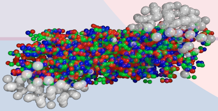
Core Corona approach, a model to explain Hyperon Global Polarization and other phenomena in semi-central heavy-ion collisions at low energy.



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MPD Cross-PWGroup meeting

Outline

Section 1

Abstract-Summary

Summary

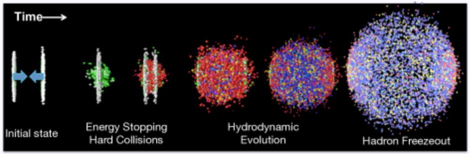
The study of heavy ion collisions at low energies is a topic of great interest since it was suggested that the onset of the hadron to quark matter transition should be expected at low energies where high baryon densities and lower temperatures dominates. To characterize the strongly interacting matter created in such type of collisions, several variables has been proposed, in particular hyperon global polarization, due to the possibility to link this observable to its vorticity, viscosity and flow, and shed light of criticality in the nuclear matter phase diagram. Recently experiments such as HADES, and BES at RHIC, measured the global polarization of Λ and $\bar{\Lambda}$ and so it has been observed that the global polarization increases as the energy of the collision decreases, being greater the effect for the Λ. To describe this behavior, we have implemented the Core-Corona model, which assumes that in non-central heavy ion collisions, hyperons come from different density regions of the created system. We found that the relative abundance between one region and another influences global polarization, and it reaches a maximum at collision energies $\sqrt{s_{NN}} \leq 10$ GeV. In this talk I will show the details of the hyperon global polarization estimation and plans to apply it to other measurements in the frame of the MPD experiment.

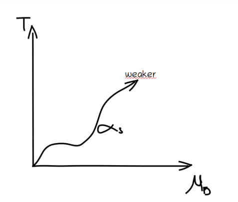
Section 2

Heavy-Ion Collisions at low energy

QGP and Heavy-Ion Collisions

- The QCD predicts a new phase of matter the QGP.
- It exist at early universe, 10 μs after the big bang.
- Heavy-Ion Collisions reproduce conditions at 10 μs after Big-Bang
- Transition from QGP to hadron gas.
- Strongly interacting matter in equilibrium is characterized by two quantities T and μ_B (or n_B)

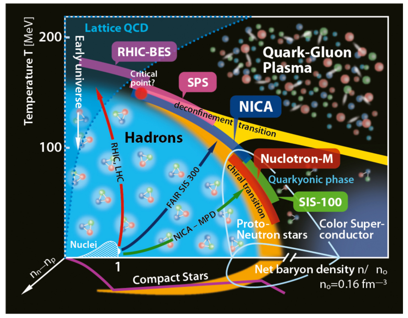




Due to asymptotic freedom, α_s diminishes with increasing the energy scale producing that interactions among strongly interacting particles will get weaker as T or μ_B are increased.

J.Phys.Conf.Ser. 50 (2006) 238-242

Exploration of the QCD Phase Diagram



MPD experiment is focus in look for new phenomena in the baryon-rich region of QCD phase diagram

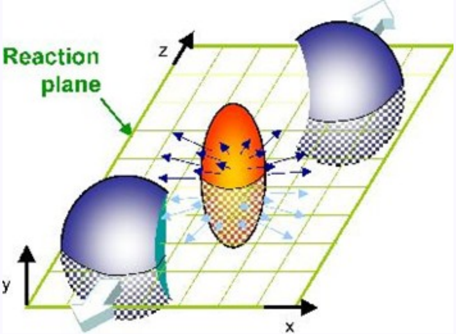
Different experiments allow us to:

- Scan different parts of the QCD phase diagram which has different characteristics
- At Low energy collisions, matter differ form that studied at SPS, RHIC or LHC, because it consist principally for baryons and few mesons and can be compressed until 3 times the density of nuclear matter for approx 10-12 fm/c
- State of matter produced in HIC has fluid properties.
- Non-central collisions have a large angular momentum and strong vortical structure.

Section 3

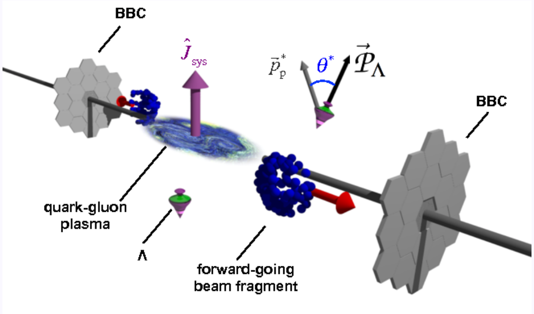
Hyperon Global Polarization and vorticity

Global Vorticity and polarization in heavy Ion collisions



- Non-central collisions have large angular momentum $L \sim 10^5 \hbar$.
- Shear forces in initial condition introduce vorticity to the QGP.
- Spin-orbit coupling: spin alignment, or polarization, along the direction of the vorticity -on average- parallel to J.

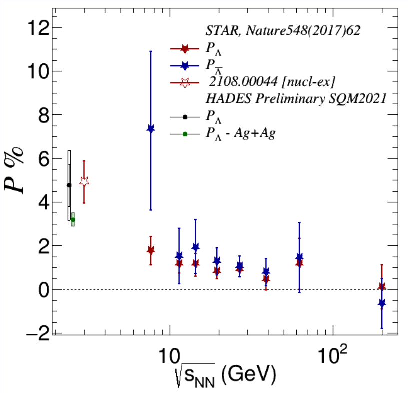
Why are we interested in measure Hyperon Global Polari



The fluid at midrapidity has a whirling substructure oriented (on average) in the direction of the total angular momentum, $\hat{J}$. [Nature 548,62-65(2017)]

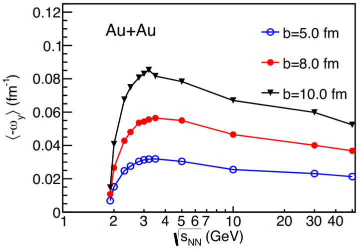
- The Λ and $\bar{\Lambda}$ polarization are linked to the properties of the medium produced in relativistic heavy-ion collisions.
- For semi-central collisions, Angular momentum can be quantified in terms of the thermal vorticity
- The global polarization can be measured using the self-analysing $\Lambda/\bar{\Lambda}$ decays.

Global Polarization as a function of energy



Energy range $\sqrt{s_{NN}} = \{2, 11\} GeV$ can be covered by ongoing/future experiments.

STAR BES-II + FXT: 3-19 GeV
HADES:2-3 GeV
NICA:4-11 GeV → MPD



Energy dependence of kinematic vorticity predicted by a transport model (UrQMD) [a]

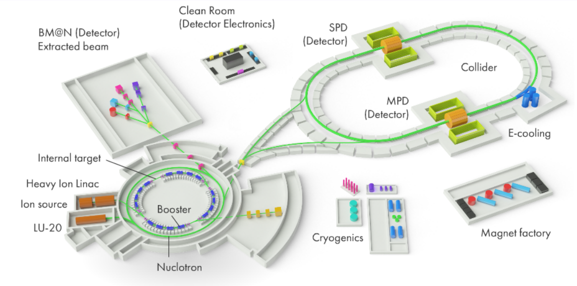
[a] X.-G. Deng et al., PRC101.064908(2020)

Section 4

MPD Experiment

NICA The Nuclotron Ion fAcility

Aimed at study of hot and dense nuclear and baryonic matter in HIC at a center of mass energies in the range $\sqrt{s_{NN}} = 4 - 11$ GeV. The average luminosity expected is $1 \cdot 10^{27}$ cm^{-2}s$^{-}1$.



