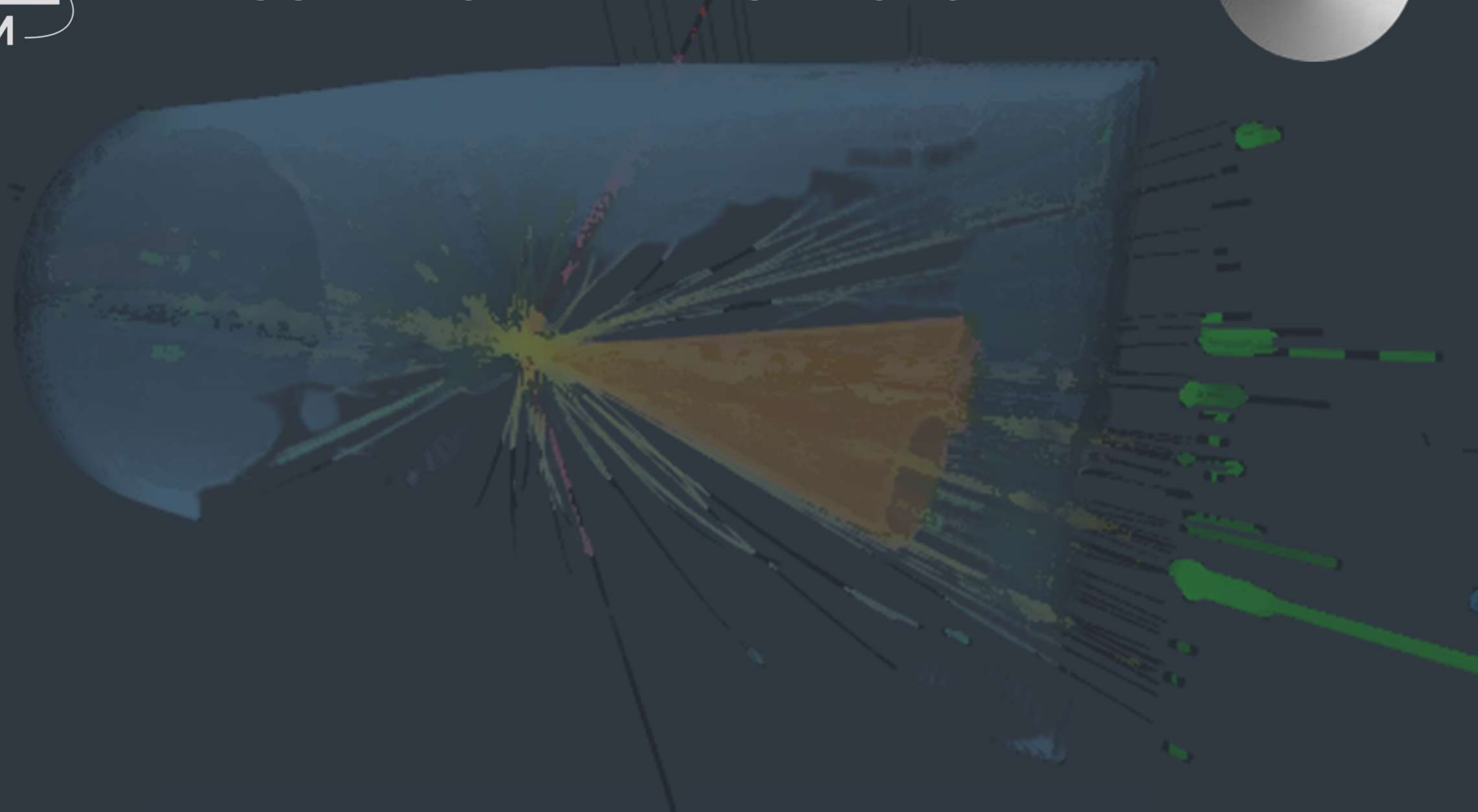
MILLICHARGED PARTICLES AT SUPER CHARM-TAU FACTORY



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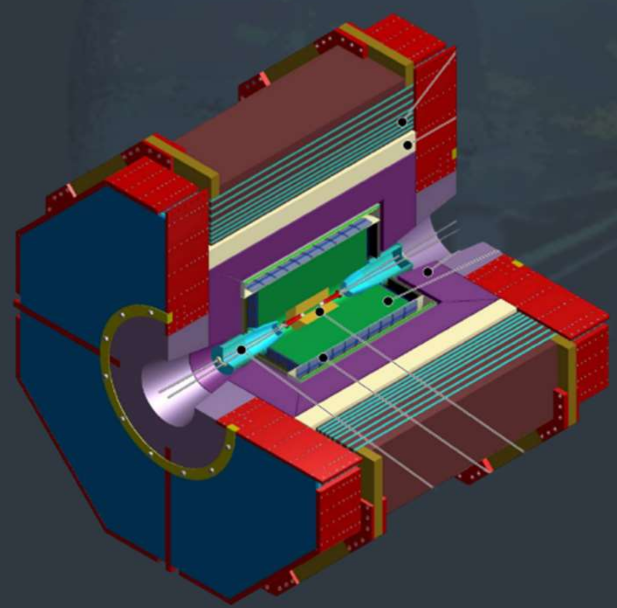
MILLICHARGED PARTICLES
Super Charm-Tau Factory

