BSM physics at future lepton colliders

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Summary

- 1. Light-by-light (LBL) scattering $\gamma\gamma \rightarrow \gamma\gamma$ of Compton backscattered (CB) photons at high energy collider CLIC. Probing anomalous quartic $\gamma\gamma\gamma\gamma$ couplings
- 2. Search for anomalous gauge couplings of $\gamma\gamma\gamma Z$ vertex in $\gamma\gamma \rightarrow \gamma Z$ collision of CB photons at the CLIC
- **3.** Study of invisible massive dark photon in γe scattering at future lepton colliders ILC, CLIC, and CEPC.
- 4. Conclusions

Light-by-light scattering of CB photons at the CLIC



Anomalous diagram for process $\gamma\gamma \rightarrow \gamma\gamma$

The goal is to examine anomalous quartic gauge couplings (QGCs) via reaction $\gamma\gamma \rightarrow \gamma\gamma$ at the CLIC

New physics contributions to 4y couplings

New charged particles via loops

ζ_i ~ Q⁴ m⁻⁴ Example: top partner

New neutral particles at tree level

$$\zeta_{i} \sim f^{-2} m^{-2}$$

Example: KK gravitons, radion (warped extra dimension)

If $f_{KK} \sim \text{TeV}$ and $m_{KK} \sim \text{few TeV}$, then $\zeta_i \sim 10^{-2} - 10^{-1}$ TeV⁻⁴



Compact Linear Collider (CLIC)



	Collision energy	(unpolarized beams)
1st stage	380 GeV	1.0 ab ⁻¹
2nd stage	1500 GeV	2.5 ab ⁻¹
3rd stage	3000 GeV	5.0 ab ⁻¹

Tests are shared to us in a site.

 λ_0 is helicity of laser photon beam, and λ_e is helicity of electron beam before Compton backscattering

 $\begin{aligned} & (\lambda_e^{(1)}, \lambda_0^{(1)}; \lambda_e^{(2)}, \lambda_0^{(2)}) = (0.8, 1; 0.8, 1) , \\ & (\lambda_e^{(1)}, \lambda_0^{(1)}; \lambda_e^{(2)}, \lambda_0^{(2)}) = (-0.8, 1; -0.8, 1) \end{aligned}$

CLIC energy stages and integrated luminosities for unpolarized and polarized electron beams

		L, fb^{-1}		
Stage	\sqrt{s}, GeV	$\lambda_e = 0$	$\lambda_e = -0.8$	$\lambda_e = 0.8$
2	1500	2500	2000	500
3	3000	5000	4000	1000

Differential cross section for LBL scattering

$$\frac{d\sigma}{d\cos\theta} = \frac{1}{128\pi s} \int_{x_{1}\min}^{x_{\max}} \frac{dx_{1}}{x_{1}} f_{\gamma/e}(x_{1}) \int_{x_{2}\min}^{x_{\max}} \frac{dx_{2}}{x_{2}} f_{\gamma/e}(x_{2})$$

$$\times \left\{ \left[1 + \xi \left(E_{\gamma}^{(1)}, \lambda_{0}^{(1)} \right) \xi \left(E_{\gamma}^{(2)}, \lambda_{0}^{(2)} \right) \right] \right.$$

$$\times \left(|M_{++++}|^{2} + |M_{++--}|^{2} \right) + \left[1 - \xi \left(E_{\gamma}^{(1)}, \lambda_{0}^{(1)} \right) \xi \left(E_{\gamma}^{(2)}, \lambda_{0}^{(2)} \right) \right]$$

$$\times \left(|M_{+-+-}|^{2} + |M_{+--+}|^{2} \right) \right\}$$

 $f_{\gamma/e}(x) = CB$ photon distribution $\xi(E_{\gamma},\lambda_0) = CB$ photon helicity in energy fraction $x = E_{\gamma}/E_e$

$$x_{1,min} = (p_t/E_e)^2$$
, $x_{2,min} = (p_t/x_1E_e)^2$, $x_{max} = 0.83$

Effective Lagrangian

Effective Lagrangian of dimension-8 operators which contribute to anomalous quartic neutral gauge couplings

