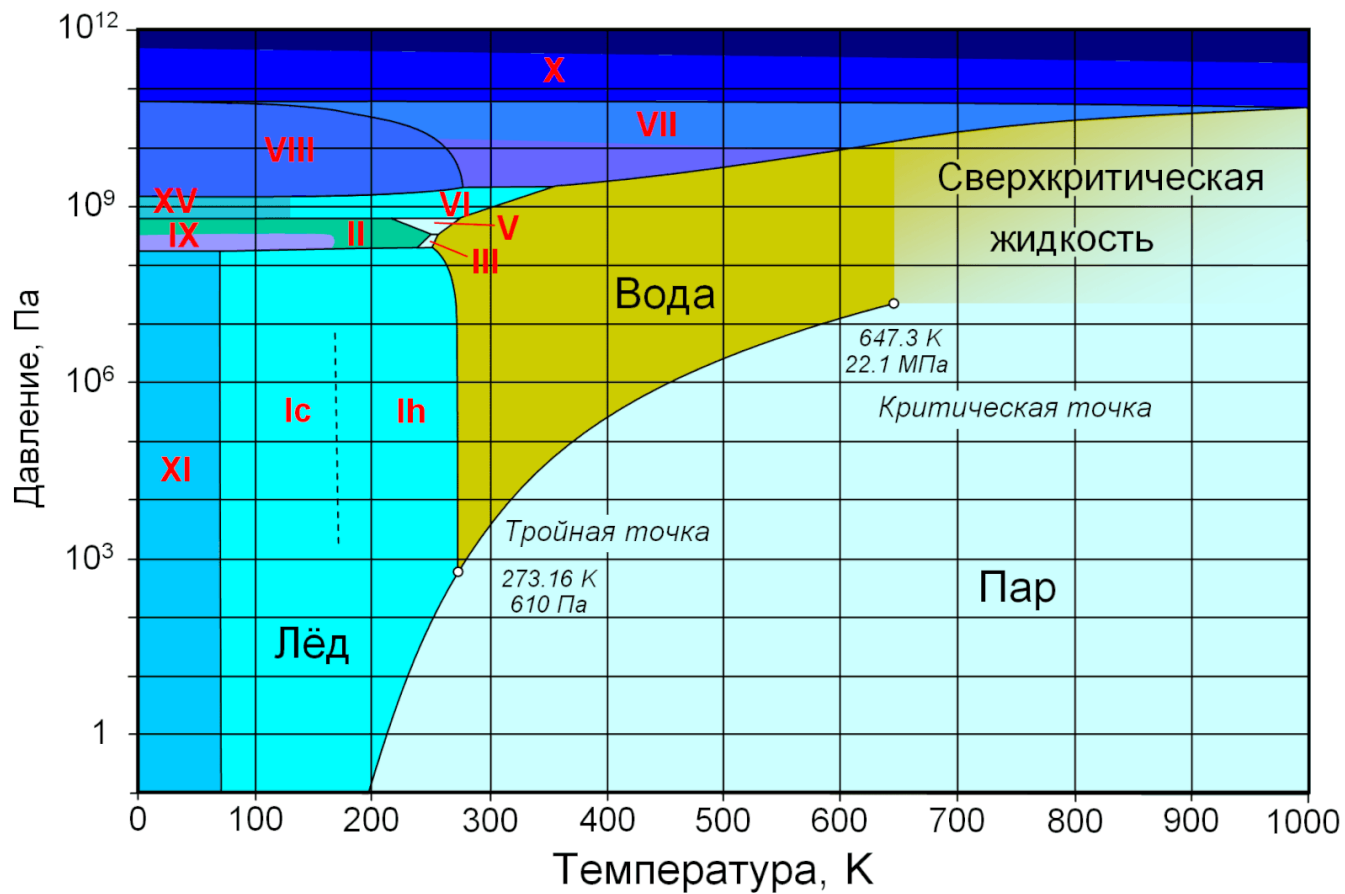


Quantum Phase Transition in the One-Dimensional Water Chain

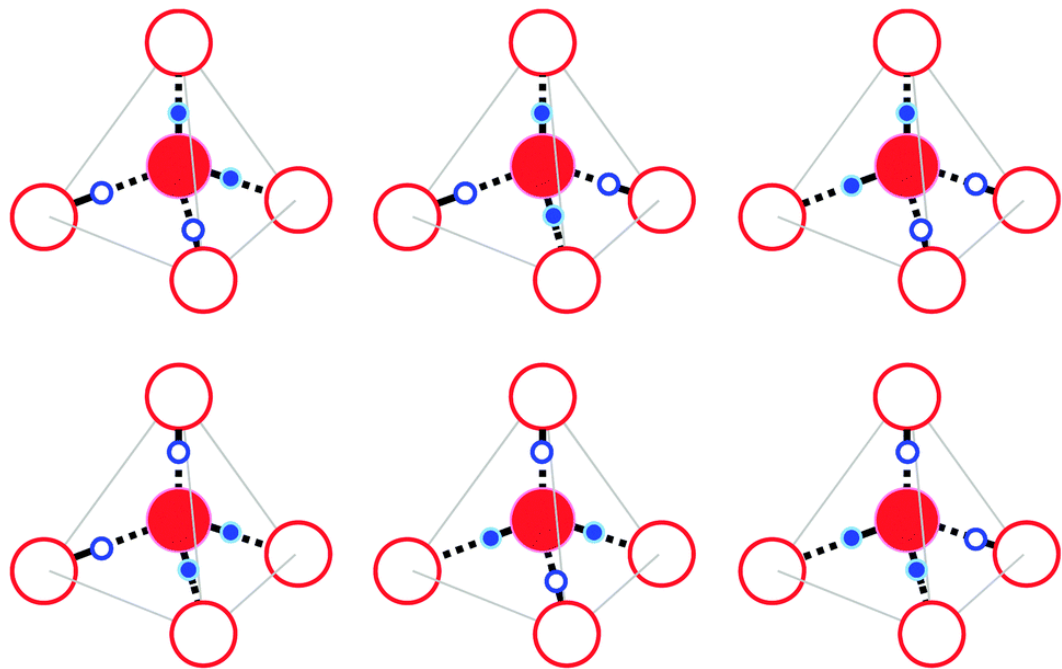