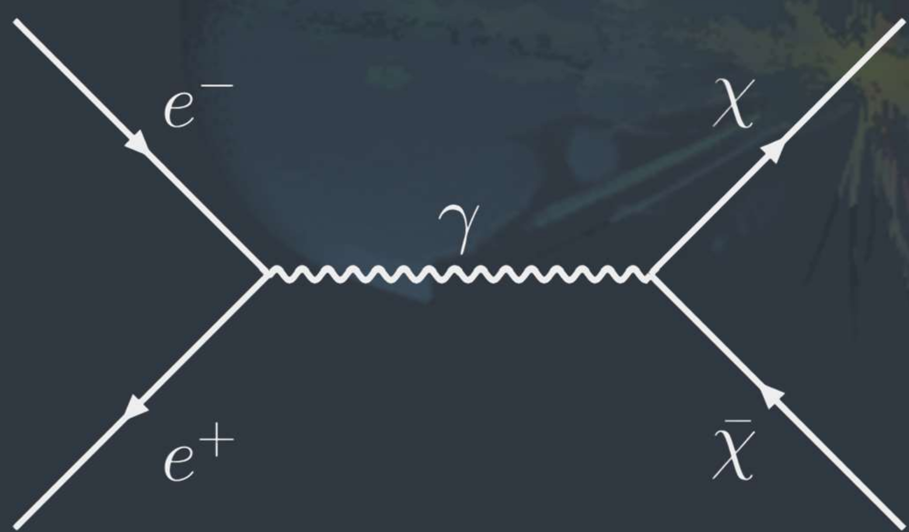
e^+e^- - collider in Sarov

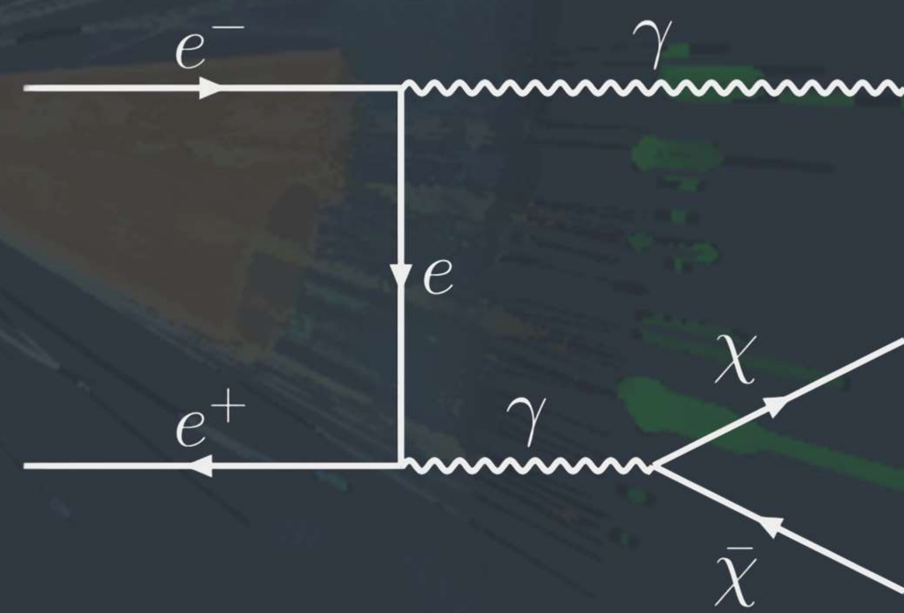


$\sqrt{s}, GeV$	L, fb^{-1}
3.097	300
3.554	50
3.686	150
3.770	300
4.110	100
4.650	100

MILLICHARGED PARTICLES
Couplings

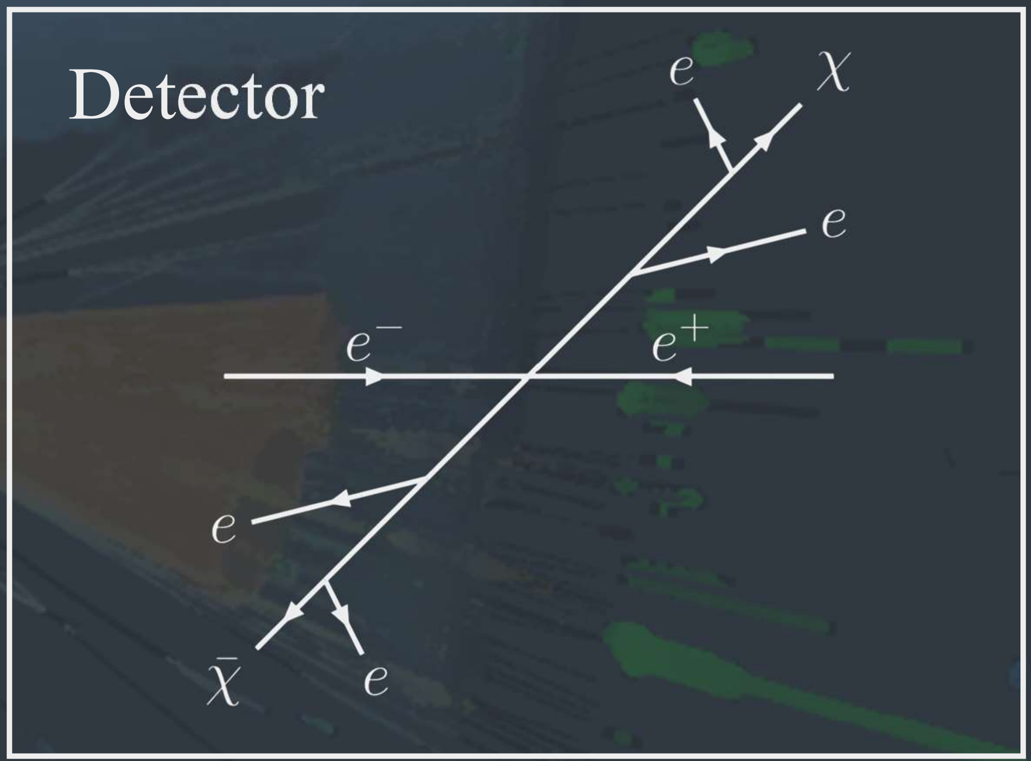
$$\mathcal{L} = \epsilon e\, A_\mu \bar{\chi} \gamma^\mu \chi \,, \qquad \epsilon = \frac{Q_\chi}{e}$$