$$\begin{aligned} \mathcal{L}_{\text{QNGC}} &= \frac{c_8}{\Lambda^4} B_{\rho\sigma} B^{\rho\sigma} B_{\mu\nu} B^{\mu\nu} + \frac{c_9}{\Lambda^4} W^a_{\rho\sigma} W^{a\rho\sigma} W^b_{\mu\nu} W^{b\mu\nu} + \frac{c_{10}}{\Lambda^4} W^a_{\rho\sigma} W^{b\rho\sigma} W^a_{\mu\nu} W^{b\mu\nu} \\ &+ \frac{c_{11}}{\Lambda^4} B_{\rho\sigma} B^{\rho\sigma} W^a_{\mu\nu} W^{a\mu\nu} + \frac{c_{13}}{\Lambda^4} B_{\rho\sigma} B^{\sigma\nu} B_{\nu\mu} B^{\mu\rho} + \frac{c_{14}}{\Lambda^4} W^a_{\rho\sigma} W^{a\sigma\nu} W^b_{\nu\mu} W^{b\mu\rho} \\ &+ \frac{c_{15}}{\Lambda^4} W^a_{\rho\sigma} W^{b\sigma\nu} W^a_{\nu\mu} W^{b\mu\rho} + \frac{c_{16}}{\Lambda^4} B_{\rho\sigma} B^{\sigma\nu} W^a_{\nu\mu} W^{a\mu\rho} \end{aligned}$$

Effective Lagrangian for anomalous yyyy couplings (in terms of physical fields)

$$L_{\rm QNGC}^{\gamma\gamma\gamma\gamma} = \varsigma_1(F_{\mu\nu}F^{\mu\nu})(F_{\rho\sigma}F^{\rho\sigma}) + \varsigma_2(F_{\mu\nu}F^{\nu\rho}F_{\rho\sigma}F^{\sigma\mu})$$

Anomalous couplings ζ_1 , ζ_2 have dimension -4

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Numerical analysis



Electron beam helicities:

 $\lambda_e = 0.8, -0.8, \text{ and } 0 \text{ (unpolarized case)}$

Cuts on the invariant mass and rapidities of final photons:

m_{γγ} > 200 GeV, |η|< 2.5

SM amplitude: (G. Gounaris et al., EPJC, 9, 673, 1999)

M_w (W-loop, dominates at Vs > 200 GeV) + M_f (fermion-loop) contributions



Total cross sections for $\gamma\gamma \rightarrow \gamma\gamma$ scattering versus minimal invariant mass of outgoing photons for $\sqrt{s} = 3$ TeV. The left, middle and right panels correspond to $\lambda_e = 0.8$, -0.8, 0. The solid curves (from top downwards): $(\zeta_1 = 10^{-13} \text{ GeV}^{-4}, \zeta_2 = 0), \ (\zeta_1 = 0, \zeta_2 = 10^{-13} \text{ GeV}^{-4}), \text{ SM.}$

Exclusion significance (δ = percentage systematic error) (Y. Zhang & J. Shen, EPJC, 80, 811, 2020)

$$S_{\text{excl}} = \sqrt{2} \left[s - b \ln \left(\frac{b + s + x}{2b} \right) - \frac{1}{\delta^2} \ln \left(\frac{b - s + x}{2b} \right) - (b + s - x) \left(1 + \frac{1}{\delta^2 b} \right) \right]^{1/2}$$

s (b) = number of signal (background) events

In the limit $\delta = 0$ $S_{\text{excl}} = \sqrt{2\left[s - b\ln\left(1 + \frac{s}{b}\right)\right]}$ $s \ll b$ $S_{\text{excl}} = \frac{s}{\sqrt{b}}$

> S_{excl} ≤ 1.645 is the region that can be excluded at 95% C.L.



95% C.L. exclusion region for anomalous couplings ζ₁, ζ₂ for unpolarized LBL scattering at the CLIC. Systematic errors are 0%, 5%, and 10%. The collision energy is 3 TeV, the integrated luminosity is 5 ab⁻¹. Couplings ζ₁, ζ₂ are in GeV⁻⁴.

