

STANDARD MODEL

RESULTS IN PP COLLISIONS

BY ATLAS & CMS



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STANDARD MODEL

❑ The SM describes the fundamental constituents of matter and their interactions

❑ The SM of particle physics is the theory describing **3 of the 4 known fundamental forces (electromagnetic, weak and strong interactions** — *excluding gravity*) in the universe and classifying all known elementary particles

❖ Quantum chromodynamics sector

❖ Electroweak sector

❖ Higgs sector







❖ Yukawa sector

$$\mathcal{L}_{\text{QCD}} = \sum_i \bar{\psi}_i \left(i\gamma^\mu (\partial_\mu \delta_{ij} - ig_s G_\mu^a T_{ij}^a) \right) \psi_j - \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a$$


$$\mathcal{L}_{\text{EW}} = \sum_\psi \bar{\psi} \gamma^\mu \left(i\partial_\mu - g' \frac{1}{2} Y_W B_\mu - g \frac{1}{2} \vec{\tau}_L \vec{W}_\mu \right) \psi - \frac{1}{4} W_{\mu\nu}^a W_{\mu\nu}^a - \frac{1}{4} B_{\mu\nu} B_{\mu\nu}$$


$$\mathcal{L}_H = \left| \left(\partial_\mu + \frac{i}{2} (g' Y_W B_\mu + g \vec{\tau} \vec{W}_\mu) \right) \varphi \right|^2 - \frac{\lambda^2}{4} (\varphi^\dagger \varphi - v^2)^2$$







$$\mathcal{L}_{\text{Yukawa}} = \bar{U}_L G_u U_R \varphi^0 - \bar{D}_L G_u U_R \varphi^- + \bar{U}_L G_d D_R \varphi^+ + \bar{D}_L G_d D_R \varphi^0 + \text{h. c.},$$


UP mass 2,3 MeV/c ² charge 2/3 spin 1/2 	CHARM 1,275 GeV/c ² 2/3 1/2 	TOP 173,07 GeV/c ² 2/3 1/2 
DOWN 4,8 MeV/c ² -1/3 1/2 	STRANGE 95 MeV/c ² -1/3 1/2 	BOTTOM 4,18 GeV/c ² -1/3 1/2 


GLUON
0
0
1


HIGGS BOSON
126 GeV/c ²
0
0
0


PHOTON
0
0
1


ELECTRON 0,511 MeV/c ² -1 1/2 	MUON 105,7 MeV/c ² -1 1/2 	TAU 1,777 GeV/c ² -1 1/2 
ELECTRON NEUTRINO <2,2 eV/c ² 0 1/2 	MUON NEUTRINO <0,17 MeV/c ² 0 1/2 	TAU NEUTRINO <15,5 MeV/c ² 0 1/2 

Z BOSON
91,2 GeV/c ²
0
0
1


W BOSON
80,4 GeV/c ²
±1
1


Interaction		Weak	Electromagnetic	Strong	
Property	Gravitational	Electroweak		Fundamental	Residual
Acts on:	Mass - Energy	Flavor	Electric charge	Color charge	Atomic nuclei
Particles experiencing:	All particles	quarks, lepton s	Electrically charged	Quarks, Gluons	Hadrons
Particles mediating:	Graviton (Not yet observed)	W ⁺ , W ⁻ and Z ⁰	γ (photon)	Gluons	Mesons
Strength at the scale of quarks:	10 ⁻⁴¹ (predicted)	10 ⁻⁴	1	60	Not applicable to quarks
Strength at the scale of protons/neutrons:	10 ⁻³⁶ (predicted)	10 ⁻⁷	1	Not applicable to hadrons	20

QUANTUM CHROMODYNAMICS (QCD)

QCD: Quantum field theory of strong interactions (C.N. Yang, R. Mills; H. Fritzsch, M. Gell-Mann, H. Leutwyler)

➤ **Interaction carried by gluons acting on quarks and gluons**

➤ **QCD-charge: colour of three types (red, blue, green)**

❖ **QCD coupling strength α_s depends on energy**

- **high energy** (= short distance or time): α_s is **small** (asymptotic freedom): **perturbative regime of QCD**
- **low energy** (= long distance or time): α_s is **large** (**confinement**): **non-perturbative regime of QCD**

❖ **The Renormalization Group Equation (RGE)**

$$\alpha_s(\mu^2) = \frac{\alpha_s(\mu_0^2)}{1 + b_0 \alpha_s(\mu_0^2) \ln \frac{\mu^2}{\mu_0^2}} = \frac{1}{b_0 \ln \frac{\mu^2}{\Lambda^2}}$$

$$b_0 = \frac{11C_A - 2n_f}{12\pi}$$

n_f : number of light quarks;

C_A : QCD colour factors

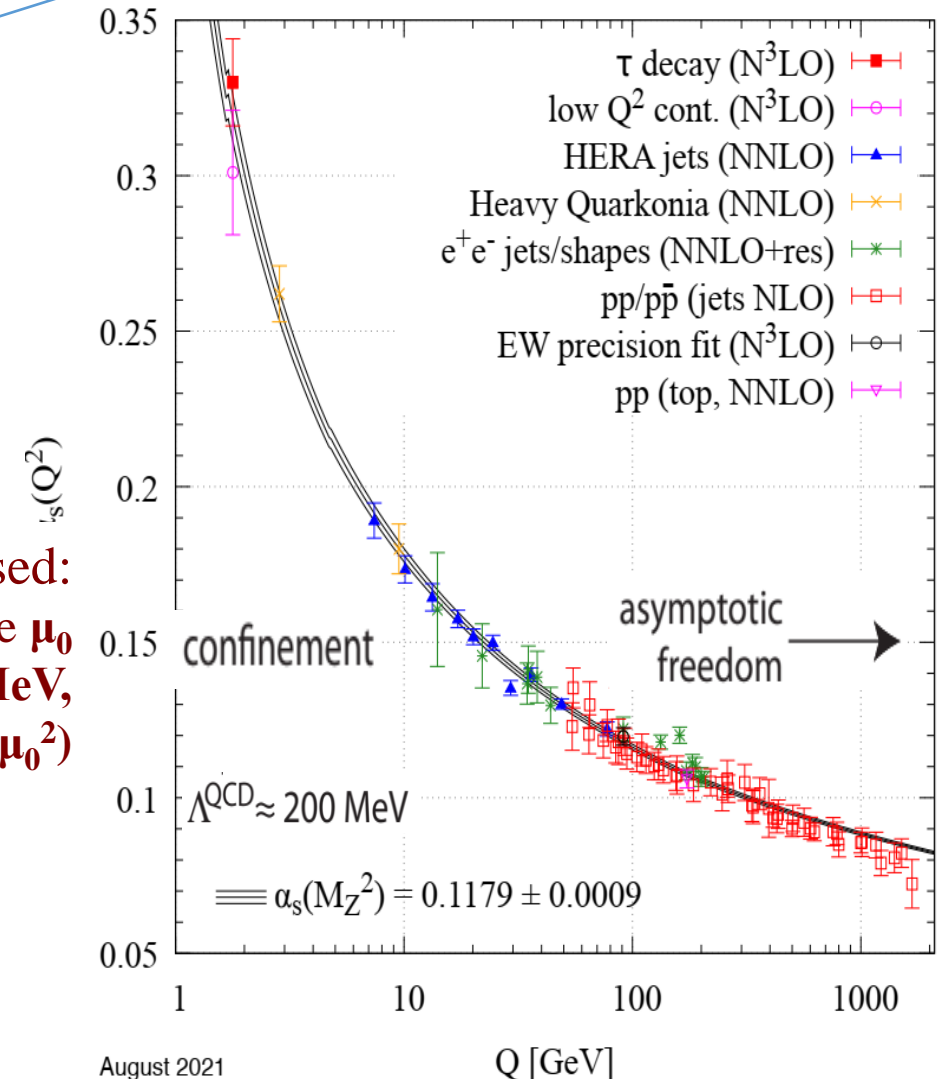
□ $\alpha_s(\mu^2)$ can be expressed:

- in terms of coupling $\alpha_s(\mu_0^2)$ at reference scale μ_0
- or by introducing non-perturbative constant $\Lambda \approx 200 \text{ MeV}$, corresponding to the divergence (Landau pole) of $\alpha_s(\mu_0^2)$

❖ **Behaviour of $\alpha_s(\mu^2)$ as a function of energy Q ($= 1/R$) determines the properties of QCD and the dynamics of quarks and gluons**

- **Large Q (small distance R): α_s small:: QCD weakly interacting; quarks and gluons asymptotically free; **regime of perturbative QCD****
- **Small $Q \approx \Lambda$ (large distance R): α_s large:: QCD strongly interacting, quarks and gluons form colour-less bound states: baryons (qqq) and mesons (qq) and do not exist as free particles, **regime of non-perturbative QCD****

$$f = f_0 + \alpha_s f_1 + \alpha_s^2 f_2 + \alpha_s^3 f_3 + \dots$$

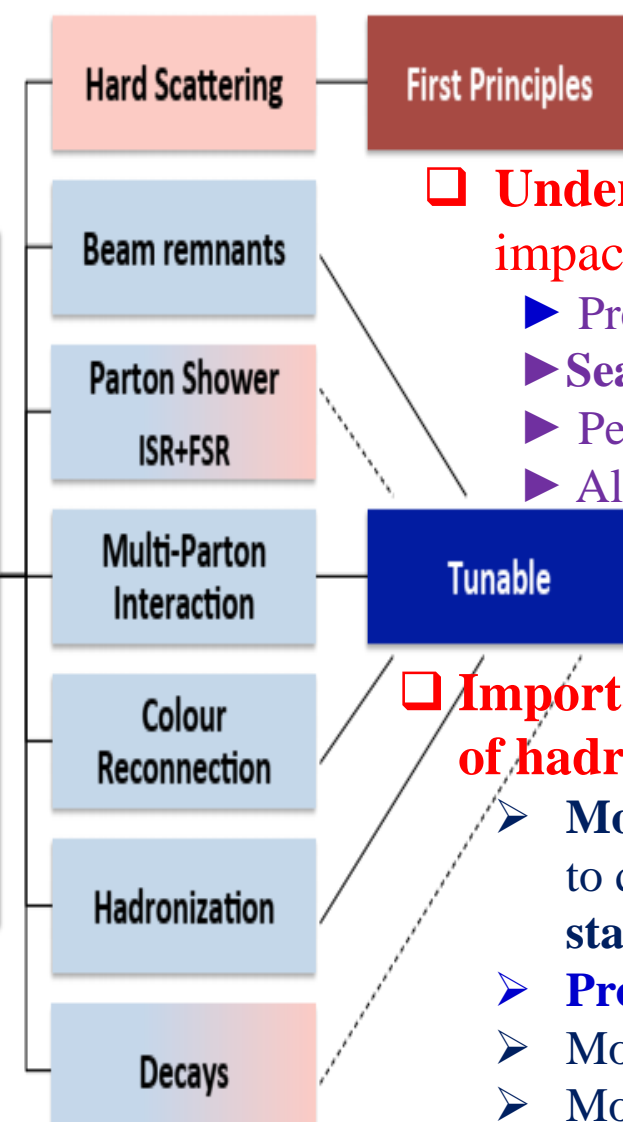


August 2021

Summary of measurements of α_s as a function of the energy scale Q

QCD MEASUREMENTS AT THE LHC

Monte Carlo Event Generator

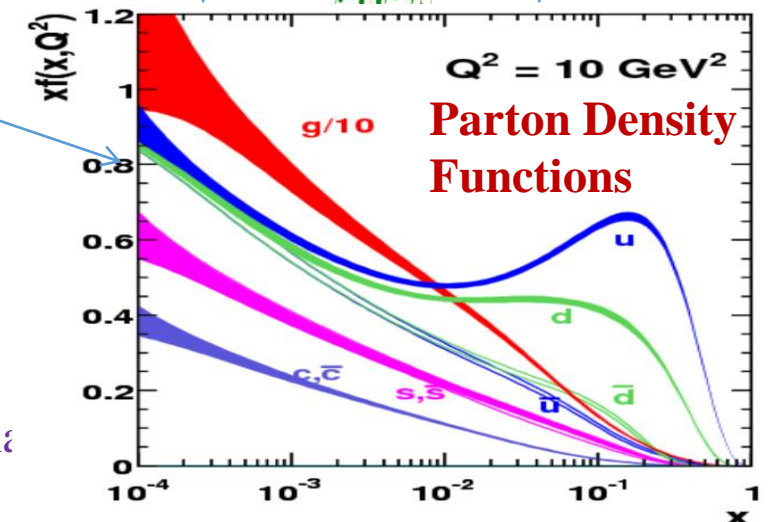
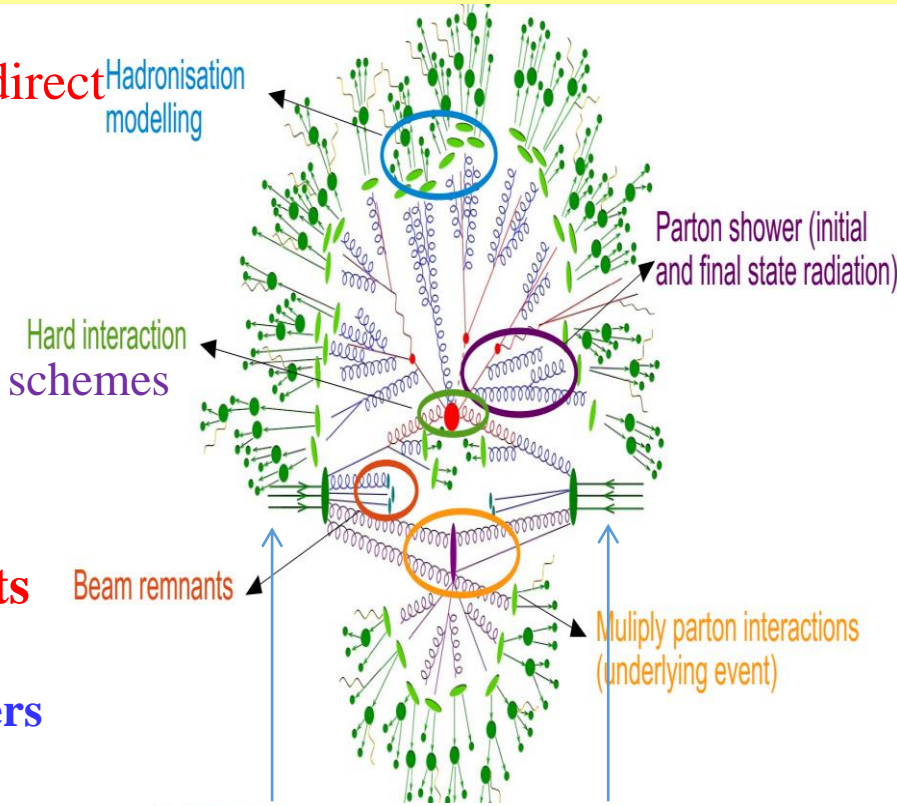


❑ Understanding of soft QCD interactions has direct impact on

- Precision measurements,
- Searches for new physics,
- Perturbative *Multi-Parton Interactions* (MPI),
- Alternative evolution equations and factorisation schemes

❑ Importance of good modelling soft QCD aspects of hadron collisions:

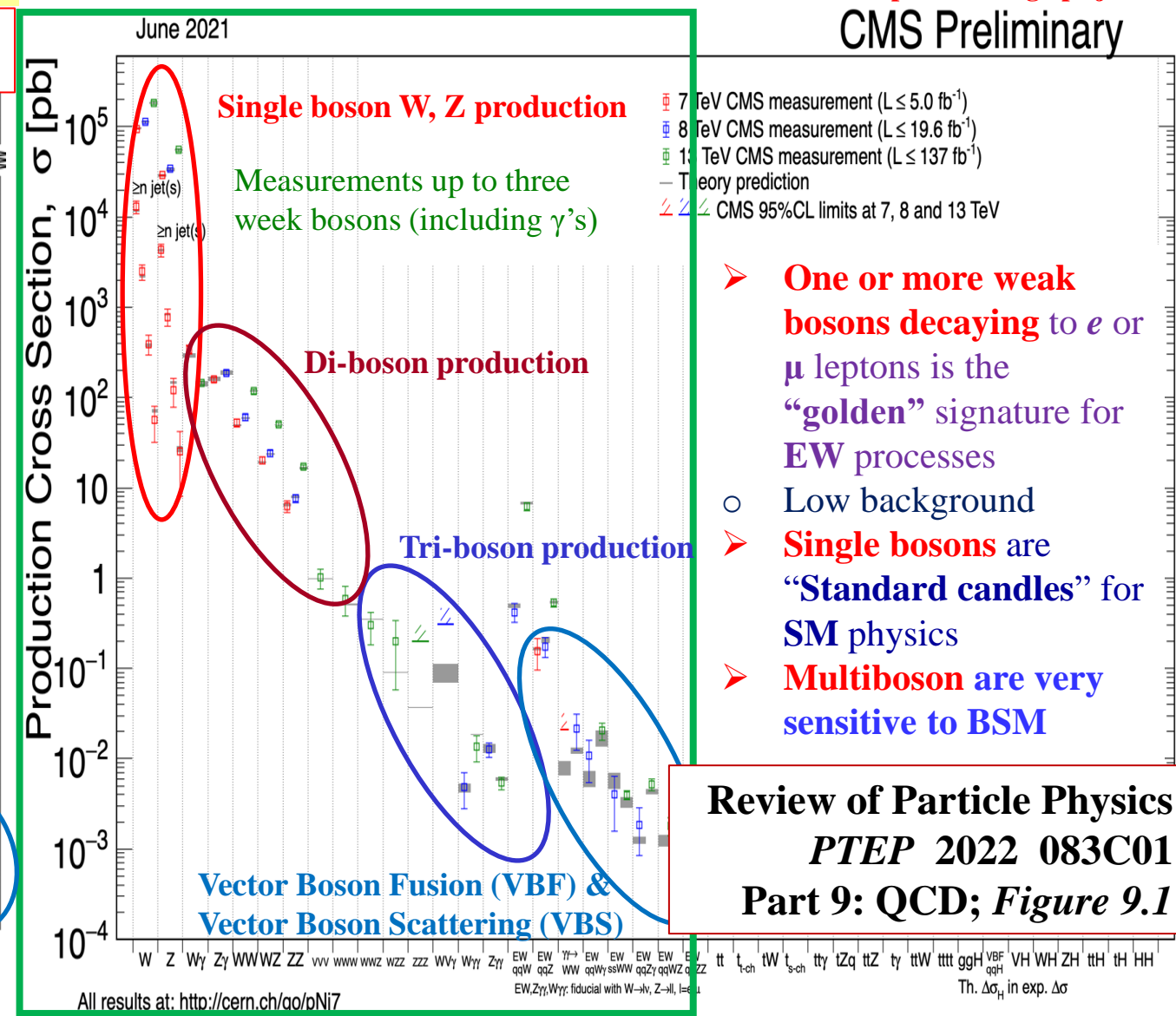
- Models developed with set of **tunable parameters** to describe the **hadron-level properties** of final states dominated by QCD
- Probing parton distribution functions (PDF)
- Modelling initial state of heavy ion collisions
- Modelling of backgrounds for electroweak