MILLICHARGED PARTICLES
Signal for direct searches

Energy loss



δ-electrons

MILLICHARGED PARTICLES
Limits. Energy loss

$$\frac{dE}{dx} = \frac{\rho K}{2}\frac{Z}{A} \times \frac{\boldsymbol{\epsilon^2}}{\boldsymbol{\beta^2}} \times ln\left(\frac{2m_e\beta^2T_{max}}{I^2}\right) \Rightarrow \frac{dE}{dx} \approx 25\ eV/cm \times \left(\frac{\epsilon}{\epsilon_0} \times \frac{\beta_0}{\beta}\right)^2$$

$$\frac{dE}{dx} > 25\ eV/cm$$

$$\frac{\epsilon}{\beta} > 0.18$$

MILLICHARGED PARTICLES
Limits. δ-electons

$$N_\delta = \frac{\rho K}{2}\frac{Z}{A} \times \frac{\boldsymbol{\epsilon}^2}{\boldsymbol{\beta}^2} \times L\left(\frac{1}{T_{min}} - \frac{1}{T_e}\right) \Rightarrow N_\delta = 28.6\left(\frac{\epsilon}{\epsilon_0} \times \frac{\beta_0}{\beta}\right)^2 \times \left(1 - \frac{1\,keV}{E_\chi}\right)$$

$$N_\delta > 2$$

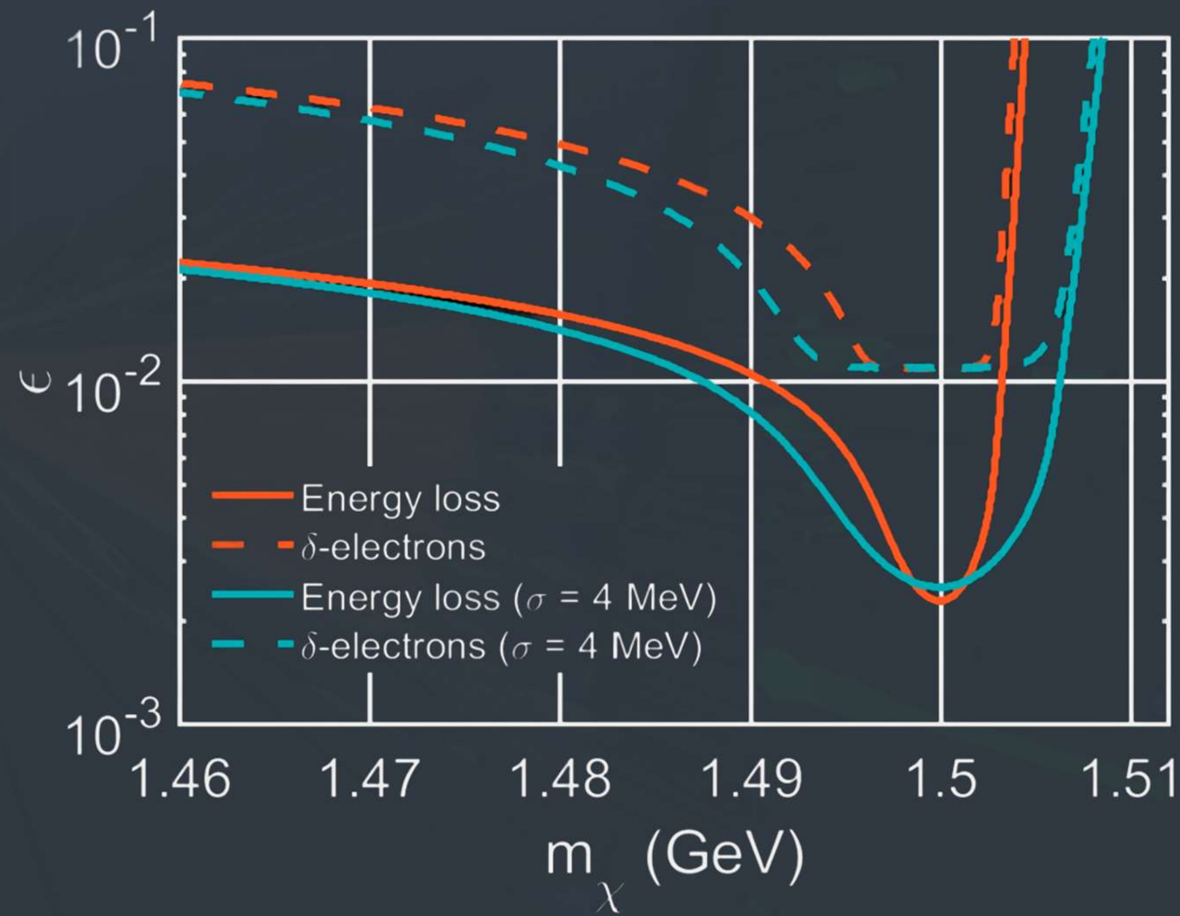
$$\frac{\epsilon}{\beta} > 0.343 \times \left(1 - \frac{1\,keV}{E_\chi}\right)^{-1/2}$$