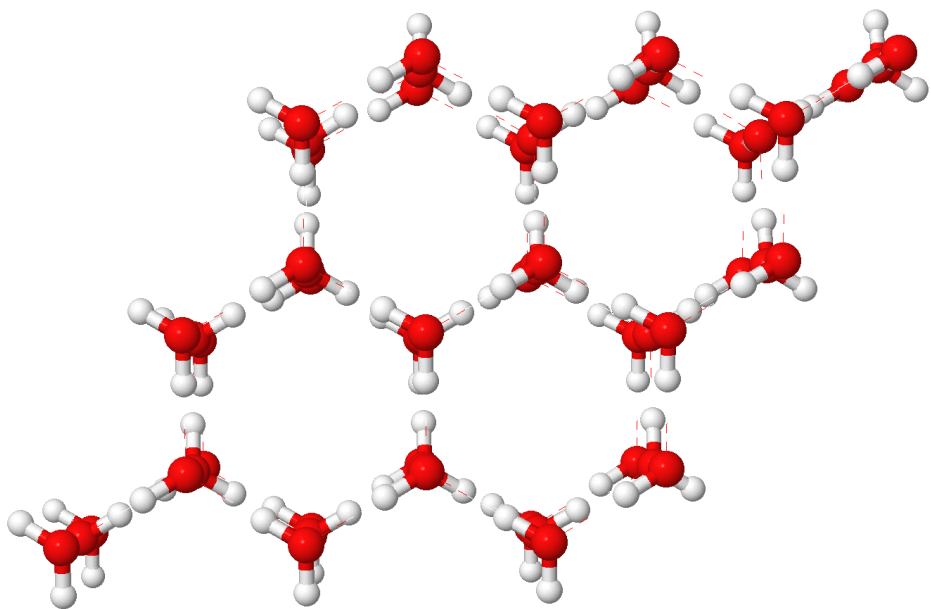
*T. Serwatka, R.G. Melko, A. Burkov, and P.-N. Roy,
Phys. Rev. Lett. 130, 026201 (2023)*