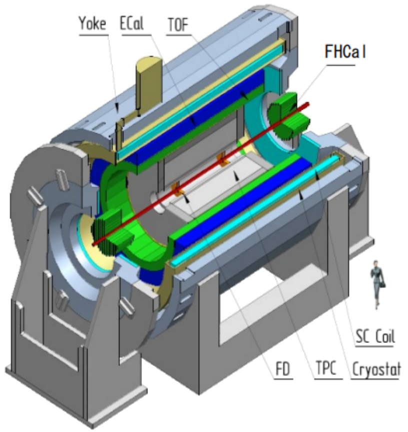
It look for sheed light on

- In-medium properties of hadrons and the nuclear matter equation of state (EoS)
- The onset of deconfinement (OD) and/or chiral symmetry restoration (CSR)
- Phase transition (PT)
- Mixed Phase (MP)
- The Critical End Point (CEP)

Multi Purpose Detector

MPD physics goals

- Hadrochemistry.
- Anisotropic flow measurements.
- Intensity interferometry.
- Fluctuations.
- Short lived resonances.
- Electromagnetic probes.

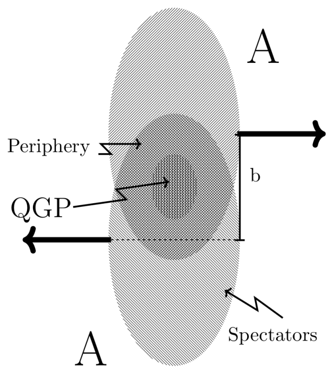


Section 5

Core-Corona Model

Core-Corona Model

In non central collisions we can identify two regions with different density: Core and Corona



- In the Corona: reactions like $p + p \rightarrow K + \Lambda + p$ Polarization is described by Lund Model, DeGrand-Miettinen model or Gluon bremsstrahlung mechanism
- In the Core: spin aligment driven by vorticity or magnetic field

Articles

- **Core meets corona: A two-component source to explain Λ and $\bar{\Lambda}$ global polarization in semi-central heavy-ion collisions** Phys.Lett.B Volume 810, (2020) 135818
- **The rise and fall of Λ and $\bar{\Lambda}$ global polarization in semi-central heavy-ion collisions at HADES, NICA and RHIC energies from the core-corona model** arXiv:2106.14379v1 [hep-ph] 28 Jun 2021

The rise and fall of Λ and $\bar{\Lambda}$ global polarization in semi-central heavy-ion collisions at HADES, NICA and RHIC energies from the core-corona model

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We compute the Λ and $\bar{\Lambda}$ global polarizations in semi-central heavy-ion collisions using the core-corona model where the source of Λ's and $\bar{\Lambda}$'s is taken as consisting of a high-density core and a less dense corona. We show that the overall properties of the polarization excitation functions can be linked to the relative abundance of Λs coming from the core versus those coming from the corona. For low collision energies, the former are more abundant whereas for higher energies the latter become more abundant. The main consequence of this reversing of the relative abundance is that both polarizations peak at collision energies $\sqrt{s_{NN}} \lesssim 10$ GeV. The exact positions and heights of these peaks depend not only on this reversal of relative abundances, but also on the centrality class, which is directly related to the QGP volume and lifetime, as well as on the relative abundances of Λs and $\bar{\Lambda}$s in the core and corona regions. The intrinsic polarizations are computed from a field theoretical approach that links the alignment of the strange quark spin with the thermal vorticity and modeling the QGP volume and lifetime using a Bjorken expansion scenario. We predict that the Λ and $\bar{\Lambda}$ global polarizations should peak at the energy range accessible to NICA and HADES.

I. INTRODUCTION

The polarization properties of Λ and $\bar{\Lambda}$ have received increasing attention over the last years due to the possibility to link this observable to the properties of the medium produced in relativistic heavy-ion collisions [1–13]. For semi-central collisions, the matter density profile in the transverse plane develops an angular momentum [14] which can be quantified in terms of the thermal vorticity [15]. When this vorticity is transferred to spin degrees of freedom, the global polarization can be measured using the self-analising $\Lambda/\bar{\Lambda}$ decays.

The Beam Energy Scan (BES) at RHIC, performed by the STAR Collaboration [16, 17] has shown a trend for the Λ and $\bar{\Lambda}$ global polarization to increase as the energy of the collision decreases and that this increase is faster for $\bar{\Lambda}$s than for Λs. In addition, the HADES Collaboration has recently provided preliminary results on the Λ global polarization in Au+Au collisions at $\sqrt{s_{NN}} = 2.42$ short-lived but intense magnetic fields [21–23] and the possibility that Λ and $\bar{\Lambda}$ align their spins with the direction of the angular momentum created in the reaction during the life-time of the evolving system [24, 25].

In a recent work [26], we have shown that when in semi-central heavy-ion collisions, the source of Λs and $\bar{\Lambda}$s is modeled as a high-density core and a less dense corona, the global polarization properties of these hyperons, as a function of the collision energy, are well described. The QGP is produced in the core when the density of participants is larger than a critical value. At the same time, this region corresponds to the low baryon density. On the other hand, the corona corresponds to the region with the overall larger baryon density. For a given impact parameter (or rather, a centrality class), the volume in the corona becomes larger for lower energies. We found that when the larger abundance of Λs compared to $\bar{\Lambda}$s coming from the corona is combined with a smaller number of Λs coming from the core, compared to those from the

Core-Corona Model: Two-component source

In heavy-ion collisions, Λ and $\bar{\Lambda}$ come from different density regions

- **Core**: Via QGP processes like

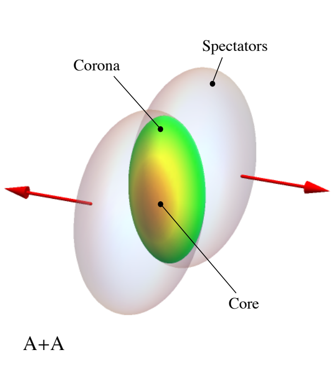
$$q\bar{q} \to s\bar{s} \text{ and } gg \to s\bar{s}$$

- **Corona**: Via $n + n$ reactions by recombination-like processes

The number of Λs can be written as: $N_\Lambda = N_{\Lambda_{QGP}} + N_{\Lambda_{REC}}$
Then the polarization given by:

$$\mathcal{P} = \frac{N^{\uparrow} - N^{\downarrow}}{N^{\uparrow} + N^{\downarrow}}$$

can be rewritten in terms of the number of Λs (or $\bar{\Lambda}$s) produced in the different density regions