MILLICHARGED PARTICLES
Non-monochromatic energy

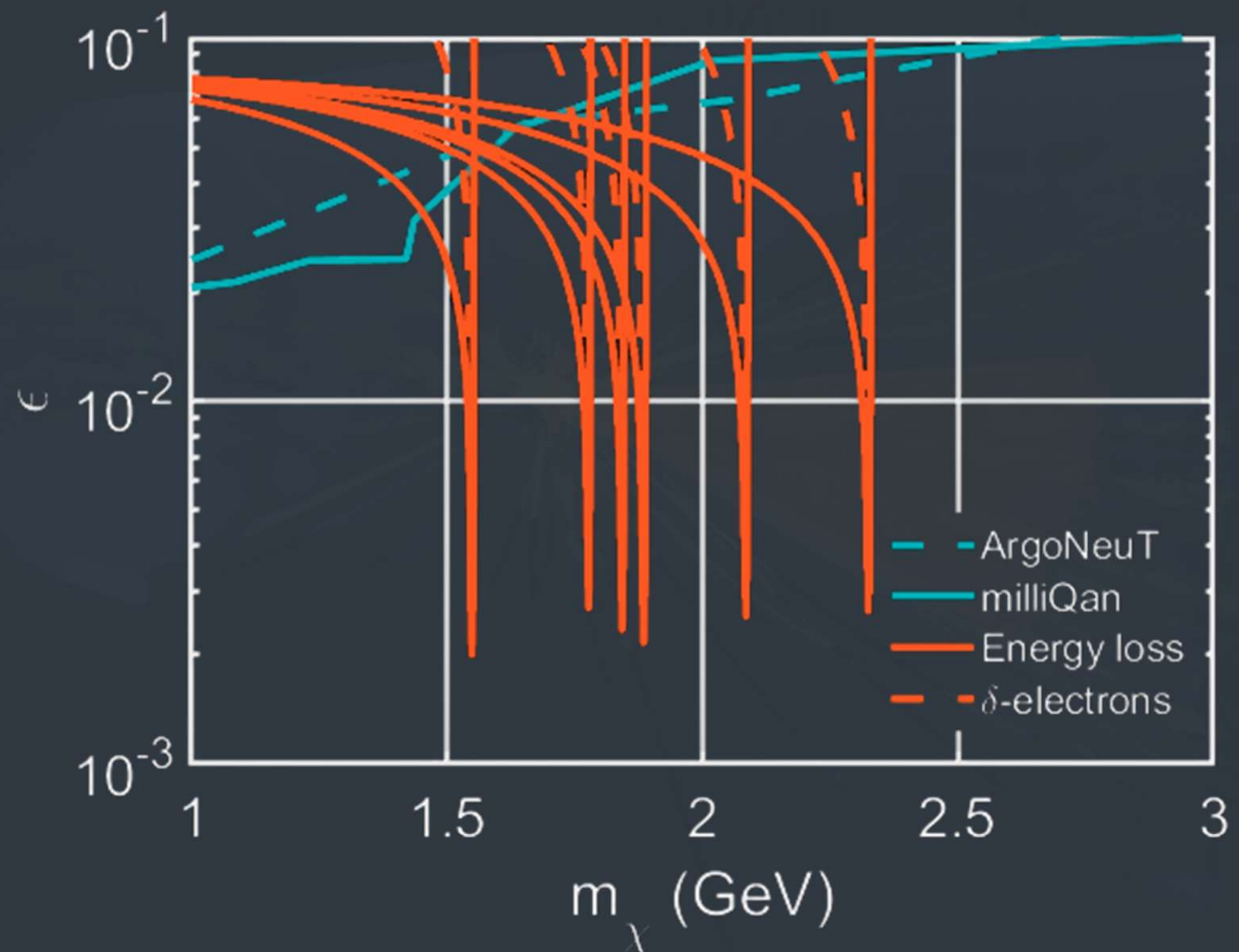
$$\frac{\mathrm{d}L}{d\sqrt{s}} = \frac{L_0}{\sqrt{2\pi}\sigma} \times exp\left(-\frac{\left(\sqrt{s}-\sqrt{s_0}\right)^2}{2\sigma^2}\right)$$