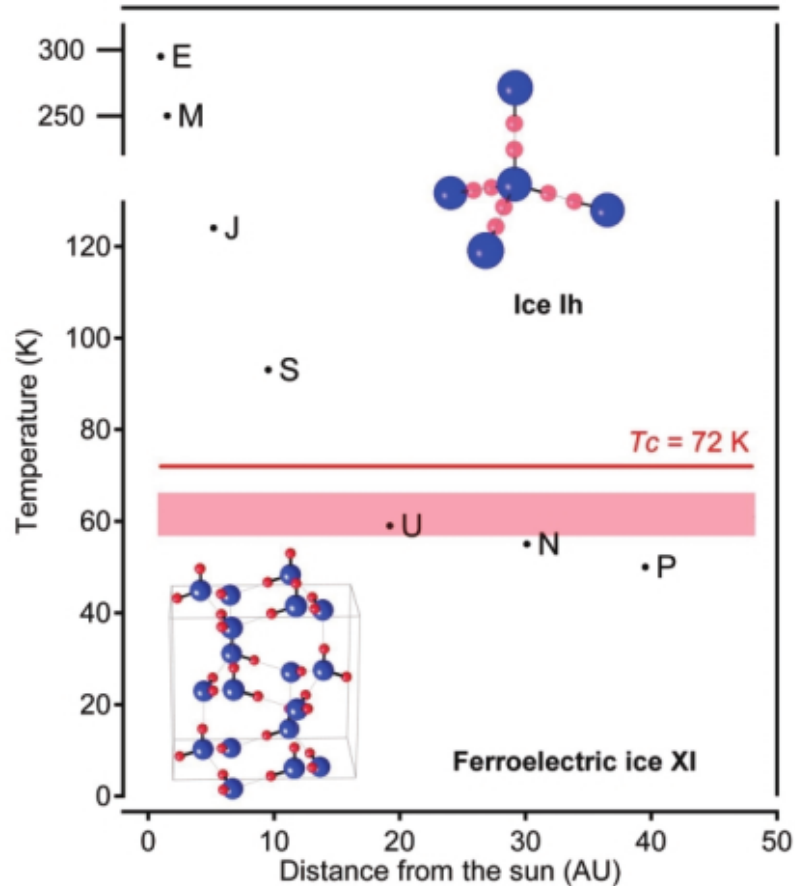
Ice



Hexagonal ice



Ferroelectric ice

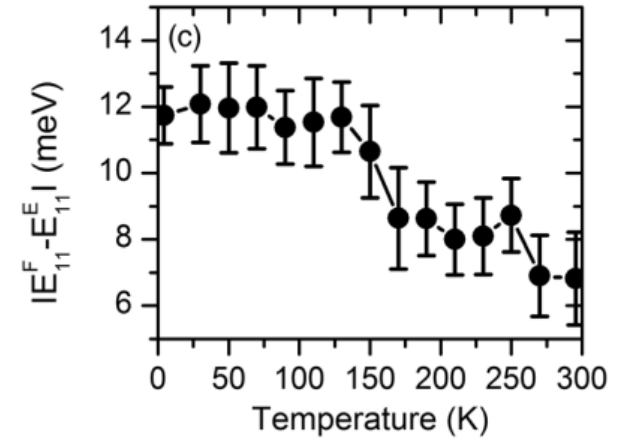
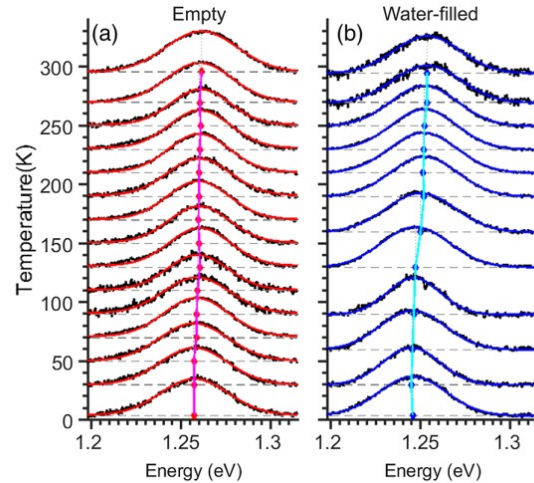
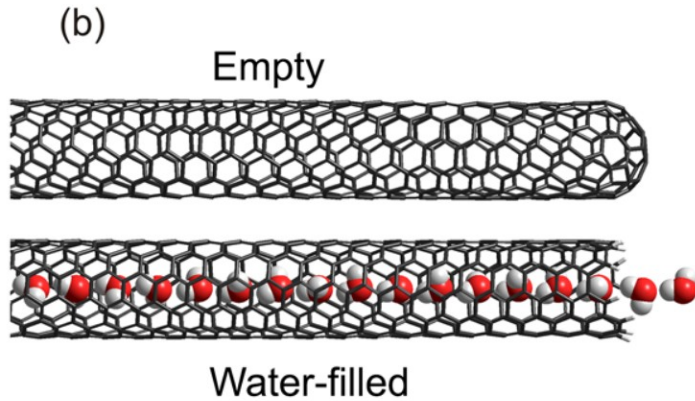


EXISTENCE OF FERROELECTRIC ICE IN THE UNIVERSE

H. FUKAZAWA,¹ A. HOSHIKAWA,¹ Y. ISHII,¹ B. C. CHAKOUMAKOS,² AND J. A. FERNANDEZ-BACA²

Received 2006 August 11; accepted 2006 October 9; published 2006 November 3

Water in porous channels



Quasiphase Transition in a Single File of Water Molecules Encapsulated in (6,5) Carbon Nanotubes Observed by Temperature-Dependent Photoluminescence Spectroscopy

Xuedan Ma,^{1,*} Sofie Cambré,^{2,†} Wim Wenseleers,² Stephen K. Doorn,¹ and Han Htoon¹

¹Center for Integrated Nanotechnologies, Materials Physics and Applications Division, Los Alamos National Laboratory, New Mexico 87545, USA

²Experimental Condensed Matter Physics Laboratory, University of Antwerp, B-2610 Antwerp, Belgium

(Received 1 May 2016; published 12 January 2017)