Phys.Lett.B 810 (2020) 135818

Rewriting Polarization

$$\mathcal{P}^{\Lambda} = \frac{(N^{\uparrow}_{\Lambda_{QGP}} + N^{\uparrow}_{\Lambda_{REC}}) - (N^{\downarrow}_{\Lambda_{QGP}} + N^{\downarrow}_{\Lambda_{REC}})}{(N^{\uparrow}_{\Lambda_{QGP}} + N^{\uparrow}_{\Lambda_{REC}}) + (N^{\downarrow}_{\Lambda_{QGP}} + N^{\downarrow}_{\Lambda_{REC}})}$$

$$\mathcal{P}^{\bar{\Lambda}} = \frac{(N^{\uparrow}_{\bar{\Lambda}_{QGP}} + N^{\uparrow}_{\bar{\Lambda}_{REC}}) - (N^{\downarrow}_{\bar{\Lambda}_{QGP}} + N^{\downarrow}_{\bar{\Lambda}_{REC}})}{(N^{\uparrow}_{\bar{\Lambda}_{QGP}} + N^{\uparrow}_{\bar{\Lambda}_{REC}}) + (N^{\downarrow}_{\bar{\Lambda}_{QGP}} + N^{\downarrow}_{\bar{\Lambda}_{REC}})}$$

After some algebra, we get:

$$\mathcal{P}^{\Lambda} = \frac{\left(\mathcal{P}^{\Lambda}_{REC} + \frac{N^{\uparrow}_{\Lambda_{QGP}} - N^{\downarrow}_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}\right)}{\left(1 + \frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}\right)}$$

$$\mathcal{P}^{\bar{\Lambda}} = \frac{\left(\mathcal{P}^{\bar{\Lambda}}_{REC} + \frac{N^{\uparrow}_{\bar{\Lambda}_{QGP}} - N^{\downarrow}_{\bar{\Lambda}_{QGP}}}{N_{\bar{\Lambda}_{REC}}}\right)}{\left(1 + \frac{N_{\bar{\Lambda}_{QGP}}}{N_{\bar{\Lambda}_{REC}}}\right)}$$

Where the polarization along the angular momentum produced in the corona is:

$$\mathcal{P}^{\Lambda}_{REC} = \frac{N^{\uparrow}_{\Lambda_{REC}} - N^{\downarrow}_{\Lambda_{REC}}}{N^{\uparrow}_{\Lambda_{REC}} + N^{\downarrow}_{\Lambda_{REC}}}$$

$$\mathcal{P}^{\bar{\Lambda}}_{REC} = \frac{N^{\uparrow}_{\bar{\Lambda}_{REC}} - N^{\downarrow}_{\bar{\Lambda}_{REC}}}{N^{\uparrow}_{\bar{\Lambda}_{REC}} + N^{\downarrow}_{\bar{\Lambda}_{REC}}}$$

Assumptions: Polarization of Λ ($\bar{\Lambda}$) from the Corona

$$\mathcal{P}^{\Lambda} = \frac{\left(\mathcal{P}^{\Lambda}_{REC} + \frac{N^{\uparrow}_{\Lambda_{QGP}} - N^{\downarrow}_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}\right)}{\left(1 + \frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}\right)}$$

$$\mathcal{P}^{\bar{\Lambda}} = \frac{\left(\mathcal{P}^{\bar{\Lambda}}_{REC} + \frac{N^{\uparrow}_{\bar{\Lambda}_{QGP}} - N^{\downarrow}_{\bar{\Lambda}_{QGP}}}{N_{\bar{\Lambda}_{REC}}}\right)}{\left(1 + \frac{N_{\bar{\Lambda}_{QGP}}}{N_{\bar{\Lambda}_{REC}}}\right)}$$

Where

$$\mathcal{P}^{\Lambda}_{REC} = \frac{N^{\uparrow}_{\Lambda_{REC}} - N^{\downarrow}_{\Lambda_{REC}}}{N^{\uparrow}_{\Lambda_{REC}} + N^{\downarrow}_{\Lambda_{REC}}}$$

$$\mathcal{P}^{\bar{\Lambda}}_{REC} = \frac{N^{\uparrow}_{\bar{\Lambda}_{REC}} - N^{\downarrow}_{\bar{\Lambda}_{REC}}}{N^{\uparrow}_{\bar{\Lambda}_{REC}} + N^{\downarrow}_{\bar{\Lambda}_{REC}}}$$

- Nucleon-nucleon scattering not enough to align the spin in the direction of the angular momentum.
- Polarization of Λ and $\bar{\Lambda}$ averages to zero.

$$\mathcal{P}^{\Lambda}_{REC} = \mathcal{P}^{\bar{\Lambda}}_{REC} = 0$$

Assumptions: Intrinsic Polarization

$$\mathcal{P}^{\Lambda} = \frac{\left(\mathcal{P}^{\Lambda}_{REC} + \frac{N^{\uparrow}_{\Lambda_{QGP}} - N^{\downarrow}_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}\right)}{\left(1 + \frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}\right)}$$

$$\mathcal{P}^{\bar{\Lambda}} = \frac{\left(\mathcal{P}^{\bar{\Lambda}}_{REC} + \frac{N^{\uparrow}_{\bar{\Lambda}_{QGP}} - N^{\downarrow}_{\bar{\Lambda}_{QGP}}}{N_{\bar{\Lambda}_{REC}}}\right)}{\left(1 + \frac{N_{\bar{\Lambda}_{QGP}}}{N_{\bar{\Lambda}_{REC}}}\right)}$$

We define **z** and **z̄** which represent the Λ and $\bar{\Lambda}$ intrinsic polarization respectively

$$N^{\uparrow}_{\Lambda_{QGP}} - N^{\downarrow}_{\Lambda_{QGP}} = \mathbf{z} N_{\Lambda_{QGP}}$$

$$N^{\uparrow}_{\bar{\Lambda}_{QGP}} - N^{\downarrow}_{\bar{\Lambda}_{QGP}} = \bar{\mathbf{z}} N_{\bar{\Lambda}_{QGP}}$$

Assumptions: The ratio $N_{\bar{\Lambda}_{REC(QGP)}}/N_{\Lambda_{REC(QGP)}}$

$$\mathcal{P}^{\Lambda} = \frac{\left(\mathcal{P}^{\Lambda}_{REC} + \frac{N^{\uparrow}_{\Lambda_{QGP}} - N^{\downarrow}_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}\right)}{\left(1 + \frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}\right)}$$

$$\mathcal{P}^{\bar{\Lambda}} = \frac{\left(\mathcal{P}^{\bar{\Lambda}}_{REC} + \frac{N^{\uparrow}_{\bar{\Lambda}_{QGP}} - N^{\downarrow}_{\bar{\Lambda}_{QGP}}}{N_{\bar{\Lambda}_{REC}}}\right)}{\left(1 + \frac{N_{\bar{\Lambda}_{QGP}}}{N_{\bar{\Lambda}_{REC}}}\right)}$$

The number of $\bar{\Lambda}$s are proportional to an energy-dependent coefficient **w**(**w**′) times the number of Λs in the corona(core).

$$N_{\bar{\Lambda}_{REC}} = \mathbf{w} N_{\Lambda_{REC}}$$
$$N_{\bar{\Lambda}_{QGP}} = \mathbf{w}' N_{\Lambda_{QGP}}$$

Λ and $\bar{\Lambda}$ global polarization

With this assumptions:

$$\mathcal{P}^{\Lambda}_{REC} = \mathcal{P}^{\bar{\Lambda}}_{REC} = 0$$

$$\begin{aligned} N^{\uparrow}_{\Lambda_{QGP}} - N^{\downarrow}_{\Lambda_{QGP}} &= z N_{\Lambda_{QGP}} \\ N^{\uparrow}_{\bar{\Lambda}_{QGP}} - N^{\downarrow}_{\bar{\Lambda}_{QGP}} &= \bar{z} N_{\bar{\Lambda}_{QGP}} \end{aligned}$$

$$\begin{aligned} N_{\bar{\Lambda}_{REC}} &= w N_{\Lambda_{REC}} \\ N_{\bar{\Lambda}_{QGP}} &= w' N_{\Lambda_{QGP}} \end{aligned}$$