$$\sigma = \sqrt{2} \times 0.01\% \ \frac{\sqrt{s}}{2}$$

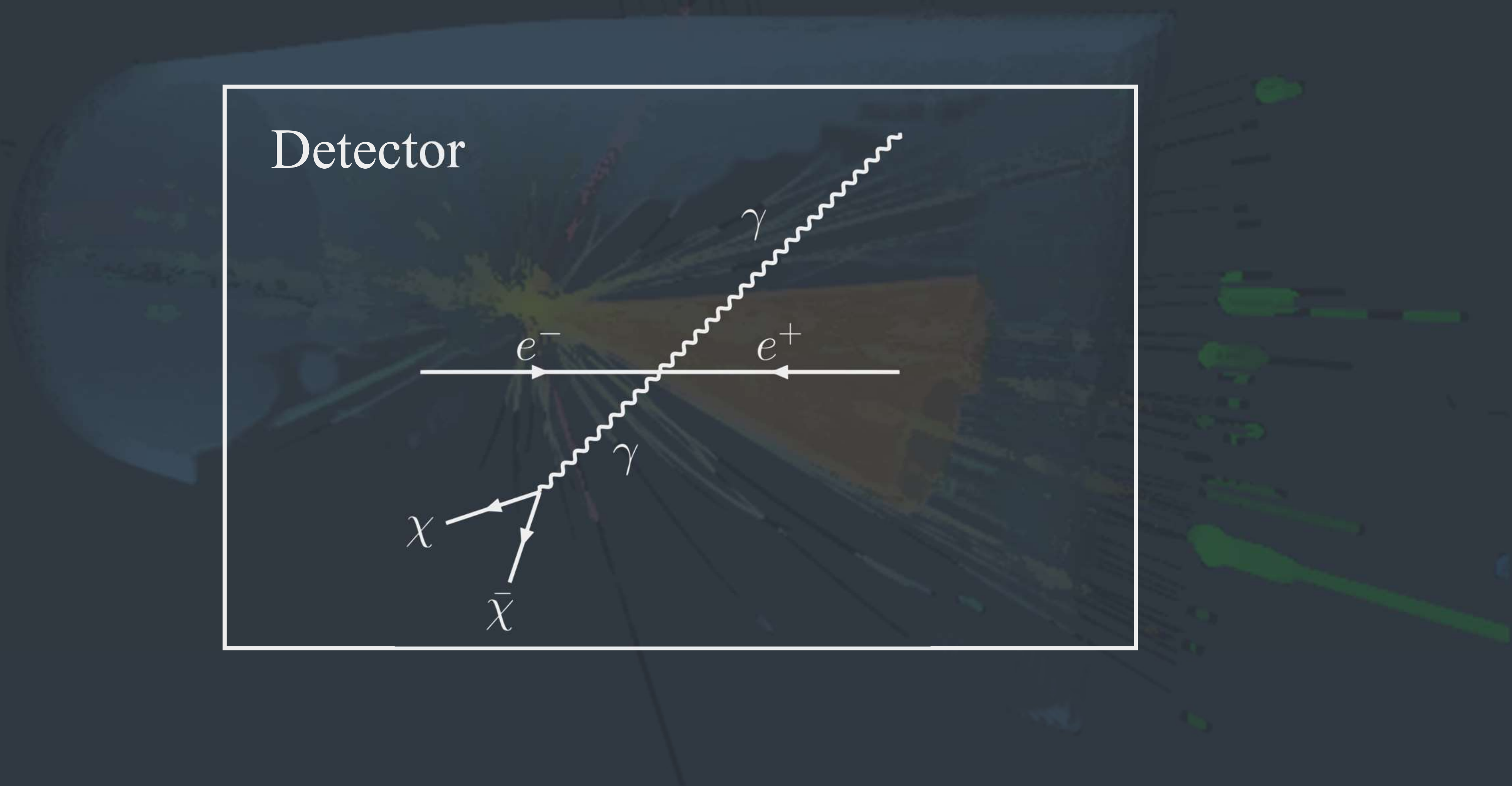
$$\sigma(e^+e^- \to \chi\bar{\chi}) \propto \varepsilon^2\beta$$



