Особенности электронного строения и химические свойства сверхтяжелых элементов 7-го и 8-го периодов

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Introduction

- Ground state electron configurations of superheavy elements (SHE)
- Relativistic effects in the electronic structure of SHE
- Electron density distribution in SHE. Electron localization function (ELF)
- Orbital collapse of the 5g-states
- Summary and Outlook

Introduction. Pyykkö's periodic table



The table is taken from P. Pyykkö Chem. Rev. 112, 371 (2012).

Introduction

In this work, the results of the electronic-structure calculations for a number of super-heavy elements (SHE) of the 7th and 8th periods with atomic numbers Z=111-170 and their lighter homologous are performed [1-3].

- Ground state electron configurations of superheavy elements (SHE)
- Ionization potentials
- Electron affinities
- One-particle electron density
- Root-mean-square radii of the valence states
- Widths of the electron-density distribution of valence shells
- Electron localization functions (ELF)
- Shannon entropy
- Quantum electrodynamics (QED) corrections
- 1. M. Y. Kaygorodov et al., Phys. Rev. A 104, 012819 (2021)
- 2. I. I. Tupitsyn et al., Opt. Spectr. 129, 1038 (2021)
- 3. M. Y. Kaygorodov et al., Phys. Rev. A 105, 062805 (2022)

Introduction. Specific features of electronic structure and chemical properties of super-heavy elements of the 7th and 8th periods

The electronic structure of SHEs is unique in several aspects:

- Strong relativistic effects cause contraction of the s- and p- orbitals
- Spin-orbital splitting of valence p-shells reaches up about 10 eV in Og (Z=118) and about 420 eV for the 7p-orbital and about 75 eV for the 8p-orbital in element with atomic number Z=165.
- As a result, due to the strong relativistic contraction, the radial distribution of the electron density of the valence $7p_{1/2}$ -shell of the Og atom starts to overlap with the outer core shells and ELF is close to 0.5 in the valence region. In Ref. [1], this effect in Og was interpreted as smearing out the valence electron density distribution and its approaching to the case of the homogeneous electron gas.
- Starting from the Z = 125 element, the 5g-shell with the large angular momentum (I = 4) is occupied with electrons.
- The effective radial potential for the 5g-electron, which includes a large centrifugal repulsive term, has two potential wells which leads to the so-called *orbital collapse*.

[1] P. Jerabek, B. Schuetrumpf, P. Schwerdtfeger, and W. Nazarewicz, Phys. Rev. Lett. 120, 053001 (2018).

In our work, we used two independent theoretical calculation methods.

• Configuration Interaction Dirac-Fock-Sturm method (CI-DFS)

At the first step, to obtain the one-electron wave functions for the occupied atomic shells, we use the Dirac-Fock method. Then the DFS orbitals are obtained by solving the DFS equations for the vacant shells. At the last step, the relativistic CI+MBPT method is used to obtain the many-electron wave functions and the total energies.

• Fock Space Coupled-Cluster method (FS-CC)

DIRAC, a relativistic ab initio electronic structure program, Release DIRAC21 (2021), http://www.diracprogram.org

FS-CC method, in contrast to the one-configuration coupled-cluster method is capable of providing not only the ground-state energy of an N-electron system, but also an important fraction of system's excitation spectrum, including ionization potentials, electron affinities, etc.

• To evaluate the QED correction we use the model QED operator approach[1]

[1] V.M. Shabaev, I.I. Tupitsyn, V.A. Yerokhin, Comput. Phys. Commun. 189, 175 (2015)

For Z=126 the configurations with the *lowest Dirac-Fcok (DF) energy* within the configuration average approximation are

Configuration	$E_{ m DF}^{ m av}$ [a.u.]
$8p^17d^16f^25g^2$	-66298.183666
$8p^{1}6f^{3}5g^{2}$.183121
$8p^{1}6f^{4}5g^{1}$.168137
$8p^26f^25g^2$.146963
$8p^17d^26f^15g^2$.114149

- $8p^26f^25g^2$ Mann et al., 1970 (DF)
- $8p^{1}7d^{1}6f^{2}5g^{2}$ Fricke et al., 1977 (DFS)
- $8p^16f^45g^1$ Umemoto and Saito, 1997 (DF+PZ SIC)
- $8p^16f^35g^2$ (0.98) Nefedov et al., 2006 (MCDF)
- $8p^17d^16f^25g^2$ Zhou et al., 2017 (DF)