Global polarization depends on the coefficients w, w', z, $\bar{z}$ and the ratio $\frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}$ that can be estimated from data or calculated.

$$\begin{aligned} \mathcal{P}^{\Lambda} &= \frac{z \frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}}{1 + \frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}} \\ \mathcal{P}^{\bar{\Lambda}} &= \frac{\bar{z} \left(\frac{w'}{w}\right) \frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}}{1 + \left(\frac{w'}{w}\right) \frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}} \end{aligned}$$

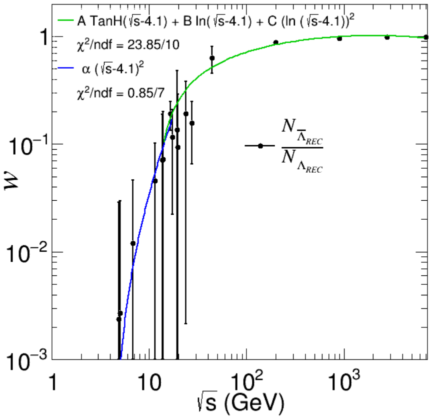
The ratio $w = N_{\bar{\Lambda}_{REC}} / N_{\Lambda_{REC}}$

Model as $p + p$ collisions

- Experimental data obtained from $p + p$ collisions at different energies[1]
- w is defined only for $\sqrt{s} > 4.1$GeV. The threshold energy for

$$p + p \rightarrow p + p + \Lambda + \bar{\Lambda}$$

- w is smaller than 1 except for energies $\sqrt{s} > 1$ TeV



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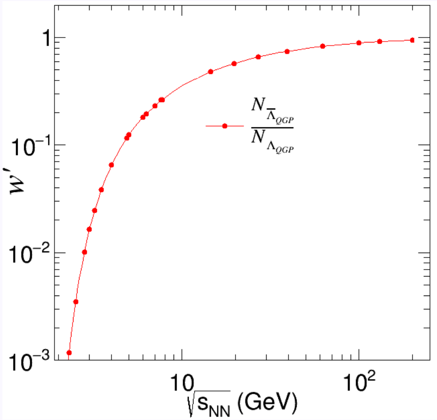
[1] M. Gazdzicki and D. Rohrich, Z. Phys. C 71 (1996) 55; V. Blobel et al. Nucl. Phys. B 69(1974), 454-492; J. W. Chapman et al., Phys. Lett. 47B (1973) 465; D. Brick et al., Nucl. Phys. B 164 (1980) 1; C. Höhne, CERN-THESIS-2003-034; J. Baechler et al. [NA35 Collaboration], Nucl. Phys. A 525 (1991) 221C; G. Charlton et al., Phys. Rev. Lett. 30 (1973) 574; F. Lopinto et al., Phys. Rev. D 22 (1980) 573; H. Kichimi et al., Phys. Rev. D 20 (1979) 37; F. W. Busser et al., Phys. Lett. 61B (1976) 309; S. Erhan, et al., Phys. Lett. 85B(1979) 447; B. I. Abelev et al. [STAR Collaboration], Phys. Rev. C 75 (2007) 064901; E. Abbas et al. [ALICE Collaboration], Eur. Phys. J. C 73 (2013) 2496

The ratio $w' = N_{\bar{\Lambda}_{QGP}}/N_{\Lambda_{QGP}}$

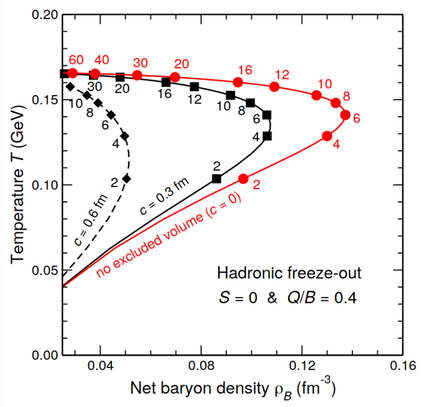
- The coefficient w' is computed as the ratio of the equilibrium distributions of $\bar{s}$ to s-quark for a given temperature T and chemical potential $\mu = \mu_B/3$ given by:

$$w' = \frac{e^{(m_s-\mu)/T} + 1}{e^{(m_s+\mu)/T} + 1}$$

where $m_s = 100$ MeV is the mass of the s-quark and T and μ_B are taken along the curve of the maximum chemical potential at freeze out.



μ_B and T at freeze out



Eur.Phys.J. 52 (2016) 218-219

At maximum freeze-out baryon density in nuclear collisions, the extracted values of T and μ_B exhibits a smooth and monotonic dependence on the collision energy

$$T(\mu_B) = 166 - 139\mu_B^2 - 53\mu_B^4$$
$$\mu_B(\sqrt{s_{NN}}) = \frac{1308}{1000 + 0.273\sqrt{s_{NN}}}$$

Phys.Rev.C 74 (2006) 047901

Production of Λ in the core and the corona

Number of Λs in the core

$$N_{\Lambda_{QGP}} = cN^2_{p_{QGP}}$$

in which

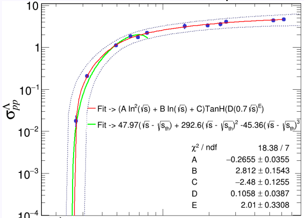
$$N_{p_{QGP}} = \int n_p(\mathbf{s,b})\theta[n_p(\mathbf{s,b}) - n_c]\, d^2 s$$

with $n_c = 3.3$ fm^{-2}, the critical density required to form the QGP [Phys.Rev.Lett. 77 (1996), 1703-1706]

Number of Λs in the corona

$$N_{\Lambda_{REC}} = \sigma^{\Lambda}_{NN} \int T_B(\mathbf{b} - \mathbf{s})\, T_A(\mathbf{s})\theta[n_c - n_p(\mathbf{s,b})]\, d^2 s$$

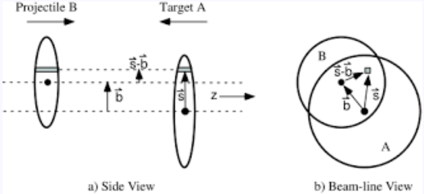
where σ^{Λ}_{NN} is obtained from experimental data

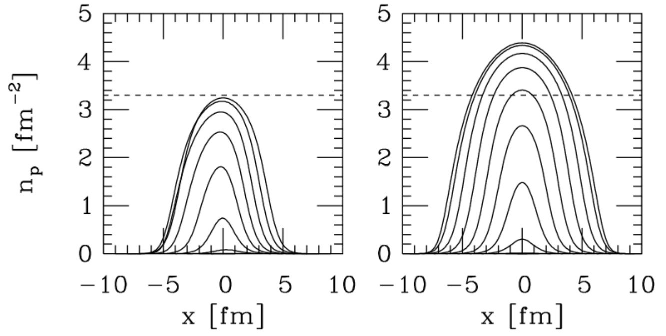


The density of participants $n_p(\mathbf{s,b}) = T_A(\mathbf{s})[1 - e^{-\sigma_{NN} T_B(\mathbf{s-b})}] + T_B(\mathbf{s-b})[1 - e^{-\sigma_{NN} T_A(\mathbf{s})}]$,
The thickness function $T_A(z,s) = \int_{\infty}^{\infty} \rho_A(z,\mathbf{s})dz$ and
the Woods-Saxon profile density $\rho_A(\mathbf{s}) = \frac{\rho_0}{1+e^{(r-R_A)/a}}$

The critical density n_c

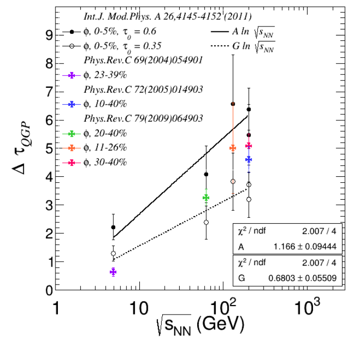
QGP formation is related with the J/Psi suppresion (Phys.Rev.Lett. 77 (1996) 1703-1706).