❑ Soft QCD measurements are:

- Crucial for success of ongoing LHC measurements and R&D for ATLAS upgrades
- Crucial for the tuning of the Monte Carlo event generator
- Essential to understand and correctly simulate any other more complex phenomena
- Ideal to study tracking performance in the “early” stage of a new data taking

CMS Preliminary



□

A TOROIDAL LHC APPARATUS (ATLAS)

Subdetector

Operational Fraction

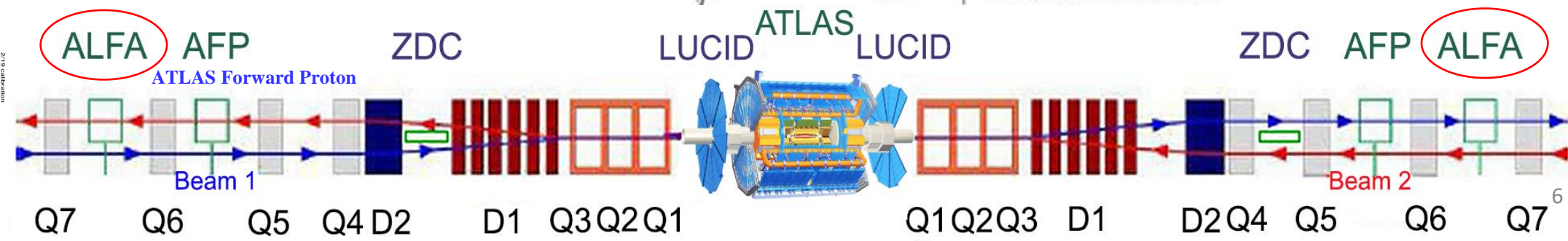
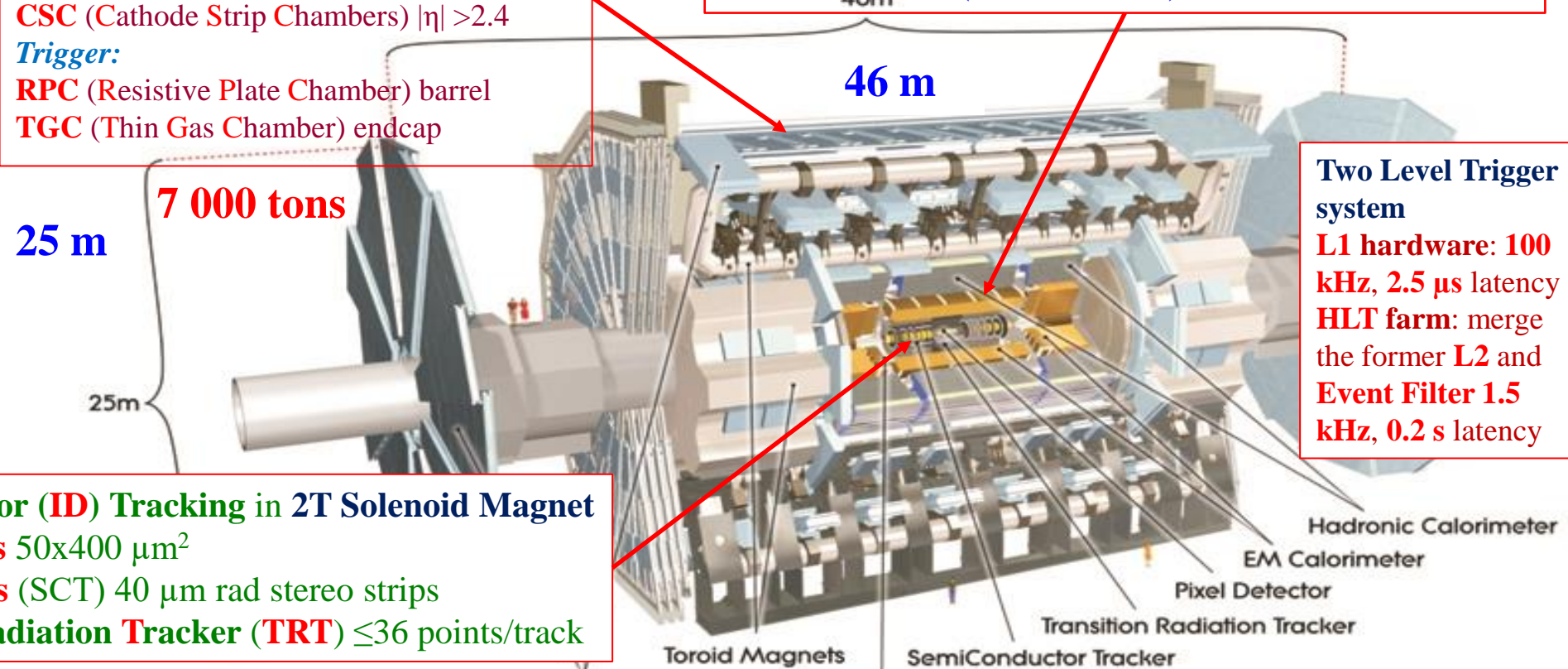
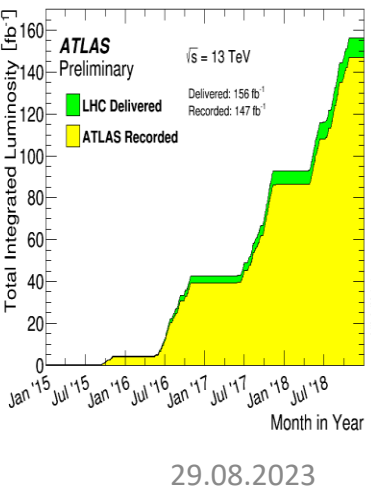
AFP	93.8%
ALFA	99.9%
CSC Cathode Strip Chambers	95.3%
Forward LAr Calorimeter	99.7%
Hadronic End-Cap Lar Cal	99.5%
LAr EM Calorimeter	100 %
LVL1 Calo Trigger	99.9%
LVL1 Muon RPC Trigger	99.8%
LVL1 Muon TGC Trigger	99.9%
MDT Muon Drift Tubes	99.7%
Pixels	97.8%
RPC Barrel Muon Chambers	94.4%
SCT Silicon Strips	98.7%
TGC End-Cap Muon Cha	99.5%
Tile Calorimeter	99.2%
TRT Transit Rad Tracker	97.2%

Air-core Muon spectrometer
 (μ Trigger/tracking and Toroid Magnets)
Precision Tracking:
 MDT (Monitored Drift Tubes)
 CSC (Cathode Strip Chambers) $|\eta| > 2.4$
Trigger:
 RPC (Resistive Plate Chamber) barrel
 TGC (Thin Gas Chamber) endcap

Longitudinally segmented Calorimeter:
 EM and Hadronic energy
 LiquidAr EM barrel and End-cap & Hadronic End-cap
 Tile calorimeter (Fe-scintillator) Hadronic barrel

Two Level Trigger system
 L1 hardware: 100 kHz, 2.5 μ s latency
 HLT farm: merge the former L2 and Event Filter 1.5 kHz, 0.2 s latency

Inner Detector (ID) Tracking in 2T Solenoid Magnet
 Silicon Pixels 50x400 μ m²
 Silicon Strips (SCT) 40 μ m rad stereo strips
 Transition Radiation Tracker (TRT) ≤ 36 points/track



COMPACT MUON SOLENOID (CMS) + TOTEM



CMS DETECTOR

Total weight : 14,000 tonnes
Overall diameter : 15.0 m
Overall length : 28.7 m
Magnetic field : 3.8 T

STEEL RETURN YOKE
12,500 tonnes

SILICON TRACKERS
Pixel ($100 \times 150 \mu\text{m}^2$) $\sim 1.9 \text{ m}^2 \sim 124\text{M}$ channels
Microstrips ($80\text{--}180 \mu\text{m}$) $\sim 200 \text{ m}^2 \sim 9.6\text{M}$ channels

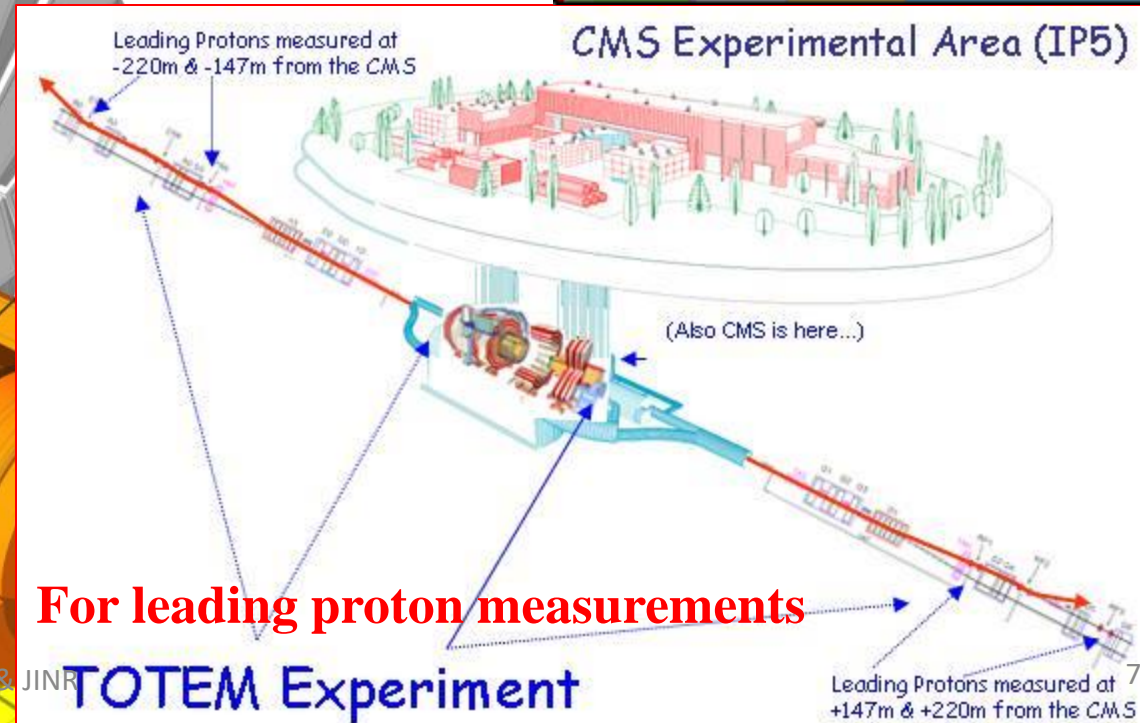
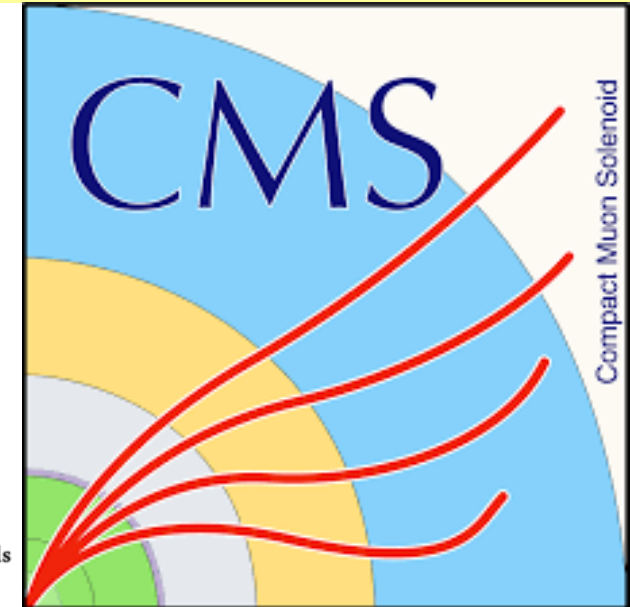
SUPERCONDUCTING SOLENOID
Niobium titanium coil carrying $\sim 18,000 \text{ A}$

MUON CHAMBERS
Barrel: 250 Drift Tube, 480 Resistive Plate Chambers
Endcaps: 540 Cathode Strip, 576 Resistive Plate Chambers

PRESHOWER
Silicon strips $\sim 16 \text{ m}^2 \sim 137,000$ channels

CRYSTAL
ELECTROMAGNETIC
CALORIMETER (ECAL)
 $\sim 76,000$ scintillating PbWO_4 crystals

HADRON CALORIMETER (HCAL)
Brass + Plastic scintillator $\sim 7,000$ channels



For leading proton measurements

TOTEM Experiment

Yuri Kulchitsky, IP NASB & JINR

□ Total Cross-sections

1. $pp \rightarrow X$

□ Diboson cross sections

1. $pp \rightarrow \gamma\gamma + X$

2. $pp \rightarrow 4l + X$

□ WW-boson cross sections

1. $pp \rightarrow (\gamma\gamma \rightarrow W^+W^-) + X$

2. $pp \rightarrow W^+W^- + \geq 1 \text{ jet} + X$

□ Measurements of the rarest processes

1. $pp \rightarrow Z(\rightarrow \nu\nu)\gamma + 2 \text{ jets} + X$

2. $pp \rightarrow Z(\rightarrow ll)\gamma + 2 \text{ jets} + X$

3. $pp \rightarrow WWW + X$

□ EW Z-boson cross sections

1. $pp \rightarrow Z + \text{jet}(\text{high-}p_T) + X$

2. $pp \rightarrow Z + \text{large } R\text{-jet} + X$

□ The b-quark fragmentation properties

1. $pp \rightarrow \text{jet}(B^\pm \rightarrow J/\psi K^\pm \rightarrow \mu^+ \mu^- K^\pm) + X$

□ Forward Proton Scattering

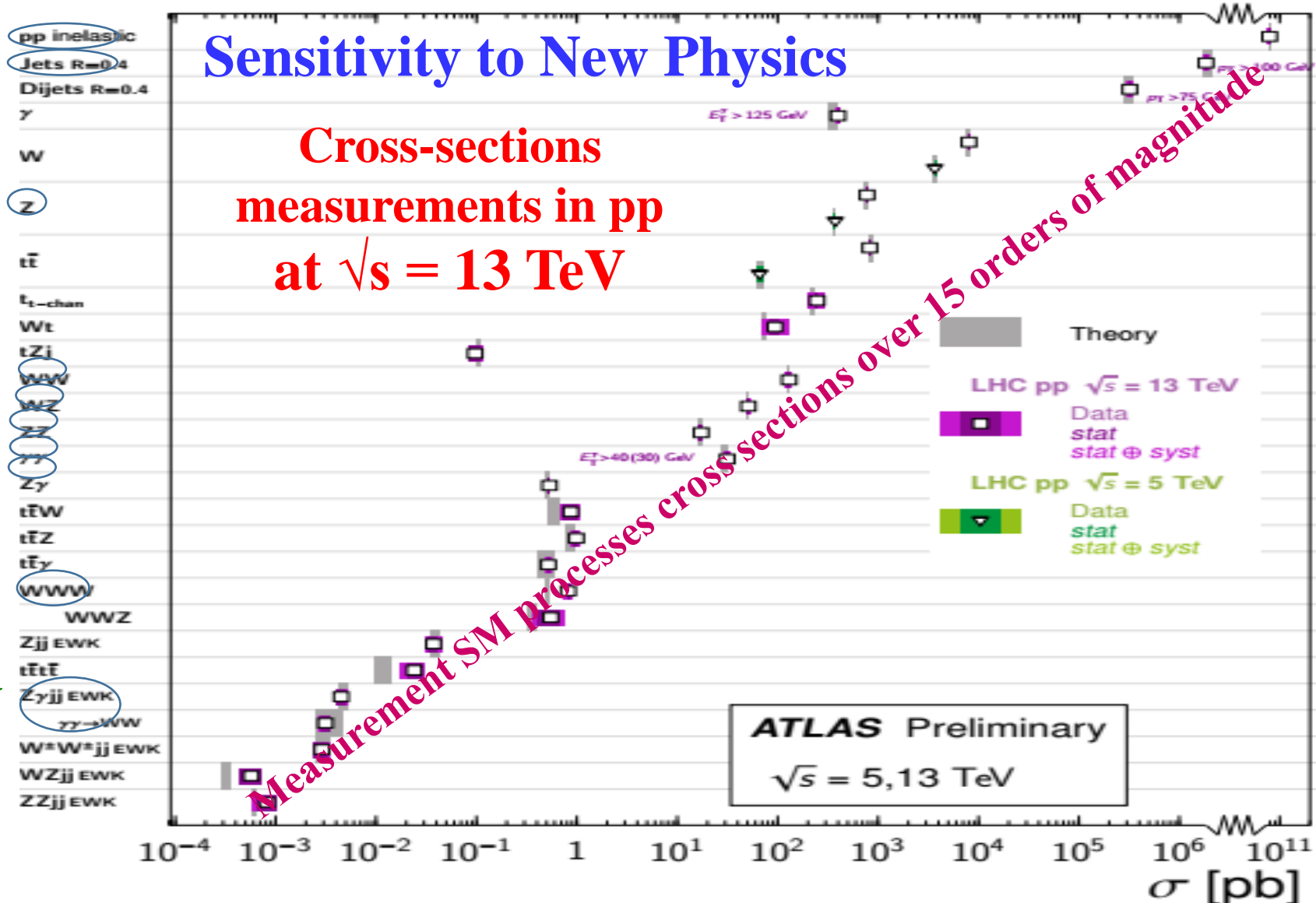
1. $pp \rightarrow p(\gamma\gamma \rightarrow l^+l^-) p^*$

□ Soft QCD: Bose-Einstein correlations

1. $pp \rightarrow h^\pm h^\pm + X$

□ Bonus: W-mass

Standard Model Production Cross Section Measurements



ATLAS: SM and fiducial production cross-section measurements in **pp** interactions at **5, 13 TeV**, corrected for branching fractions, compared to the corresponding **theoretical expectations**

