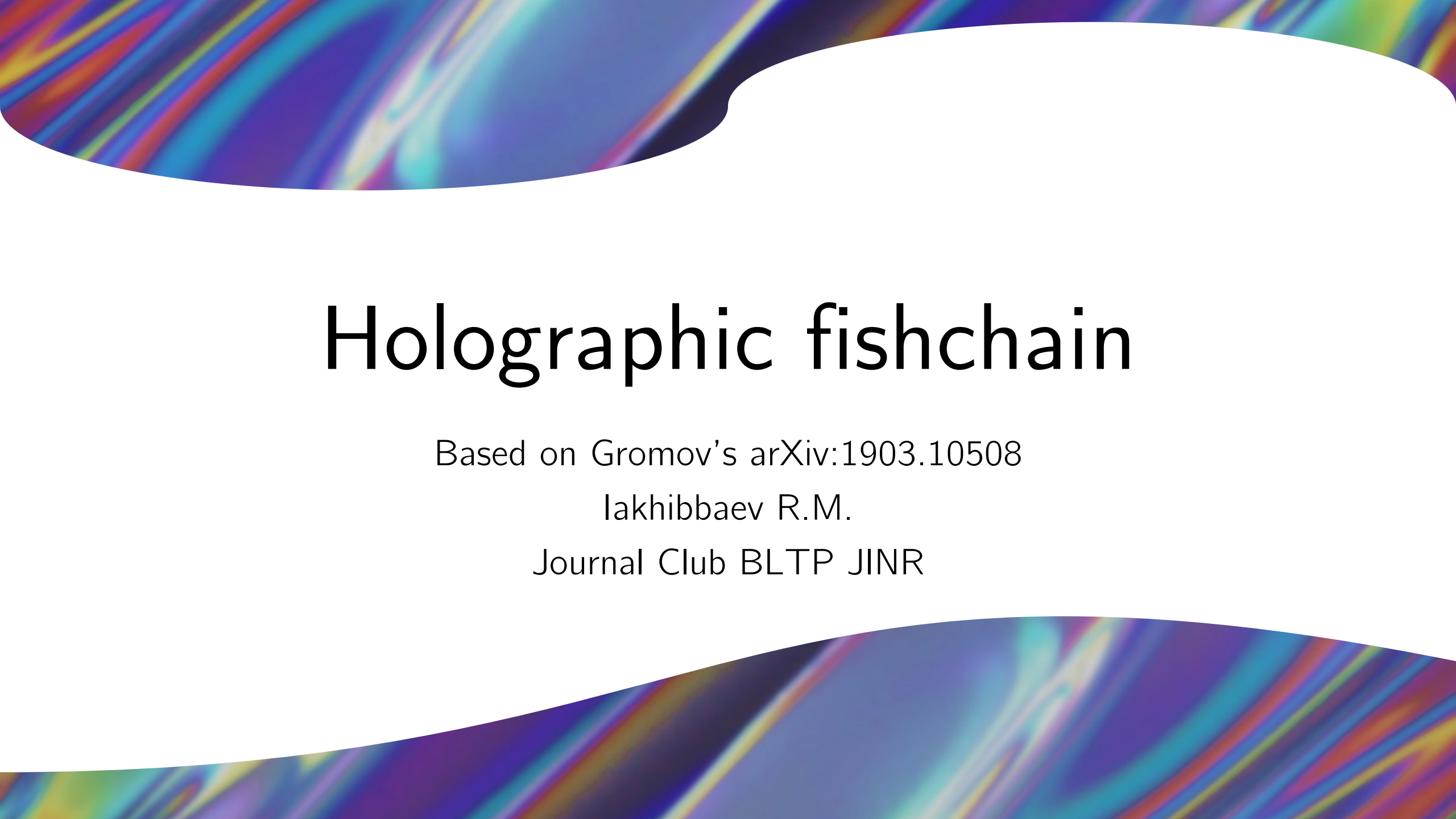




Holographic fishchain

Based on Gromov's [arXiv:1903.10508](https://arxiv.org/abs/1903.10508)

Iakhibbaev R.M.