Таблица:	Ground	state	electron
configurat	ions		

	Core: [Rn] $5f^{14}$	
Z	Conf.	J
111	$6d^9 7s^2$	2.5
112	$\mathrm{6d}^{10}$ $\mathrm{7s}^2$	0.0
113	$6d^{10}$ $7s^27p^1$	0.5
114	$6d^{10}$ $7s^27p^2$	0.0
118	$6d^{10}$ $7s^27p^6$	0.0
119	$6d^{10}\ 7s^27p^68s^1$	0.5
120	$\mathrm{6d}^{10}~\mathrm{7s}^{2}\mathrm{7p}^{6}\mathrm{8s}^{2}$	0.0

	Core: [Og] 8s ²					
Z	Conf.	J	Ref[1]	Ref[2]		
125	$5g^16f^27d^18p^1$	8.5	$5g^16f^38p^1$	$5g^16f^28p^2$		
126	$5g^26f^27d^18p^1$	10	$5g^26f^27d^18p^1$	$5g^2 6f^3 8p^1$		
127	$5g^36f^27d^18p^1$	13.5	$5g^36f^28p^2$	$5g^36f^28p^2$		
	Core: [Og] 8s ² 8p ² _{1/2}					
144	$5g^{18}6f^17d^3$	4.0	$5g^{18}6f^17d^3$	$5g^{17}6f^17d^3$		
145	$5g^{18}6f^37d^2$	6.5	$5g^{18}6f^37d^2$	$5g^{18}6f^37d^2$		
162	$5g^{18}6f^{14}7d^8$	4.0	$5g^{18}6f^{14}7d^8$	$5g^{18}6f^{14}7d^79s^1$		
163	$5g^{18}6f^{14}7d^9$	2.5	$5g^{18}6f^{14}7d^9$	$5g^{18}6f^{14}7d^89s^1$		
164	$5g^{18}6f^{14}7d^{10}$	0.0	$5g^{18}6f^{14}7d^{10}$	$5g^{18}6f^{14}7d^99s^1$		
165	$5g^{18}6f^{14}7d^{10}9s^{1}$	0.5	$5g^{18}6f^{14}5g^{18}7d^{10}9s^{1}$	—		
166	$5g^{18}6f^{14}7d^{10}9s^2$	0.0	$5g^{18}6f^{14}5g^{18}7d^{10}9s^2$	—		

Таблица: Ground state electron configurations

[1] B. Fricke and G. Soff, Atomic Data and Nuclear Data Tables 19, 83 (1977).

[2] V.I. Nefedov, M. Trzhaskovskaya, Dokl. Phys. Chem. 408, 149 (2006).

Z	7p _{1/2}	7p _{3/2}	$\Delta_{\rm SO}$
114	10.4	4.5	5.9
116	14.3	6.1	8.2
118	20.1	8.3	11.8
	$8p_{1/2}$	$8p_{3/2}$	
125	5.3	2.4	2.9
144	13.8	2.7	11.1
145	16.2	2.7	13.5
164	69.9	3.6	66.3
165	79.7	5.0	74.7
166	90.1	6.6	83.5

Таблица: Spin-orbit splitting $\Delta_{\rm SO}$ [eV]





Electron density distribution



Radial electron density $4\pi\rho(r)r^2$ plot for Og in both relativistic and non-relativistic approximations.

There is no visible peak in the valence shell region. Electron density alone does not show any valence shell structure

The electron localization function is defined by [4]

$$\mathrm{ELF}(\boldsymbol{r}) = \left(1 + \left[\frac{D(\boldsymbol{r})}{D_0(\boldsymbol{r})}\right]^2\right)^{-1},\qquad(1)$$

where

$$D(\boldsymbol{r}) = \frac{1}{2} \left[\tau - \frac{1}{4} \frac{|\boldsymbol{\nabla}\rho(\boldsymbol{r})|^2}{\rho(\boldsymbol{r})} \right], \quad \text{and} \quad \tau = \sum_{i,\sigma} |\boldsymbol{\nabla}\varphi_{i\sigma}(\boldsymbol{r})|^2.$$
(2)

Here ρ is total density and τ the kinetic energy density.

 $D_0(r)$ corresponds to a uniform electron gas (Thomas-Fermi) kinetic energy density

$$D_0(\mathbf{r}) = \frac{3}{10} (3\pi^2)^{2/3} \rho^{5/3}(\mathbf{r}) \,. \tag{3}$$

The ELF values lie by definition between zero and one. Small values are typical for the region between two electron shells. In a homogeneous electron gas $\mathsf{ELF} = 0.5$.