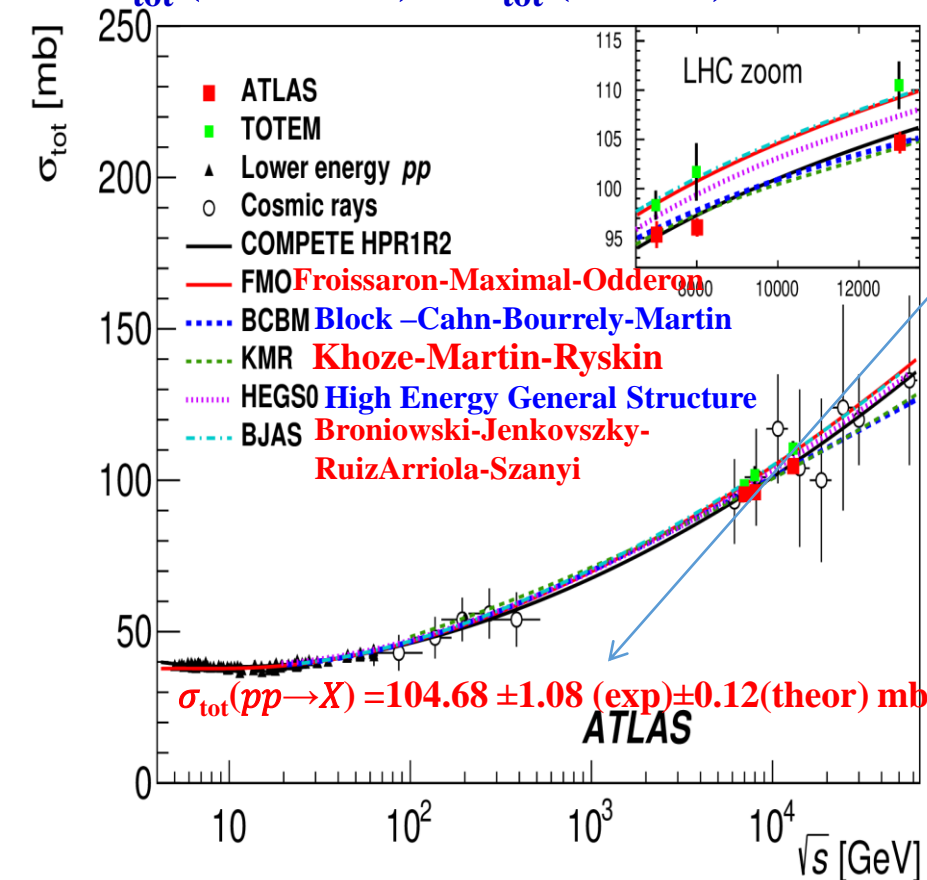
- **ATLAS** measurements of elastic scattering can be linked to other processes occurring on hadronic pp-interactions
- **Calculation the total pp cross section σ_{tot} & the ρ -ratio**
- The **ρ -ratio** determines the complex phase between the **Coulomb** and the **nuclear amplitudes**

❑ σ_{tot} (TOTEM) > σ_{tot} (ALFA) on 5.8 mb (2.2σ)

$$\rho = \frac{\text{Re}[f_{el}(t)]}{\text{Im}[f_{el}(t)]} \Big|_{t \rightarrow 0}$$

where $f_{el}(t)$ is the **elastic-scattering amplitude**, t stands for the four-momentum transfer in the reaction

➤ $\rho = 0.0978 \pm 0.0085$ (exp) ± 0.0064 (th)

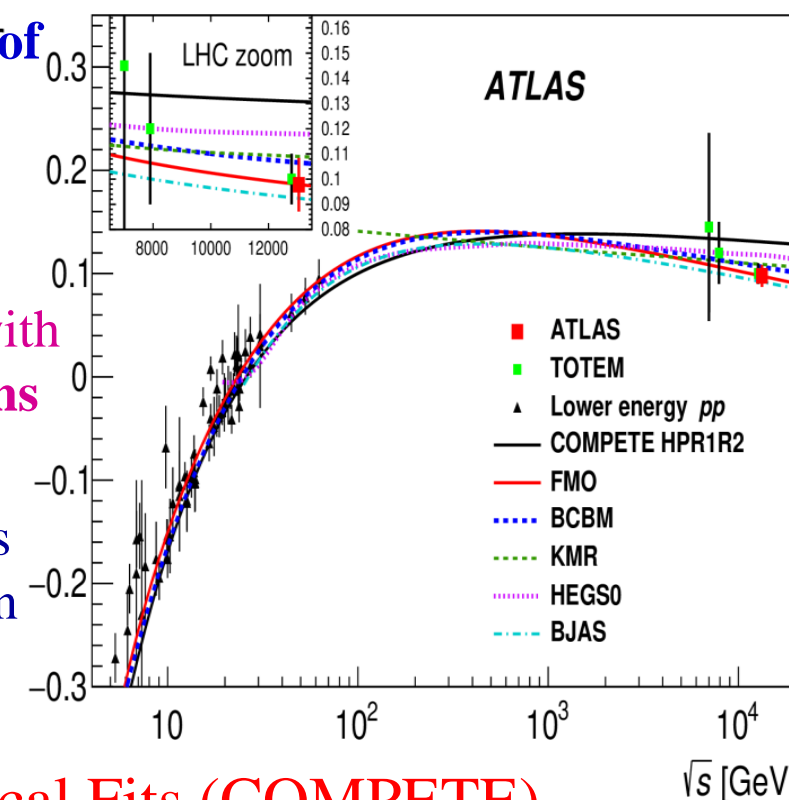


➤ The most precise measurement of the total cross section

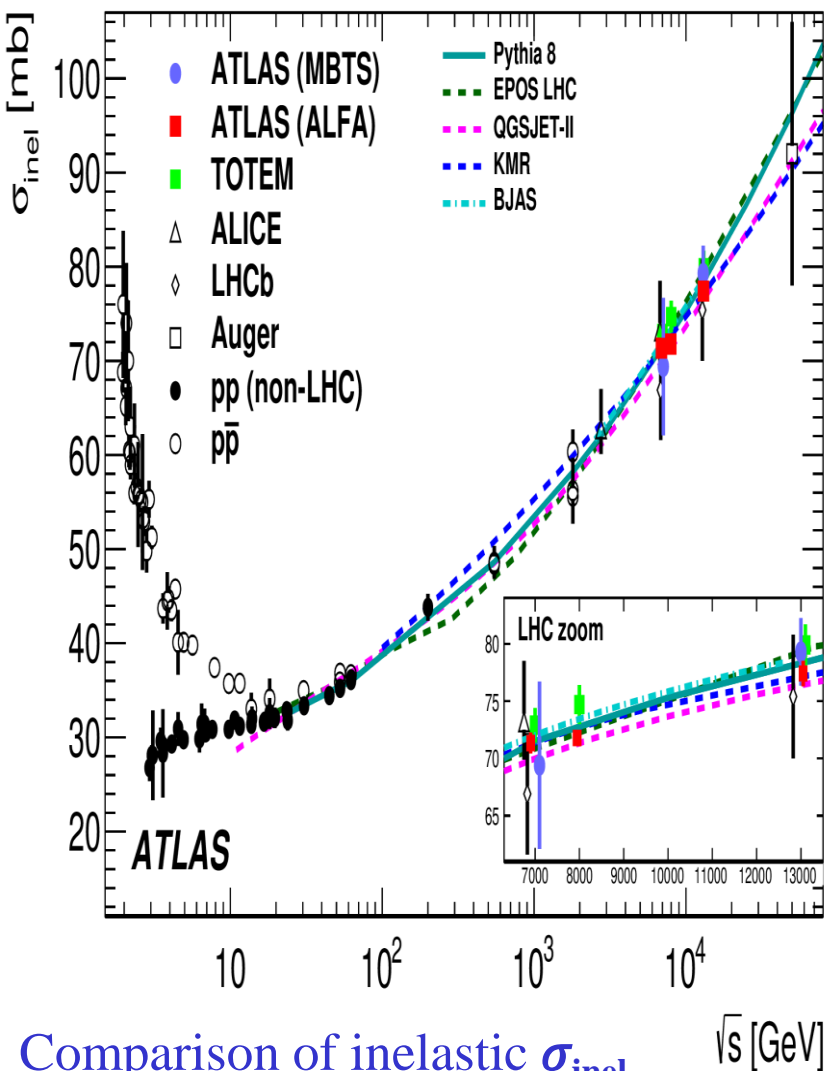
❑ The confirms the trend of σ_{tot} increase with \sqrt{s}

❑ The ρ -ratio is in disagreement with pre-LHC theoretical expectations

❑ σ_{tot} and ρ together with other measurements of these parameters compared to **model predictions** in dependence of \sqrt{s}



σ_{tot} and ρ incompatible with Community-Standard Semi-Empirical Fits (COMPETE) indicating **Odderon** exchange or a slowdown of σ_{tot} rise at high \sqrt{s}



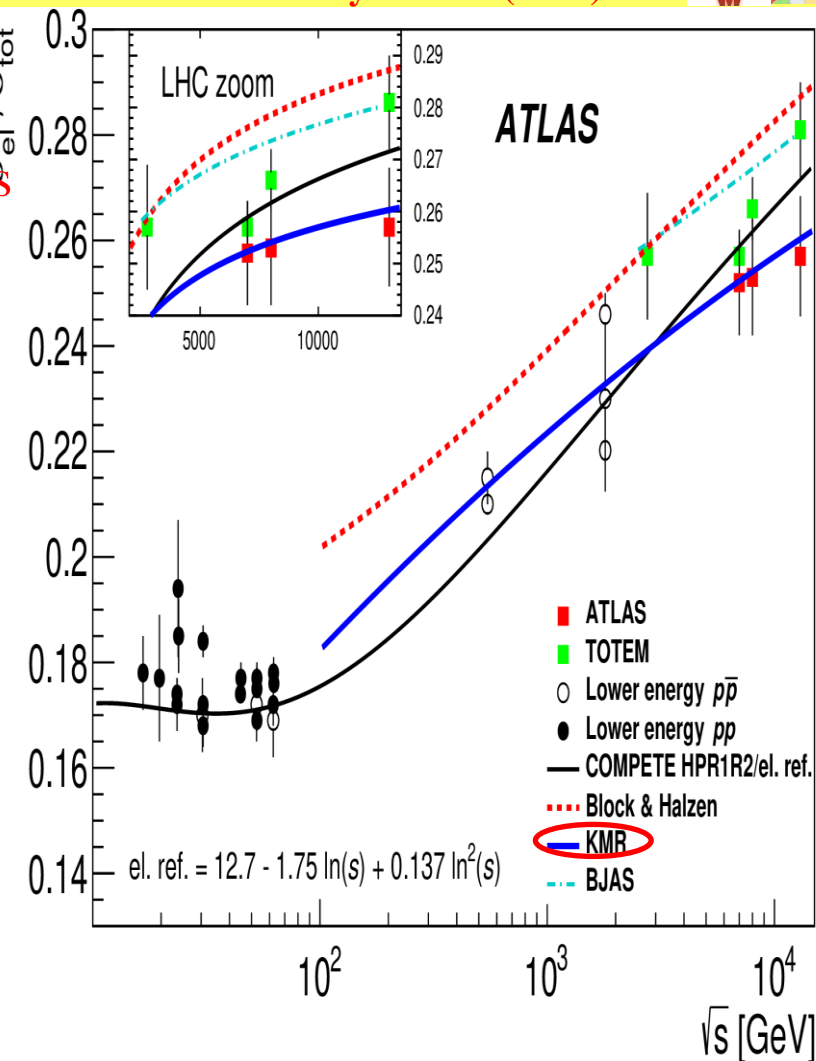
Comparison of inelastic σ_{inel} measurements with other published measurements and model predictions as a function of the \sqrt{s}

➤ **Total inelastic cross section σ_{tot} is in agreement with previous ATLAS measurements using MBTS detectors**

❑ **The ratio $\sigma_{\text{inel}}/\sigma_{\text{tot}}$, a measure of the opaqueness of the proton, continues to grow slowly with \sqrt{s} , and its evolution is well described by the Khoze-Martin-Ryskin (KMR) model. (This is a two-channel eikonal model with few parameters and it uses all available high-energy data for ρ and σ_{tot} , as well as the corresponding differential elastic cross sections, and also all available measurements of low-mass diffraction.)**

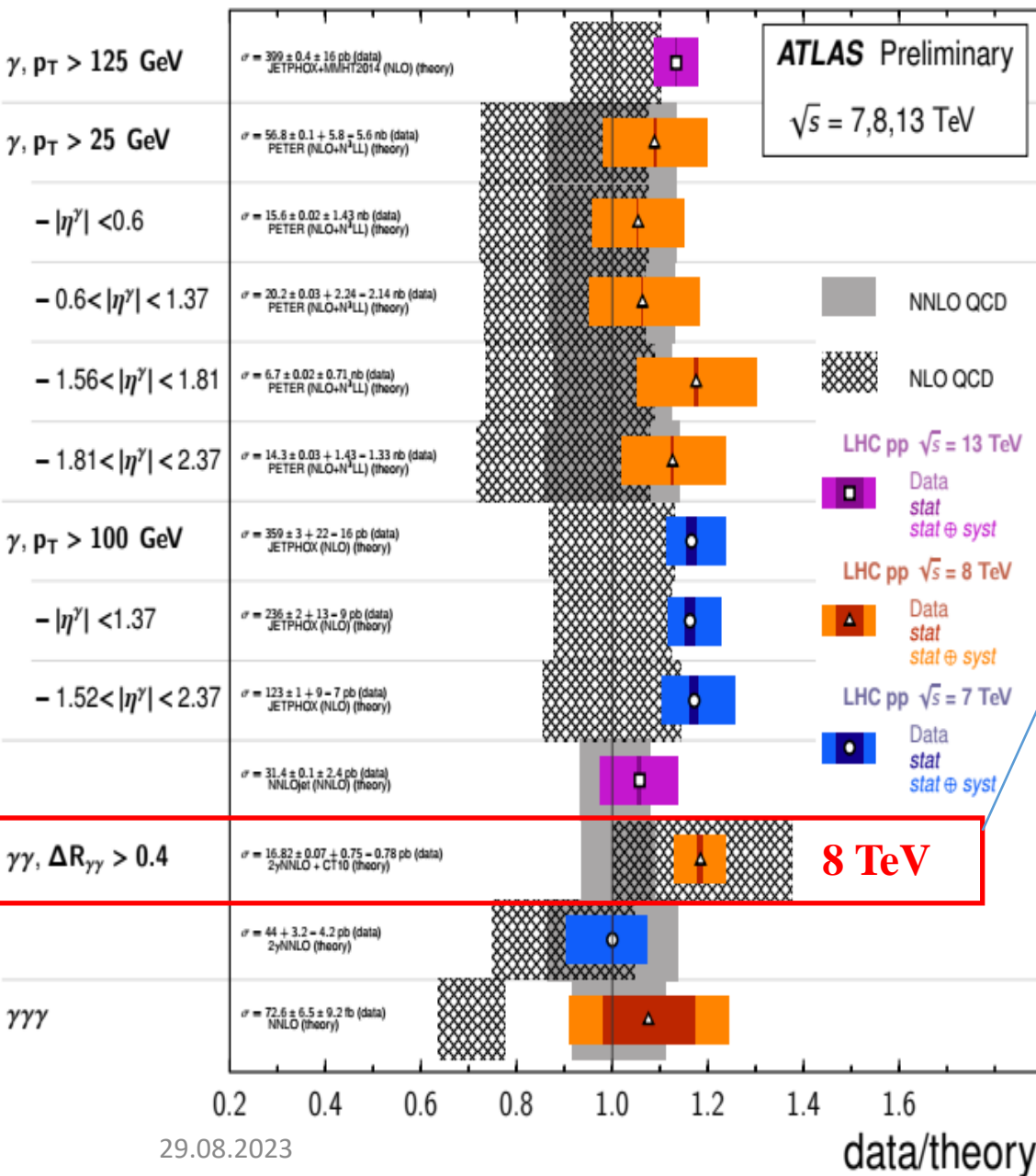
❑ **The measurement remains far from probing the black-disc limit, i.e. a totally opaque proton**

➤ **Ratio of elastic to total cross section $\sigma_{\text{el}}/\sigma_{\text{tot}}$ in tension with the values from TOTEM and lower energies**



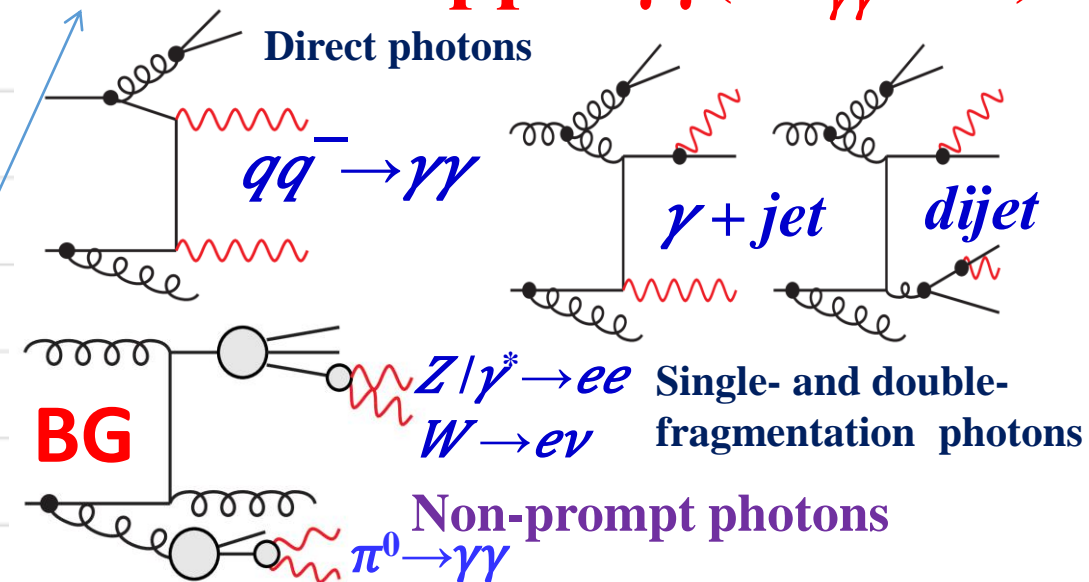
Measurements of the ratio $\sigma_{\text{el}}/\sigma_{\text{tot}}$ at different \sqrt{s} compared to model predictions and for illustrative purposes the COMPETE prediction of σ_{tot} divided by a conventional parameterization of the elastic cross section

Inclusive Photon Fiducial Cross Section Measurements


 $\int \mathcal{L} dt$
[fb⁻¹]

Re

Cross section of $pp \rightarrow \gamma\gamma (\Delta R_{\gamma\gamma} > 0.4) + X$



Selection	Detector level	Particle level
Photon kinematics	$p_{T,\gamma_{1(2)}} > 40$ (30) GeV, $ \eta_\gamma < 2.37$ excluding $1.37 < \eta_\gamma < 1.52$	
Photon identification	tight	stable, not from hadron decay
Photon isolation	$E_{T,\gamma}^{\text{iso},0.2} < 0.05 \cdot p_{T,\gamma}$	$E_{T,\gamma}^{\text{iso},0.2} < 0.09 \cdot p_{T,\gamma}$
Diphoton topology		$N_\gamma \geq 2, \Delta R_{\gamma\gamma} > 0.4$