Exclusion limits on anomalous couplings for the LBL scattering at the CLIC with energy 3.0 TeV

Helicity		0	-0.8	0.8
Luminosity, fb^{-1}		5000	4000	1000
8	$\delta = 0\%$	6.85×10^{-16}	8.82×10^{-16}	8.73×10^{-16}
$ \zeta_1 , \text{GeV}^{-4} (\zeta_2 = 0)$	$\delta = 5\%$	1.90×10^{-15}	2.48×10^{-15}	1.56×10^{-15}
	$\delta = 10\%$	2.63×10^{-15}	3.37×10^{-15}	2.12×10^{-15}
	$\delta = 0\%$	1.43×10^{-15}	1.85×10^{-15}	1.82×10^{-15}
$ \zeta_2 , \text{GeV}^{-4} (\zeta_1 = 0)$	$\delta = 5\%$	3.99×10^{-15}	5.12×10^{-15}	3.28×10^{-15}
	$\delta = 10\%$	5.53×10^{-15}	7.10×10^{-15}	4.46×10^{-15}

(S.Inan & A.K., EPJC, 81, 664, 2021)

Bounds on anomalous couplings for LHC and HL-LHC

LHC, L=300 fb⁻¹: |ζ₁| < 1.5*10⁻¹⁴ GeV⁻⁴, |ζ₂| < 3.0*10⁻¹⁴ GeV⁻⁴

(S.Fichet et al., JHEP 02, 165, 2015)

HL-LHC, L=3000 fb⁻¹: |ζ₁| < 7.0*10⁻¹⁵ GeV⁻⁴, |ζ₂| < 1.5*10⁻¹⁴ GeV⁻⁴

(S.Fichet et al., Phys. Rev. D 89, 114004, 2014)

vZ production in photon-photon scattering

Anomalous diagram for process $\gamma\gamma \rightarrow \gamma Z$



Effective Lagrangian for anomalous yyyZ couplings

$$L_{\text{QNGC}}^{\gamma\gamma\gamma\overline{Z}} = g_1(F^{\rho\mu}F^{\alpha\nu})(\partial_{\rho}F_{\mu\nu}Z_{\alpha}) + g_2(F^{\rho\mu}F_{\mu\nu})(\partial_{\rho}F_{\alpha\nu}Z^{\alpha})$$

Anomalous couplings g₁, g₂ have dimension -4

Feynman rule for anomaly vertex yyyZ

$$P^{\mu\nu\rho\alpha} = \mathcal{P}\left[g_1 [(p_1 \cdot p_2)(p_2 \cdot p_3)g^{\mu\nu}g^{\rho\alpha} - (p_1 \cdot p_3)p_2^{\mu}p_1^{\nu}g^{\rho\alpha} - (p_1 \cdot p_3)p_1^{\nu}p_2^{\alpha}g^{\mu\rho} + p_2^{\mu}p_1^{\nu}p_1^{\alpha}g^{\alpha} \right] + g_2 [-(p_1 \cdot p_2)(p_1 \cdot p_3)g^{\mu\alpha}g^{\nu\rho} + (p_2 \cdot p_3)p_1^{\nu}p_1^{\alpha}g^{\mu\rho} - (p_2 \cdot p_3)p_1^{\nu}p_1^{\beta}g^{\mu\alpha} + (p_2 \cdot p_3)p_1^{\nu}p_2^{\alpha}g^{\mu\rho} + 2(p_2 \cdot p_3)p_2^{\mu}p_1^{\rho}g^{\nu\alpha} - (p_1 \cdot p_3)p_2^{\rho}p_1^{\alpha}g^{\mu\nu} + p_3^{\mu}p_1^{\nu}p_2^{\rho}p_1^{\alpha} \right] \}$$

P- permutations (symmetrization with respect to photon's momenta and indices)

 $M_{\lambda_1\lambda_2\lambda_3\lambda_4}(p_1, p_2, p_3) = P_{\mu\nu\rho\alpha}(p_1, p_2, p_3) \,\varepsilon_{\mu}^{\lambda_1}(p_1) \varepsilon_{\nu}^{\lambda_2}(p_2) \varepsilon_{\rho}^{*\lambda_3}(p_3) \varepsilon_{\alpha}^{*\lambda_4}(p_4)$



24 helicity amplitudes proportional to g₁ 24 helicity amplitudes proportional to g₂ (S. Inan & A.K., JHEP, 10, 121, 2021)

SM amplitudes are taken from G. Gounaris et al., EPJC, 10, 499, 1999



Total cross sections for $\gamma\gamma \rightarrow \gamma Z$ scattering versus minimal invariant mass of outgoing photons for Vs = 3 TeV. The left, middle and right panels correspond to $\lambda_e = 0.8$, -0.8, 0. Solid curves correspond to $(g_1 = 10^{-13} \text{ GeV}^{-4}, g_2 = 0)$, SM.