MILLICHARGED PARTICLES
Direct



MILLICHARGED PARTICLES
Signal with missing energy



MILLICHARGED PARTICLES
Background

Reducible

$e^+e^- \rightarrow \gamma\ e^+e^-\ (\mu^+\mu^-, \gamma\gamma)$

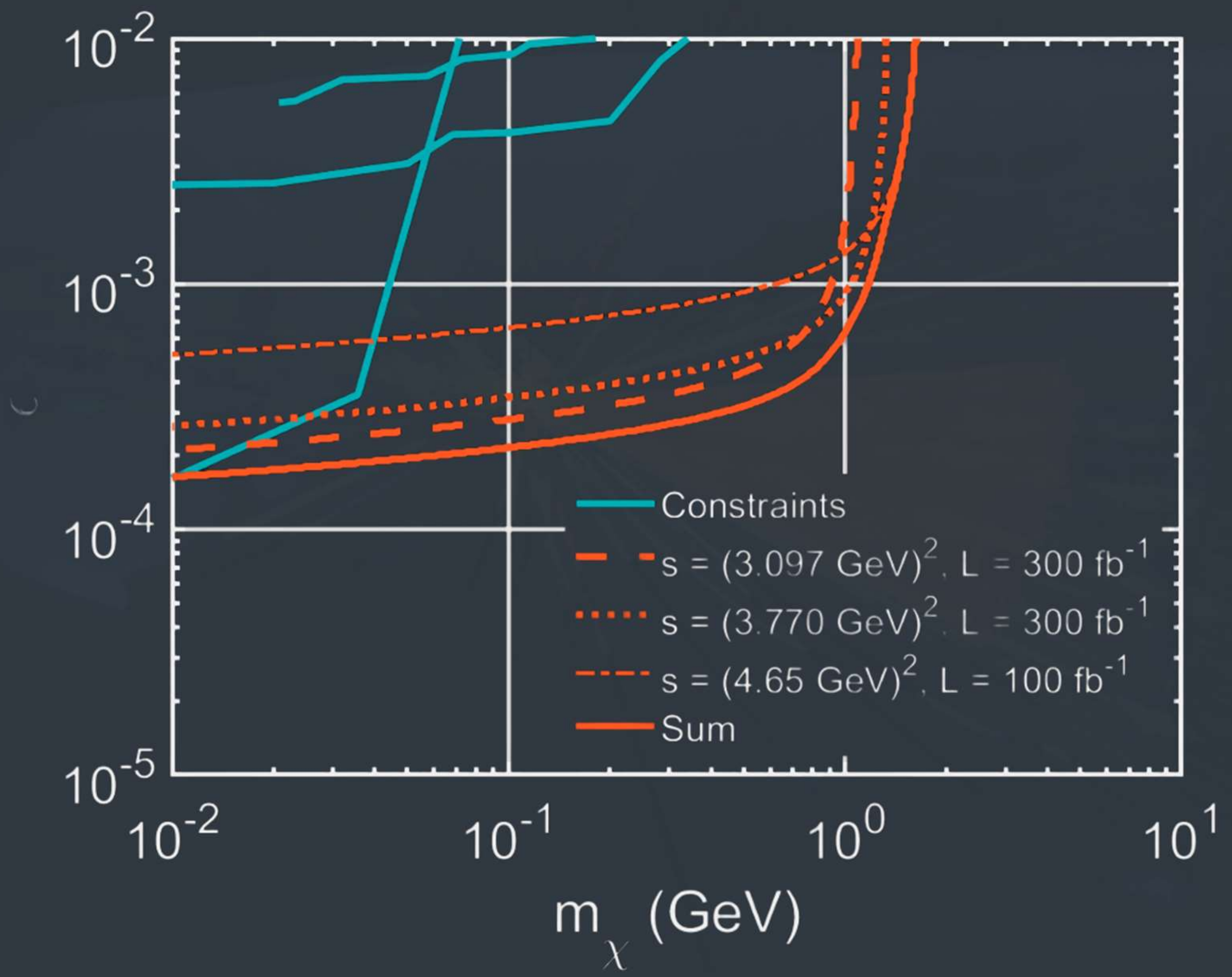
$E_\gamma > E_{\gamma min}(cos\theta_\gamma)$

Irreducible

$e^+e^- \rightarrow \gamma\nu\nu$

e^- - beam polarization

MILLICHARGED PARTICLES
Invisible



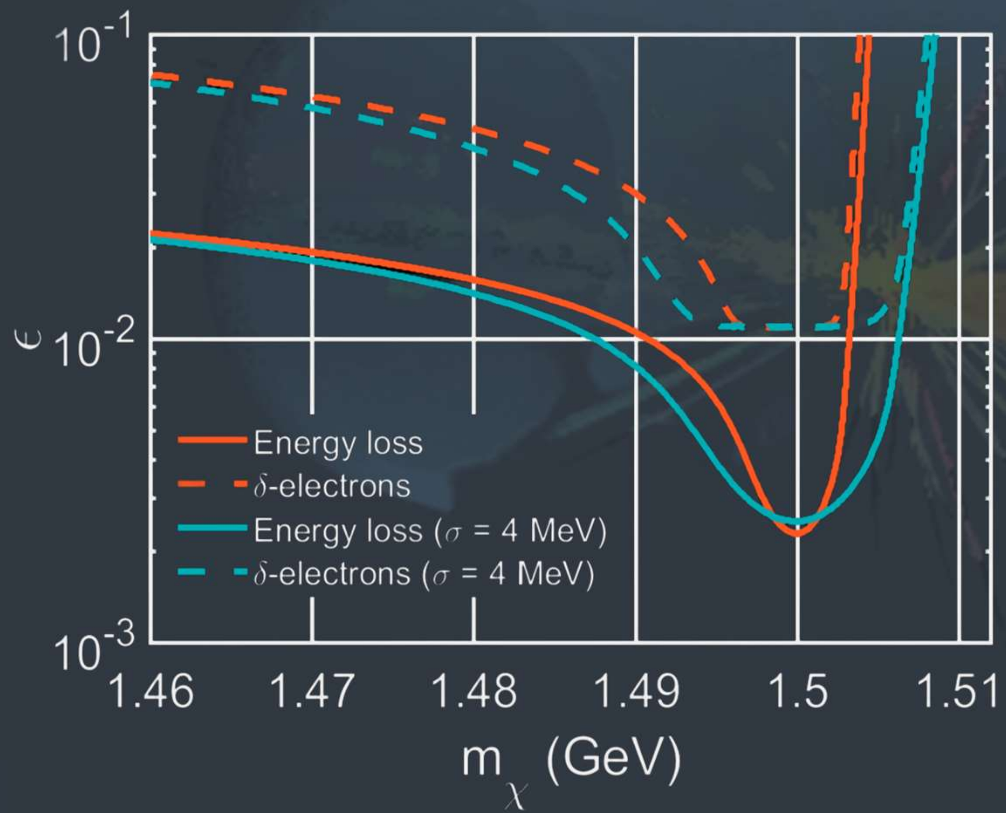
MILLICHARGED PARTICLES
Ways to increase sensitivity