The density of participants $n_p(s)$, for s along the direction of the impact parameter, for various values of the impact parameter: $b = 0, 2, 4 \cdot \cdot \cdot$ fm. left: S−U collision; right: Pb−Pb collision. The horizontal dashed line corresponds to the largest density achieved in the S−U system, $n_p = 3.3$fm^{-2}.

Number of Λs and $\bar{\Lambda}$s as a function of energy

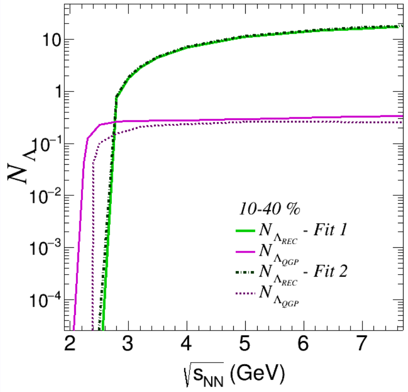


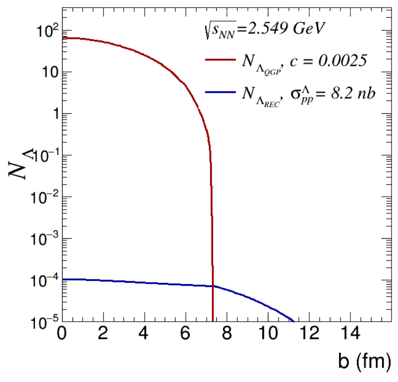
$N_{\Lambda_{QGP}}$ and $N_{\Lambda_{REC}}$ as a function of the collision energy for impact parameters $b = 0, 4, 7$ fm.

- At small b, Λ particle production is dominated by the core region.
- For peripheral collisions, Λ particle production is dominated by the corona region $\rightarrow$ relevant for vorticity and polarization studies.

Core-corona model introduces a critical density of participants n_c above which the core can be produced. For peripheral collisions is difficult to achieve this critical density n_c, even for the largest collision energies.

Λs in the Core and Corona





At low energies $N_{\Lambda_{QGP}}$ depends on σ_{NN}, different parametrizations impact on the strenght of polarization

σ_{NN} affects the ratio $N_{\Lambda QGP}/N_{\Lambda_{REC}}$ and the value of b at which the ratio is smaller than 1

Intrinsic Polarization

Intrinsic polarization is given by:

$$z = 1 - e^{-\Delta\tau_{QGP}/\tau}$$

and

$$\bar{z} = 1 - e^{-\Delta\tau_{QGP}/\bar{\tau}}$$

in terms of the relaxation times τ and $\bar{\tau}$ and the QGP life-time $\Delta\tau_{QGP}$
The relaxation time can be computed as the inverse of the interaction rate

$$\tau \equiv 1/\Gamma$$

given by

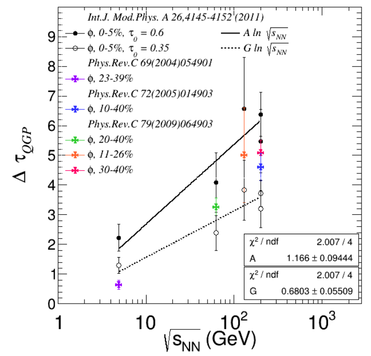
$$\Gamma = V \int \frac{d^3p}{2\pi^3}\Gamma(p_0), \text{ with } V = \pi R^2 \Delta\tau_{QGP}$$

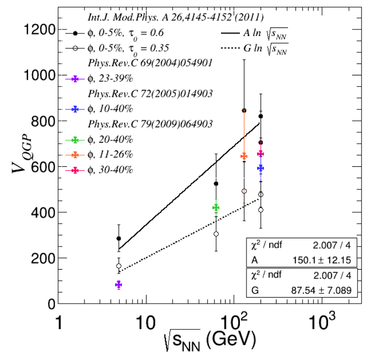
where V is the volume of the core region[2], related with the QGP life-time $\Delta\tau_{QGP}$ in the scenario of a Bjorken expansion

$$\Delta\tau_{QGP} = \tau_f - \tau_0 = \tau_0[\left(\frac{T_0}{T_f}\right)^3 - 1]$$

[2] Phys.Rev.D 102,056019 (2020)

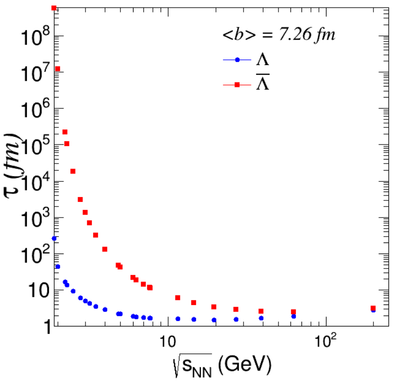
Volume and QGP life-time

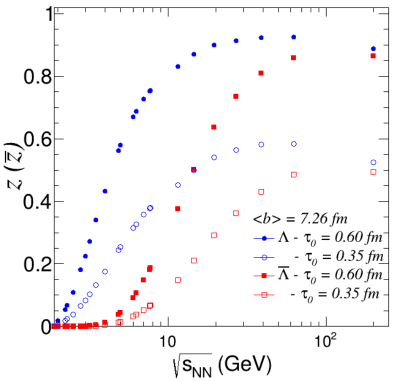




T_0 is estimated from p_T of ϕ mesons, $\tau_0 = 0.35 - 0.60$ fm to incorporate the effect of collision centrality, and T_f is taken as the value along the maximum chemical potential curve at freeze-out

Relaxation time and intrinsic polarization as a function o





For energies below the Λ production threshold energy, the τ and $\bar{\tau}$ increase dramatically, as expected, since the interaction rate should vanish below these energies.