Journal Club BLTP JINR

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1. Introduction and motivation
2. Fishnet biscalar model
3. From fishnet to fishchain
4. Test
5. Integrability
6. Final remarks (quantization and so on)



From N=4 SYM to fishnet

- N=4 SYM considered to be integrable
- How to “simplify” the theory and study particular integrable parts of it? Twisting!

$$\text{Tr}(\Phi_1[\Phi_2, \Phi_3]) \rightarrow \text{Tr}(\Phi_1\Phi_2\Phi_3e^{i\pi\beta} - \Phi_1\Phi_3\Phi_2e^{-i\pi\beta})$$

$$\mathcal{L} = N_c \text{Tr} \left[-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} D^\mu \phi_i^\dagger D_\mu \phi^i + i \bar{\psi}_{\dot{\alpha} A} D^{\dot{\alpha} \alpha} \psi_\alpha^A \right] + \mathcal{L}_{\text{int}}$$

$$\begin{aligned} \mathcal{L}_{\text{int}} = N_c g \text{Tr} & \left[\frac{g}{4} \{ \phi_i^\dagger, \phi^i \} \{ \phi_j^\dagger, \phi^j \} - g e^{-i\epsilon^{ijk} \gamma_k} \phi_i^\dagger \phi_j^\dagger \phi^i \phi^j \right. \\ & - e^{-\frac{i}{2} \gamma_j^-} \bar{\psi}_j \phi^j \bar{\psi}_4 + e^{+\frac{i}{2} \gamma_j^-} \bar{\psi}_4 \phi^j \bar{\psi}_j + i \epsilon_{ijk} e^{\frac{i}{2} \epsilon_{jkm} \gamma_m^+} \psi^k \phi^i \psi^j \\ & \left. - e^{+\frac{i}{2} \gamma_j^-} \psi_4 \phi_j^\dagger \psi_j + e^{-\frac{i}{2} \gamma_j^-} \psi_j \phi_j^\dagger \psi_4 + i \epsilon^{ijk} e^{\frac{i}{2} \epsilon_{jkm} \gamma_m^+} \bar{\psi}_k \phi_i^\dagger \bar{\psi}_j \right] \end{aligned}$$

From N=4 SYM to fishnets

- More simplifications:

DS-limit: strong twist + weak coupling

$$\mathcal{L}_{\text{int}} = N_c \text{Tr} \left[\xi_1^2 \phi_2^\dagger \phi_3^\dagger \phi^2 \phi^3 + \xi_2^2 \phi_3^\dagger \phi_1^\dagger \phi^3 \phi^1 + \xi_3^2 \phi_1^\dagger \phi_2^\dagger \phi^1 \phi^2 + i\sqrt{\xi_2 \xi_3} (\psi^3 \phi^1 \psi^2 + \bar{\psi}_3 \phi_1^\dagger \bar{\psi}_2) \right. \\ \left. + i\sqrt{\xi_1 \xi_3} (\psi^1 \phi^2 \psi^3 + \bar{\psi}_1 \phi_2^\dagger \bar{\psi}_3) + i\sqrt{\xi_1 \xi_2} (\psi^2 \phi^3 \psi^1 + \bar{\psi}_2 \phi_3^\dagger \bar{\psi}_1) \right].$$

- MORE simplifications:

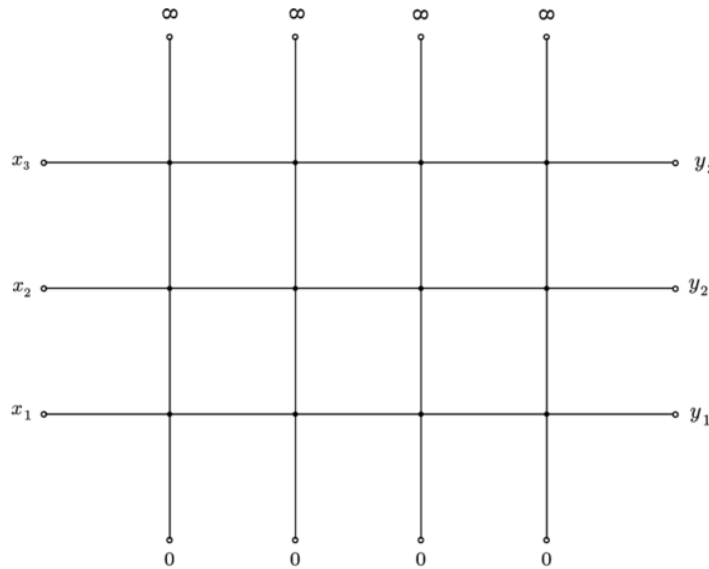
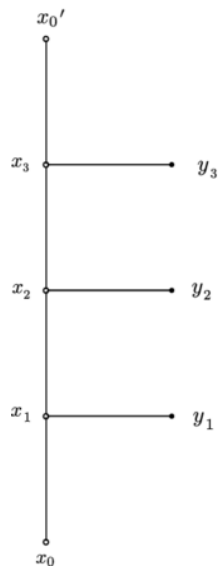
- $\xi_1 = \xi_2 = \xi_3$ beta-deformed SYM
- $\xi_1 = 0$ χ 0-CFT
- $\xi_1 = \xi_2 = 0, \xi_3 \equiv \xi \neq 0$ biscalar fishnet **(the simplest case)**

$$\mathcal{L}_{4d} = N \text{tr} \left(|\partial \phi_1|^2 + |\partial \phi_2|^2 + (4\pi)^2 \xi^2 \phi_1^\dagger \phi_2^\dagger \phi_1 \phi_2 \right)$$

Biscalar fishnet theory

- Conformal, non-unitary, solvable in the planar limit

$$\mathcal{L}_{4d} = N \operatorname{tr} \left(|\partial\phi_1|^2 + |\partial\phi_2|^2 + (4\pi)^2 \xi^2 \phi_1^\dagger \phi_2^\dagger \phi_1 \phi_2 \right)$$



- Feynman rules:

$$\langle \phi_1(x_i) \phi_1^\dagger(y_i) \rangle = \frac{1}{(2\pi)^2 (x_i - y_i)^2}$$

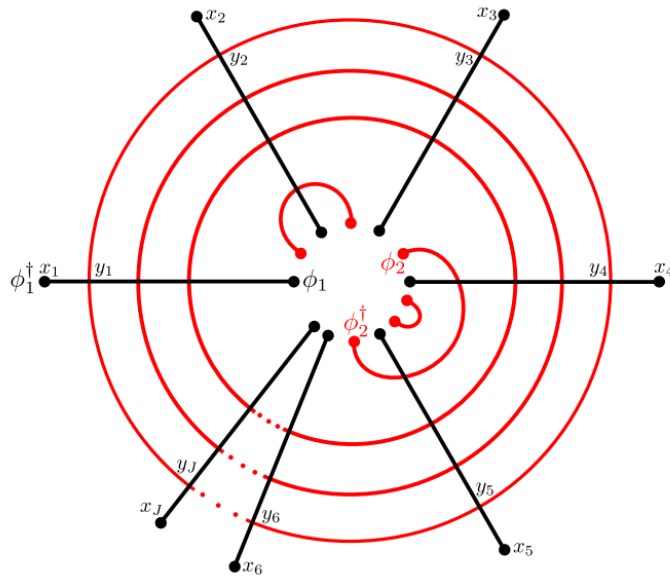
$$\langle \phi_2(y_i) \phi_2^\dagger(y_{i+1}) \rangle = \frac{1}{(2\pi)^2 (y_i - y_{i+1})^2}$$

$$(4\pi)^2 \xi^2 \phi_1^\dagger \phi_2^\dagger \phi_1 \phi_2$$

Biscalar fishnet theory

- Feynman diagram (“wheel-graph”) for local operator:

$$\mathcal{O}_J = \phi_1^J (\phi_2^\dagger \phi_2)^n$$



- Graph-building operator:

$$\hat{B}(\{\vec{y}_i\}_{i=1}^J, \{\vec{x}_j\}_{j=1}^J) = \prod_{i=1}^J \frac{(4\pi)^2 \xi^2}{(2\pi)^2 (\vec{y}_i - \vec{y}_{i+1})^2 (2\pi)^2 (\vec{x}_i - \vec{y}_i)^2}$$