Water in porous channels

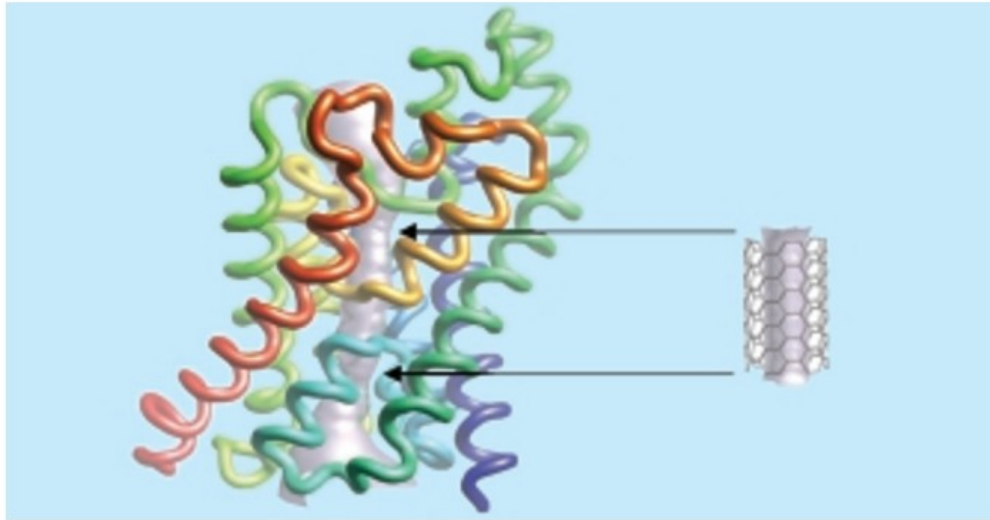
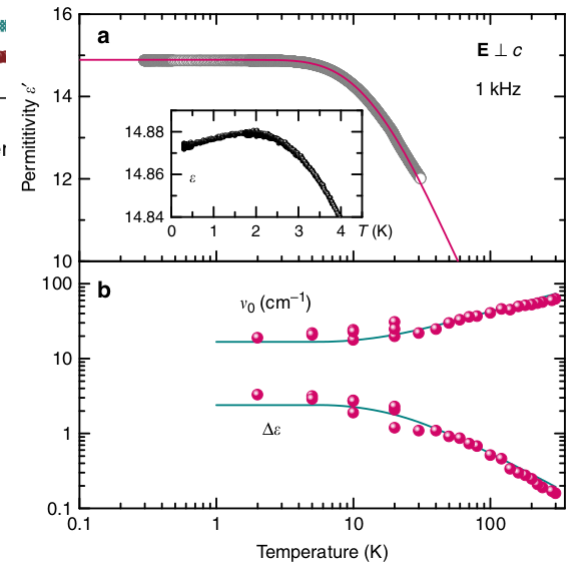
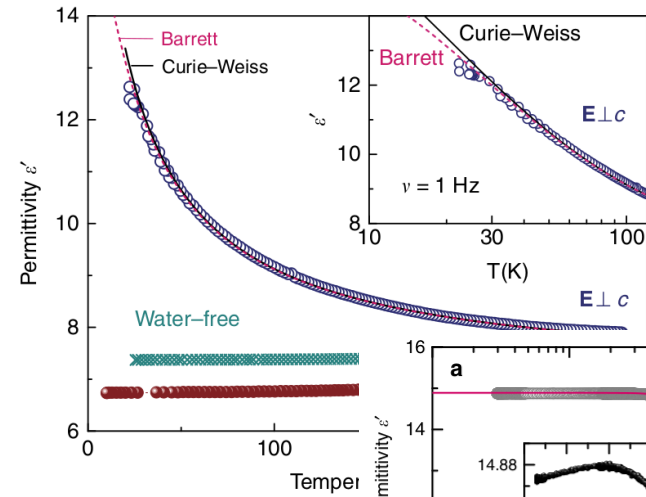
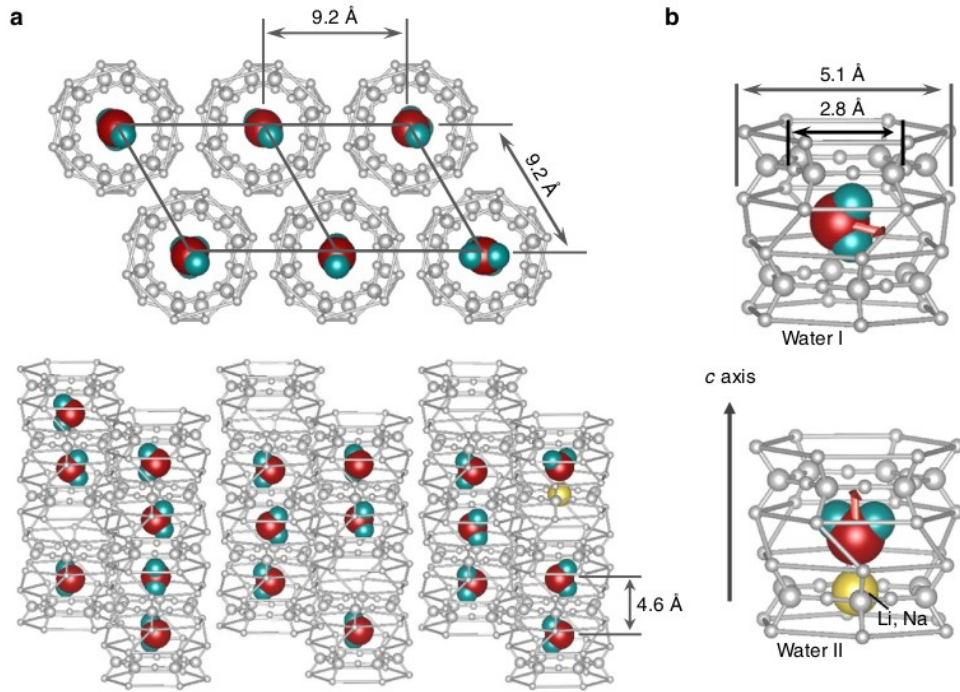