Overview of the event selection at the detector level and at the particle level (fiducial phase space)

Process	Event fraction [%]
$\gamma\gamma$ signal	60.4 ± 3.0
γj	20.0 ± 1.3
$j\gamma$	10.1 ± 1.1
$j j$	6.3 ± 1.2
Electron	2.6 ± 0.1
$\gamma\gamma$ pile-up	0.6 ± 0.4

Status: February 2022

Diphoton production fiducial & differential cross sections

Low-invariant mass diphoton event

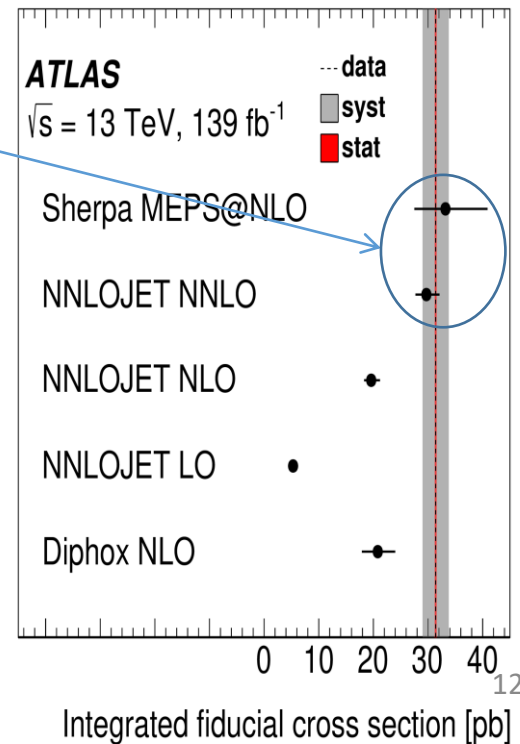
$$\Delta R_{\gamma\gamma} > 0.4$$

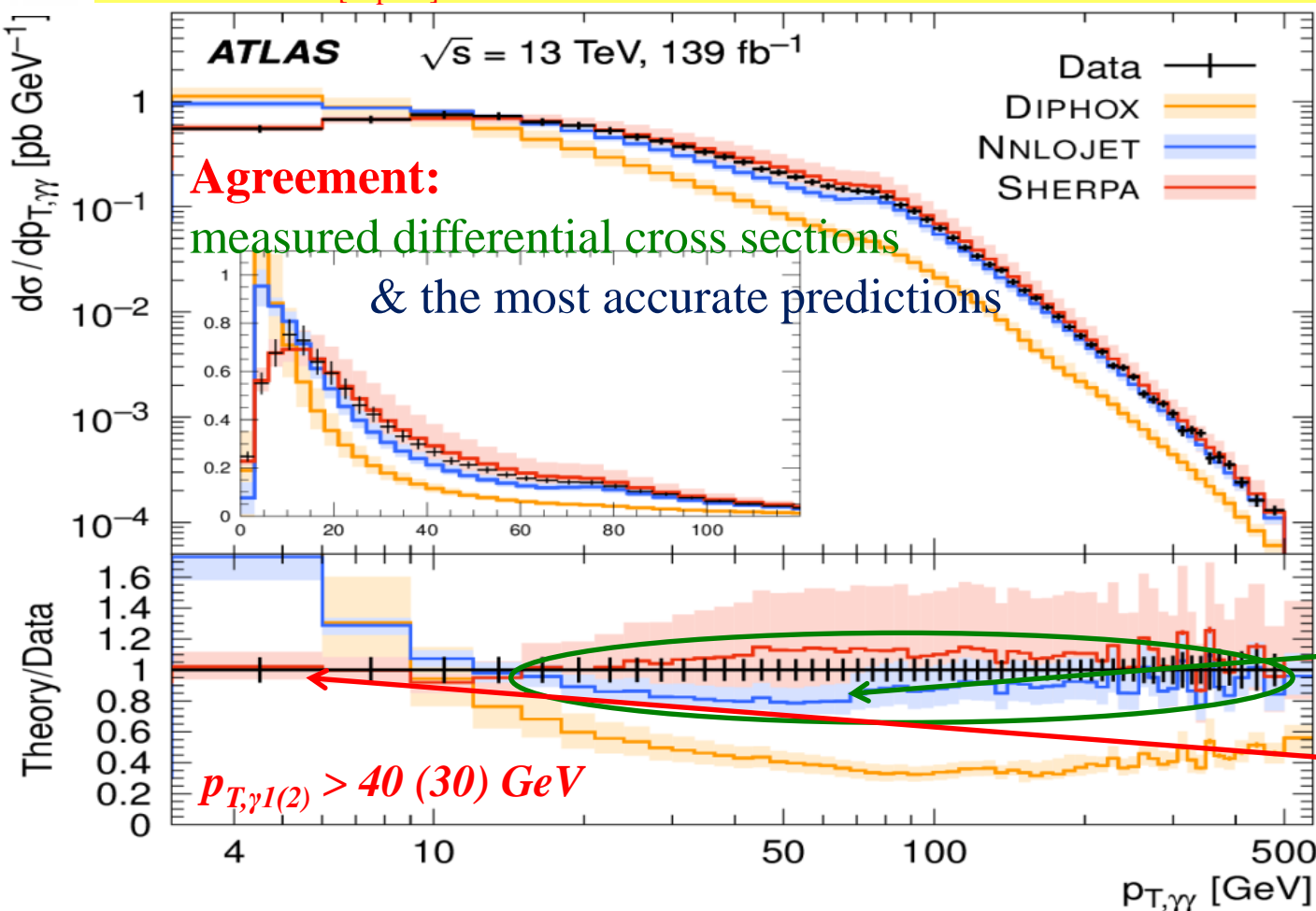
The measured integrated cross section compatible with the NNLO predictions and multileg-merged calculations

Fiducial cross section [pb]	$\sigma_{\gamma\gamma}$	\pm unc.
SHERPA MEPS@NLO	33.2	+7.7 -5.6
NNLOJET NNLO	29.7	+2.4 -2.0
NLO	19.6	+1.6 -1.3
LO	5.3	+0.5 -0.5
DIPHON NLO	20.8	+3.2 -2.9
Data	31.4	2.4

$$\sigma_{\gamma\gamma} = 31.4 \pm 0.1 \text{ (stat)} \pm 2.4 \text{ (syst) pb}$$

A comparison of the measured cross section with the theoretical predictions shows the **importance of higher-order QCD contributions** even for such an **inclusive $\gamma\gamma$ measurement**





Higher-order perturbative QCD

corrections are crucial for $pp \rightarrow \gamma\gamma + X$

- ❑ In the **Leading Order (LO)** the $\gamma\gamma$ momenta balance each other perfectly, resulting in a $p_{T,\gamma\gamma}$ of zero. This obviously doesn't match the measured distribution, which shows that most of the γ -pairs have $p_{T,\gamma\gamma}$ values different from 0
- ❑ In the **Diphox NLO** and the **Nnlojet NNLO** physicists include the **production of the extra quarks and gluons** in the prediction.
- ❑ The **Nnlojet NNLO** prediction (**blue line**) gives a **good description** of the measured values at **high- $p_{T,\gamma\gamma}$**
- ➔ The **Sherpa** (red line) combine calculations with additional techniques, which involve the simulations of arbitrarily many quark and gluon emissions, especially relevant at low energies

❖ **Sherpa** provide the **best description** of the **entire measured distribution**, including the **low- $p_{T,\gamma\gamma}$**

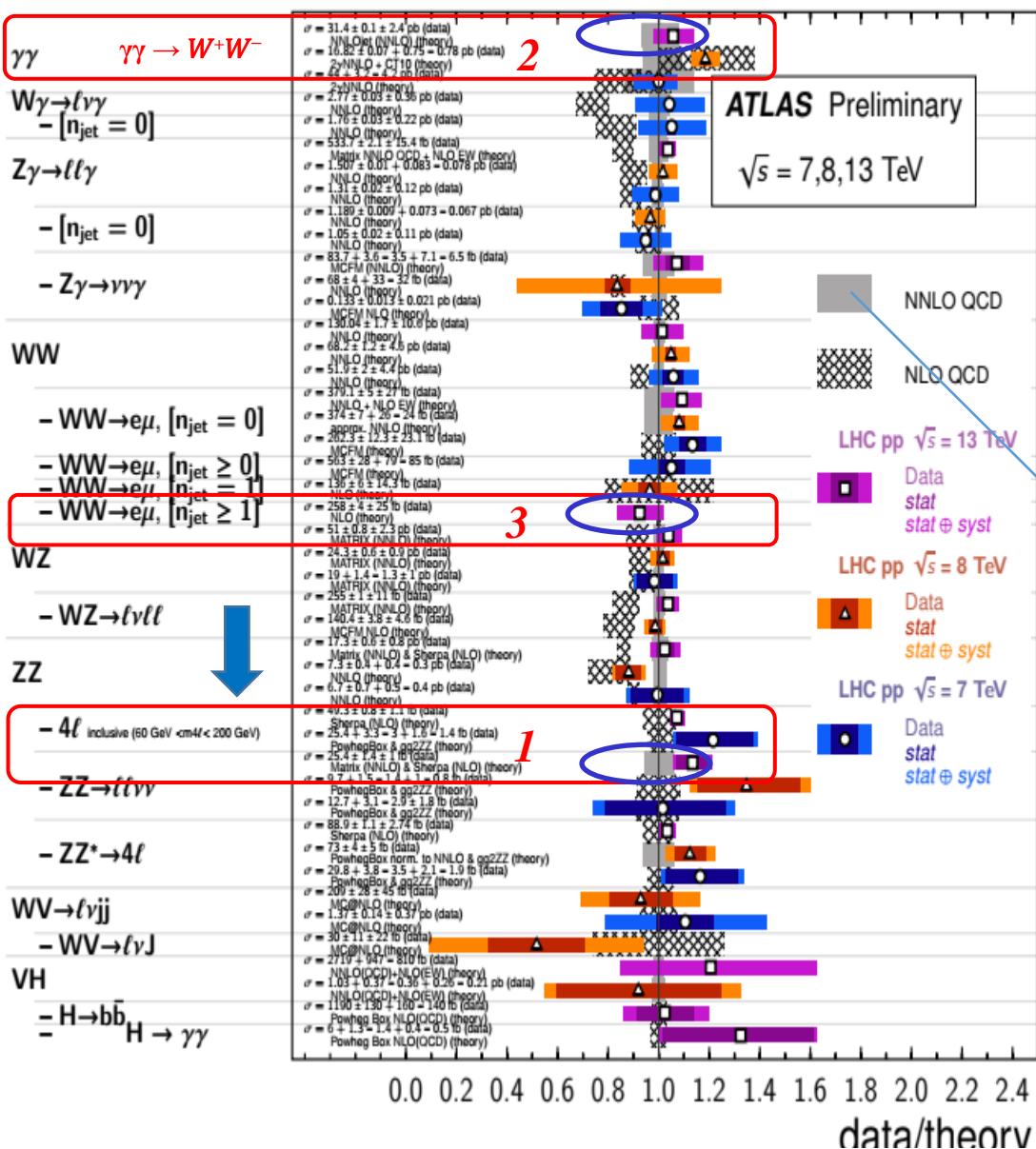
Differential cross sections measured as functions of $p_{T,\gamma\gamma}$ (also as functions of $\phi_{\gamma\gamma}^*, \pi - \Delta\phi_{\gamma\gamma}, a_{T,\gamma\gamma}$) compared with the predictions from **Diphox NLO**, **Nnlojet NNLO** and **Sherpa MEPS@NLO**. At the bottom, the ratio of the prediction to the data is shown. **Uncertainty bars** on the data represent the **total uncertainty**, while **uncertainty bands** on the predictions represent perturbative scale (**statistical**) uncertainties.

DIBOSON CROSS SECTION MEASUREMENTS

<https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsCombined>

Diboson Cross Section Measurements

Status: February 2022


 $\int \mathcal{L} dt$
 $[\text{fb}^{-1}]$

Referen

139	JHEP 11 (2021) 1
20.2	PRD 95 (2017) 1
4.9	JHEP 01, 086 (2017)
4.6	PRD 87, 112003 (2013)
4.6	PRD 87, 112003 (2013)
36.1	JHEP 03 (2020) 1
20.3	PRD 93, 112002 (2016)
20.3	PRD 87, 112003 (2013)
4.6	PRD 87, 112003 (2013)
36.1	JHEP 12 (2018) 1
20.3	PRD 93, 112002 (2016)
4.6	PRD 87, 112003 (2013)
36.1	EPJC 79 (2019) 1
20.3	PLB 763, 114 (2016)
4.6	Phys. Rev. D 87, 1408.5243 (2013)
36.1	EPJC 79 (2019) 1
20.3	JHEP 09 (2016) 1
4.6	PRD 87, 112001 (2013)
4.6	PRD 91, 052005 (2015)
20.3	PLB 763, 114 (2016)
139	ATL-COM-Phys-2019-001
36.1	EPJC 79 (2019) 1
20.3	PRD 93, 082004 (2016)
4.6	EPJC 72 (2012) 2
36.1	EPJC 79 (2019) 1
20.3	PRD 93, 082004 (2016)
36.1	PRD 97 (2018) 0
20.3	JHEP 01, 099 (2017)
4.6	JHEP 03, 128 (2017)
139	JHEP 07 (2021) 1
4.6	JHEP 03, 128 (2017)
36.1	JHEP 10 (2019) 1
20.3	JHEP 10 (2019) 1
4.6	JHEP 03, 128 (2017)
139	JHEP 07 (2021) 1
20.3	PLB 753, 552-57 (2016)
4.6	JHEP 03, 128 (2017)
20.2	EPJC 77 (2017) 1
4.6	JHEP 01, 049 (2017)
36.1	JHEP 12 (2017) 1
20.3	JHEP 12 (2017) 1
139	ATLAS-CONF-2019-010
139	ATLAS-CONF-2019-010

May 2021

CMS Preliminary

CMS measurements

vs. NNLO (NLO) theory

7 TeV CMS measurement (stat,stat+syst)

8 TeV CMS measurement (stat,stat+syst)

13 TeV CMS measurement (stat,stat+syst)

 $\gamma\gamma$

2

 $1.06 \pm 0.01 \pm 0.12$ 5.0 fb^{-1} $W\gamma$, (NLO th.) $1.16 \pm 0.03 \pm 0.13$ 5.0 fb^{-1} $W\gamma$, (NLO th.) $1.01 \pm 0.00 \pm 0.05$ 137 fb^{-1} $Z\gamma$, (NLO th.) $0.98 \pm 0.01 \pm 0.05$ 5.0 fb^{-1} $Z\gamma$, (NLO th.) $0.98 \pm 0.01 \pm 0.05$ 19.5 fb^{-1} $WW+WZ$ $1.01 \pm 0.13 \pm 0.14$ 4.9 fb^{-1} WW $1.07 \pm 0.04 \pm 0.09$ 4.9 fb^{-1} WW

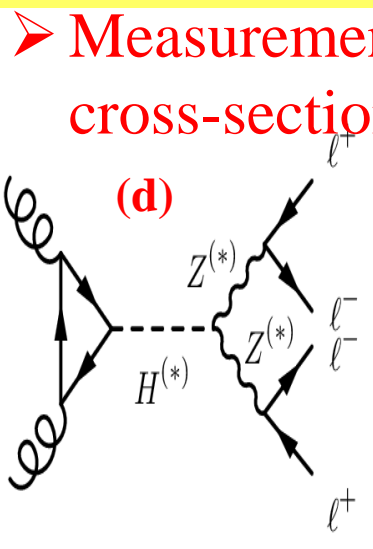
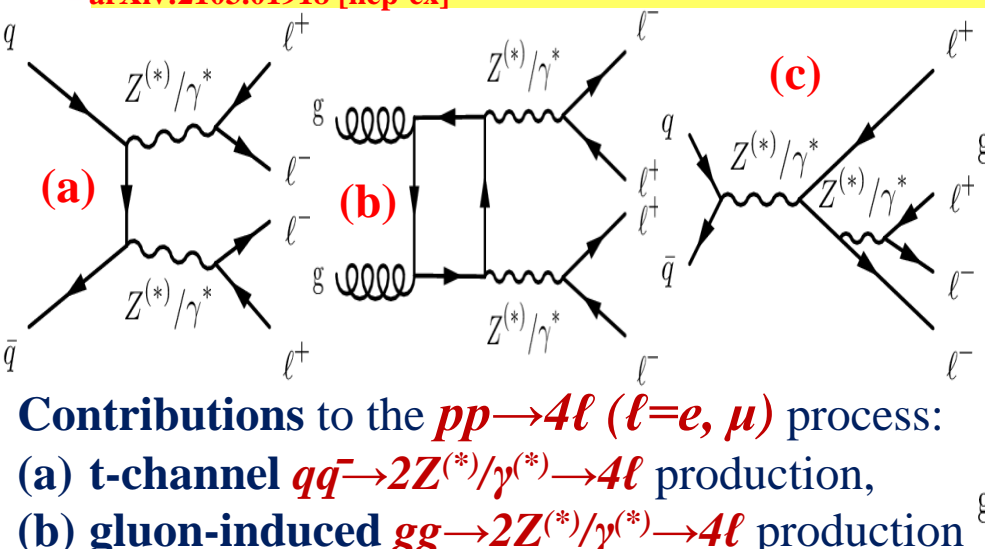
3

 $1.00 \pm 0.02 \pm 0.08$ 19.4 fb^{-1} WW $1.00 \pm 0.01 \pm 0.06$ 35.9 fb^{-1} WZ $1.05 \pm 0.07 \pm 0.06$ 4.9 fb^{-1} WZ $1.02 \pm 0.04 \pm 0.07$ 19.6 fb^{-1} WZ $1.00 \pm 0.02 \pm 0.03$ 137 fb^{-1} ZZ $0.97 \pm 0.13 \pm 0.07$ 4.9 fb^{-1} ZZ $0.97 \pm 0.06 \pm 0.08$ 19.6 fb^{-1} ZZ $1.04 \pm 0.02 \pm 0.04$ 137 fb^{-1}

All results at:

Production Cross Section Ratio: $\sigma_{\text{exp}} / \sigma_{\text{theo}}$

□ Diboson cross section ratio comparison to theory: Theory predictions updated to latest NNLO calculations.