- All graph summation:

$$1 + \hat{B} + \hat{B}^2 + \dots = \frac{1}{1 - \hat{B}}$$

- 2J-correlation function:

$$\langle x_1, \dots, x_J | \frac{1}{1 - \hat{B}} | y_1, \dots, y_J \rangle$$

Exact correlation function for J=2

- Conformal symmetry fixes form of eigenvectors (CFT wave function) of the graph-building operator

$$\Psi_{\Delta,S}(x_1, x_2) = \frac{1}{x_1^{\Delta-S} x_2^{\Delta-S} x_{12}^{S+2-\Delta}} (\text{tensor})^S$$

- Eigenvalue of the graph-building operator

$$\hat{B}_{\Delta,S} \Psi_{\Delta,S}(x_1, x_2) = E_{\Delta,S} \Psi_{\Delta,S}(x_1, x_2)$$

- So the correlation function can be represented as

$$\begin{aligned} & \langle x_1, \dots, x_J | \frac{1}{1 - \hat{B}_{\Delta,S}} | y_1, \dots, y_J \rangle = \\ & = \sum_S \int d\Delta \frac{1}{1 - E_{\Delta,S}} \langle x_1, \dots, x_J | \Psi_{\Delta,S} \rangle \langle \Psi_{\Delta,S} | y_1, \dots, y_J \rangle \end{aligned}$$

- Poles given by:

$$E_{\Delta,S} = 1$$

Exact correlation function for J=2

- 4 pt correlation function can be written in the next form

$$\mathcal{G}(x_1, x_2; y_1, y_2) = \sum_{\Delta, S} \underset{\substack{\uparrow \\ \text{structure constant}}}{C_{\Delta, S}} \overset{\text{conformal block}}{g_{\Delta, S}(u, v)}$$

- Spectral equation and its solutions:

$$(\nu^2 + S^2/4)(\nu^2 + (S+2)^2/4) = \xi^4 \quad \longrightarrow \quad \begin{aligned} \Delta_2(S) &= 2 + \sqrt{(S+1)^2 + 1 - 2\sqrt{(S+1)^2 + 4\xi^4}} \\ \Delta_4(S) &= 2 + \sqrt{(S+1)^2 + 1 + 2\sqrt{(S+1)^2 + 4\xi^4}} \end{aligned}$$

- Classical limit:

Spectrum:

$$\Delta = \sqrt{S^2 \pm 4\xi^2}$$

4pt correlation function:

$$\mathcal{G} \sim e^{-2i\xi\sqrt{\theta^2 + \rho^2}}$$

$$\begin{aligned} u &= \frac{4}{(\cos \theta - \cosh \rho)^2} \\ v &= \frac{(\cos \theta + \cosh \rho)^2}{(\cos \theta - \cosh \rho)^2} \end{aligned}$$

Derivation of dual

- Wave functions corresponding to dilatation operator are at the residues:

$$(\hat{B} - 1)\Psi = 0$$

Derivation of dual

- Wave functions corresponding to dilatation operator are at the residues:

$$\begin{array}{ccc}
 (\hat{B} - 1)\Psi = 0 & & \\
 \vec{p}_i = -i \frac{\partial}{\partial \vec{x}_i} & \longrightarrow & \prod_i \vec{p}_i^2 \\
 \prod_{i=1}^J \frac{\xi^2/\pi^2}{(\vec{y}_i - \vec{y}_{i+1})^2 (\vec{x}_i - \vec{y}_i)^2} & \longrightarrow & \square_i \frac{1}{(2\pi)^2 (x_i - y_i)^2} = \delta(x_i - y_i)
 \end{array}$$

$$\mathcal{H} \circ \Psi(x_i) = 0 \quad \mathcal{H} = \prod_{i=0}^J p_i^2 - \prod_{i=0}^J \frac{4\xi^2}{(x_i - x_{i+1})^2}$$

Quantum system with time-reparametrization symmetry?

