1^1=(1/3)(T-T_0^0)$. Thus,

$$T = T[\text{part}] + T[\text{cr}]$$

$$T[\text{part}] = \varepsilon - 3p$$

$$T[\text{cr}] = \ddot{\lambda}_1 (8\beta' R_0^0 - 4\beta' R - 12\gamma' R)$$

$$-4\dot{\lambda}_1 \left(\beta' \frac{\dot{a}}{a} (R + 2R_0^0) + 6\gamma' \dot{R} + 9\gamma' \frac{\dot{a}}{a} R\right) - 12\lambda_1 \gamma' (\ddot{R} + 3\frac{\dot{a}}{a} R)$$

$$T_0^0 = T_0^0[\text{part}] + T_0^0[\text{cr}], \quad T_0^0[\text{part}] = \varepsilon$$

$$T_0^0[\mathbf{cr}] = 8\gamma' \dot{\lambda}_1 \frac{\dot{a}}{a} R_0^0 - 4(\beta' + 3\gamma') \dot{\lambda}_1 \frac{\dot{a}}{a} R$$
$$-\gamma' \lambda_1 \left(12 \frac{\dot{a}}{a} \dot{R} + R(4R_0^0 - R) \right)$$

How about the equations of motion for the cosmological perfect fluid? We are left with only one equation plus the law of the particle creation, namely

$$\left\{ \begin{array}{l} \dot{\lambda}_1 = -\frac{\varepsilon + p}{n} \\ \frac{(na^3)}{a^3} = \Phi(\mathrm{inv}) \end{array} \right.$$

$$\Phi(\text{inv}) = -\frac{4}{3}\beta' R_0^0 (2R_0^0 - R) + \gamma' R^2.$$

Creation of the universe from nothing

Empty from the very beginning Vacuum is not absolute, but physical May or may not it persists?

"Pregnant" vacuum

$$\Phi(\text{inv}) = 0, \quad |\beta'| + |\gamma'| \neq 0$$

$$\frac{4}{3}\beta' R_0^0 (2R_0^0 - R) = \gamma' R^2$$

$$G^0[\text{part}] = T_0^0[\text{part}] = T[\text{part}] = 0$$

$$n = 0$$

$$\dot{\lambda}_1 = -\frac{\varepsilon + p}{n} \quad \Rightarrow$$

Non-dust matter

$$\dot{\lambda}_1 = 0 \Rightarrow \lambda_1 = const$$

2 Dust matter

$$\dot{\lambda}_1 = -\phi(0) = const \Rightarrow \lambda_1 = -\phi(0)(t - t_0)$$



Non-dust pregnancy I

General case: β' , $\gamma' \neq 0$

$$\begin{split} \lambda_1 &= const \\ R &= \xi R_0^0 \quad \Rightarrow \\ &(3\gamma'\xi^2 + 4\beta'(\xi - 2))R_0^0 = 0 \\ &- 6(\gamma - \gamma'\lambda_1)\dot{R} = 0 \\ &- 12(\gamma - \gamma'\lambda_1)\frac{(\dot{R}a^3)}{a^3} = 0 \\ &- (\gamma - \gamma'\lambda_1)\left\{12\frac{\dot{a}}{a}\dot{R} + R(R - 4R_0^0)\right\} = 0 \end{split}$$
 Let, first, $\gamma \neq \gamma'\lambda_1 \quad \Rightarrow \quad \dot{R} = 0$ Either $R = 0 \quad \Rightarrow \quad R_0^0 = 0 \quad \Rightarrow \quad \text{Milne universe}$ Or $\xi = 4 \quad \Rightarrow \quad \text{de Sitter for } \beta + 6\gamma' = 0$

Non-dust pregnancy II

General case: β' , $\gamma' \neq 0$

$$\gamma = \gamma' \lambda_1$$

It is not a special condition, but the solution λ_1

$$\dot{a}^2 + k = C_0 a^{\frac{4}{\xi - 2}}$$

Instability and so on...

Dust pregnancy

$$\lambda_1 = -\phi(0)(t - t_0) \quad \Rightarrow \quad \dot{\lambda}_1 = -\phi < 0, \quad \ddot{\lambda}_1 = 0$$

Surely, there exists the solution with $R = R_0^0 = 0$, i. e., Milne universe.

It can be shown that the only other solution is just the de Sitter universe, with $\xi = 4$ (R, $R_0^0 = const$) for $\beta' + 6\gamma' = 0$.

If it is not so, the universe emerged from the vacuum foam, like Aphrodite, immediately starts to produce dust particles !!!

The End