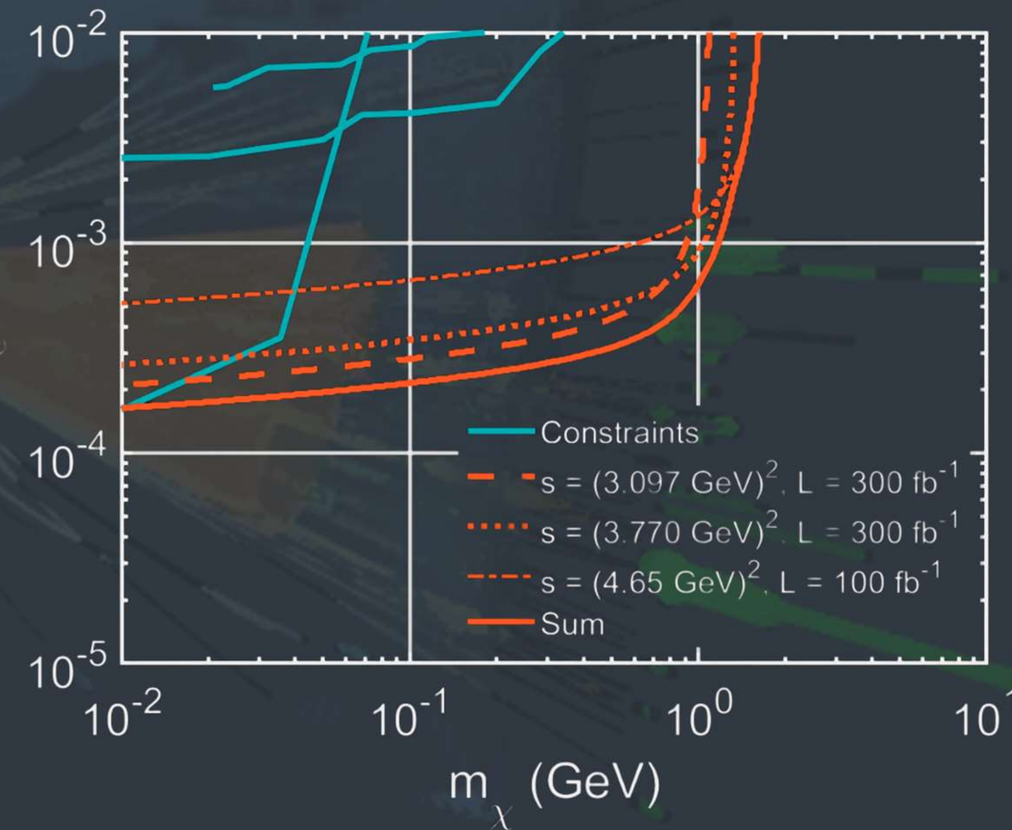
- Greater luminosity
  For direct detection energy limits remain. For invisible signal we also get larger background but we can apply additional cuts to reduce it

- Monochromatic beam
  More particles with small β but suppressed production $\sigma \propto \epsilon^2 \beta$

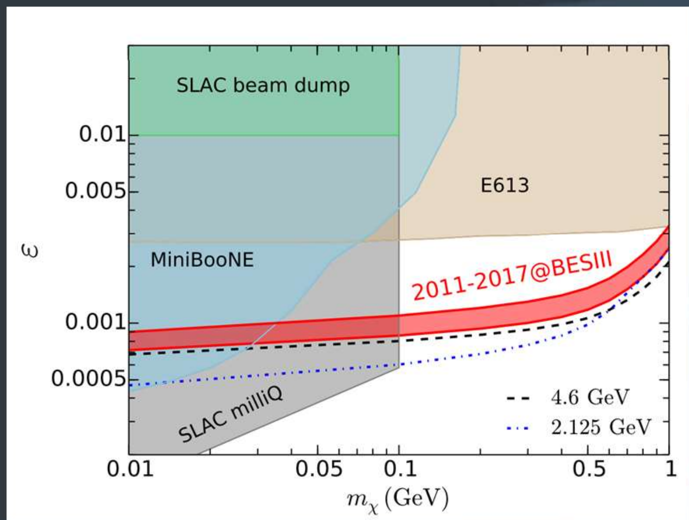
- Other gases
  To increase weaken the limits on ϵ/β but it alters the zero background condition