Symmetries, EOMs, constraints

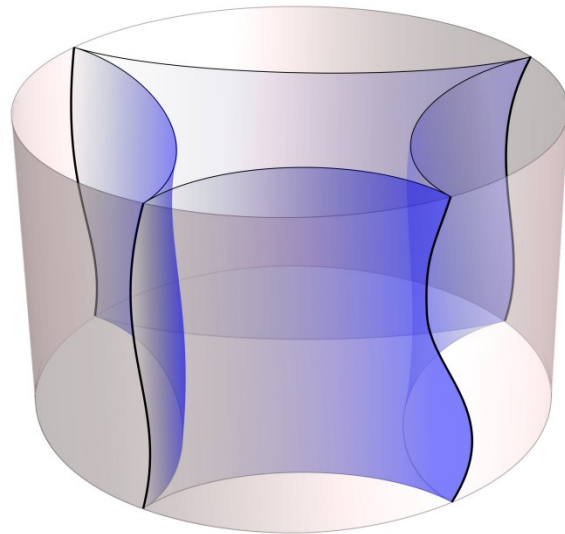
- Lagrangian at the classical level

$$L = \frac{2J-1}{2^{\frac{2J}{2J-1}}} \left(\frac{1}{\gamma} \prod_{i=1}^J \vec{x}_i^2 \right)^{\frac{1}{2J-1}} + \gamma \prod_{i=1}^J \frac{4\xi^2}{(\vec{x}_i - \vec{x}_{i+1})^2} \quad S = \xi \int L dt = 2J\xi \int \left(\prod_{i=1}^J \frac{\dot{\vec{x}}_i^2}{(\vec{x}_i - \vec{x}_{i+1})^2} \right)^{\frac{1}{2J}} dt$$

As it is well known the conformal group in 4D coincides with the group of rotations in $R(1,5)$, so we embed actions above to 6D space

$$x_i^\mu \rightarrow X_i^M / X_i^+, \quad X_i^2 = 0, \quad X^+ = X_i^0 + X_i^{-1}$$

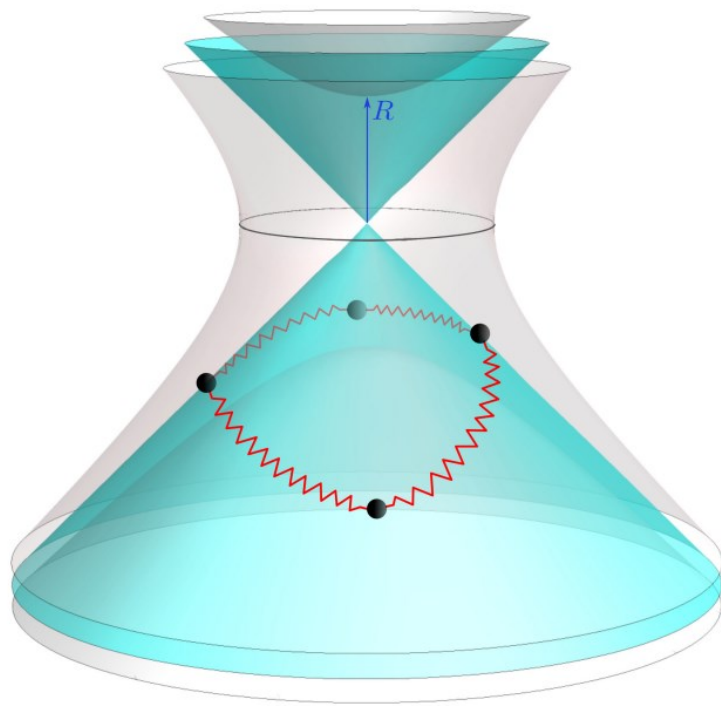
Fishnet action on projective lightcone:
$$L = 2J \left(\prod_{i=1}^J \frac{\dot{X}_i \cdot \dot{X}_i}{-2X_i \cdot X_{i+1}} \right)^{\frac{1}{2J}}$$