Figure 1 Pore comparisons. Left, a biological pore (grey), as found in the bacterial water and glycerol channel, GlpF. The pore is formed by the side chains of the protein backbone (seen here as coloured spaghetti). Right, a similar pore in a narrow carbon nanotube, such as that for which Hummer *et al.*¹ carried out simulations of water conduction. It is hoped that nanotubes can be used in simulations

Water in porous channels



ARTICLE

Received 9 Mar 2016 | Accepted 8 Aug 2016 | Published 30 Sep 2016

DOI: 10.1038/ncomms12842

OPEN

Incipient ferroelectricity of water molecules confined to nano-channels of beryl

B.P. Gorshunov^{1,2,3}, V.I. Torgashev⁴, E.S. Zhukova^{1,2,3}, V.G. Thomas⁵, M.A. Belyanchikov¹, C. Kadlec⁶, F. Kadlec⁶, M. Savinov⁶, T. Ostapchuk⁶, J. Petzelt⁶, J. Prokeška⁷, P.V. Tomas^{8,9}, E.V. Pestrjakov¹⁰, D.A. Fursenko⁵, G.S. Shakurov¹¹, A.S. Prokhorov^{1,2}, V.S. Gorelik¹², L.S. Kadyrov¹, V.V. Uskov¹, R.K. Kremer¹³ & M. Dressel³

Water in porous channels

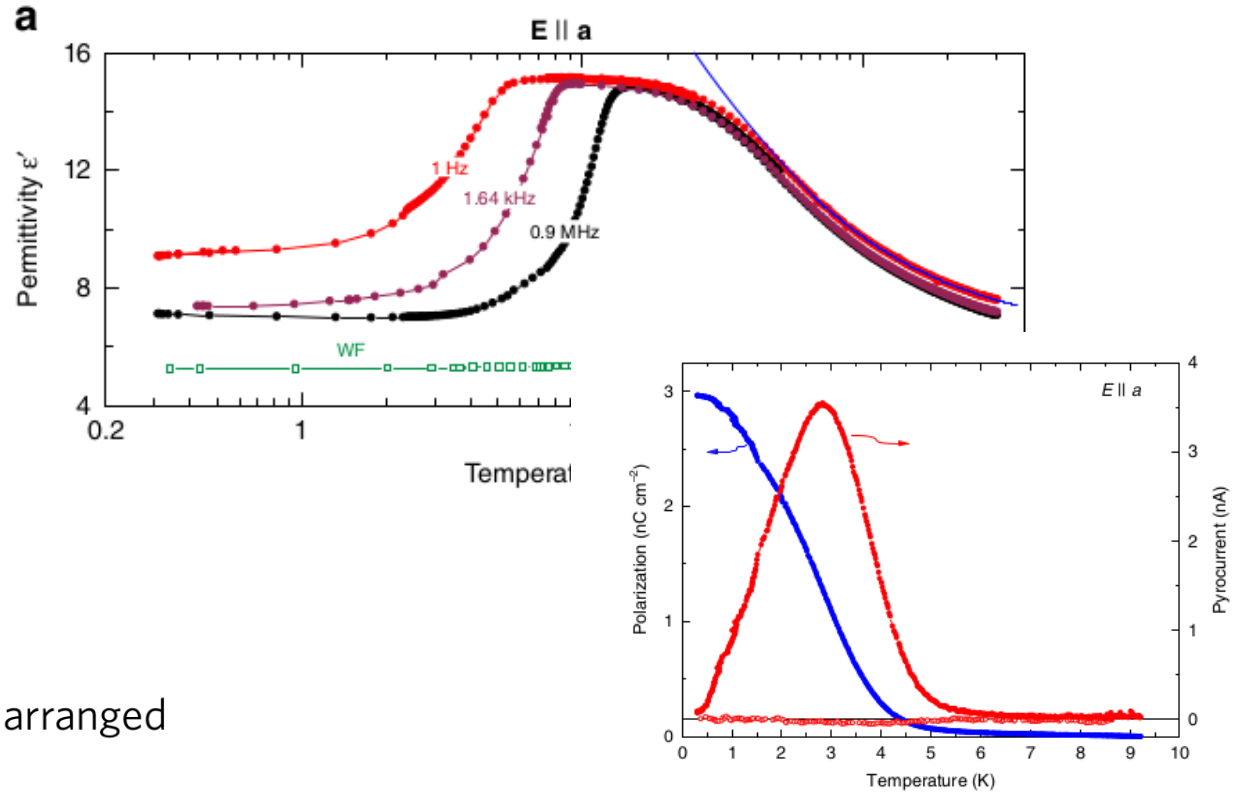
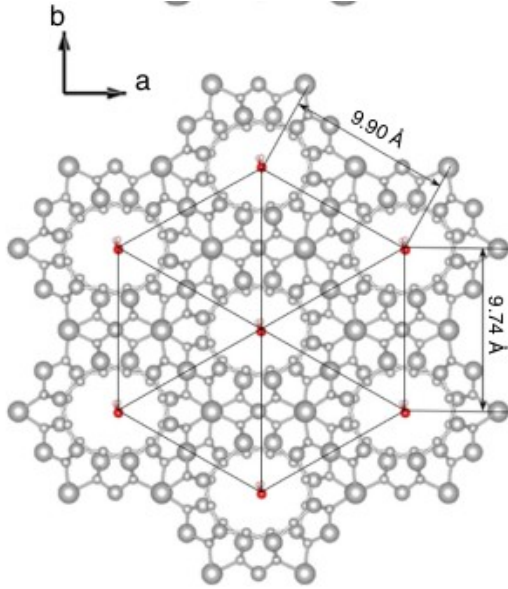