95% C.L. exclusion region for anomalous couplings g_1 , g_2 for unpolarized $\gamma\gamma \rightarrow \gamma Z$ scattering at the CLIC. Systematic errors are 0%, 5%, and 10%. The collision energy is 3 TeV, the integrated luminosity is 5 ab⁻¹. Couplings g_1 , g_2 are in GeV⁻⁴.

Exclusion limits on anomalous couplings at the CLIC for energy 3.0 TeV. Polarized case

λ_e		0	-0.8	0.8
L, fb^{-1}		5000	4000	1000
$ g_1 , \mathrm{GeV}^{-4}$ $(g_2 = 0)$	$\delta = 0\%$	5.98×10^{-15}	7.14×10^{-15}	5.13×10^{-15}
	$\delta = 5\%$	1.33×10^{-14}	1.73×10^{-14}	7.79×10^{-15}
	$\delta = 10\%$	1.85×10^{-14}	2.39×10^{-14}	1.04×10^{-14}
$ g_2 , \mathrm{GeV}^{-4}$ $(g_1 = 0)$	$\delta = 0\%$	5.18×10^{-15}	6.62×10^{-15}	5.19×10^{-15}
	$\delta = 5\%$	1.16×10^{-14}	1.60×10^{-14}	7.87×10^{-15}
	$\delta = 10\%$	1.62×10^{-14}	2.21×10^{-14}	1.05×10^{-14}

(S.Inan & A.K., JHEP, 10, 121, 2021)

Bounds on anomalous couplings for the LHC and HL-LHC

SM expectation: $B(Z \rightarrow \gamma \gamma \gamma) = 5.41^{*}10^{-10}$

B(Z → γγγ) < 0.8*10⁻⁵ (LEP) **B**(Z → γγγ) < 2.2*10⁻⁶ (ATLAS)

 $if_i < 1.3*10^{-9} \text{ GeV}^{-4}$ where $f_i = g_i/8$

From $\gamma\gamma \rightarrow \gamma Z$ at LHC: LHC, L=300 fb⁻¹: $|f_i| < 1.0*10^{-13} \text{ GeV}^{-4}$ HL-LHC, L=3000 fb⁻¹: $|f_i| < 7.8*10^{-14} \text{ GeV}^{-4}$

(C. Baldenegro et al., JHEP 06, 142, 2017)

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Invisible dark photons in ve scattering at future lepton colliders ILC, CLIC, CEPC

We consider dark matter (DM) scenario in which no DM are charged under SM gauge group

Lightest stable DM particles can only interact with SM through exchange of massive vector mediator, dark photon (DP) A'

DP kinematically mixes with SM U(1)_Y hypercharge gauge field (kinematic mixing portal)

Such mixing can be generated by loops of massive particles charged under both $U(1)_{\gamma}$ and secluded U(1)' groups



Gauge Lagrangian $\mathcal{L}_{gauge} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} \bar{F}'_{\mu\nu} \bar{F}'^{\mu\nu} - \frac{\varepsilon}{2c_W} \bar{F}'_{\mu\nu} B^{\mu\nu}$ $B_{\mu\nu}, \ \overline{F}'_{\mu\nu}$ are field strength tensors of $U(1)_Y, U(1)'$

ε – kinetic mixing parameter

Interaction Lagrangian (after diagonalizaion of gauge and DP fields)

$$\mathcal{L}_{\rm int} = eJ_{\mu}A^{\mu} - \varepsilon eJ_{\mu}A^{\prime\mu} + \varepsilon e^{\prime}t_{W}J_{\mu}^{\prime}Z_{\mu} + e^{\prime}J_{\mu}^{\prime}A^{\prime\mu} + \mathcal{L}_{A^{\prime}\chi}$$