MILLICHARGED PARTICLES
Ways to increase sensitivity

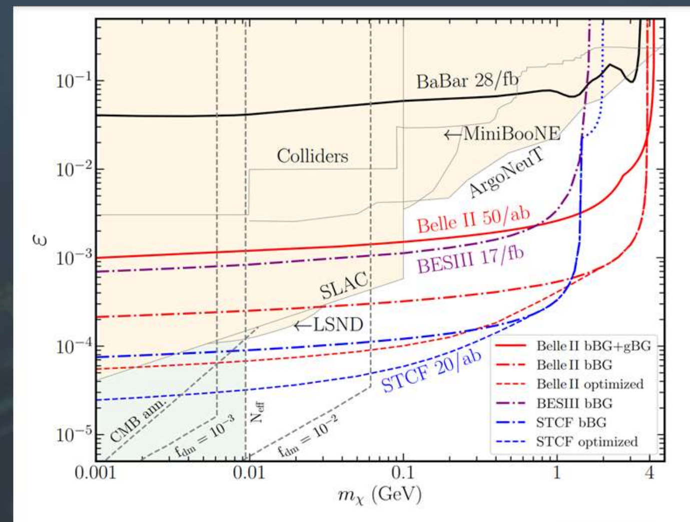




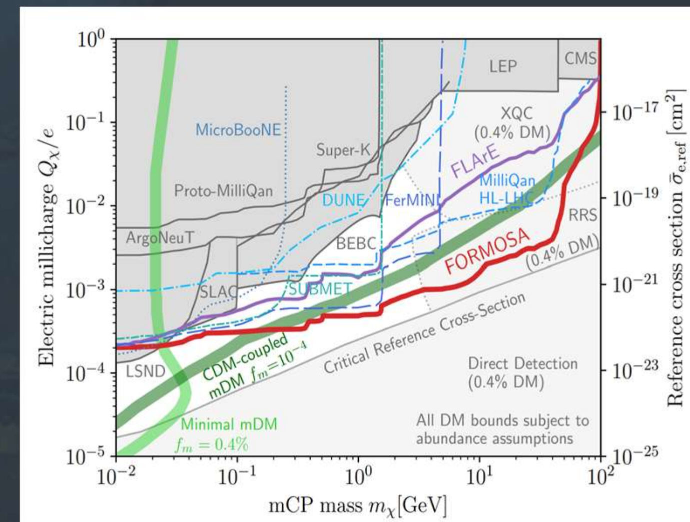
MILLICHARGED PARTICLES
Other studies



Zuowei Liu, Yu Zhang. Probing millicharge at BESIII. Phys. Rev. D 99, 015004 (2019)

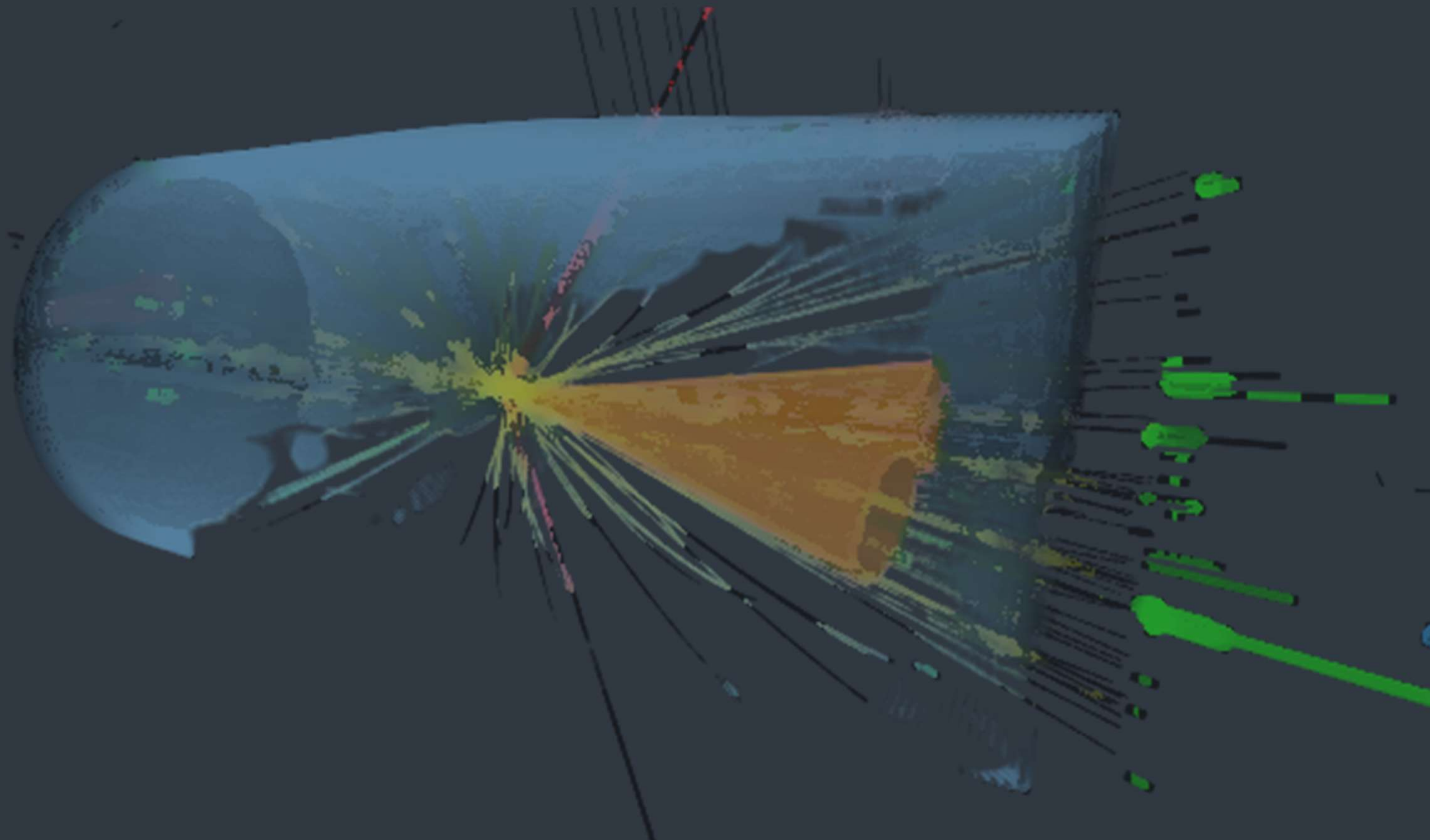


Jinhan Liang, Zuowei Liu, Yue Ma, Yu Zhang. Millicharged particles at electron colliders. Phys. Rev. D 102, 015002 (2020)



Jonathan L. Feng et al. The Forward Physics Facility at the High-Luminosity LHC. Contribution to Snowmass 2021

MILLICHARGED PARTICLES



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MILLICHARGED PARTICLES
Initial State Radiation

$$\sigma_{corr} = \int_0^{x_{max}} \sigma\big((1-x)s\big)\, H(x,s)$$

$$E_\gamma < 25\ MeV$$