





# (ANTI-)STRANGENESS PRODUCTION IN HEAVY ION COLLISIONS

**Pierre Moreau** 

for the PHSD group

Strangeness in Quark Matter 2015, Russia



Introduction HIC BES Ingredients

Conclusion

### From NICA to LHC, passing by FAIR and RHIC...



- Explore the QCD phase diagram and properties of hadrons at high temperature or high baryon density
- Phase transition from hadronic to partonic matter
- Goal: Study the properties of strongly interacting matter under extreme conditions from a microscopic point of view

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Realization: dynamical many-body transport approach

- Explicit parton-parton interactions, explicit phase transiton from hadronic to partonic degrees of freedom
- Transport theory: off-shell transport equations in phase-space representation based on Kadanoff-Baym equations for the partonic and hadronic phase



#### **Parton-Hadron-String-Dynamics (PHSD)**

W.Cassing, E.Bratkovskaya, PRC 78 (2008) 034919; NPA831 (2009) 215; W.Cassing, EPJ ST 168 (2009) 3

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Dynamical Quasi-Particle Model (DQPM)

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The QGP phase is described in terms of interacting quasiparticles with Lorentzian spectral functions: 

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$$\rho_i(\omega,T) = \frac{4\omega\Gamma_i(T)}{(\omega^2 - \mathbf{p}^2 - M_i^2(T))^2 + 4\omega^2\Gamma_i^2(T)} \qquad (i = q, \bar{q}, g)$$

Properties of quasiparticles are fitted to the lattice QCD results:



PHSU

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Masses and widths of partons depend on the temperature of the medium

Peshier, Cassing, PRL 94 (2005) 172301; Cassing, NPA 791 (2007) 365: NPA 793 (2007)



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Conclusion

### Stages of a collision in PHSD



- String formation in primary NN collisions
- String decays to pre-hadrons (baryons and mesons)



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## Stages of a collision in PHSD



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- Formation of a QGP state if  $\epsilon > \epsilon_c = 0.5 \text{ GeV.fm}^{-3}$
- Dissolution of new produced secondary hadrons into massive colored quarks and mean-field energy

 $B 
ightarrow qqq~(\bar{q}\bar{q}\bar{q}),~m
ightarrow q\bar{q}~+~U_q$ 

- DQPM define the properties (masses and widths) of partons  $m_q(\epsilon)$   $\Gamma_q(\epsilon)$
- ... and mean-field potential at a given local energy density  $\boldsymbol{\epsilon}$ 
  - $U_q(\epsilon)$

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## Stages of a collision in PHSD



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- Propagation of partons, considered as dynamical quasiparticles, in a self-generated mean-field potential from the DQPM
- EoS of partonic phase: ,crossover' from Lattice QCD fitted by DQPM



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### Stages of a collision in PHSD



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#### Stages of a collision in PHSD



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### Stages of a collision in PHSD



#### Hadron-string interactions – off-shell HSD

 Elastic and inelastic collisions between baryons (B), mesons (m) and resonances (R)

Distribution of hadron collisions as a function of time and collisional energy:



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#### Stages of a collision in PHSD

t = 0.1 fm/c



Conclusion

#### Stages of a collision in PHSD

t = 1.63549 fm/c



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Conclusion

#### Stages of a collision in PHSD

t = 2.06543 fm/c







Conclusion

#### Stages of a collision in PHSD

t = 3.20258 fm/c







Conclusion

#### Stages of a collision in PHSD

t = 5.56921 fm/c







Conclusion

#### Stages of a collision in PHSD

t = 8.06922 fm/c





Au + Au  $\sqrt{s_{NN}}$  = 200 GeV b = 2.2 fm – Section view

- Baryons (559)
- Antibaryons (139)
- Mesons (2686)
- Quarks (2628)
- Gluons (442)



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#### Stages of a collision in PHSD

t = 10.5692 fm/c







Conclusion

#### Stages of a collision in PHSD







Conclusion

#### Stages of a collision in PHSD





#### Partonic energy fraction in central A+A

- At SPS, only a small part of the initial energy is converted into the QGP phase
- At top RHIC energies, the QGP phase at midrapidity contains roughly 90% of the energy

Time evolution of the partonic energy fraction for different energies:





#### Transverse mass spectra (PHSD – HSD)

- With the HSD model, the high-pT spectra is not described properly especially at high energies where the parton energy fraction is major
- At low SPS energies, the difference is less visible since the partonic phase is not predominant



Transverse mass spectra for pions and kaons at different energies:

#### Central Pb+Pb – SPS energies

Central Au+Au – RHIC

W. Cassing & E. Bratkovskaya, NPA 831 (2009) 215; E. Bratkovskaya, W. Cassing, V. Konchakovski, O. Linnyk, NPA856 (2011) 162

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#### Au-Au at Top RHIC energies

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- At high energies, particles and antiparticles are produced in quasi-equal quantities at midrapidity whatever the centrality of the collision
- Anti-baryon absorption at low pT is visible



Conclusion

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#### Au-Au at BES @ RHIC energies

At low energies, a clear difference appears between the production of particles and antiparticles, and also between positively and negatively charged mesons



Production at midrapidity dN/dy:



Conclusion

#### Rapidity spectra

- At high energies, the hadrons produced at midrapidity come mostly from the QGP phase
- At high rapidity, particles are more produced than antiparticles due to the high baryon density
- At low energies, the stopping of initial nucleons induces a high baryon density even at midrapidty which favors the production of baryons compared to antibaryons



#### Rapidity spectras:

Conclusion

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#### Beam energy scan study

#### Production at midrapidity as a function of the collisional energy:

- By decreasing the collisional energy, the more the composition of produced particles is conditioned by the composition of the initial state
- At the highest energies, the composition of produced particles at midrapidity is conditionned by the plasma composition



Conclusion

#### Strange baryon production

D Multi-step production of multi-strange baryons  $\Xi$  and Ω:

Hyperon (Y =  $\Lambda$ ,  $\Sigma$ ) production

 $\pi + N \leftrightarrow K + Y$  $N + N \leftrightarrow N + Y + K$  $N + \bar{K} \leftrightarrow Y + \pi$ 



Strange baryon production

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**D** Multi-step production of multi-strange baryons  $\Xi$  and Ω:

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- At low energies, the production of multistrange baryons comes mostly from the hadronic processes
- QGP phase plays an important role in the production of anti-multi-strange baryons



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Ingredients for multi-strange baryon production

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Implementation of charge-channel decomposition of strangeness exchange reactions

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5 :  $d\sigma/d\Omega$  ( $\overline{K^0}$ + $\Sigma^+ \rightarrow \pi^+$ + $\Xi^0$ ) [mb.sr<sup>-1</sup>]

1.0

0.0

- Taken from the coupled channel approach based on a SU(3)-invariant hadronic Lagrangian from
- C.H. Li & C.M. Ko, Nucl.Phys. A712 (2002) 110-130

[GeV] v<sub>0</sub>s-√s



 $\bar{K} + \Lambda / \Sigma \leftrightarrow \Xi + \pi$ 



Conclusion





2.5

2.0

1.5 1.0 0.5

(Anti-)strangeness production in heavy ion collisions

#### Rapidity spectra - Channel decomposition

- □ (Anti-)strange baryons are dominantly produced by the QGP phase at mid-rapidity
- Production of strange baryons through hadronic processes are favoured at high rapidity regions where the fragments of initial nuclei are situated
- Even at high energies, the production of Ξ through strange exchange reactions is about 30% of the total production



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#### Ingredients for multi-strange baryon production

- Implementation of hyperon-hyperon reactions
- Taken from the coupled channel approach based on a SU(3)-invariant hadronic Lagrangian from
- F. Li, L. Chen, C.M. Ko, and S.H. Lee,



Born diagrams for the reactions  $\Lambda\Lambda \to N\Xi$ ,  $\Lambda\Sigma \to N\Xi$ , and  $\Sigma\Sigma \to N\Xi$ 



Note: this cross-section is not considered in the results presented here

By Alessia Palmese

#### Phys. Rev. C 85, 064902



#### Conclusion

- By decreasing the collisionnal energy, more differences appear between the production of particle and antiparticle, even at midrapidity regions
- Partonic and hadronic processes are both necessary to approach the yield of multi-strange baryons (work in progress)
- Cross sections from the DQPM at finite chemical potential may also play a significant role at low collisional energy



Signatures from a chiral phase transition are under study



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