Production of DP in γe⁻ collision with subsequent invisible decay



χ – dark matter particle

We assume that $\mathcal{B}(A' \rightarrow \chi bar(\chi)) = 1$

SM background $\gamma e^- \rightarrow e^- + v + bar(v)$



International Linear Collider (ILC)





Tutometed

	Collision energy	(unpolarized beams)
1st stage	250 GeV	2.0 ab ⁻¹
2nd stage	500 GeV	4.0 ab ⁻¹
3rd stage	1000 GeV	5.4 ab ⁻¹

(ILC Technical Design Report, arXiv:1306.6327, 1903.01629)

Circular Electron Positron Collider (CEPC)





	Collision energy	Integrated luminosity (unpolarized beams)
1st stage	90 GeV	16 ab ⁻¹
2nd stage	180 GeV	2.6 ab ⁻¹
3rd stage	240 GeV	5.6 ab ⁻¹

(CEPC Study Group, arXiv:1809.00285, 1811.10545)



Differential cross sections for $\gamma e^- \rightarrow e^- + E_{miss}$ scattering at the collider CEPC. DP mass is $m_{A'} = 50$ GeV, mixing parameter is $\epsilon = 0.1$. Black curves are SM predictions.



Differential cross sections for $\gamma e^- \rightarrow e^- + E_{miss}$ scattering at the collider ILC. DP mass is $m_{A'} = 100$ GeV, mixing parameter is $\epsilon = 0.1$. Black curves are SM predictions.



Differential cross sections for $\gamma e^- \rightarrow e^- + E_{miss}$ scattering at the collider CLIC. DP mass is $m_{A'} = 200$ GeV, mixing parameter is $\epsilon = 0.1$. Black curves are SM predictions.

(S.Inan & A.K., EPJC, 82, 592, 2022)



Excluded bounds at 95% C.L. on DP mass $m_{A'}$ and mixing parameter ϵ for invisible DP production in $\gamma e^- \rightarrow A' e^$ *unpolarized* scattering at future lepton colliders

 $p_{e,t} > 10 \text{ GeV}, |\eta_e| < 2.5, |m_{A'} - m_{miss}| < 10 \text{ GeV}$



Exclusion limits on mixing parameter ε for massive DP going into invisible states (A. Filippi & M. De Napoli, Rev. Phys. 5, 100042, 2020)

Conclusions

- We have studied anomalous quartic neutral gauge couplings in unpolarized and polarized LBL scattering of CB photons at 3 TeV CLIC
- Exclusion regions for these couplings are obtained. Our best constraints are $\zeta_1 = 6.9*10^{-4}$ TeV⁻⁴ and $\zeta_2 = 1.4*10^{-3}$ TeV⁻⁴, one order of magnitude stronger than the HL-LHC bounds derived for integrated luminosity of 3 ab⁻¹
- Sensitivity to anomalous couplings of the $\gamma\gamma\gamma Z$ vertex is examined using CB photon scattering $\gamma\gamma \rightarrow \gamma Z$ at 3 TeV CLIC. For both couplings, g_1 and g_2 , our best bound is 5*10⁻³ TeV⁻⁴, while the HL-LHC bounds are $(1-0.7)*10^{-1}$ TeV⁻⁴

- For the first time, production of the massive DP in γe scattering at future lepton colliders ILC, CLIC and CEPC is studied
- Invisible decay mode of DP is addressed. Wide range, 1 GeV 1 TeV, of DP mass m_{A'} is considered. We have obtained exclusion regions in the plane (m_{A'}, ε), where ε is the kinetic mixing parameter
- In low mass region, 1-10 GeV, our bounds on ε for the 90 CEPC are several times stronger than BaBar limit (~ 10⁻³). The CEPC bounds for 160 and 240 GeV, as well as bounds for the 250 ILC are close to the BaBar constraints

Thank you for your attention



Possible scenario of future colliders



Anomalous yyyy vertex

The first evidence of $\gamma\gamma \rightarrow \gamma\gamma$ process was observed by ATLAS & CMS in Pb-Pb collisions

(ATLAS Collab., Nat. Phys., 13, 852, 2017) (CMS Collab., Phys. Rev. Lett. 123, 052001, 2019)



Inverse laser-Compton scattering process





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Helicity of CB photon

$$\xi(E_{\gamma},\lambda_0) = \frac{\lambda_0(1-2r)[1-x+1/(1-x)] + \lambda_e r \zeta [1+(1-x)(1-2r)^2]}{1-x+1/(1-x) - 4r(1-r) - \lambda_e \lambda_0 r \zeta (2r-1)(2-x)},$$

where $x = E_{\gamma}/E_e$, $r = x/\zeta(1-x)$, $\zeta = 4E_e E_0/m_e^2$, m_e being the electron mass.

