#### Classical Patterns in Quantum Rainbow Channeling

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# Classical patterns in the quantum rainbow channeling of high energy electrons

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# Introduction

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# The Channeling Effect



Figure 1: (a) Scattering on atomic string. (a) Axial channel potential. (c) Schematic representation of the channeling effect.

Potential of the planar chennel (220) is  $V(x) = \sum_{m=-N}^{N} V_{220}^{\text{th}}(x - md_{220})$   $V_{220}^{\text{th}}(x) = -\frac{7e^2 a_{\text{s}}}{2\sqrt{2}\varepsilon_0 d_{220}^2} \sum_{l=1}^{3} \frac{\alpha_l}{\beta_l} e^{\frac{\beta_l^2 \sigma_{\text{th}}^2}{2a_{\text{s}}^2}}$  $\times \left\{ e^{-\beta_l \frac{|x|}{a_{\text{s}}}} \operatorname{erfc}\left(\frac{\beta_l \sigma_{\text{th}}}{\sqrt{2}a_{\text{s}}} - \frac{|x|}{\sqrt{2}\sigma_{\text{th}}}\right) + e^{\beta_l \frac{|x|}{a_{\text{s}}}} \operatorname{erfc}\left(\frac{\beta_l \sigma_{\text{th}}}{\sqrt{2}a_{\text{s}}} + \frac{|x|}{\sqrt{2}\sigma_{\text{th}}}\right) \right\}.$ 

The relativistic mass, critical angle and the dynamical thickness are

$$m_{\rm r} = \gamma m_{\rm e}, \ \Theta_{\rm c} = \sqrt{\frac{V(d_{220}/2)}{E_k} \frac{1 + \frac{E_k}{m_{\rm e}c^2}}{1 + \frac{E_k}{2m_{\rm e}c^2}}}, \ \Lambda = \frac{L}{2\pi} \sqrt{\frac{\partial_x^2 V(0)}{2E_k} \frac{1 + \frac{E_k}{m_{\rm e}c^2}}{1 + \frac{E_k}{2m_{\rm e}c^2}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}, \ \Lambda = \frac{L}{2\pi\sqrt{2\pi\sqrt{2}}} \sqrt{\frac{\partial_x^2 V(0)}{2E_k} \frac{1 + \frac{E_k}{m_{\rm e}c^2}}{1 + \frac{E_k}{2m_{\rm e}c^2}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}, \ \Lambda = \frac{L}{2\pi\sqrt{2\pi\sqrt{2}}} \sqrt{\frac{2\pi\sqrt{2}}{2E_k} \frac{1 + \frac{E_k}{2\pi\sqrt{2}}}{1 + \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}, \ \Lambda = \frac{L}{2\pi\sqrt{2\pi\sqrt{2}}} \sqrt{\frac{2\pi\sqrt{2}}{2E_k} \frac{1 + \frac{E_k}{2\pi\sqrt{2}}}{1 + \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}, \ \Lambda = \frac{L}{2\pi\sqrt{2\pi\sqrt{2}}} \sqrt{\frac{2\pi\sqrt{2}}{2E_k} \frac{1 + \frac{E_k}{2\pi\sqrt{2}}}{1 + \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}, \ \Lambda = \frac{L}{2\pi\sqrt{2\pi\sqrt{2}}} \sqrt{\frac{2\pi\sqrt{2}}{2E_k} \frac{1 + \frac{E_k}{2\pi\sqrt{2}}}{1 + \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}, \ \Lambda = \frac{L}{2\pi\sqrt{2\pi\sqrt{2}}} \sqrt{\frac{2\pi\sqrt{2}}{2E_k} \frac{1 + \frac{E_k}{2\pi\sqrt{2}}}{1 + \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}, \ \Lambda = \frac{L}{2\pi\sqrt{2}} \sqrt{\frac{2\pi\sqrt{2}}{2E_k} \frac{1 + \frac{E_k}{2\pi\sqrt{2}}}{1 + \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}, \ \Lambda = \frac{L}{2\pi\sqrt{2}} \sqrt{\frac{2\pi\sqrt{2}}{2E_k} \frac{1 + \frac{E_k}{2\pi\sqrt{2}}}}{1 + \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}, \ \Lambda = \frac{L}{2\pi\sqrt{2}} \sqrt{\frac{2\pi\sqrt{2}}{2E_k} \frac{1 + \frac{E_k}{2\pi\sqrt{2}}}}{1 + \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}_{+ \frac{E_k}{2\pi\sqrt{2}}}}_{+ \frac{$$

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# The Classical Channeling of Electrons

# Hamilton's equations are solved by Runge-Kutta method of 4th order

$$\frac{\mathrm{d}\theta_x}{\mathrm{d}t} = -\frac{\partial_x V(x)}{m_\mathrm{r} k_z}, \quad \Rightarrow \theta_x^{(n+1)} = \theta_x^{(n)} + \sum_s c_s F_s \left(\theta_x^{(n)}, -\frac{\partial_x V(x^{(n)})}{m_\mathrm{r} k_z}\right),$$

$$\frac{\mathrm{d}x}{\mathrm{d}x} = h \theta_s \quad \Rightarrow \sigma_s^{(n+1)} = \sigma_s^{(n)} + \sum_s \sigma_s F_s \left(\sigma_s^{(n)}, -\frac{\partial_x V(x^{(n)})}{m_\mathrm{r} k_z}\right)$$

$$\frac{\mathrm{d}x}{\mathrm{d}t} = k_z \theta_x, \qquad \rightarrow x^{(n+1)} = x^{(n)} + \sum_s c_s F_s\left(x^{(n)}, k_z \theta_x^{(n)}\right),$$

with the initial conditions  $x(0) = x^{(0)} = b$ , and  $\theta_x(0) = \theta_x^{(0)} = 0$ .





