

The processes $\tau \rightarrow (a_1\pi, K_1\pi, K_1K)\nu_\tau$ and
 $e^+e^- \rightarrow (a_1\pi, K_1K)$ in the NJL model

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BLTP, JINR

Nambu–Jona-Lasinio model

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- ▶ M. K. Volkov, C. Weiss, Phys. Rev. D 56, 221 (1997)

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- ▶ $\tau \rightarrow (a_1\pi, K_1\pi, K_1K)\nu_\tau$
- ▶ $e^+e^- \rightarrow (a_1\pi, K_1K)$

Lagrangian of the NJL model

$$\begin{aligned} \Delta L_{int} = & \bar{q} \left\{ \sum_{i=0,\pm} \left[iA_\pi \gamma^5 \lambda_i^\pi \pi^i + iA_K \gamma^5 \lambda_i^K K^i + \frac{1}{2} \gamma^\mu \lambda_i^\rho (A_\rho \rho_\mu^i + B_\rho \rho_\mu^{\prime i}) \right. \right. \\ & + \frac{1}{2} \gamma^\mu \lambda_i^K (A_{K^*} K_\mu^{*i} + B_{K^*} K_\mu^{\prime *i}) + \frac{A_{K_1}}{2} \gamma^\mu \gamma^5 \lambda_i^K K_{1\mu}^i + \frac{A_{a_1}}{2} \gamma^\mu \gamma^5 \lambda_i^\rho a_{1\mu}^i \left. \right] \\ & \left. + \frac{1}{2} \gamma^\mu \lambda^\omega (A_\omega \omega_\mu + B_\omega \omega_\mu') + \frac{1}{2} \gamma^\mu \lambda^\phi (A_\phi \phi_\mu + B_\phi \phi_\mu') \right\} q \end{aligned}$$

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$$A_M = \frac{1}{\sin(2\theta_M^0)} \left[g_M \sin(\theta_M + \theta_M^0) + g_M' f_M(k_\perp^2) \sin(\theta_M - \theta_M^0) \right],$$

$$B_M = \frac{-1}{\sin(2\theta_M^0)} \left[g_M \cos(\theta_M + \theta_M^0) + g_M' f_M(k_\perp^2) \cos(\theta_M - \theta_M^0) \right]$$

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$$f_M(k_\perp^2) = 1 + d_M k_\perp^2$$

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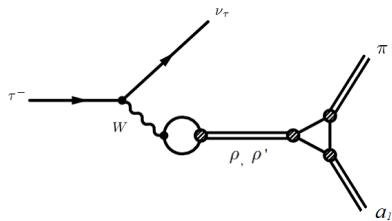
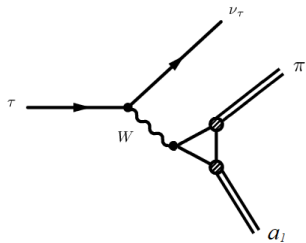
$$A_M = \frac{1}{\sin(2\theta_M^0)} \left[g_M \sin(\theta_M + \theta_M^0) + g_M' f_M(k_\perp^2) \sin(\theta_M - \theta_M^0) \right],$$

$$B_M = \frac{-1}{\sin(2\theta_M^0)} \left[g_M \cos(\theta_M + \theta_M^0) + g_M' f_M(k_\perp^2) \cos(\theta_M - \theta_M^0) \right]$$

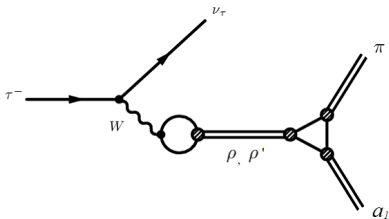
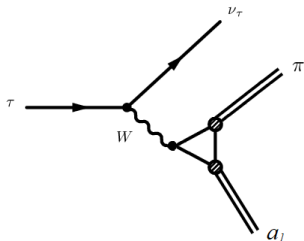
$$f_M(k_\perp^2) = 1 + d_M k_\perp^2$$

$$I_{n_1 n_2}^{f^m} = -i \frac{N_c}{(2\pi)^4} \int \frac{f^m(k_\perp^2)}{(m_u^2 - k^2)^{n_1} (m_s^2 - k^2)^{n_2}} \Theta(\Lambda^2 - k_\perp^2) d^4 k$$

Decay $\tau \rightarrow a_1 \pi \nu_\tau$

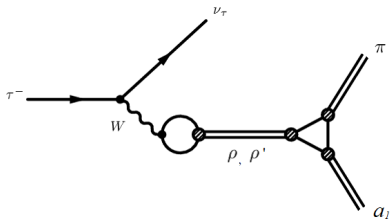
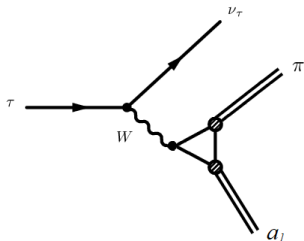