The superscript (*) refers to a particle that can be either on-shell or off-shell, whereas * indicates that it is always off-shell

➤ Measurements of 4ℓ differential & integrated fiducial cross-sections in events with $2(e^+e^-)$ or $2(\mu^+\mu^-)$ pairs

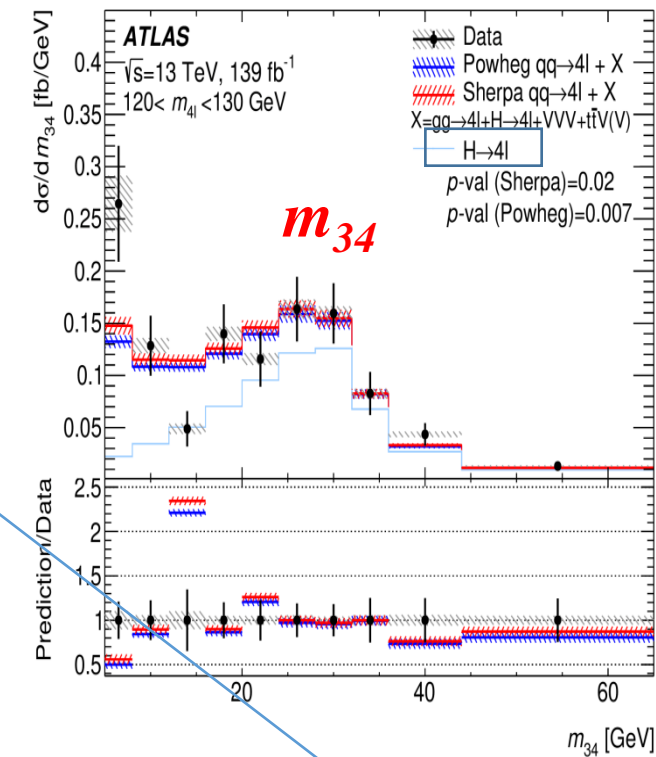
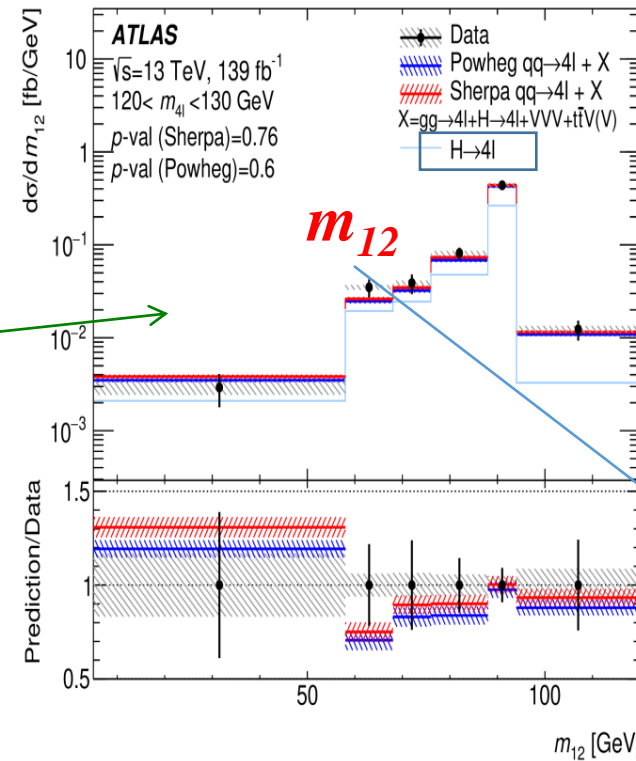
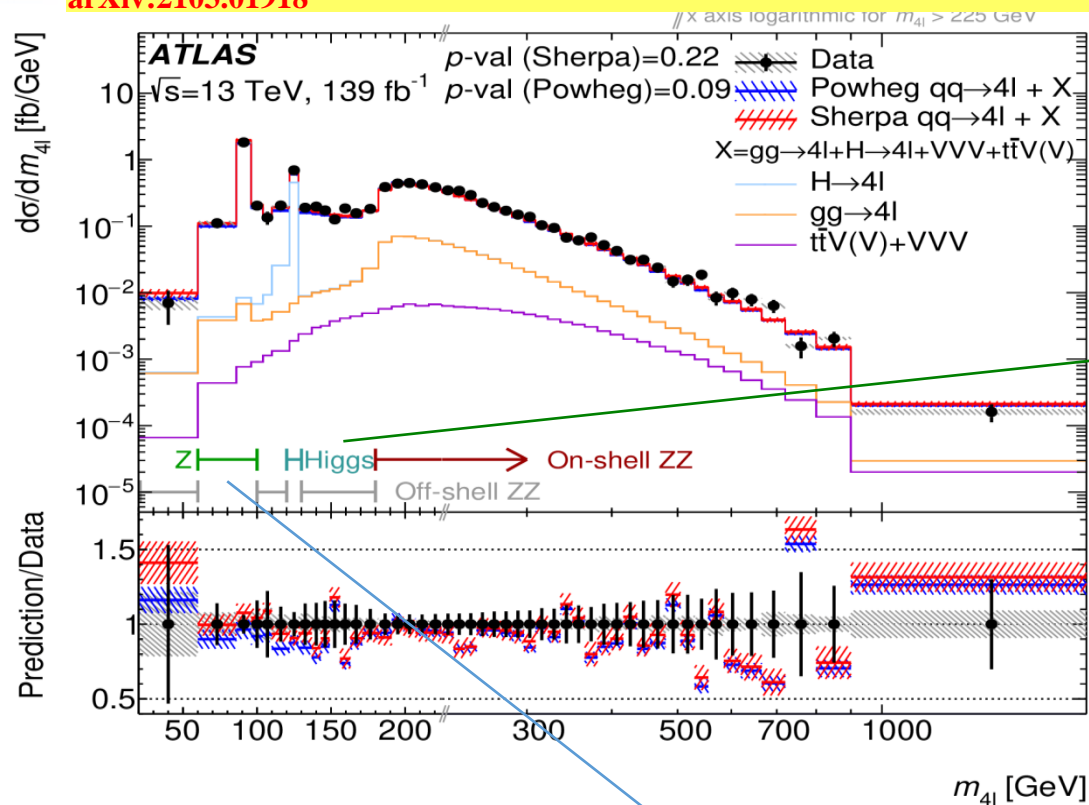
- ❑ The 4ℓ final state has contributions from interesting SM processes: 1. Single **Z** boson production, 2. **Higgs** boson production, 3. On-shell **ZZ** production.
- ❑ Sensitive to New Physics / BSM contributions: 1. Modifications to the couplings of the **Higgs** or gauge boson, 2. Possible **4-fermion** interactions, 3. Models with leptonic decays of Z bosons or 4. new particles.

➤ Fiducial cross-sections in fb in the full fiducial phase space and in the following regions of $m_{4\ell}$:

1. $Z \rightarrow 4\ell$ ($60 < m_{4\ell} < 100$ GeV)
2. $H \rightarrow 4\ell$ ($120 < m_{4\ell} < 130$ GeV)
3. off-shell ZZ ($20 < m_{4\ell} < 60$ GeV; $100 < m_{4\ell} < 120$ GeV; $130 < m_{4\ell} < 180$ GeV)
4. on-shell ZZ ($180 < m_{4\ell} < 2000$ GeV)

➤ Two predictions are shown for the $q\bar{q} \rightarrow 4\ell$ process simulated with Sherpa & Powheg+Pythia8

	Region				
	Full	$Z \rightarrow 4\ell$	$H \rightarrow 4\ell$	Off-shell ZZ	On-shell ZZ
Measured	88.9	22.1	4.76	12.4	49.3
fiducial cross-section [fb]	± 1.1 (stat.)	± 0.7 (stat.)	± 0.29 (stat.)	± 0.5 (stat.)	± 0.8 (stat.)
	± 2.3 (syst.)	± 1.1 (syst.)	± 0.18 (syst.)	± 0.6 (syst.)	± 0.8 (syst.)
	± 1.5 (lumi.)	± 0.4 (lumi.)	± 0.08 (lumi.)	± 0.2 (lumi.)	± 0.8 (lumi.)
	± 3.0 (total)	± 1.3 (total)	± 0.35 (total)	± 0.8 (total)	± 1.3 (total)
SHERPA	86 ± 5	23.6 ± 1.5	4.57 ± 0.21	11.5 ± 0.7	46.0 ± 2.9
POWHEG + PYTHIA8	83 ± 5	21.2 ± 1.3	4.38 ± 0.20	10.7 ± 0.7	46.4 ± 3.0



Differential cross-section as a function of $m_{4\ell}$

- ❖ The measured data are compared with the SM prediction using *Sherpa* (red) or *Powheg+PYTHIA* (blue) models of the $qq \rightarrow 4\ell$ contribution
- ❖ The p -value is the probability for the χ^2 , with k degrees of freedom.

➤ The SM predictions agree well within uncertainties over $m_{4\ell}$ spectrum

❖ Differential cross-section as a function of m_{12} & m_{34} in the four $m_{4\ell}$ regions:: for the contribution from Higgs production.

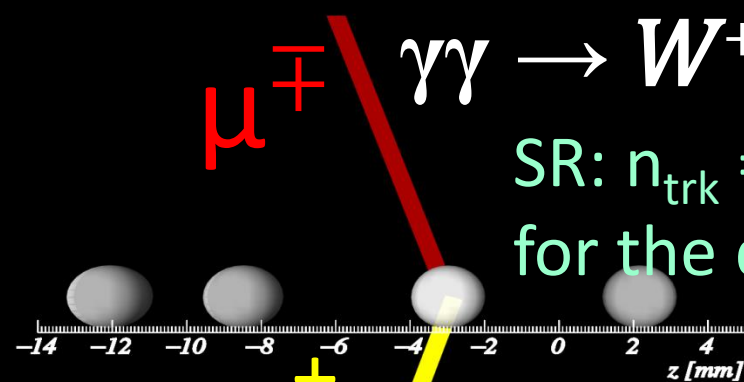
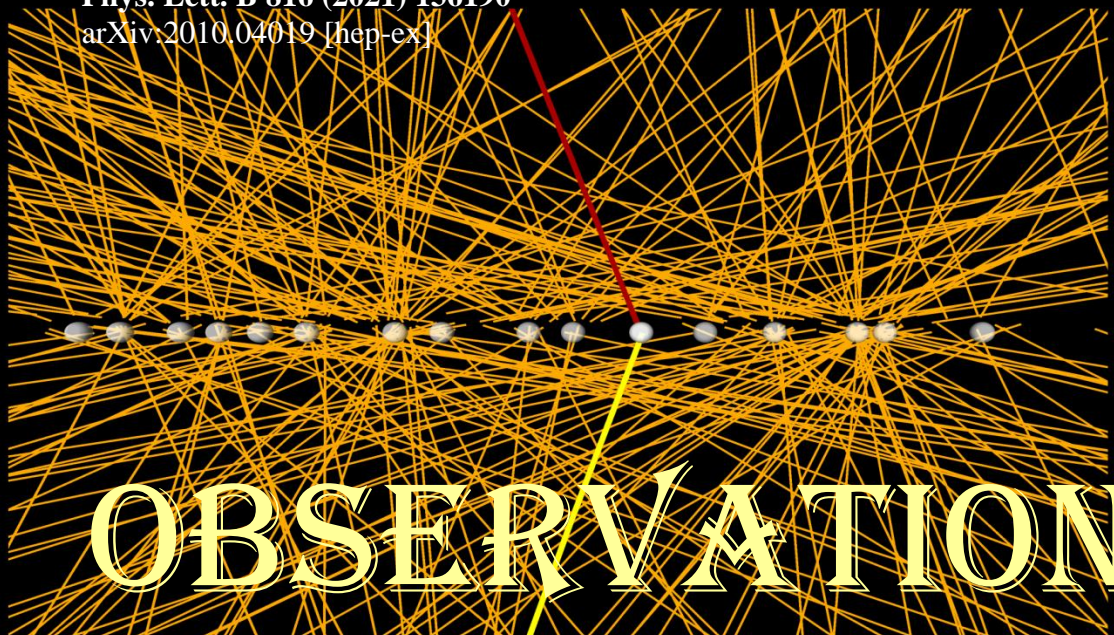
The same-flavour, opposite-charge pair with an invariant mass closest to m_Z is selected as the primary pair in the event with m_{12} .

❑ The region dominated by $Z \rightarrow 4\ell$ production is used to extract the most precise measurement of the $Z \rightarrow 4\ell$ branching fraction to date,

$$B_{Z \rightarrow 4\ell} = [4.41 \pm 0.13 (\text{stat}) \pm 0.23 (\text{syst}) \pm 0.09 (\text{theory}) \pm 0.12 (\text{lumi})] \times 10^{-6}$$

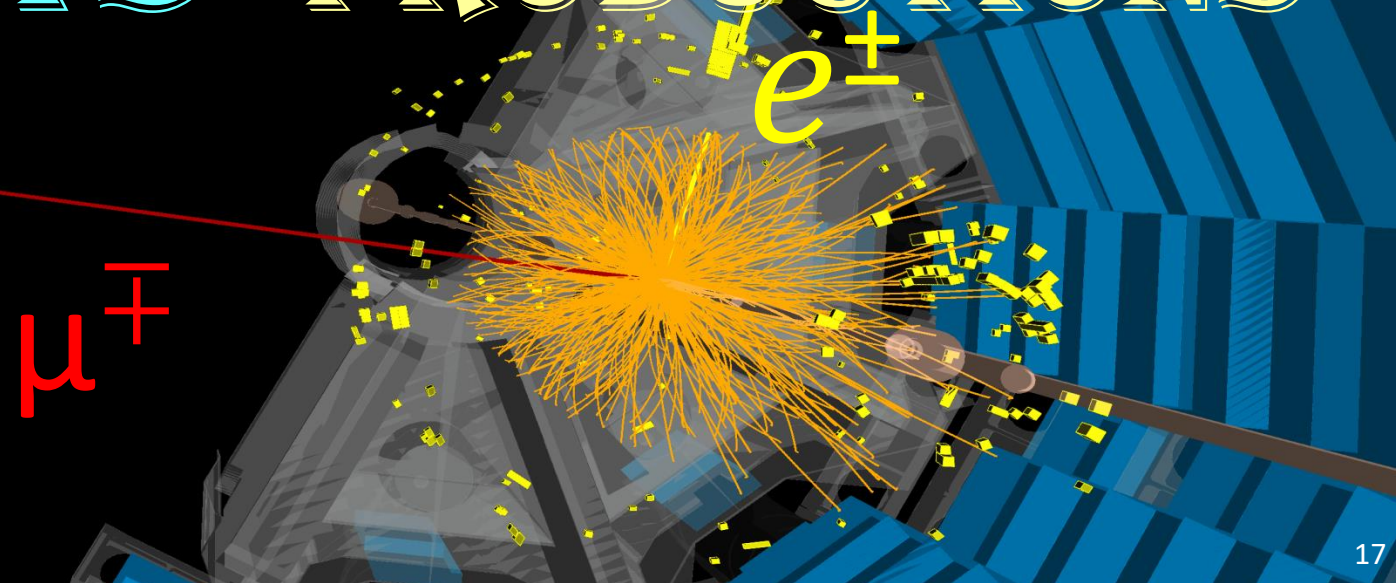
This is consistent with previous exp. and with the **Powheg** prediction

$$B_{Z \rightarrow 4\ell} = [4.50 \pm 0.01] \times 10^{-6}$$



μ^\pm $\gamma\gamma \rightarrow W^+W^- \rightarrow e^\pm \nu \mu^\mp \nu$
SR: $n_{\text{trk}} = 0$ & $p_{\text{T}}^{e\mu} > 30$ GeV
for the dilepton system

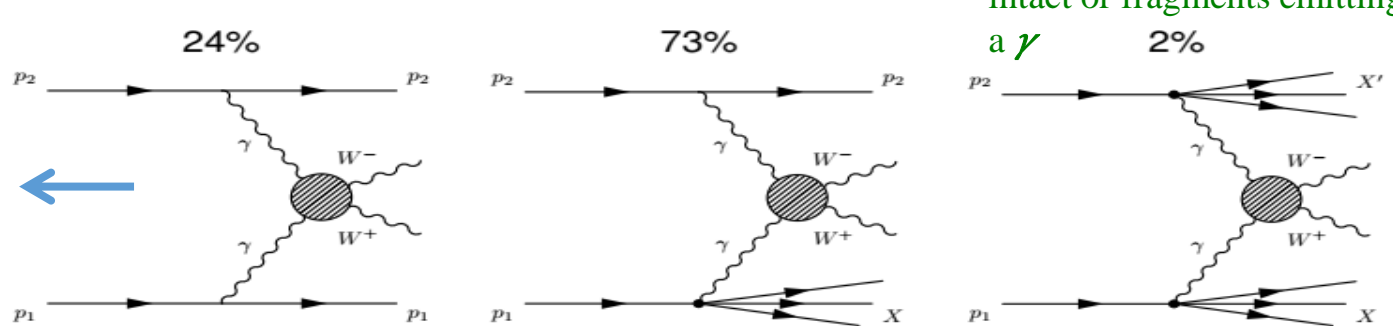
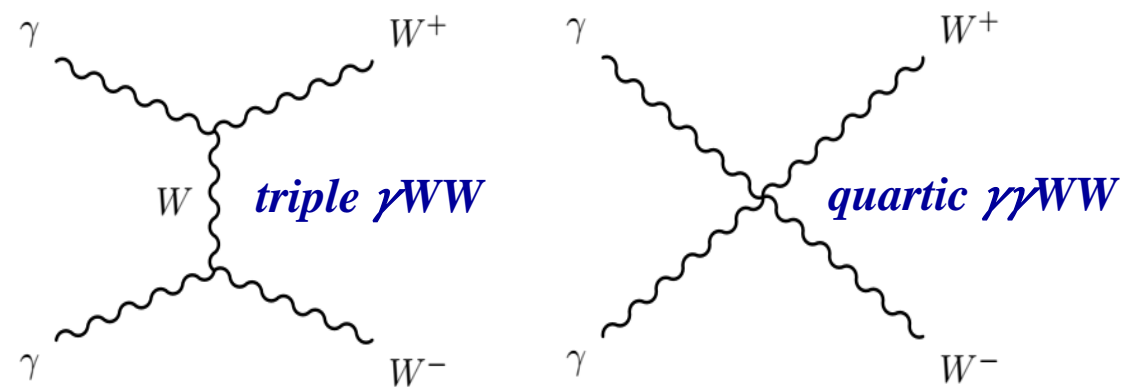
OBSERVATION OF $\gamma\gamma \rightarrow W^+W^-$ & $W^+W^- \rightarrow \geq 1$ JETS PRODUCTIONS



2. CROSS-SECTION $PP(\gamma\gamma) \rightarrow P^{(*)} W^+ W^- P^{(*)}$, $W^+ W^- \rightarrow e^\pm \nu \mu^\mp \nu$ #1

➤ The observation of **photon-induced production of W -boson pairs**, $\gamma\gamma \rightarrow W^+ W^- \rightarrow e^\pm \nu \mu^\mp \nu$