Spectrum of CB photons

$$f_{\gamma/e}(x) = \frac{1}{g(\zeta)} \Big[1 - x + \frac{1}{1 - x} - \frac{4x}{\zeta(1 - x)} + \frac{4x^2}{\zeta^2(1 - x)^2} \\ + \lambda_0 \lambda_e r \zeta(1 - 2r)(2 - x) \Big],$$

where
$$g(\zeta) = g_1(\zeta) + \lambda_0 \lambda_e g_2(\zeta) ,$$

$$g_1(\zeta) = \left(1 - \frac{4}{\zeta} - \frac{8}{\zeta^2} \right) \ln(\zeta + 1) + \frac{1}{2} + \frac{8}{\zeta} - \frac{1}{2(\zeta + 1)^2} ,$$

$$g_2(\zeta) = \left(1 + \frac{2}{\zeta} \right) \ln(\zeta + 1) - \frac{5}{2} + \frac{1}{\zeta + 1} - \frac{1}{2(\zeta + 1)^2} .$$



Differential cross sections for the process $\gamma\gamma \rightarrow \gamma\gamma$ versus invariant mass of outgoing photons at $\forall s = 3$ TeV. The left, middle and right panels correspond to $\lambda_e = 0.8$, -0.8, 0. Solid curves (from top downwards): $(\zeta_1 = 10^{-13} \text{ GeV}^{-4}, \zeta_2 = 0), (\zeta_1 = 0, \zeta_2 = 10^{-13} \text{ GeV}^{-4}), \text{ SM}.$

Discovery significance (\delta = percentage systematic error)

$$S_{\rm dis} = \sqrt{2} \left[(s+b) \ln \left(\frac{(s+b)(1+\delta^2 b)}{b+\delta^2 b(s+b)} \right) - \frac{1}{\delta^2} \ln \left(1 + \frac{\delta^2 s}{1+\delta^2 b} \right) \right]^{1/2}$$

In the limit $\delta = 0$

$$S_{\text{dis}} = \sqrt{2\left[(s+b)\ln\left(1+\frac{s}{b}\right)-s\right]}$$

 $s << b$
 $S_{\text{dis}} = \frac{s}{\sqrt{b}}$

S_{dis} ≥ 5 as discovery region



95% C.L. exclusion region for anomalous couplings ζ₁, ζ₂ for unpolarized LBL scattering at the CLIC. Systematic errors are 0%, 5%, and 10%. The collision energy is 1.5 TeV, the integrated luminosity is 2.5 ab⁻¹. Couplings ζ₁, ζ₂ are in GeV⁻⁴.

(S.Inan & A.K., EPJC, 81, 664, 2021)

Bounds on anomalous couplings for LHC

(S.Fichet et al., JHEP 02, 165, 2015)



Anomalous yyyZ vertex

Another effective Lagrangian for anomalous yyyZ couplings

$$L_{\gamma\gamma\gamma Z} = \widetilde{g}_1(F_{\mu\,\nu}F^{\mu\,\nu})(F_{\rho\sigma}Z^{\rho\sigma}) + \widetilde{g}_2(F_{\mu\,\nu}\widetilde{F}^{\mu\,\nu}F_{\rho\sigma}\widetilde{Z}^{\rho\sigma})$$

(C. Baldenegro et al., JHEP 06, 142, 2017)