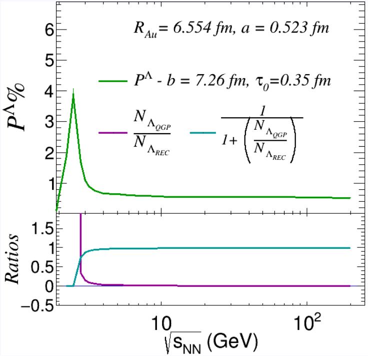
Section 6

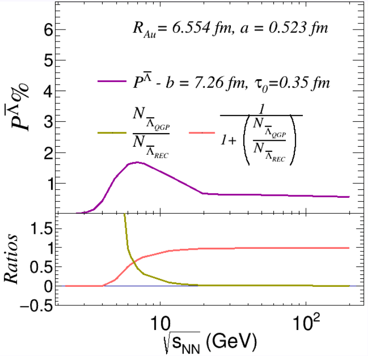
Excitation function for the Global Λ and $\bar{\Lambda}$ Polarization

The ratios describing the polarization

$$\mathcal{P}^{\Lambda}/z = \frac{\frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}}{1+\frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}}$$

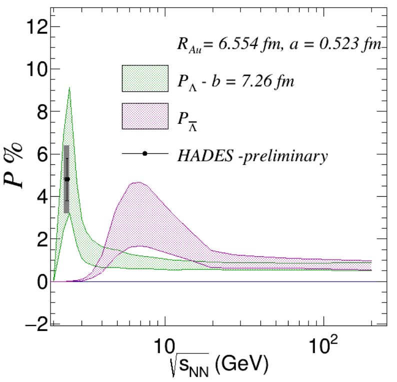
$$\mathcal{P}^{\bar{\Lambda}}/\bar{z} = \frac{\left(\frac{w'}{w}\right)\frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}}{1+\left(\frac{w'}{w}\right)\frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}}$$

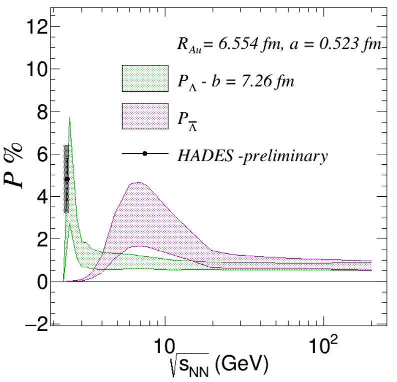




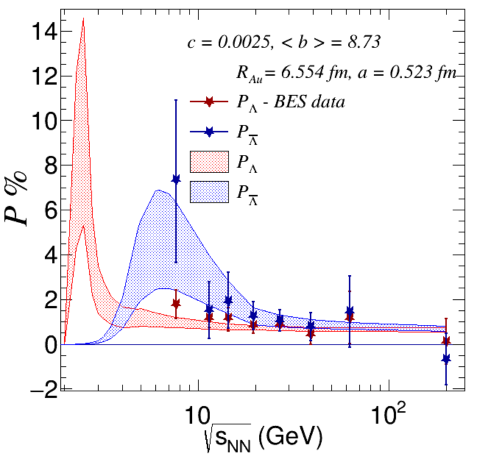
Λ and $\bar{\Lambda}$ polarization in Au+Au at HADES centrality

10 - 40 %





Λ and $\bar{\Lambda}$ in Au+Au at BES centrality

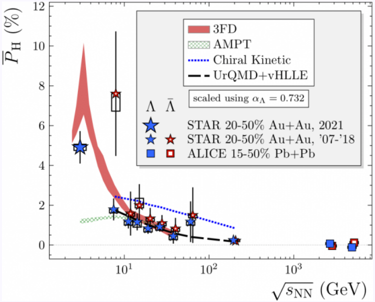


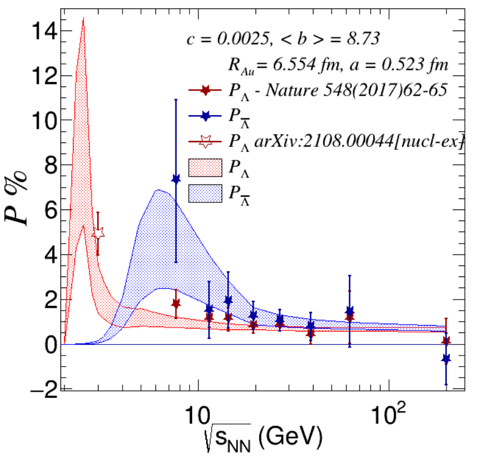
20 - 50 %

- Similar trend to the case of the analysis with smaller centrality.
- Magnitude of global polarization increases for a larger centrality as a consequence of the angular velocity increase

More Recent Results - $Au + Au$ at $\sqrt{s_{NN}} = 3$ GeV

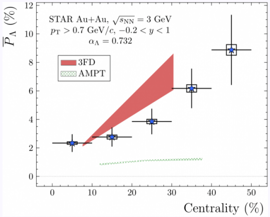
Recently STAR collaboration presents:
arXiv:2108.00044v2 [nucl-ex]

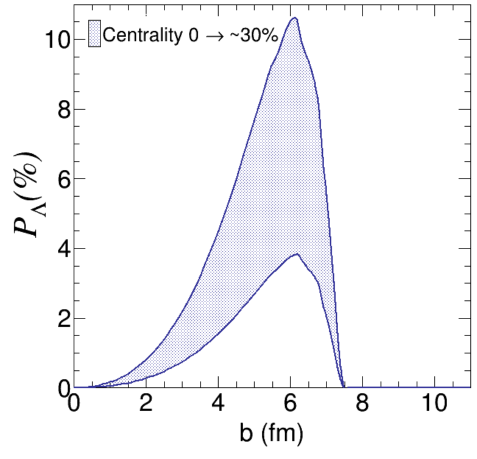




Our previous prediction fits the new value.
3FD model predicts also a peak.

Λ polarization as a function of centrality for $Au + Au$ at $\sqrt{s_{NN}} = 3$ GeV

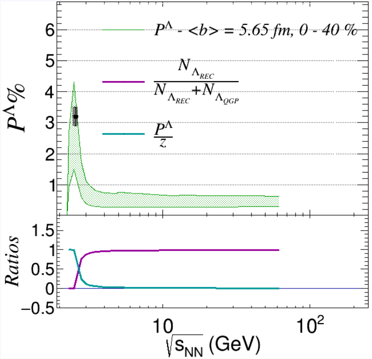


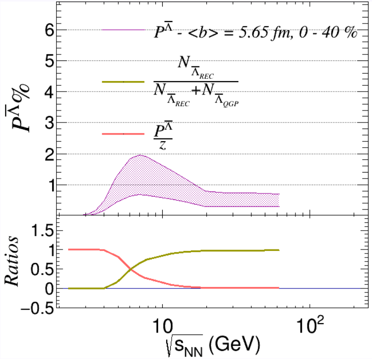


Polarization for more peripheral collisions goes to zero, as the critical density n_c of the system is not achieved, vanishing the number of Λs from the core.

Λ and $\bar{\Lambda}$ polarization in $Ag + Ag$ collisions

Similar trend to Au+Au



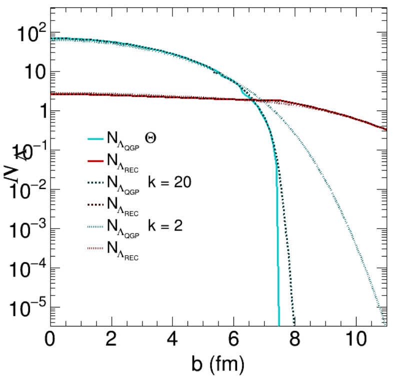


Due the system's size, the minimum critical density n_c to produce QGP is barely achieved for non-central collisions.