[4] A. D. Becke and K. E. Edgecombe, J. Chem. Phys. v.92, 5397 (1990).

$$D(\boldsymbol{r}) = \sum_{\lambda=1,2} \left[W^{\lambda}(r) T^{\lambda}(\boldsymbol{r}) - \frac{1}{8} \frac{|\boldsymbol{\nabla} \rho^{\lambda}(r)|^2}{\rho(r)} \right],$$
(4)

where $\rho(r)$ – total electron density

$$\rho(r) = \sum_{\lambda=1,2} \rho^{\lambda}(r), \qquad \rho^{\lambda}(r) = \frac{4\pi}{r^2} \sum_{a} q_a \begin{cases} P_a^2(r), & \lambda = 1, \\ Q_a^2(r), & \lambda = 2. \end{cases}$$
(5)

 $T^{\lambda}(r)$ in formula (4) is the relativistic analogue of the non-relativistic kinetic energy density

$$T^{\lambda}(\boldsymbol{r}) = \sum_{a} q_{a} t_{a}^{\lambda}(\boldsymbol{r}), \qquad t_{a}^{\lambda}(\boldsymbol{r}) = \frac{1}{2} \frac{1}{2j_{a}+1} \sum_{\mu_{a},\sigma} |\boldsymbol{\nabla}\phi_{a\mu_{a}}^{\lambda}(\boldsymbol{r},\sigma)|^{2}$$
(6)

and $W^{\lambda}(r)$ is a weight function that has the form

$$W^{\lambda}(r) = \frac{\rho^{\lambda}(r)}{\rho(r)} \,. \tag{7}$$

I.I Tupitsyn et al., Optics and Spectroscopy, 130, 1022 (2022)



"Spin-orbit splitting in the 7p electronic shell becomes so large ($\sim 10 \text{ eV}$) that Og is expected to show uniform-gas-like behavior in the valence region" [5].

[5] Jerabek et al., PRL 120, 053001 (2018).





Z164 Electron configuration: $[Og]5g^{18}6f^{14}7d^{10}$.

One-electron energies ε and mean radii of the core electrons

	ε [keV]	<r> [a.u.]</r>
1s 1/2	-770.6	0.0028
2p 1/2	-343.0	0.0038
2s 1/2	-196.3	0.0130
2p 3/2	-66.2	0.0279

 mc^2 =510.7 keV

Orbital collapse. Two solutions



$$E_{\text{inner}} = -64846.2788 \text{ a.u.}$$

$$E_{\text{outer}} = -64846.0878 \text{ a.u.}$$
(9)

Orbital collapse. Two solutions. La (Z=57)



1. I.M. Band and V.I. Fomichev, Phys.Letters A, 75, 178 (1980)

2. J.-R Connerade and R.C. Kamatak, Handbook on the Physics and Chemistry of Rare Earths, v. 28, p. 1 (2000)

Thank You for Attention.

Таблица: Comparison of contributions to the Electron Affinities (EA) of Og (Z=118) atom ([Rn]5f¹⁴6d¹⁰7s²7p⁶) calculated with the CI-DFS and FSCC methods in the present work [1] with results of Refs. [2-5] (eV). In Ref. [4] a combination of CI with the many-body perturbation theory is used.

4)
10)
3)
6)
- 4 1 3 6

- 1. M. Y. Kaygorodov et al., PRA, 104, 012819, (2021).
- 2. E. Eliav, U. Kaldor, Y. Ishikawa, and P. Pyykkö, PRL 77, 5350 (1996).
- 3. I. Goidenko, L. Labzowsky, E. Eliav, U. Kaldor, and P. Pyykkö, PRA 67, 020102 (2003).
- 4. B. G. C. Lackenby, V. A. Dzuba, and V. V. Flambaum, PRA 98, 042512 (2018).
- 5. Y Guo, et al., arXiv:2107.02164 [physics], (2021)

According to the formula of K. Shannon [1], the amount of information is defined as:

$$S = -\sum_{i=1}^{N} p_i \ln p_i \tag{12}$$

where N is the number of random events, p_i is the probability of the i-th event and

$$\sum_{i=1}^{N} p_i = 1.$$
 (13)

The minimum value of S is reached for the deterministic event when one of the probabilities of p_i is 1, and the rest are zero. In this case, $S = S_{\min} = 0$.