❑ This is **unique process**: it only involves diagrams with **self-couplings** of the **EW gauge bosons** $p^{(*)}$ final-state proton stays intact or fragments emitting a γ



Diagrams for the exclusive $\gamma\gamma \rightarrow W^+ W^-$ production representing the

I. elastic process

II. single-dissociation where one initial proton dissociates (SD)

III. double-dissociation where both protons fragment (DD)

- ❖ Test of the **EW sector** of the **SM**
- ❖ Direct access to **triple γWW** and **quartic $\gamma\gamma WW$** interactions, $O(a_{EM}^2)$

n_{trk} $p_T^{e\mu}$	Signal region		Control regions			
	SR	$n_{\text{trk}} = 0$	CR	CR	$1 \leq n_{\text{trk}} \leq 4$	CR
	> 30 GeV		< 30 GeV	> 30 GeV	< 30 GeV	
$\gamma\gamma \rightarrow WW$	174 ± 20		45 ± 6	95 ± 19	24 ± 5	
$\gamma\gamma \rightarrow \ell\ell$	5.5 ± 0.3		39.6 ± 1.9	5.6 ± 1.2	32 ± 7	
Drell–Yan	4.5 ± 0.9		280 ± 40	106 ± 19	4700 ± 400	
$qq \rightarrow WW$ (incl. gg and VBS)	101 ± 17		55 ± 10	1700 ± 270	970 ± 150	
Non-prompt	14 ± 14		36 ± 35	220 ± 220	500 ± 400	
Other backgrounds	7.1 ± 1.7		1.9 ± 0.4	311 ± 76	81 ± 15	
Total	305 ± 18		459 ± 19	2460 ± 60	6320 ± 130	
Data	307		449	2458	6332	

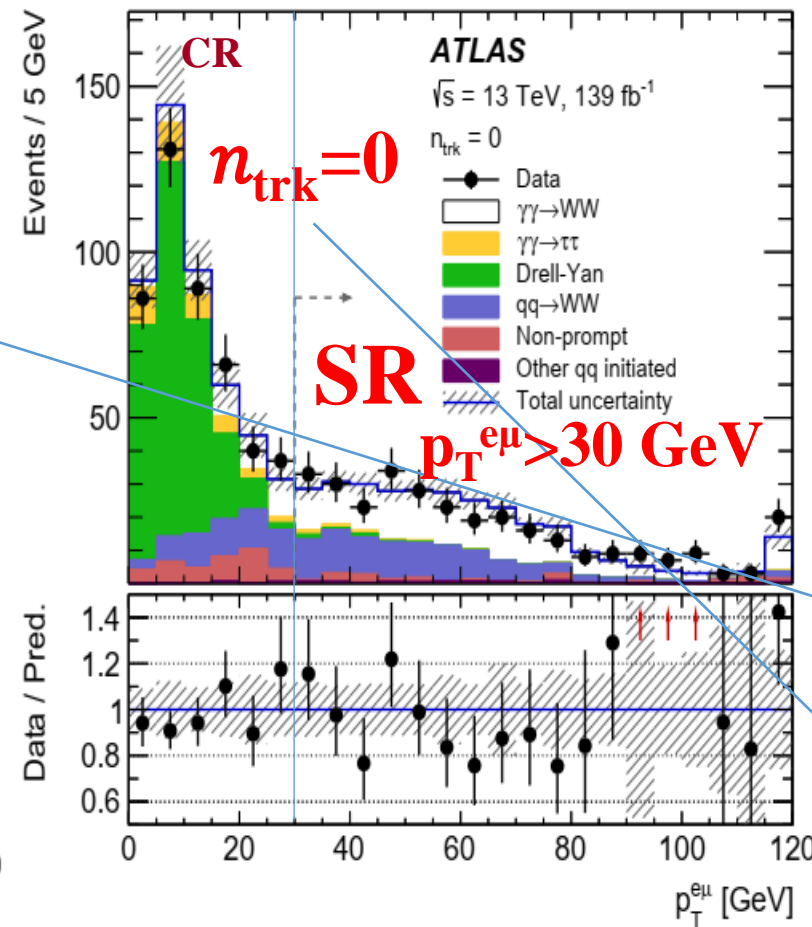
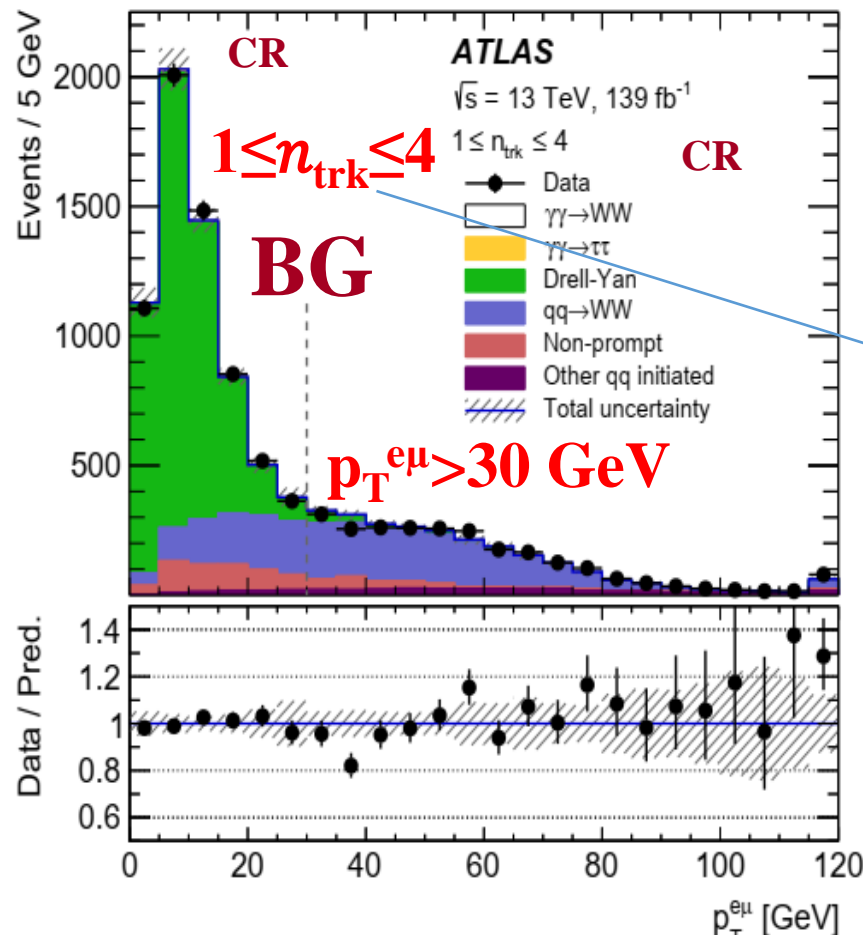
Agreement

Run 1 **evidence** of this process has turned into **observation** in Run2 at 8 TeV

✓ ATLAS: Phys. Rev. D 94 (2016) 032011

✓ CMS: JHEP 08 (2016) 119

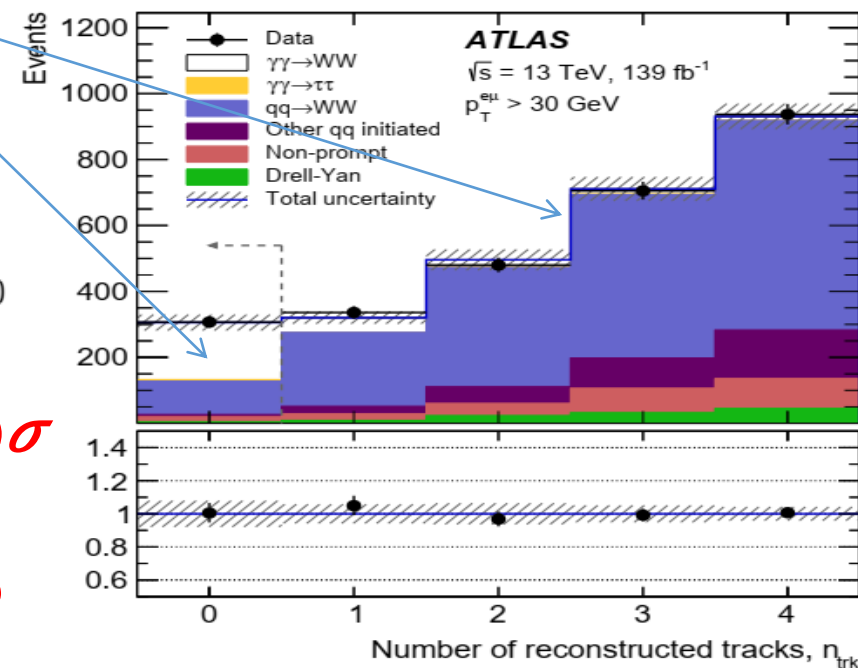
Summary of the data event yields, the predicted signal and **BG** event yields in the **Signal Region (SR)** and **Control Regions (CR)** as obtained after the fit



The **interaction vertex** is reconstructed from the **leptons** as the weighted average z -position of the tracks extrapolated to the beam line:

$$z_{\text{vtx}}^{\ell\ell} = \frac{z_{\ell_1} \sin^2 \theta_{\ell_1} + z_{\ell_2} \sin^2 \theta_{\ell_2}}{\sin^2 \theta_{\ell_1} + \sin^2 \theta_{\ell_2}}$$

The weighted mean **z -position** of the tracks extrapolated to the beam line. n_{trk} : number of tracks in a window $\Delta z = \pm 1 \text{ mm}$ around $z_{\text{vtx}}^{\ell\ell}$ excluding the tracks from leptons



The background-only hypothesis is rejected with a significance of **$8.4 (5)\sigma$**

□ The **WW observation** of **photon-induced** production in pp collisions (evidence – previously reported) **$\sigma^{\text{fid}} = 3.13 \pm 0.31 (\text{stat}) \pm 0.28 (\text{syst}) \text{ fb}$**

✓ The result is in agreement with the theoretical predictions

$\sigma_{\text{theo}} = 4.3 \pm 1.0 (\text{scale}) \pm 0.1 (\text{PDF}) \text{ fb}$ by MG5_aMC@NLO+Pythia8

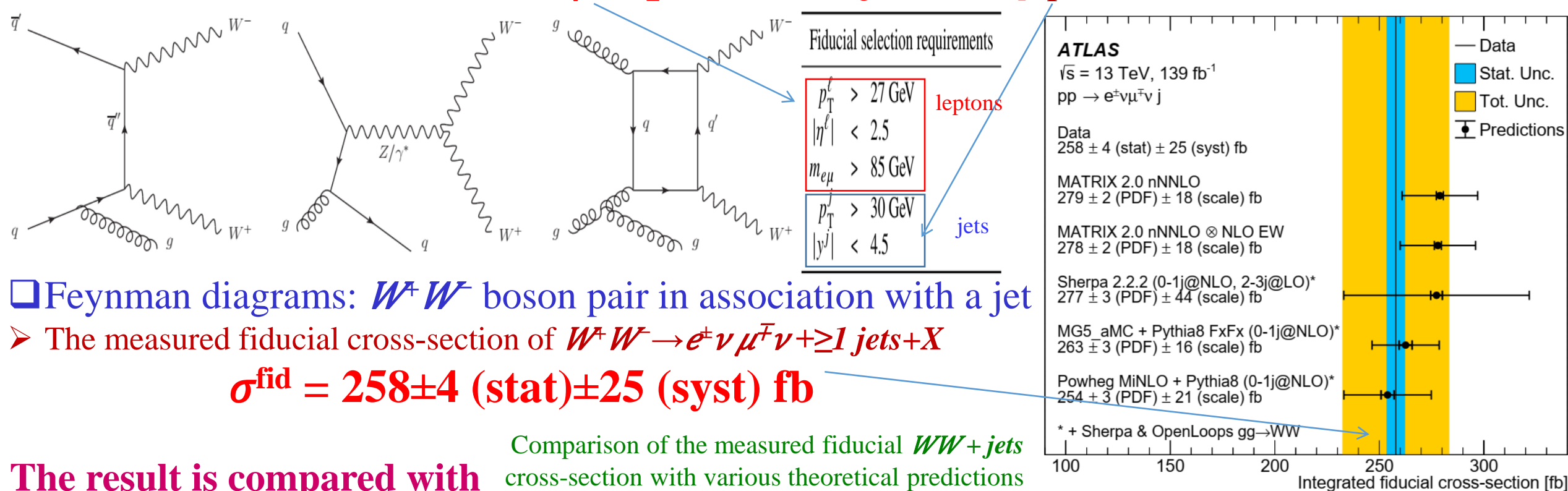
Without additional tracks from hadronic activity and from close-by pileup interactions

3. CROSS-SECTION $pp \rightarrow W^+ W^- + \geq 1 \text{ JETS} + X$ #1

➤ **Fiducial & differential cross-section measurements of $W^+ W^-$ production with ≥ 1 jets**

➤ It is **sensitive** to the properties of **EW-boson self-interactions**; provide a test of **pQCD & EW theory**

□ Events are selected with **one $e^\pm \mu^\mp$ pair & ≥ 1 jets** with $p_T^{\text{jet}} > 30 \text{ GeV}$ and $|y^{\text{jet}}| < 4.5$



□ Feynman diagrams: $W^+ W^-$ boson pair in association with a jet

➤ The measured fiducial cross-section of $W^+ W^- \rightarrow e^\pm \nu \mu^\mp \bar{\nu} + \geq 1 \text{ jets} + X$

$$\sigma^{\text{fid}} = 258 \pm 4 \text{ (stat)} \pm 25 \text{ (syst) fb}$$

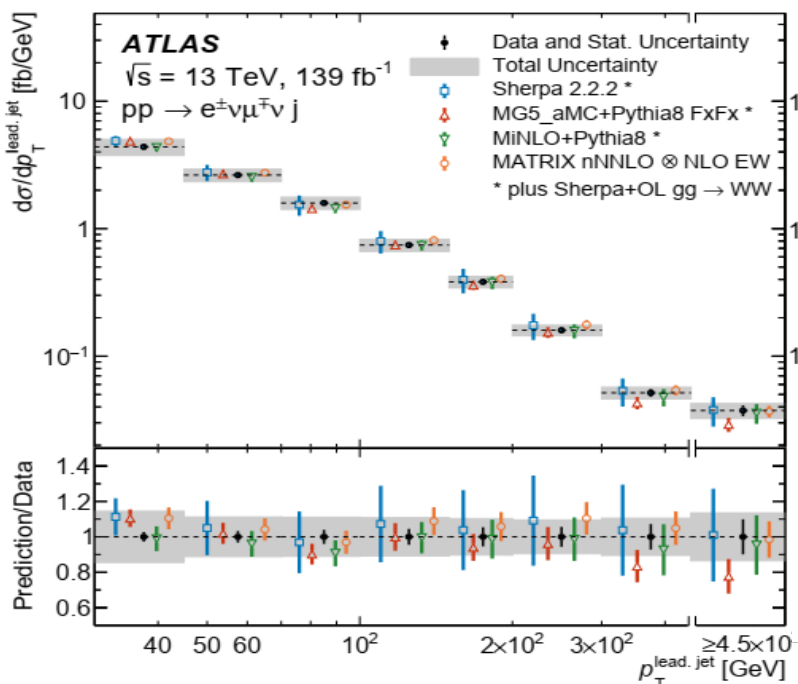
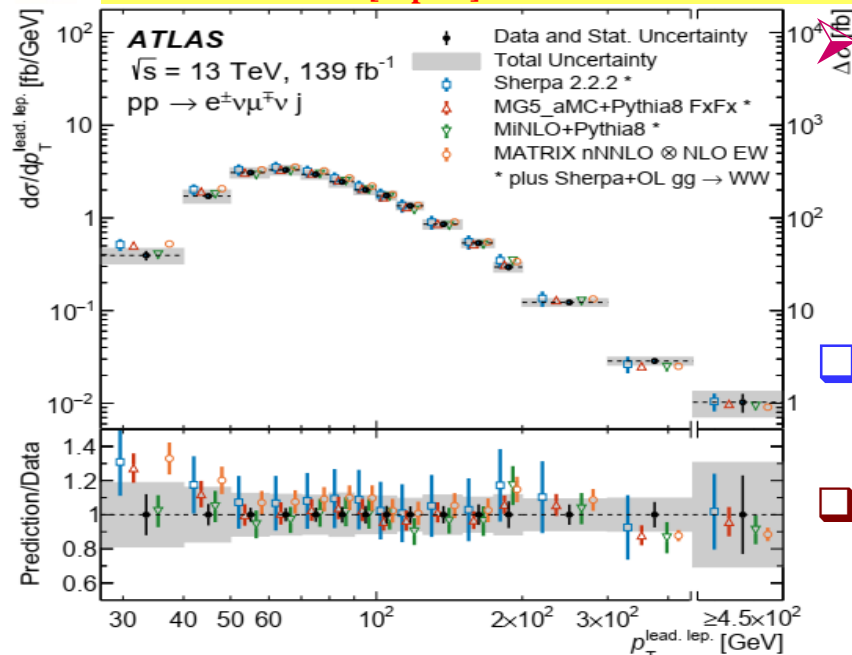
Comparison of the measured fiducial $WW + \text{jets}$ cross-section with various theoretical predictions

The result is compared with

- fixed-order parton-level prediction from **MATRIX 2.0** that is accurate to **NNLO (NLO)** for $qq^- \rightarrow WW (Wgg \rightarrow W)$ product
- prediction that additionally accounts for **EW** corrections to $WW + \text{jet}$ production: calculated with **Sherpa 2.2.2+OpenLoops**
- predictions from **Sherpa 2.2.2, MadGraph5_aMC@NLO+Pythia8** with **FxFx** merging, and **Powheg MiNLO+Pythia8**, which are all supplemented by a **Sherpa 2.2.2+OpenLoops $gg \rightarrow WW$ LO+PS** prediction