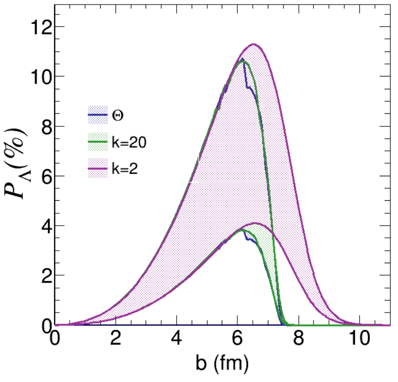
A change in the critical densisty n_c

The number of Λ's is dependent on the critical density $n_c = 3.3$ fm^{-2}.

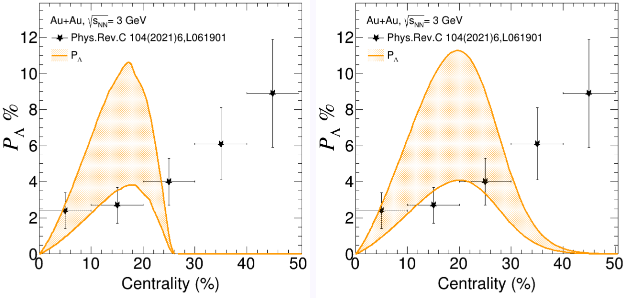
- A change in the value of the n_c, allow the QGP formation for higher b.
- Change $\theta(x) \to \frac{1}{1+2e^{-2kx}}$ with $k = \{2, 20\}$.
- Higher $k \to \theta(x)$



Polarization as a function of centrality



The polarization increases for 20-30% centrality bin, and is different from zero for 30-40% however does not describe data at higher bins of centrality



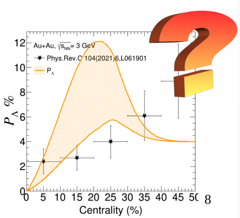
What about the contribution of P_{REC}?

- Which is the effect of transverse Λ polarization in the corona?
- The polarization in pp collisions is not zero[3].
  - At $\sqrt{s} = 19.6\text{GeV} \to \mathcal{P} = -0.25 \pm 0.26$
  - At $\sqrt{s} = 53\text{GeV} \to \mathcal{P} = -0.34 \pm 0.07$
  - At $\sqrt{s} = 62\text{GeV} \to \mathcal{P} = -0.40 \pm 0.10$

$$\mathcal{P}^{\Lambda} = \frac{\mathcal{P}^{\Lambda}_{REC} + z\frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}}{1 + \frac{N_{\Lambda_{QGP}}}{N_{\Lambda_{REC}}}}$$

In more peripheral collisions transverse polarization is not diluted by rescattering within QCD medium and can be measured at MPD[4].

Which is the value w.r.t. the angular momentum?



8

Adding an arbitrary value of $\mathcal{P}^{\Lambda}_{REC}$ Data could be described?

3 4

[3] PoS HEP2005 (2006) 122, V. Blobel et al., Nucl. Phys. B122 (1977) 429, Phys. Rev.,D11:2405, 1975

[4] Nazarova, et. al. Phys. of Part. and Nuclei Lett., 2021, Vol. 18, No. 4, pp. 429–438

Polarization from Corona region

In pp collisions, transverse polarization is measured with respect to production plane. Polarization along vector $\hat{n}$ perpendicular to the plane defined by the beam $\hat{p}_{beam}$ and Λ directions, p_Λ

$$\hat{n} \equiv \frac{\bar{p}_{beam} \times \bar{p}_\Lambda}{|\bar{p}_{beam} \times \bar{p}_\Lambda|}$$

Assuming the beam direction parallel to $\hat{z}$, we can express $\hat{n}$ like:

$$\hat{n} = \frac{1}{p_{T_\Lambda}}(-p_{y_\Lambda}, p_{x_\Lambda}, 0)$$

Local Polarization projected along angular momentum 1

Assuming that in pp collisions, polarization $\mathcal{P}_T$ is only different from zero along $\hat{n}$; for Λ's in the corona, the contribution to global polarization can be measured by:

$$\frac{dN}{d\Omega} = \frac{N}{4\pi}\left(1 + \alpha\mathcal{P}_T \cos\sigma^*\right)$$

where σ^* is the angle between $\hat{n}$ and the direction of the angular momentum $\hat{L} = \hat{b} \times \hat{p}_{beam} = (\sin\Psi_{RP}, -\cos\Psi_{RP}, 0)$.
Then $cos\sigma^*$ is given by:

$$\cos\sigma^* = \hat{n}\cdot\hat{L} = \frac{1}{p_{T_\Lambda}}\left(-p_{y_\Lambda}\sin\Psi_{RP} - p_{x_\Lambda}\cos\Psi_{RP}\right)$$

Substituting

$$\begin{aligned} p_{x_\Lambda} &= p_\Lambda \sin\theta_\Lambda \cos\phi_\Lambda \\ p_{y_\Lambda} &= p_\Lambda \sin\theta_\Lambda \sin\phi_\Lambda \\ p_{T_\Lambda} &= p_\Lambda \sin\theta_\Lambda \end{aligned}$$

Local Polarization projected along angular momentum 2

then

$$\begin{aligned}\cos\sigma^* &= -\sin\phi_\Lambda \sin\Psi_{RP} - \cos\phi_\Lambda \cos\Psi_{RP} \\ &= -\cos\left(\phi_\Lambda - \Psi_{RP}\right)\end{aligned}$$

angular distribution can be rewritten like:

$$\frac{dN}{d\Omega} = \frac{N}{4\pi}\left(1 - \alpha\mathcal{P}_T \cos\left(\phi_\Lambda - \Psi_{RP}\right)\right)$$

Considering $d\Omega = \sin\theta\, d\theta\, d\phi$ and integrating w.r.t. $d\theta$?

$$\begin{aligned}\frac{dN}{d\phi} &= \int_0^\pi \left[\frac{N}{4\pi}\left(1 - \alpha\mathcal{P}_T \cos\left(\phi_\Lambda - \Psi_{RP}\right)\right)\right] \sin\theta\, d\theta \\ &= \frac{N}{2\pi}\left(1 - \alpha\mathcal{P}_T \cos\left(\phi_\Lambda - \Psi_{RP}\right)\right)\end{aligned}$$

Local Polarization projected along angular momentum 3

Calculating the mean angular distribution $\langle \cos(\phi_\Lambda - \Psi_{RP}) \rangle$

$$\langle \cos(\phi_\Lambda - \Psi_{RP}) \rangle = -\frac{\alpha \mathcal{P}_T}{2}$$

The transverse polarization projected along angular momentum should be

$$\mathcal{P}_T = \frac{-2\langle \cos(\phi_\Lambda - \Psi_{RP}) \rangle}{\alpha}$$

that differs from the global polarization given by $\mathcal{P}_\Lambda = -\frac{8\langle \sin(\phi_p - \Psi_{RP}) \rangle}{\pi\alpha}$
There is a similarity between this expression and directed flow for Λs
Can we measure this contribution experimentally?

Perspectives of study at MPD with UrQMD

The UrQMD generator implements an **hybrid model** that includes an ideal fluid-dynamic evolution for the hot an dense stage[a].

The **fluid-dynamic evolution** is carried out by the SHASTA (SHarp and Smooth Transport Algorithm)[b]

Implementation of EoS that includes a **deconfinement plus a chiral phase transition**, through a smooth crossover between a chiral hadronic model and an interacting constituent quark model [c].