The maximum value of S is reached for an equally probable distribution $p_i=1/N.$ Then, $S_{\rm max}=\ln(N)$

$$0 \le S \le \ln(N) \tag{14}$$

For continuous distribution

$$S = -\int d\boldsymbol{r} \ln(\rho(\boldsymbol{r})) \,\rho(\boldsymbol{r}) \,, \qquad \int d\boldsymbol{r} \,\rho(\boldsymbol{r}) = 1 \,. \tag{15}$$

[1] C.E. Shannon, Bell Syst. Tech. J. 27, 379; 623 (1948).



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Group 1. Ionization Potentials and RMS, ns¹

n	Z	$\epsilon(d_{3/2})$	$\epsilon(d_{5/2})$	$\epsilon(ns)$	$\operatorname{RMS}(d_{5/2})$	$\operatorname{RMS}(d_{5/2})$	$\mathrm{RMS}(ns)$
4	29	-13.276	-12.906	-6.661	0.991	1.002	3.261
5	47	-14.312	-13.646	-6.453	1.371	1.396	3.451
6	79	-13.430	-11.667	-7.937	1.536	1.619	3.061
7	111	-14.082	-11.047	-11.593	1.708	1.868	2.588
8	163	-14.396	-6.310	-60.791	1.745	2.282	1.401
4	30	-20.977	-20.537	-8.126	0.875	0.883	2.848
5	48	-20.089	-19.290	-7.658	1.254	1.274	3.089
6	80	-17.689	-15.637	-8.926	1.431	1.499	2.843
7	112	-15.313	-12.030	-12.275	1.643	1.795	2.499
8	164	-15.875	-7.138	-65.815	1.680	2.161	1.359

Таблица: One-electron energies and RMS valence shell radii of atoms with a serial number Z=163 and Z=164 and its homologous

For atom Z=164 RMS(5g) $\simeq 0.32$ a.u.

Dirac-Fock-Sturm method (DFS)

Within the DCB approximation, the many-electron wave function $\Psi(JM)$ can be obtained in the form of an expansion in terms of the configuration-state functions (CSF) $\Phi_I(JM)$:

$$\Psi(JM) = \sum_{I} C_{I}^{JM} \Phi_{I}(JM) \,. \tag{16}$$

Each function $\Phi_I(JM)$ is a an eigenfunction of the operators \hat{J}^2 ın \hat{J}_z and there is a linear combination of Slater determinants

$$\sum_{K} H_{KI} C_{K}^{JM} = E_{I}(J) C_{I}^{JM} , \qquad (17)$$

where $H_{KI} = \langle \Phi_K \mid \hat{H}_{\text{DCB}} \mid \Phi_I \rangle$.

The Slater determinants are constructed form the one-electron wave functions ψ_i obtained by the Dirac–Fock method in the basis of Dirac–Fock–Sturm (DFS) orbitals φ_k :

$$\psi_i = \sum_k u_{ki} \,\varphi_k \,. \tag{18}$$

The one-electron DFS functions, obtained by numerically solving the Dirac–Fock (DF) integro-differential equations for the occupied in the ground and low-lying excited states.

For virtual (high-lying vacant) one-electron states, the functions φ_k are obtained by numerically solving the Dirac–Fock–Sturm equations

$$\left[\hat{h}_{\rm DF} - \varepsilon_0\right] \varphi_k = \mu_k W(r) \varphi_k \,, \tag{19}$$

where $\hat{h}_{\rm DF}$ — is the Dirac–Fock operator, ε_0 — is the reference one-electron energy, and W(r) is a positive weight function tending to zero at infinity.

We note that all the DFS orbitals have approximately the same characteristic size and the same asymptotic at infinity, determined by the reference energy ε_0 :

$$\varphi_k(r) \xrightarrow[r \to \infty]{} C_k \exp(-\sqrt{2\varepsilon_0} r).$$
 (20)

The function tending to a constant for $r \to 0$ was chosen as the weight function W (r):

$$W(r) = \frac{1 - \exp(-(\alpha r)^2)}{(\alpha r)^2} \,. \tag{21}$$

Sturmian orbitals. Hydrogen



The one-electron QED potential V^{QED} can be consider as the sum of *Vacuum Polarization* (VP) potential V^{VP} and *Self-Energy* (SE) potential V^{SE}

$$V^{\text{QED}} = V^{\text{VP}} + V^{\text{SE}} \tag{22}$$

With a good accuracy VP contribution can be presented as the sum of the local Uehling and Wichmann-Kroll potentials

$$V^{\rm VP} = V_{\rm Uehl} + V_{\rm WK} \tag{23}$$

The calculation of the self-energy is complicate and time-consuming part of the QED corrections.