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# The Quantum Channeling of Electrons

Channeling of quantum particles is governed by the Schrödinger equation

$$i\hbar\partial_t\psi(x,t) = -\frac{\hbar^2}{2m_{\rm r}}\partial_x^2\psi(x,t) + V(x)\psi(x,t). \label{eq:phi}$$

Propagator  $\hat{U}(t)=\exp[-i\hat{H}t/\hbar]$  is expanded in the basis of the Chebyshev polynomials  $T_n$ 

$$\psi(t) = \hat{U}(t)\psi_0 \approx e^{-\frac{i}{\hbar}\bar{E}t} \sum_{n=0}^N a_n(\alpha)T_n\left(-\frac{i}{\hbar}\hat{\mathcal{H}}\right)\psi_0,$$
$$a_n(\alpha) = (2-\delta_{n0})i^n J_n(\alpha), \quad \alpha = \frac{\Delta Et}{2\hbar}, \quad \hat{\mathcal{H}} = \frac{2}{\Delta E}(\hat{H}-\bar{E}).$$

where  $\Delta E=E_M-E_m,\,\bar{E}=(E_M+E_m)/2.$  At the boundaries of computational domain absorptive potential is used

$$V_{\rm abs}(x) = \begin{cases} i \frac{V_0}{\cosh^2[\alpha(|x| - x_{\rm max})]}, & |x| \ge x_{\rm max}, \\ 0, & |x| < x_{\rm max}. \end{cases}$$

Main iteration of the method becomes

$$\psi_n = 2\hat{\mathcal{H}}\psi_{n-1} - \psi_{n-2}, \rightarrow \psi_n = e^{-\gamma} \left( 2\hat{\mathcal{H}}\psi_{n-1} - e^{-\gamma}\psi_{n-2} \right),$$

where  $\gamma(x) = -iV_{\rm abs}(x)$ .

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# The Electron's Self-Interference

The initial electron state is represented by a Gaussian wave packet

$$\psi_b(x) = \sqrt{\mathcal{N}\left(b, \frac{1}{2k_z\Omega}\right)}$$

Electron state in the angular representation is given by the integral

$$\varphi_b(\theta_x) = \sqrt{\frac{k_z}{2\pi}} \int \psi_b(x) e^{-ik_z \theta_x x} \,\mathrm{d}x, \quad \varphi_b \sim \mathrm{FFTW}\left[\psi_b\right].$$



Figure 3: Evolution of the 255-MeV quantum state having  $b = d_{220}/3$  in the (a) spatial and (b) angular representations. The dashed lines show boundaries of (220) channels. The dot-dashed lines show the corresponding classical trajectories.

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# The Probability Densities of Electron Beam

The probability densities of ensemble of noninteracting particles

$$\rho_x(x) = \frac{1}{d_{220}} \int_{-d_{220}/2}^{d_{220}/2} |\psi_b(x)|^2 \,\mathrm{d}b, \quad \rho_{\theta_x}(\theta_x) = \frac{1}{d_{220}} \int_{-d_{220}/2}^{d_{220}/2} |\varphi_b(x)|^2 \,\mathrm{d}b.$$

# reflects a classical caustic pattern



Figure 4: The probability density of the 255-MeV electron beam for  $\Omega = 0.25\Theta_c$  in (a) spatial and (b) angular representations. Full and dashed lines show classical caustics and boundaries of (220) channels, respectively.

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# The Probability Densities of Electron Beam

or canonical diffraction patterns, depending on the value of the angular divergence of the electron beam  $\Omega.$ 





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# Channeling in the Phase Space

Wigner function in the  $(x, \theta_x)$  phase space is defined by the following integral

$$W(x,\theta_x) = \frac{k_z}{2\pi} \int \psi_b^{\dagger} \left(x - \frac{\xi}{2}\right) \psi_b \left(x + \frac{\xi}{2}\right) e^{-ik_z \theta_x \xi} d\xi,$$
$$W \sim \text{CZT} \left[\psi_b^{\dagger} \psi_b, -k_z \Theta_c, k_z \Theta_c\right].$$





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# Conslusions

- The classical dynamics electrons can be efficiently simulated using Runge-Kutta method of the 4th order.
- Distributions of the channeled beam can obtained by easily parallelized Monte-Carlo simulation.
- The quantum dynamics can be obtained accurately using Chebyshev method of global propagation.
- Efficiency of the calculation can be increased by the Task parallelism where the master thread executes the Chebyshev iteration, while slaves accumulate results to get time evolution of the quantum state.
- The dynamics in the phase space can be obtained efficiently if Wigner transform is implemented using CZT transform.
- The results show that quantum-classical transition can emerge on the level of the ensemble without the need for wave packets to transform into mass points.



Pafnuty Chebyshev (4. May 1821 - 26. December 1894).



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