A **core-corona like separation** mechanism for the initial state of the fluid evolution. Quark density cut in η intervals for select particles in the fluid-dynamical evolution [d].

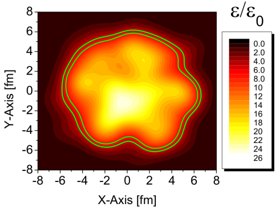


FIG. 1. (Color online) Contour plot of the local rest frame energy density in the transverse plane ($z = 0$) of a central ($b = 0$) collision of Pb+Pb at $E_{lab} = 40A$ GeV. The energy density is normalized to the ground state energy density ($\epsilon_0 \approx$ 145 MeV/fm^3). The two green lines correspond to lines of a constant energy density of $\epsilon/\epsilon_0 = 5$ and 7.

[a] Phys. Rev. C 78 (2008) 044901

[b] Nucl.Phys.A595(1995)346, Nucl. Phys.A595(1995)383

[c] Phys.Rev.C84,045208(2011)

[d] Phys.Rev.C84,024905(2011)

Another studies with core-corona approach

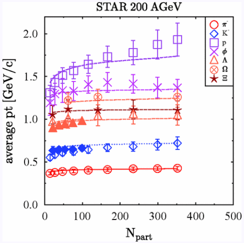
Core-Corona approach has been used to explain data from different experiments

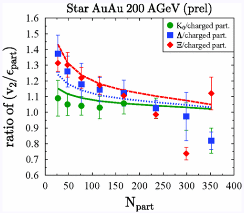
- Centrality Dependence of Strangeness Enhancement in Ultrarelativistic Heavy Ion Collisions - a Core-Corona Effect[5].
- Is the centrality dependence of the elliptic flow v_2 and of the average $\langle p_T \rangle$ more than a Core-Corona Effect?[6]

Observable $O(N_{part})$ as a function of centrality

$$O(N_{part}) = f_{core} O_{core} + (1 - f_{core}) O_{corona}$$

$core$ - central collisions
$corona$ - pp collisions
f_{core} - fraction of core nucleons

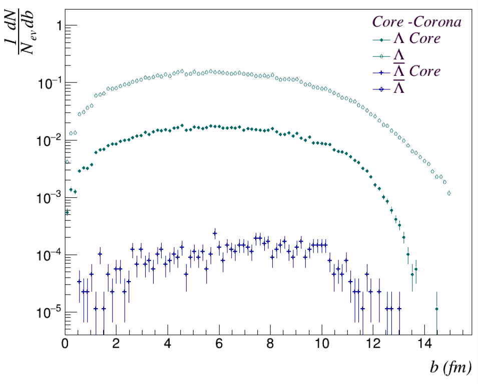




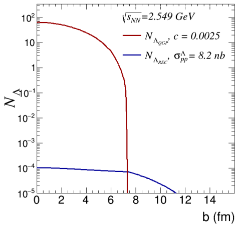
[5] Phys.Lett.B 673(2009)19-23
[6] Phys.Rev.C 82(2010)034906

Core-Corona dN/db

From pure MC ~ 89k events



Λ abundances in different regions differs from critical density - Glauber calculation



However we can use it to estimate the contribution to Global polarization

Implementation in MPD

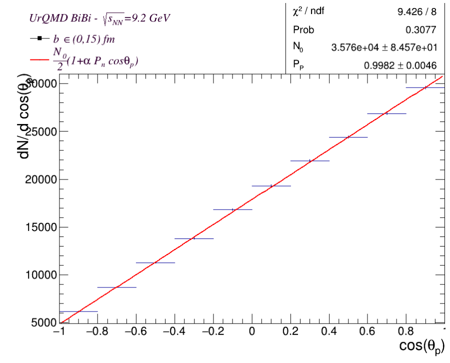
As a first attempt assign an arbitrary local polarization → 40% only to corona particles.

- Class MpdUrQMDGenerator → modified to read parent process type.
- Assign a fixed local polarization value to Λs in the corona in the $\hat{n}$ direction
- Transfer polarization to decay particles by the same procedure developed to Hyperon Global Polarization transfer in PHSD (PWG2 - E. Nazarova, and V. Voronyuk, [7])
- Λ reconstruction and measurement of Hyperon Global Polarization with MCTracks.

[7] https://indico.jinr.ru/event/3202/contributions/17250/attachments/12925/21604/Nazarova_26_07_22.pdf

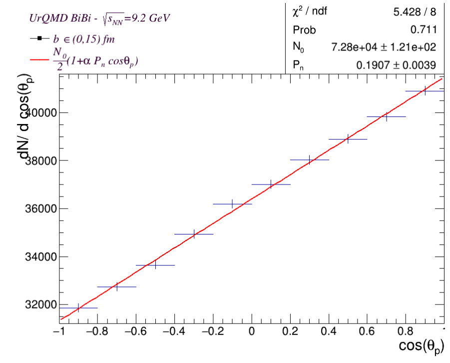
Polarization transfer

All Λs in the corona are polarized, projection along its own $\mathcal{P}$ is consistent with 100%

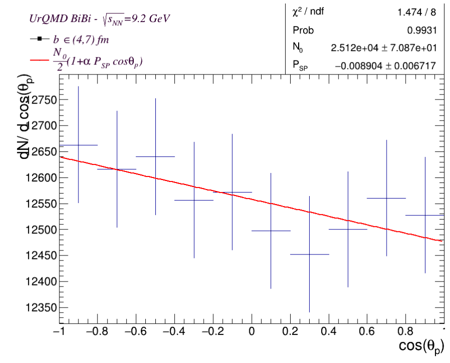


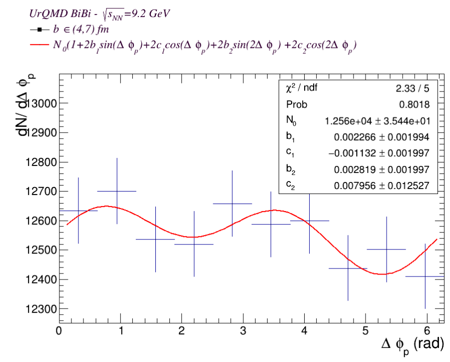
sample ~ 36k events

Local Polarization measured in all the Λ sample, smaller than 40% due to core contribution



Polarization for Λ produced at $b \in (4,7) fm$





$$\mathcal{P} \sim -0.009$$

$$\mathcal{P} = -\frac{8}{\pi \alpha_\Lambda} \frac{b_1}{R_{SP}} \rightarrow \sim -0.010$$

Does local polarization contribute to global polarization?

Section 7

Summary

Summary

- The description of global polarization has been shown with the core-corona model, which describes the experimental data at energy ranges of the HADES, STAR and NICA experiments.
- It has been shown that local polarization could contribute to global polarization with a fraction of the value of transverse polarization measured in pp collisions
- UrQMD has been proposed to simulate both the hydrodynamic phase of the core and cascade transport of the corona and separate Λ contribution.
- Mpdroot has been used to simulate the decay and transport of polarization in the charged decay of Λ.
- It has been shown that local polarization could contribute to global polarization, however, the results are inconclusive, due to the size of the sample used and the uncertainties of the calculation.

Work in progress and future plans

- Increase the size of the sample to repeat measurements and verify results.
- Get polarization with reconstructed tracks.
- Contribute to the implementation of core separation to another measurements within MPD.