We suppose that there exist one-electron self-energy (SE) operator $\hat{\Sigma}$ which can be directly included into the Dirac-Coulomb-Breit (DCB) many-electron calculations. The approximation to the self-energy operator $\hat{\Sigma}$ based on the expansion (1) is given by

$$V^{\rm SE} = \sum_{i,k=1}^{n} |\psi_i^{(0)}\rangle \Sigma_{ik} \langle \psi_k^{(0)}|, \quad \text{where} \quad \Sigma_{ij} = \langle \psi_i^{(0)}|\hat{\Sigma}|\psi_j^{(0)}\rangle.$$
(24)

The SE operator have to be localized in a small region whose size is on the order of Compton wavelength (1/137 a.u.)

The hydrogen-like wave functions should be replaced by the properly localized so-called projected functions $\phi_i.$ Then

$$V^{\rm SE} = \sum_{i,k=1}^{n} |\phi_i\rangle B_{ik} \langle \phi_k | , \qquad (25)$$

The matrix elements B_{ik} are chosen so that the matrix elements of the model operator $V_{ik}^{\rm SE}$ calculated with hydrogen like wave functions $\psi_i^{(0)}$ have to be equal to matrix elements Q_{ik} of the symmetrized exact one-loop energy-dependent SE operator $\Sigma(\varepsilon)$

$$\langle i \mid V^{\rm SE} \mid k \rangle = Q_{ik} = \frac{1}{2} \left[\Sigma_{ik} + \Sigma_{ki} \right]$$
(26)

Then

$$B_{ik} = \sum_{j,l=1}^{n} (D^{-1})_{ji} Q_{jl} (D^{-1})_{lk}, \qquad D_{ik} = \langle \phi_i | \psi_k^{(0)} \rangle$$
(27)

This approximation to the SE operator is also not very successful. If any strongly localized function is orthogonal to the finite set of projection functions ϕ_i then the corresponding SE contribution will be equal to zero.

To overcome this problem we introduce the local potential $V_{\rm loc}^{\rm SE}$ in (4) by the following way

$$V^{\rm SE} = V_{\rm loc}^{\rm SE} + \sum_{i,k=1}^{n} |\phi_i\rangle \Delta B_{ik} \langle \phi_k |, \qquad (28)$$

where

$$\Delta B_{ik} = \sum_{j,l=1}^{n} (D^{-1})_{ji} \, \Delta \Sigma_{jl} \, (D^{-1})_{lk} \,, \tag{29}$$

and $\Delta \Sigma_{ik} = \Sigma_{ik} - \langle \psi_i^{(0)} | V_{\text{loc}}^{\text{SE}} | \psi_k^{(0)} \rangle$.

At the present time, the matrix Σ is constructed in the interval $3 \le Z \le 125$. However, A. Malyshev is finishing work on expanding the region of Z up to Z=170.

This approximation to the SE operator was used in our papers [1-3]:

V.M. Shabaev, I.I. Tupitsyn, V.A. Yerokhin, Phys. Rev. A, 88, 012513 (2013)
 V.M. Shabaev, I.I. Tupitsyn, V.A. Yerokhin, Computer Phys. Comm., 189, 175 (2015)
 I.I. Tupitsyn, M.G. Kozlov, M.S. Safronova, V.M. Shabaev, and V.A. Dzuba, PRL, 117, 253001 (2016)

Contribution	I_1	EA
FSCC-SD	10.94(6)	1.92(6)
FSCC-T	-0.585	-0.076
Gaunt (FSCC)	0.026(1)	0.024(2)
Retard (CI)	-0.002(1)	-0.002(1)
Freq (CI)	-0.0006(3)	0.0001(3)
QED (CI)	0.021(3)	0.019(3)
Total	10.40(6)	1.89(6)
Eliav <i>et al.</i>	10.60	1.565
Lackenby <i>et al.</i>	11.175	
Hangele <i>et al.</i>	11.871	2.5077

Таблица: Contributions to the first ionization potential I_1 and electron affinity EA of Rg (Z=111) atom (eV). Configuration: $6d^97s^2$

- 1. E. Eliav et al., PRL, 73, 3203, (1994).
- 2. B. G. C. Lackenby et al., PRA, 101, 012514 (2020).
- 3. T. Hangele et al, J. Chem. Phys., 136, 214105 (2012).