Decay $\tau \rightarrow a_1 \pi \nu_\tau$



$$\mathcal{M}_{\tau \rightarrow a_1 \pi \nu_\tau} = -iG_F V_{ud} 4m_u g_\pi L_\mu^{\text{weak}} \left[I_{20}^{a_1} g_{\mu\nu} + I_{20}^{a_1 \rho} \frac{C_\rho}{g_\rho} \frac{g_{\mu\nu} s - p_\mu p_\nu}{M_\rho^2 - s - i\sqrt{s}\Gamma_\rho} \right. \\ \left. + I_{20}^{a_1 \rho'} \frac{C_{\rho'}}{g_{\rho'}} \frac{g_{\mu\nu} s - p_\mu p_\nu}{M_{\rho'}^2 - s - i\sqrt{s}\Gamma_{\rho'}} \right] \epsilon_\nu^*(p_{a_1})$$

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$$C_\rho = \frac{1}{\sin(2\theta_\rho^0)} \left[\sin(\theta_\rho + \theta_\rho^0) + R_\rho \sin(\theta_\rho - \theta_\rho^0) \right],$$

$$C_{\rho'} = \frac{-1}{\sin(2\theta_\rho^0)} \left[\cos(\theta_\rho + \theta_\rho^0) + R_\rho \cos(\theta_\rho - \theta_\rho^0) \right]$$

Decay $\tau \rightarrow K_1 \pi \nu_\tau$

$$K_1(1270) = K_{1A} \sin \alpha + K_{1B} \cos \alpha,$$

$$K_1(1400) = K_{1A} \cos \alpha - K_{1B} \sin \alpha$$

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$$L_B = \frac{g_B}{2} \sum_{j=0,\pm} K_{1B}^{\mu j} \left(\bar{q} \lambda_j^K \gamma^5 \overleftrightarrow{\partial}_\mu q \right)$$

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$$\mathcal{M}_{\tau^- \rightarrow K_1(1270)^- \pi^0 \nu_\tau} = -G_F V_{us} g_\pi L_\mu^{\text{weak}} [iT_{K_{1A}} + T_{K_{1B}}]^{\mu\nu} \epsilon_\nu^*(p_{K_1})$$

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$$T_{K_{1A}}^{\mu\nu} = 2m_s \sin \alpha \left[I_{11}^{K_1} g^{\mu\nu} + I_{11}^{K_1 K^*} \frac{C_{K^*}}{g_{K^*}} \frac{g^{\mu\nu} s - p^\mu p^\nu}{M_{K^*}^2 - s - i\sqrt{s}\Gamma_{K^*}} \right. \\ \left. + I_{11}^{K_1 K^{*'}} \frac{C_{K^{*'}}}{g_{K^{*'}}} \frac{g^{\mu\nu} s - p^\mu p^\nu}{M_{K^{*'}}^2 - s - i\sqrt{s}\Gamma_{K^{*'}}} \right]$$

$$T_{K_{1B}}^{\mu\nu} = g_B \cos \alpha \left[\left(I_{10}^{K_1} - m_s^2 I_{11}^{K_1} \right) g^{\mu\nu} + \left(I_{10}^{K_1 K^*} - m_s^2 I_{11}^{K_1 K^*} \right) \frac{C_{K^*}}{g_{K^*}} \frac{g^{\mu\nu} s - p^\mu p^\nu}{M_{K^*}^2 - s - i\sqrt{s}\Gamma_{K^*}} \right. \\ \left. + \left(I_{10}^{K_1 K^{*'}} - m_s^2 I_{11}^{K_1 K^{*'}} \right) \frac{C_{K^{*'}}}{g_{K^{*'}}} \frac{g^{\mu\nu} s - p^\mu p^\nu}{M_{K^{*'}}^2 - s - i\sqrt{s}\Gamma_{K^{*'}}} \right]$$