➤ **The data agree well with all MC predictions**

3. CROSS-SECTION $PP \rightarrow W^+ W^- + \geq 1 \text{ JETS} + X$ #2



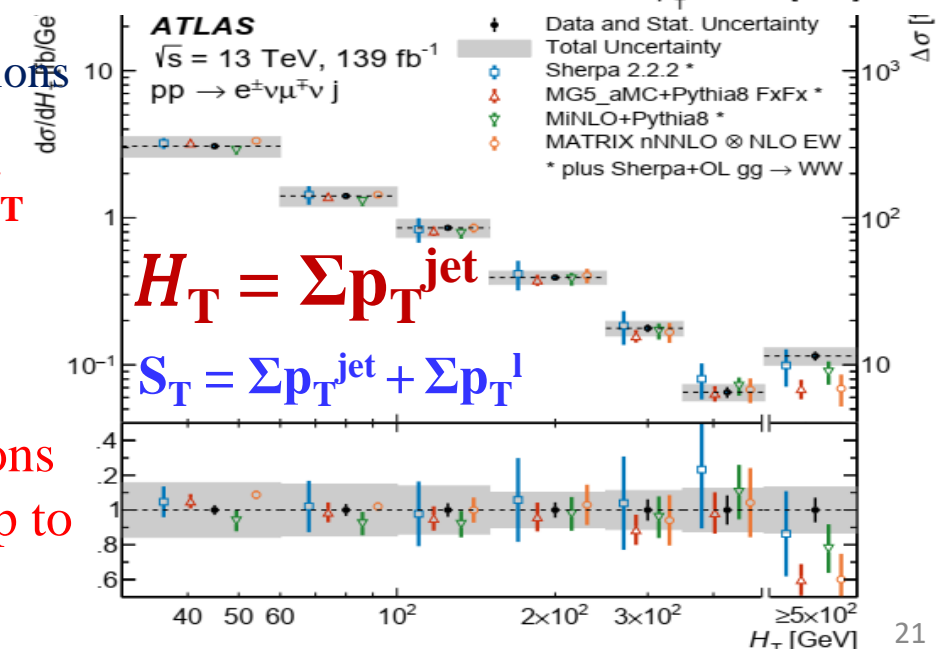
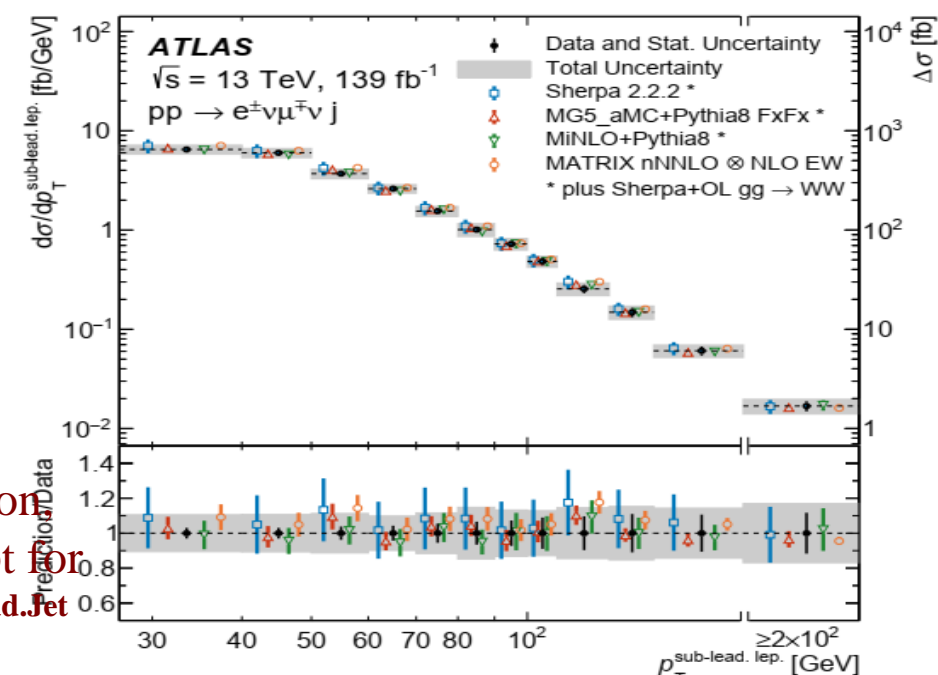
Measured fiducial cross-sections of WW +jets for:
 $p_T^{\text{lead.lep.}}$, $p_T^{\text{sub-lead.lep.}}$,
 $p_T^{\text{lead.Jet}}$, H_T

□ All predictions give an **excellent description** of the observed data

□ For the nominal **Sherpa 2.2.2** prediction values of χ^2/ndf are below **one**, except for the $m_{e\mu}$ distribution measured for $p_T^{\text{lead.Jet}} > 200 \text{ GeV}$, for which the value is **1.4**

□ Comparisons of the remaining predictions with the data yield similar χ^2 values, except for the jet multiplicity, H_T & S_T distributions, where, for the **highest multiplicities and energies**, small discrepancies exist: data / predictions

➤ The **data agree well with predictions in all differential distributions**, up to the highest p_T and for ≤ 5 jets



RAREST ELECTROWEAK

$Z(\rightarrow LL)\gamma JJ$ & $Z(\rightarrow \nu\nu)\gamma JJ$

MET PRODUCTION

jet

γ

$Z(\rightarrow \nu\nu)\gamma jj$ candidate event:

- $p_T^\gamma = 64$ GeV,
- $E_T^{\text{miss}} = 264$ GeV
- $m_{jj} = 2.6$ TeV

jet

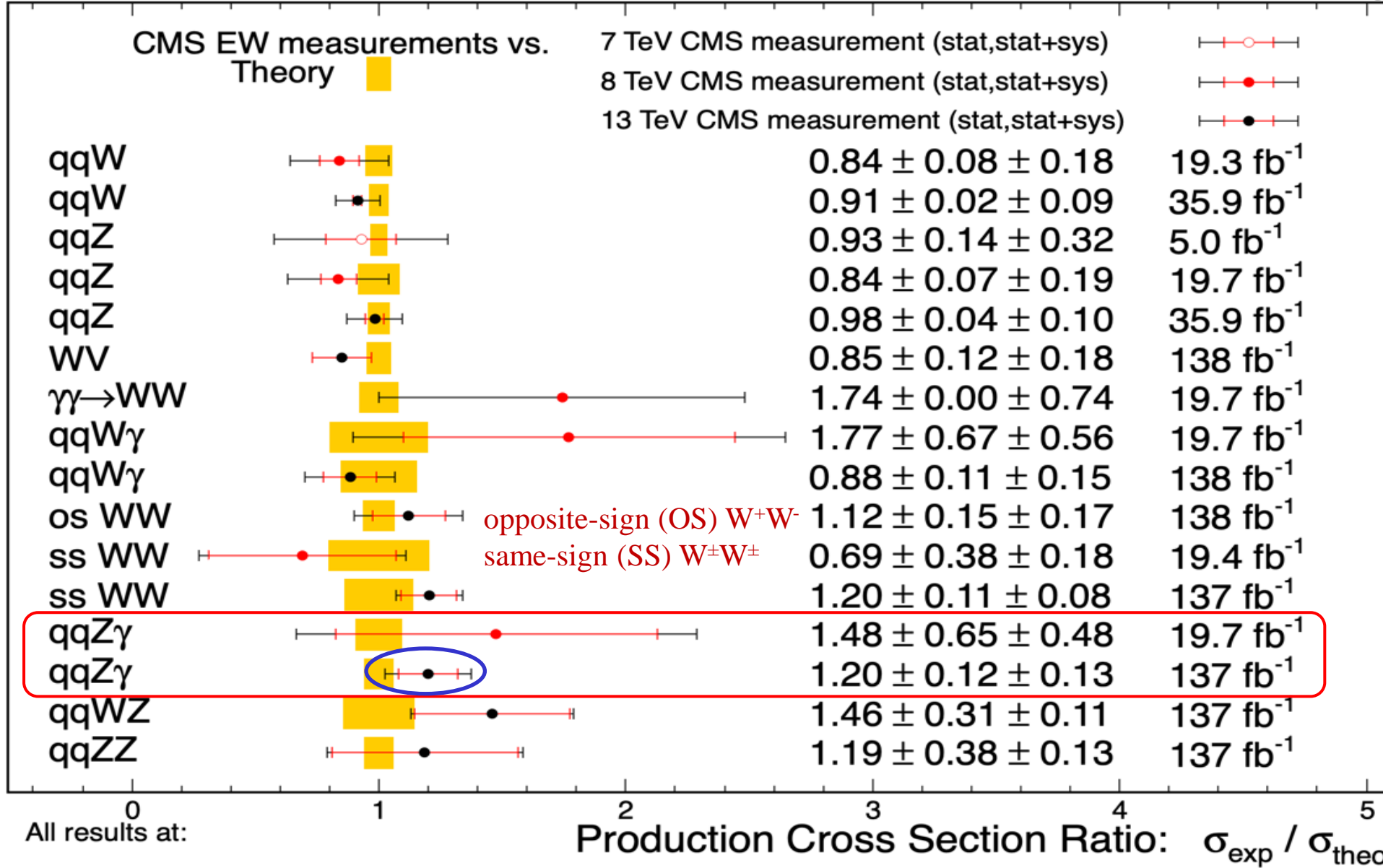
$Z(\rightarrow \nu\nu)\gamma JJ$

MET

jet

May 2022

CMS Preliminary



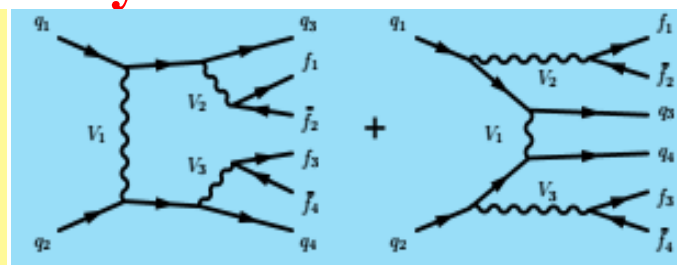
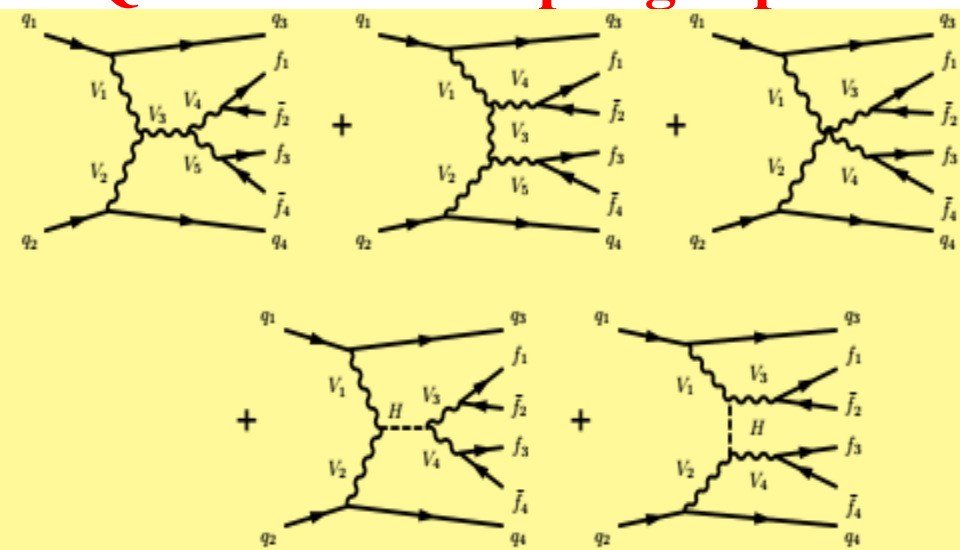
Summary of the cross sections of pure Electroweak (EW) interactions among the gauge bosons presented as a ratio compared to theory

JHEP 06 (2020) 076

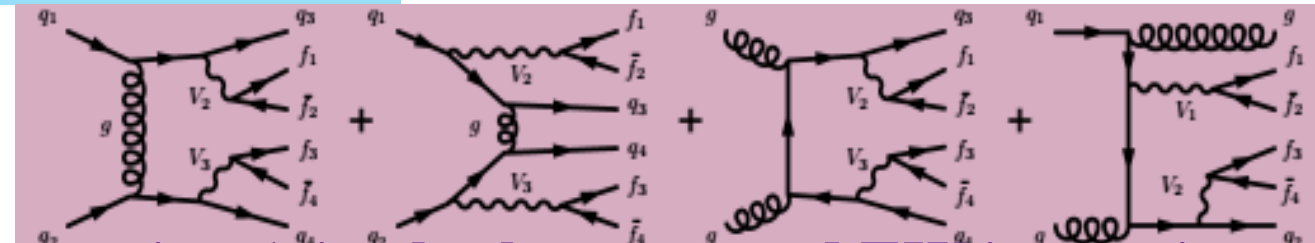
Eur. Phys. J. C 82 (2021) 105

➤ Vector boson scattering (VBS) processes ($VV \rightarrow VV$ with $V = W/Z/\gamma$)

➤ Quartic EW coupling experimentally accessible in electroweak production of $VVjj$



□ Purely EW interactions without self interactions



□ Processes involving both strong and EW interactions

□ Purely EW interactions involving only cubic and quartic self interactions

$Z\gamma$ production

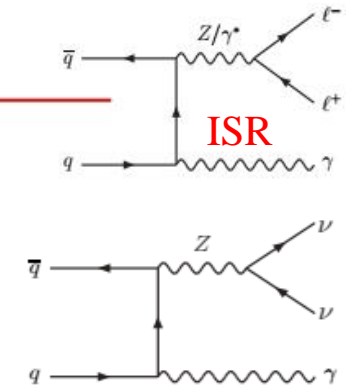
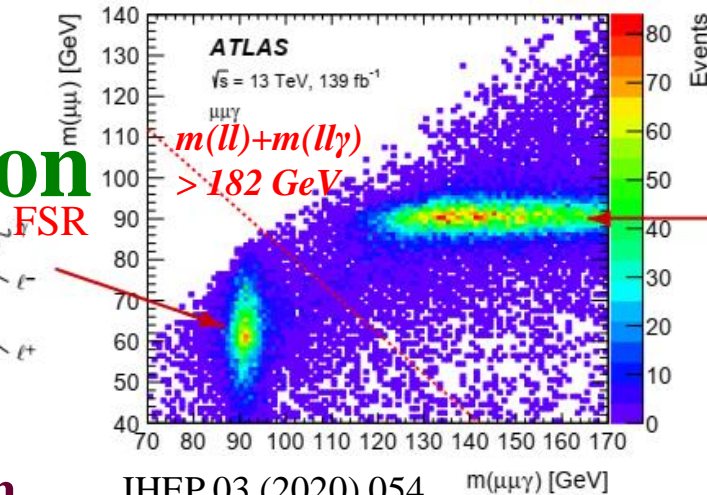
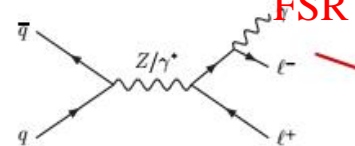
The $Z \rightarrow ll$ & $Z \rightarrow \nu\nu$ decays are studied

$Z \rightarrow ll$ large contributions from FSR off of leptons

□ Increase sensitivity to EW couplings, by removing FSR events

➤ Selection on the masses of the ll and $ll\gamma$ systems: $m_{ll} + m_{ll\gamma} > 2m_Z$

➤ Using $Z \rightarrow \nu\nu$ decays



Two-dimensional distribution of $m(ll)$ and $m(ll\gamma)$ for events satisfying all $\mu^+\mu^-\gamma$ selection criteria except that on the sum of $m(ll)$ and $m(ll\gamma)$. The photon is emitted from an initial-state quark

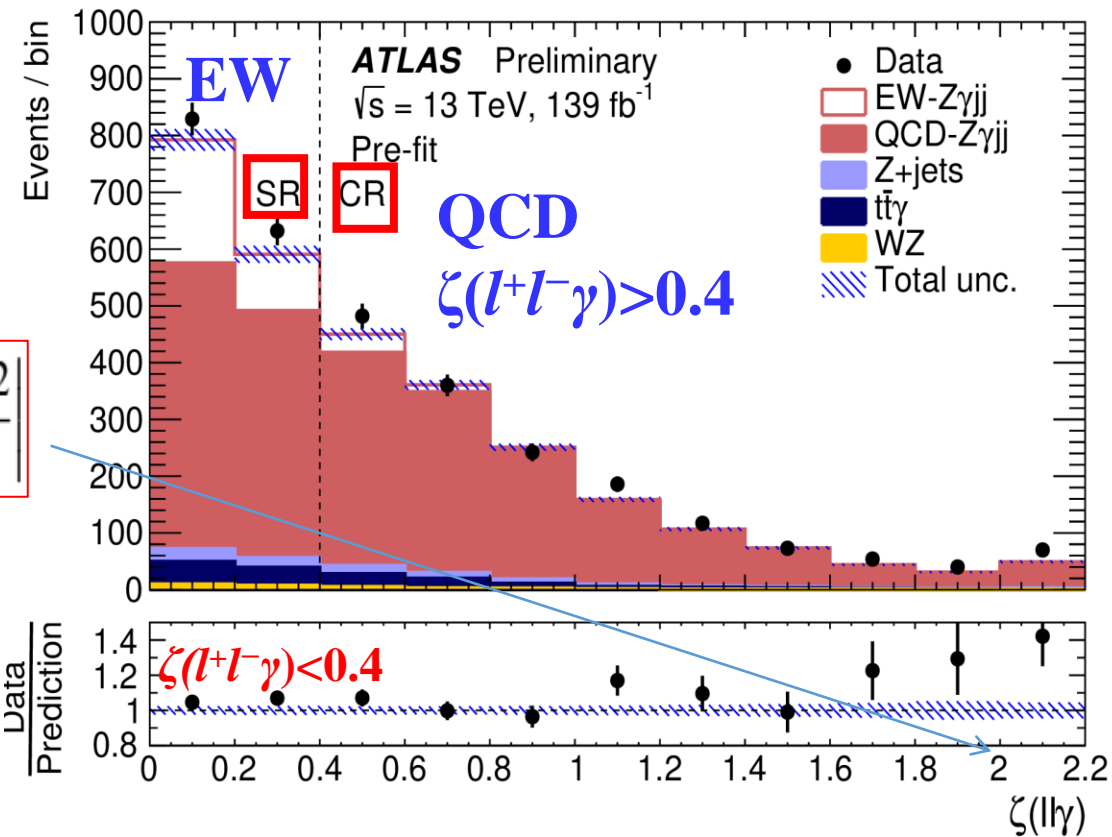
- The cross-section of the **EW production**: sensitivity to the **gauge boson self-interactions**
- ❑ Improved constraints probe scales of **new physics** in the **multi-TeV range** and provide a way to look for **signals of new physics in a model-independent way**

Lepton	$p_T^\ell > 20, 30(\text{leading}) \text{ GeV}, \eta_\ell < 2.47$ $N_\ell \geq 2$
Photon	$E_T^\gamma > 25 \text{ GeV}, \eta_\gamma < 2.37$ $E_T^{\text{cone20}} < 0.07 E_T^\gamma$ $\Delta R(\ell, \gamma) > 0.4$
Jet	$p_T^{\text{jet}} > 50 \text{ GeV}, y_{\text{jet}} < 4.4$ VBS signature $ \Delta y > 1.0$ $m_{jj} > 150 \text{ GeV}$ remove jets if $\Delta R(\gamma, j) < 0.4$ or if $\Delta R(\ell, j) < 0.3$
Event	$m_{\ell\ell} > 40 \text{ GeV}$ $m_{\ell\ell} + m_{\ell\ell\gamma} > 182 \text{ GeV}$ Strong QCD-$Z\gamma jj$ separation $\zeta(\ell\ell\gamma) < 0.4$ $N_{\text{jets}}^{\text{gap}} = 0$