Fig. 4 Pyrocurrent and polarization indicate a phase transition at $T_0 = 3$ K. Temperature dependences of pyrocurrent (red) measured for $\mathbf{E} \parallel \mathbf{a}$ polarization while heating in zero electric field after cooling hydrous cordierite crystal in external electric field 8 kV cm^{-1} . The temperature-dependent polarization (blue) was calculated from pyrocurrent. Open dots correspond to the pyrocurrent measured during heating in zero electric field

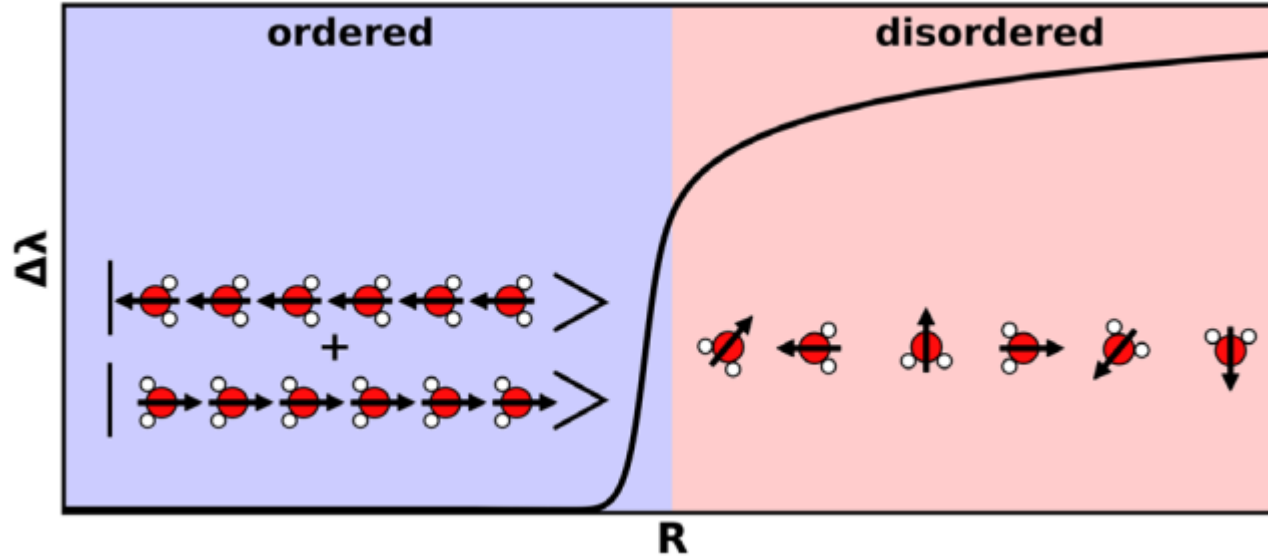
<https://doi.org/10.1038/s41467-020-17832-y>

OPEN

Dielectric ordering of water molecules arranged in a dipolar lattice

M. A. Belyanchikov^{1,2}, M. Savinov², Z. V. Bedran¹, P. Bednyakov², P. Proschek³, J. Prokleska³, V. A. Abalmasov⁴, J. Petzelt², E. S. Zhukova¹, V. G. Thomas^{5,6}, A. Dudka⁷, A. Zhugayevych⁸, A. S. Prokhorov^{1,9}, V. B. Anzin^{1,9}, R. K. Kremer¹⁰, J. K. H. Fischer^{11,12}, P. Lunkenheimer¹¹, A. Loidl¹¹, E. Uykur¹³, M. Dressel^{1,13} & B. Gorshunov^{1,2}