Numerical estimations

Decay mode	[1]	[2]	NJL[3]	
$\tau \rightarrow a_1(1260)^- \pi^0 \nu_\tau$	6.9 ± 6.3		0.14 ± 0.02	$\times 10^{-3}$
	6.1 ± 5.9			$\times 10^{-3}$
$\tau \rightarrow a_1(1260)^0 \pi^- \nu_\tau$	6.8 ± 6.1	1.3	0.13 ± 0.02	$\times 10^{-3}$
	5.9 ± 5.7			$\times 10^{-3}$
$\tau \rightarrow K_1(1270)^- \pi^0 \nu_\tau$	0.8 ± 0.2		3.59 ± 0.61	$\times 10^{-6}$
$\tau \rightarrow K_1(1270)^0 \pi^- \nu_\tau$	1.4 ± 0.5	21	6.84 ± 1.16	$\times 10^{-6}$
$\tau \rightarrow K_1(1400)^- \pi^0 \nu_\tau$	1.1 ± 0.1		0.27 ± 0.04	$\times 10^{-6}$
$\tau \rightarrow K_1(1400)^0 \pi^- \nu_\tau$	2.1 ± 0.2	4.1	0.49 ± 0.08	$\times 10^{-6}$
$\tau \rightarrow K_1(1270)^- K^0 \nu_\tau$	2.8 ± 1.9		4.25 ± 0.72	$\times 10^{-9}$
$\tau \rightarrow K_1(1270)^0 K^- \nu_\tau$	13 ± 8.9		6.79 ± 1.15	$\times 10^{-9}$

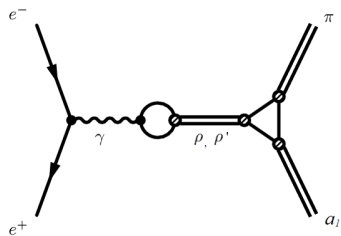
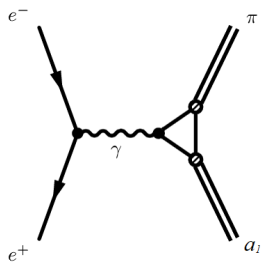
[1] G. Calderon, J. H. Munoz and C. E. Vera, *Phys. Rev. D* **87** (2013) no.11, 114011

[2] L. R. Dai, L. Roca and E. Oset, *Phys. Rev. D* **99** (2019) no.9, 096003

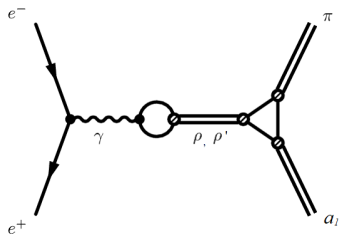
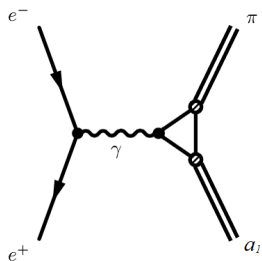
[3] M. K. Volkov, A. A. Pivovarov and K. Nurlan, *arXiv:2205.02810 [hep-ph]* —

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Process $e^+e^- \rightarrow a_1\pi$

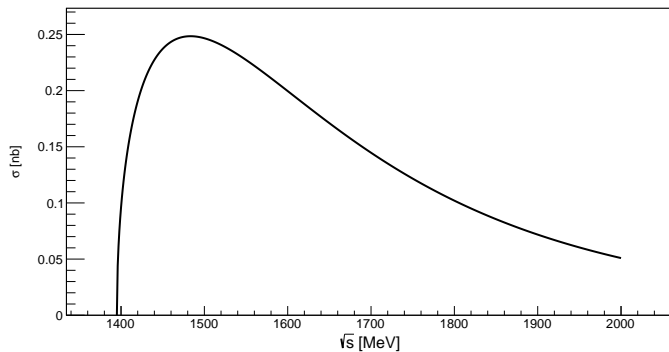


Process $e^+e^- \rightarrow a_1\pi$



$$\begin{aligned} \mathcal{M}_{e^+e^- \rightarrow a_1\pi} = & \frac{16\pi\alpha_{em}}{s} L_\mu^{em} \left[g^{\mu\nu} I_{20}^{a_1} + \frac{C_\rho}{g_\rho} \frac{g_{\mu\nu}s - p_\mu p_\nu}{M_\rho^2 - s - i\sqrt{s}\Gamma_\rho} I_{20}^{\rho a_1} \right. \\ & \left. + \frac{C_{\rho'}}{g_{\rho'}} \frac{g_{\mu\nu}s - p_\mu p_\nu}{M_{\rho'}^2 - s - i\sqrt{s}\Gamma_{\rho'}} I_{20}^{\rho' a_1} \right] e_\nu^*(a_1) \end{aligned}$$

Cross section of the process $e^+e^- \rightarrow a_1\pi$



Process $e^+e^- \rightarrow K_1K$

$$\mathcal{M}_{e^+e^- \rightarrow K_1(1270)K} = \frac{8\pi\alpha_{em}}{s} L_\mu^{em} \left[T_\gamma + T_{\rho+\rho'} + T_{\omega+\omega'} + e^{i\pi} T_{\phi+\phi'} \right]^{\mu\nu} e_\nu^*(p_{K_1})$$