Relations between two sets of coupling constants

$$g_1 = 8(\tilde{g}_2 - \tilde{g}_1), \quad g_2 = 8\tilde{g}_2$$

Unpolarized case

$$\sum_{\lambda_1\dots\lambda_4} |M_{\lambda_1\lambda_2\lambda_3\lambda_4}|^2 = \frac{1}{4} [g_1^2(3A+2B) - 4g_1g_2(A+B) + 4g_2^2(A+B)]$$

$$A = s^2 t^2 + t^2 u^2 + u^2 s^2 , \quad B = stu \, m_Z^2$$



Differential cross sections for the process $\gamma\gamma \rightarrow \gamma Z$ versus invariant mass of outgoing photons at $\sqrt{s} = 3$ TeV. The left, middle and right panels correspond to $\lambda_e = 0.8$, -0.8, 0. Solid curves (from top downwards): (g₁ = 10⁻¹³ GeV⁻⁴, g₂ = 0), SM.



95% C.L. exclusion region for anomalous couplings g_1 , g_2 for unpolarized $\gamma\gamma \rightarrow \gamma Z$ scattering at the CLIC. Systematic errors are 0%, 5%, and 10%. The collision energy is 1.5 TeV, the integrated luminosity is 2.5 ab⁻¹. Couplings g_1 , g_2 are in GeV⁻⁴. (S.Inan & A.K., JHEP, 10, 121, 2021)

Partial-wave expansion of helicity amplitude

(M. Jacob & G. Wick, Ann. Phys. 7, 404, 1959; ibid 281, 774 2000)

$$M_{\lambda_1\lambda_2\lambda_3\lambda_4}(s,\theta,\varphi) = 16\pi \sum_J (2J+1)\sqrt{(1+\delta_{\lambda_1\lambda_2})(1+\delta_{\lambda_3\lambda_4})}$$
$$\times e^{i(\lambda-\mu)\phi} d^J_{\lambda\mu}(\theta) T^J_{\lambda_1\lambda_2\lambda_3\lambda_4}(s)$$

 $d_{\lambda\mu}^{J}$ = Wigner's (small) d-function

$$T^{J}_{\lambda_{1}\lambda_{2}\lambda_{3}\lambda_{4}}(s) = \frac{1}{32\pi} \frac{1}{\sqrt{(1+\delta_{\lambda_{1}\lambda_{2}})(1+\delta_{\lambda_{3}\lambda_{4}})}} \int_{-1}^{1} M_{\lambda_{1}\lambda_{2}\lambda_{3}\lambda_{4}}(s,z) d^{J}_{\lambda\mu}(z) dz$$

Partial-wave unitarity bound

$$\left|T^{J}_{\lambda_{1}\lambda_{2}\lambda_{3}\lambda_{4}}(s)\right| \leq 1$$

Invisible dark photons

Diagonalization of gauge fields W³, B and DP field A'

$$\begin{pmatrix} W_{\mu}^{3} \\ B_{\mu} \\ \bar{A}'_{\mu} \end{pmatrix} = \begin{pmatrix} c_{W} & s_{W} - s_{W}\varepsilon \\ -s_{W} & c_{W} - c_{W}\varepsilon \\ t_{W}\varepsilon & 0 & 1 \end{pmatrix} \begin{pmatrix} Z_{\mu} \\ A_{\mu} \\ A'_{\mu} \end{pmatrix}$$

CEPC as a yy factory

(Ch. Zhang & W. Chou, Photon Beam Workshop, Padova, Italy, 2017)



FCC-ee as a yy factory

(R. Alesan et al., IPAC 2015, Richmond, VA, USA, 2015)



γγ Higgs factory can be realized by back-scattering two counter-propagating electron bunches off laser pulses

(S.Inan & A.K., EPJC, 82, 592, 2022)



Excluded bounds at 95% C.L. on DP mass m_{A'} and mixing parameter ε for invisible DP production in collision of *unpolarized* photon and electron whose *polarization is* 80%

 $p_{e,t} > 10 \text{ GeV}, |\eta_e| < 2.5, |m_{A'} - m_{inv}| < 10 \text{ GeV}$



Existing bounds (shaded regions) for DP in visible decays (A. Filippi & M. De Napoli, Rev. Phys. 5, 100042, 2020)



CMS limit on mixing parameter for DP decaying into muon pair

(CMS Collab., Phys. Rev. Lett. 124, 131802, 2020)

DP at future ep colliders LHeC and FCC-he

(CMS Collab., Phys. Rev. D 101, 015020, 2020)