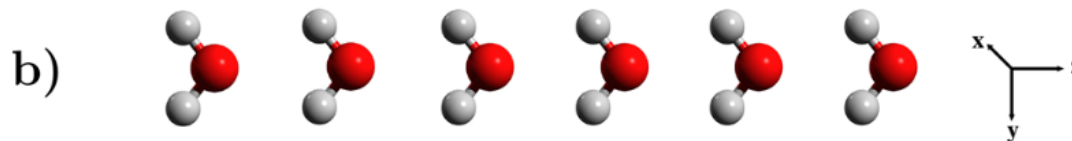
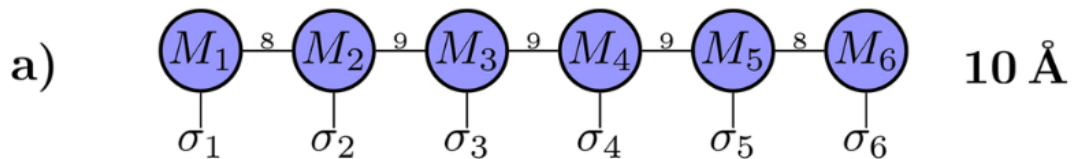
Phase transition simulation



Computational details

$$|\Psi\rangle = \sum_{\sigma_1, \dots, \sigma_N} C_{\sigma_1, \dots, \sigma_N} |\sigma_1 \dots \sigma_N\rangle,$$

$$\begin{aligned} C_{\sigma_1, \dots, \sigma_N} &= \sum_{\{\alpha\}} M_{1\alpha_1}^{(1),\sigma_1} M_{\alpha_1\alpha_2}^{(2),\sigma_2} \dots M_{\alpha_{N-1}1}^{(N),\sigma_N} \\ &= \mathbf{M}^{(1),\sigma_1} \mathbf{M}^{(2),\sigma_2} \dots \mathbf{M}^{(N),\sigma_N}, \end{aligned}$$



Schmidt values

The Schmidt spectrum $\{\lambda_l\}$ of the system's quantum state $|\Psi\rangle$ is obtained by a Schmidt decomposition:

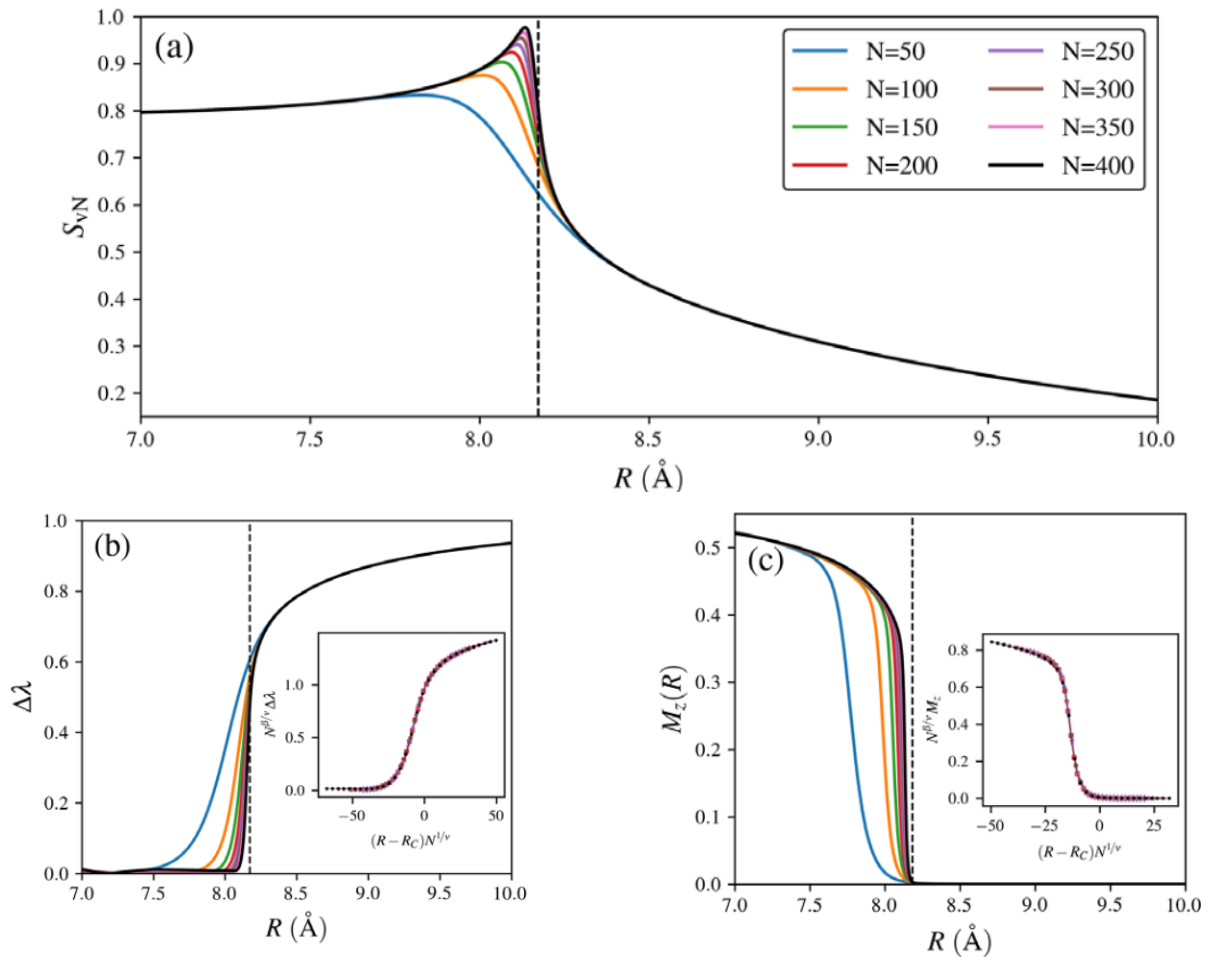
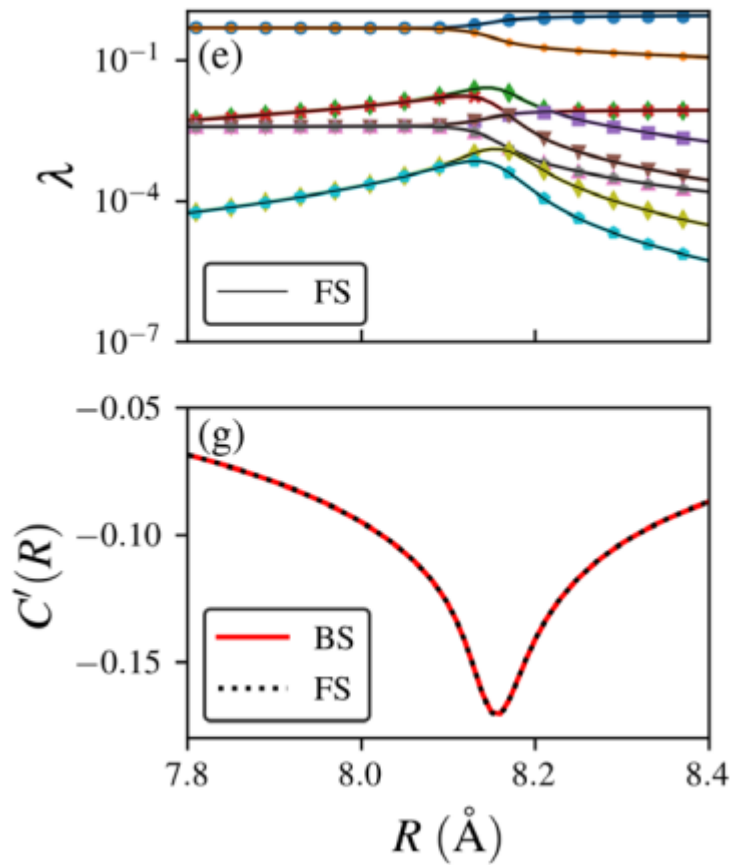
$$|\Psi\rangle = \sum_l \sqrt{\lambda_l} |\phi_A^l\rangle \otimes |\phi_B^l\rangle \quad (2)$$

where the complete Hilbert space is decomposed in a bipartite tensor product $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$ with orthonormal sets $\{|\phi_A^l\rangle\}$ and $\{|\phi_B^l\rangle\}$ for region A and B, respectively. The Schmidt coefficients are nothing else but the eigenvalues of the reduced density matrices ρ_A (and ρ_B) for the chosen partition and are directly connected to the von-Neumann entanglement entropy,

$$S_{\text{vN}}(\rho_A) = -\text{tr}(\rho_A \log \rho_A) \quad (3)$$

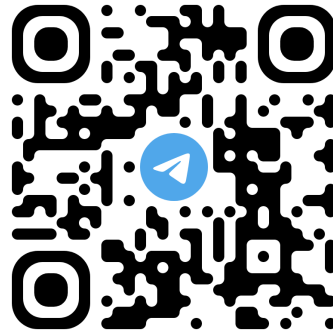
$$= -\sum_l \lambda_l \ln(\lambda_l). \quad (4)$$

Quantum phase transition



Вступайте в чат молодых ученых ЛТФ! /Join
the BLTP young scientists discussion!

Telegram link:



Группа для обсуждений и новостей, касающихся молодежи в
ЛТФ/News, discussions and giveaways

-Павел Максимов, представитель ЛТФ в
Объединении молодых ученых и специалистов
ОИЯИ