Process $e^+e^- \rightarrow K_1K$

$$\mathcal{M}_{e^+e^- \rightarrow K_1(1270)K} = \frac{8\pi\alpha_{em}}{s} L_\mu^{em} \left[T_\gamma + T_{\rho+\rho'} + T_{\omega+\omega'} + e^{i\pi} T_{\phi+\phi'} \right]^{\mu\nu} e_\nu^*(p_{K_1})$$

$$T_{(\gamma)\mu\nu} = i(m_s + m_u) I_{11}^{K_1K} \sin \alpha g_{\mu\nu} + g_B g_K \cos \alpha \left[\frac{2}{3} F(u, s) + \frac{1}{3} F(s, u) \right] g_{\mu\nu}$$

$$T_{(V+V')\mu\nu} = r_V \left[\frac{C_V}{g_V} \frac{g_{\mu\nu} s - p_\mu p_\nu}{M_V^2 - s - i\sqrt{s}\Gamma_V} (i \sin \alpha G_{VK_{1A}K} + \cos \alpha G_{VK_{1B}K}) \right. \\ \left. + \frac{C_{V'}}{g_{V'}} \frac{g_{\mu\nu} s - p_\mu p_\nu}{M_{V'}^2 - s - i\sqrt{s}\Gamma_{V'}} (i \sin \alpha G_{V'K_{1A}K} + \cos \alpha G_{V'K_{1B}K}) \right]$$

Process $e^+e^- \rightarrow K_1K$

$$\mathcal{M}_{e^+e^- \rightarrow K_1(1270)K} = \frac{8\pi\alpha_{em}}{s} L_\mu^{em} \left[T_\gamma + T_{\rho+\rho'} + T_{\omega+\omega'} + e^{i\pi} T_{\phi+\phi'} \right]^{\mu\nu} e_\nu^*(p_{K_1})$$

$$T_{(\gamma)\mu\nu} = i(m_s + m_u) I_{11}^{K_1K} \sin\alpha g_{\mu\nu} + g_B g_K \cos\alpha \left[\frac{2}{3} F(u, s) + \frac{1}{3} F(s, u) \right] g_{\mu\nu}$$

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$$F(u, s) = I_{10} - \left((m_s - m_u)^2 + m_u^2 \right) I_{11} - 2m_u^2(m_s - m_u) I_{21}$$

$$G_{VK_{1A}K} = (m_s + m_u) I_{11}^{VK_1K}$$

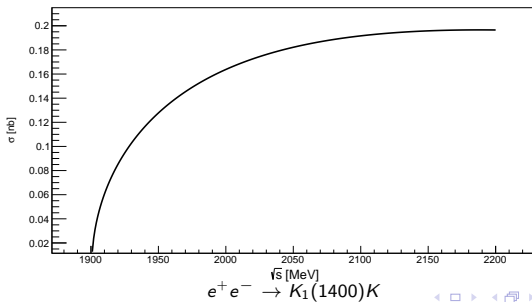
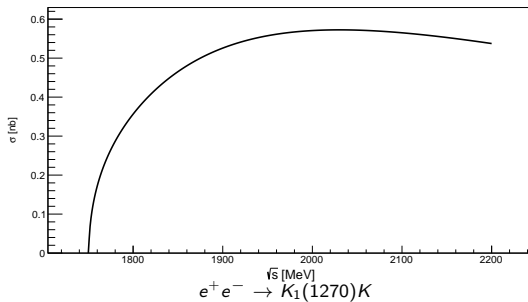
$$G_{V'K_{1A}K} = (m_s + m_u) I_{11}^{V'K_1K}$$

$$G_{\rho K_{1B}K} = g_B \left[I_{10}^{\rho K} - \left((m_s - m_u)^2 + m_u^2 \right) I_{11}^{\rho K} - 2m_u^2(m_s - m_u) I_{21}^{\rho K} \right]$$

$$G_{\rho' K_{1B}K} = g_B \left[I_{10}^{\rho' K} - \left((m_s - m_u)^2 + m_u^2 \right) I_{11}^{\rho' K} - 2m_u^2(m_s - m_u) I_{21}^{\rho' K} \right]$$

$$r_\rho = 1/2 \quad r_\omega = 1/6 \quad r_\phi = 1/3$$

Cross section of $e^+e^- \rightarrow (K_1(1270), K_1(1400))K$



Summary

- ▶ The processes $\tau \rightarrow (a_1\pi, K_1\pi, K_1K)\nu_\tau$ were calculated in the framework of the NJL model
- ▶ The comparison of the obtained results with the results of other theoretical works was carried out
- ▶ The processes $e^+e^- \rightarrow (a_1\pi, K_1K)$ were calculated in the framework of the NJL model
- ▶ The mixing of the states $K_1(1270)$ and $K_1(1400)$ was taken into account