Symmetries of the action:

1. Conformal symmetry
2. Time-reparametrisation symmetry
3. Time-dependent scaling
4. Translation along the chain
5. Gauge symmetry

Symmetries, EOMs, constraints



Virasoro-like constraints

$$\dot{X}_k^2 = 2 \prod_i (-X_i \cdot X_{i+1}) \equiv \mathcal{L}$$

Polyakov-type action:

$$L = - \sum_i \left[\frac{\dot{X}_i^2}{2\alpha_i} + \eta_i X_i^2 + \gamma \prod_k (-X_k \cdot X_{k+1})^{-\frac{1}{J}} \right]$$

EOM

$$\ddot{X}_i = 2\eta_i X_i - \frac{\mathcal{L}}{2} \left(\frac{X_{i+1}}{X_{i+1} \cdot X_i} + \frac{X_{i-1}}{X_i \cdot X_{i-1}} \right)$$

$$q_i^{MN} = 2\dot{X}_i^{[M} X_i^{N]} \quad j_i^{MN} = 2 \frac{X_{i-1}^{[M} X_i^{N]}}{X_{i-1} \cdot X_i} \quad \dot{q}_i = \frac{\mathcal{L}}{2} (j_{i+1} - j_i) ,$$

Nodes or SO(1,5)-charge and current density

Test (J=2 classical solutions)

- Fishnet:

Classical spectrum:

$$\Delta^2 = S^2 \mp 4\xi^2$$

4pt correlation function:

$$\mathcal{G} \sim \exp \left(-2i\xi \sqrt{\rho^2 + \theta^2} \right)$$

$$u = \frac{4}{(\cos \theta - \cosh \rho)^2}$$

$$v = \frac{(\cos \theta + \cosh \rho)^2}{(\cos \theta - \cosh \rho)^2}$$

- Fishchain: define two 6D null-vectors

$$X_{1,2} = \frac{r}{\sqrt{2}} (\cosh s, \sinh s, \pm \cos \phi, \mp \sin \phi, 0, 0)$$

- Conserved charges:

$$S = \mp 4\xi\tau, \quad s = -i\tau\Delta/\xi,$$

$$\phi_1 = \tau S_1/\xi, \quad d\tau = \frac{dt}{r^2(t)}$$

Solution:

$$S^2 - \Delta^2 = \pm 4\xi^2$$

Integrability

- Zero curvature condition (like Toda chain):

$$\dot{\mathbb{L}}_i = \mathbb{V}_{i+1} \cdot \mathbb{L}_i - \mathbb{L}_i \cdot \mathbb{V}_i \qquad 0 = j_{i+1} q_i^2 - j_i q_i^2 + \text{EOM}$$

- a pair of spacelike and timelike connections (Lax pairs):

$$\mathbb{L}_i^{\mathbf{6}} = u^2 + u q_i + \frac{q_i^2}{2} \ , \qquad \mathbb{V}_i^{\mathbf{6}} = \frac{j_i}{u} \frac{\mathcal{L}}{2}$$

- Integral of motions will be given

$$\mathbb{T}(u) = \text{tr} (\mathbb{L}_1(u) \dots \mathbb{L}_J(u))$$

$$\dot{\mathbb{T}}(u) = 0$$

Final remarks and conclusion

Conclusion:

- The fishchain can be derived explicitly from fishnet model
- It's conformal, integrable. Also the model reminds string theory
- It is dual to fishnets and live on lightcone

Further development

- One can quantise the model. Quantum fishchain lives in AdS5!
- Fishchain theory is applicable on gamma-deformed N=4 SYM (at least on U(1)-sector with some magnons)
- Open fishchain(open fishchain without periodical boundary conditions)

Open problems

- Fishchain for N=4 SYM?
- Continuum limit?
- Wilson lines?

Thanks!