Contribution	$I_{2}^{J=0}$	$I_{2}^{J=2}$	$I_{2}^{J=3}$	$I_2^{J=4}$	I_1
FSCC-SD	23.32(4)	22.53(4)	22.69(4)	22.46(4)	11.93(4)
FSCC-T	-0.143(2)	-0.238(2)	-0.159(2)	-0.302(2)	-0.006(2)
Gaunt (FSCC)	-0.100(3)	-0.031(3)	-0.074(3)	0.025(2)	0.027(2)
Gaunt (CI)				0.029(2)	0.031(2)
Retard (CI)				-0.001(1)	-0.002(1)
Freq (CI)				-0.001(1)	-0.0006(2)
QED (CI)				0.025(3)	0.022(3)
Total	23.07(4)	22.26(4)	22.45(4)	22.20(4)	11.97(4)
Eliav <i>et al.</i>	23.28	22.54	22.68	22.49	11.97
Nash					11.675
Yu et al.				21.98	11.73
Lackenby et al.				22.84	12.14
Hangele <i>et al.</i>				21.989	11.353

Таблица: Contributions to the ionization potentials (IP) of Cn atom (Z = 112 element), in eV.

- 1. E. Eliav et al., PRA, 52,2765, (1995).
- 2. T. Hangele et al, J. Chem. Phys., 136, 214105 (2012).
- 3. B. G. C. Lackenby et al., PRA, 101, 012514 (2020).
- 4. C. S. Nash et al., J. Phys. Chem. A, 109, 3493, (2005).
- 5. Y. J. Yu et al., EPJ D, 44, 51, (2007).

Contribution	I_3	I_2	I_1	EA
FSCC-SD	33.42(3)	23.91(3)	7.38(1)	0.66(1)
FSCC-T	-0.060(24)	0.008(3)	0.049(19)	0.023(9)
Gaunt (CI)	-0.060(6)	-0.049(6)	-0.055(6)	-0.040(10)
Retard (CI)	0.004(1)	0.003(1)	0.004(1)	0.003(1)
Freq (CI)	-0.005(1)	-0.005(1)	-0.003(1)	-0.002(1)
QED (CI)	-0.088(6)	-0.078(8)	-0.001(1)	0.007(4)
Total	33.21(4)	23.79(3)	7.37(2)	0.65(1)
Dzuba <i>et al.</i>	33.5	23.6	7.37	
Eliav <i>et al.</i>	33.47	23.96	7.306	0.68(5)
Pershina <i>et al.</i>			7.420	
Hangele <i>et al.</i>		23.6271	7.2779	

Таблица: Contributions to the ionization potentials (IP) and electron affinity (EA) of Nh atom (Z = 113 element), in eV.

- 1. V. A. Dzuba, PRA, 93, 032519 (2016).
- 2. E. Eliav et al., PRA, 53, (1996).
- 3. V. Pershina et al., J. Chem. Phys., 131, 084713, (2009).
- 4. T. Hangele et.al., J. Chem. Phys., 136, 214105, (2012).

Contribution	I_4	I_3	I_2	I_1
SD	46.19(4)	35.73(4)	16.92(4)	8.66(4)
Т	-0.074(25)	-0.019(6)	0.054(18)	0.036(12)
Gaunt	-0.076(2)	-0.069(2)	-0.068(2)	-0.044(2)
Retard	0.006(2)	0.004(2)	0.006(2)	0.005(2)
Freq	-0.006(1)	-0.006(1)	-0.004(1)	-0.003(1)
QED	-0.105(5)	-0.093(5)	-0.004(2)	-0.003(3)
Total	45.94(5)	35.55(4)	16.90(4)	8.65(4)
Dzuba <i>et al.</i>			17.00	8.37
Landau <i>et al.</i>	46.272	35.739	16.871	8.539
Yu et al.	46.57	35.82	17.22	8.28
Nash				8.529
Hangele <i>et al.</i>		35.3826	16.1105	7.2601

Таблица: Contributions to the ionization potentials (IP) of FI atom Z = 114 element, in eV.

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[†]Based upon ¹²C. () indicates the mass number of the longest-lived isotope

For the most precise values and uncertainties visit claaw.org and pml.nist.gov/data NIST SP 966 (July 2019)