Selection criteria applied at particle level

$$\zeta(\ell\ell\gamma) = \left| \frac{y_{\ell\ell\gamma} - (y_{j1} + y_{j2})/2}{y_{j1} - y_{j2}} \right|$$

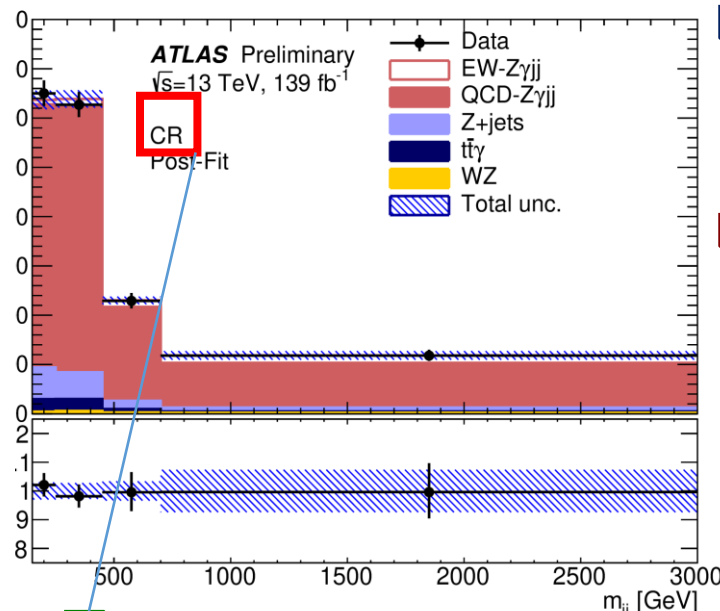
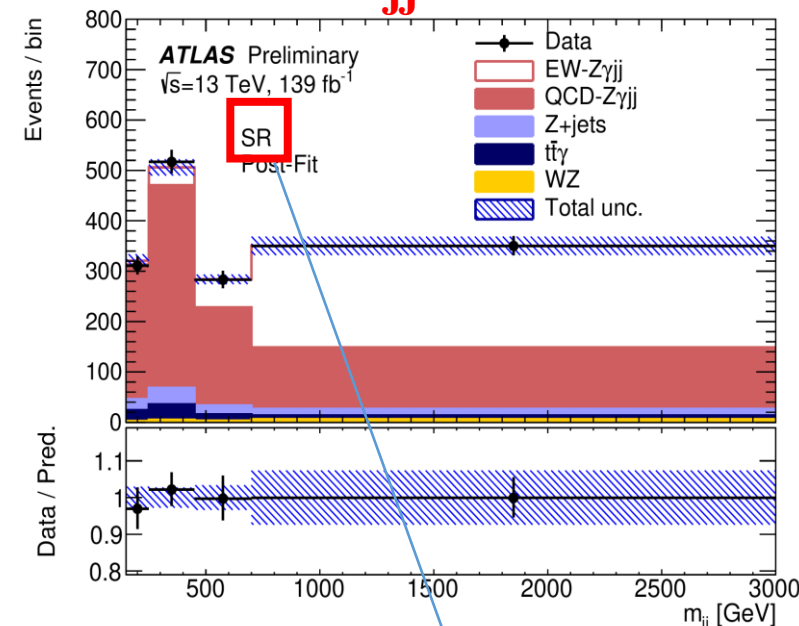
The **centrality** of the $l^+l^-\gamma$ system relative to the tagging jets (j_1, j_2)



- ❑ Events selected with high m_{jj} and high $|\Delta y_{jj}|$
- ❑ **Centrality** used to **control background** from Strong QCD - $Z\gamma jj$ production
- ❑ Background from **misidentified photons** estimated in data, background from $t\bar{t}\gamma$ validated in data

Centrality distributions, $\zeta(l^+l^-\gamma)$, in **Signal Region**, $\zeta(l^+l^-\gamma) < 0.4$, and **Control Region**, $\zeta(l^+l^-\gamma) > 0.4$, before the fit to extract the **EW- $Z\gamma jj$** component is performed. The sum of the signal and the various backgrounds is shown.

➤ Post-fit m_{jj} distributions in Signal Region (SR) & Control Region (CR)



□ To minimise dependence on theory modelling, the **high- $\zeta_{ll\gamma}$** is only used to **constrain the m_{jj} distribution**

□ The normalisation parameters of the **$QCD-Z\gamma jj$** background, **constrained by data in the SR & CR** are

$$\beta_{Z\gamma\text{-strong, CR}} = 1.00^{+0.18}_{-0.16}$$

$$\beta_{Z\gamma\text{-strong, SR}} = 1.06^{+0.17}_{-0.16}$$

□ **Background-only hypothesis rejected with 10σ (11σ)**

➤ Measured and theoretical **$Z\gamma$ -EW** cross sections:

$$\sigma_{EW}^{\text{exp}} = 4.49 \pm 0.40 \text{ (stat)} \pm 0.42 \text{ (syst) fb}$$

$$\sigma_{EW}^{\text{theor}} = 4.73 \pm 0.01 \text{ (stat)} \pm 0.15 \text{ (PDF)}^{+0.23}_{-0.22} \text{ (scale) fb}$$

(from MG5_aMC@NLO+Pythia8 at LO)

➤ Measured and theoretical **$Z\gamma$ -EW+QCD** cross sections:

$$\sigma_{EW+QCD}^{\text{exp}} = 20.6 \pm 0.6 \text{ (stat)}^{+1.2}_{-1.1} \text{ (syst) fb}$$

$$\sigma_{EW+QCD}^{\text{theor}} = 20.4 \pm 0.1 \text{ (stat)} \pm 0.2 \text{ (PDF)} \pm 2.2 \text{ (scale) fb}$$

(from MG5_aMC@NLO+Pythia8)

Summary of the observed number of events in data, N_{obs} , background (N_{Z+jets} , $N_{t\bar{t}\gamma}$, N_{WZ}), **EW- $Z\gamma jj$** signal ($N_{EW-Z\gamma jj}$) and **QCD $Z\gamma jj$** ($N_{QCD-Z\gamma jj}$) after the fit.

Sample	SR	CR
$N_{EW-Z\gamma jj}$	300 ± 36	55 ± 7
$N_{QCD-Z\gamma jj}$	987 ± 55	1352 ± 60
$N_{t\bar{t}\gamma}$	72 ± 11	59 ± 9
N_{WZ}	17 ± 3	14 ± 3
N_{Z+jets}	85 ± 30	143 ± 43
Total	1461 ± 38	1624 ± 40
N_{obs}	1461	1624

Agreement

Agreement

EW $Z(\rightarrow \nu\nu)\gamma jj$ selection

Observable	Requirements
N_{jet} with $p_T > 25$ GeV	≥ 2
$ \eta(j_{1,2}) $	< 4.5
$p_T(j_1)$ [GeV]	> 60
$p_T(j_2)$ [GeV]	> 50
$\Delta R(j, \ell)$	> 0.4
$ \Delta\eta_{jj} $	> 3.0
C_3	< 0.7
m_{jj} [TeV]	> 0.5
truth- E_T^{miss} [GeV]	> 150
$\Delta\phi(\text{truth-}\vec{E}_T^{\text{miss}}, j_i)$	> 1.0
$p_T(\gamma)$ [GeV]	$> 15, < 110$
$ \eta(\gamma) $	< 2.37
$E_T^{\text{cone20}}/E_T^\gamma$	< 0.07
$\Delta R(\gamma, \text{jet-or-}\ell)$	> 0.4
C_γ	> 0.4
$\Delta\phi(\text{truth-}\vec{E}_T^{\text{miss}}, \gamma)$	> 1.8
N_ℓ with $p_T > 4$ GeV and $ \eta < 2.47$	0

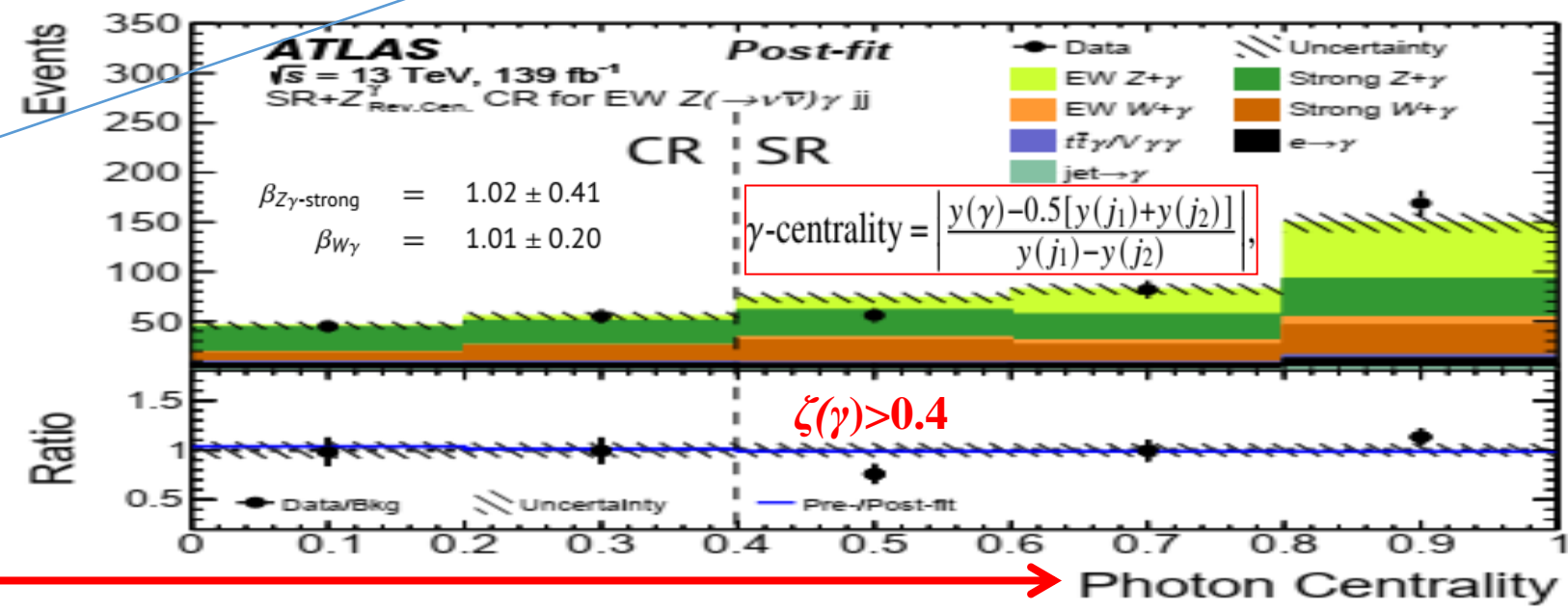
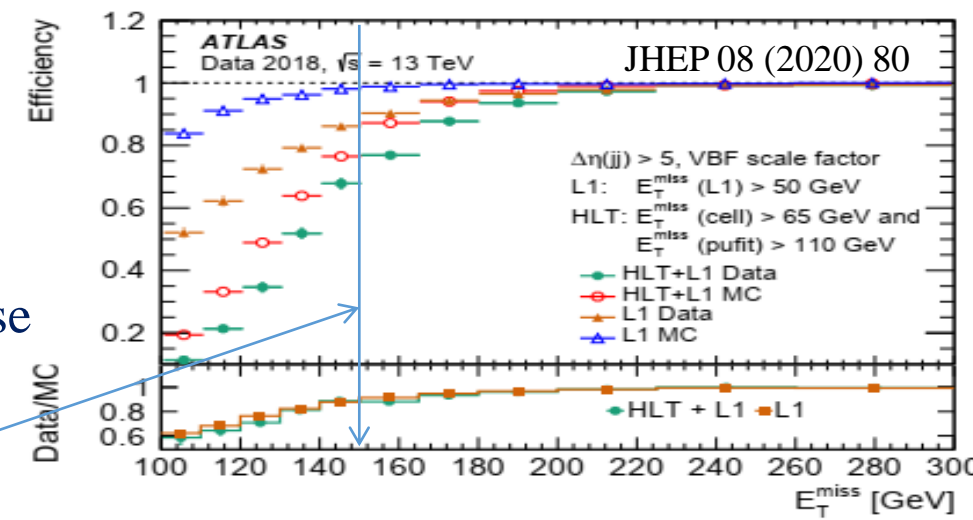
Vector boson scattering (VBS) signature

strong- Zjj separation

reject Wjj

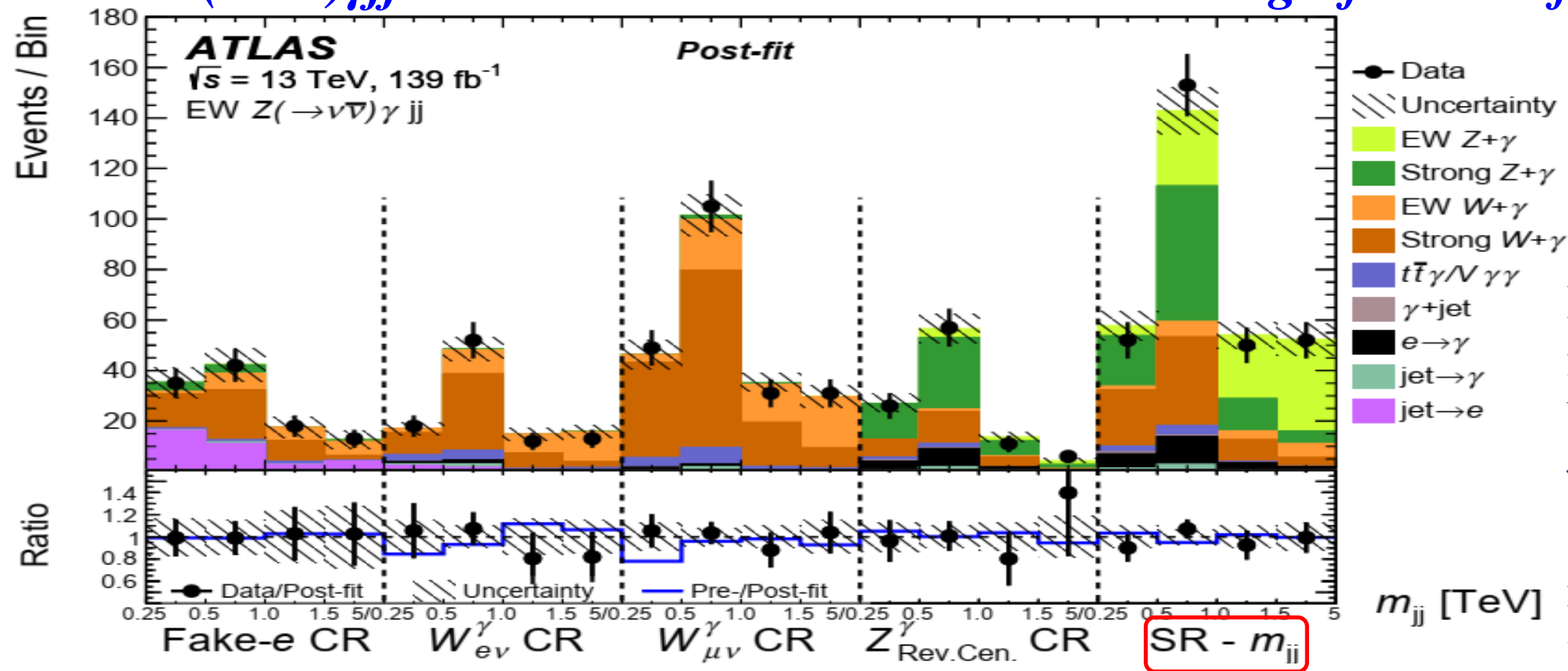
➤ Using the E_T^{miss} to trigger $Z(\rightarrow \nu\nu)\gamma$ events

❑ This usually doesn't impact sensitivity to **BSM** effects, as these are expected at high energy



Similar strategy to select **VBS** events; control strong $Z\gamma jj$ production
Additional BGs, mainly $W\gamma jj$, compared to $Z(\rightarrow ll)\gamma$ in data

➤ EW $Z(\rightarrow \nu\nu)\gamma jj$ is observed in this final state with a *significance of 5.2 (5.1) σ*



Post-fit results
for m_{jj} SR & CR
bins in the EW Z
 γ +jets cross-
section
measurement
with the $\mu_{Z\gamma EW}$
signal
normalization
floating

□ Measured and theoretical $Z\gamma$ -EW cross sections:

$$\sigma_{EW}^{meas.} = 1.31 \pm 0.20 (stat) \pm 0.20 (syst) fb$$

$$\sigma_{EW}^{theo.} = 1.27 \pm 0.01 (stat) \pm 0.17 (scale) \pm 0.03 (PDF) fb$$

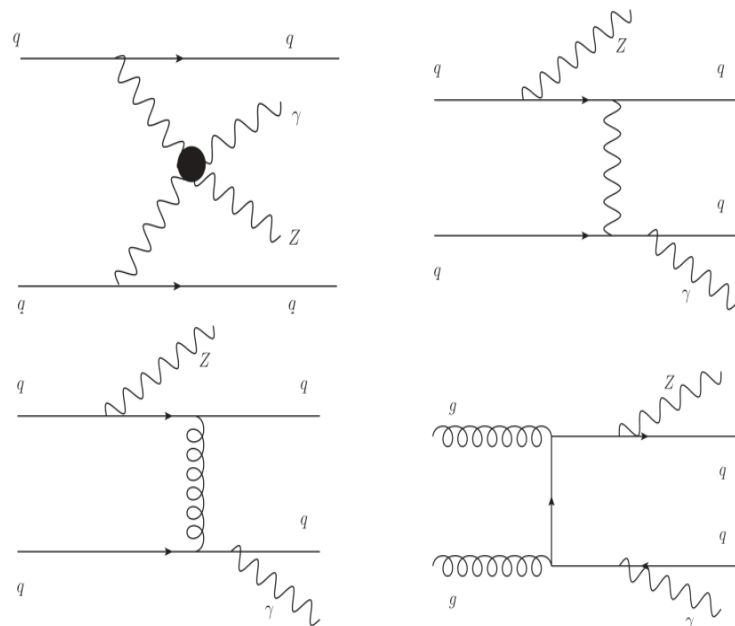
(from MG5_aMC@NLO+Pythia8 at LO, rescaled by 0.3% to VBFNLO)

➤ Largest sources of uncertainties in jet energy scale/resolution is 7.6% and in $V\gamma$ +jets modelling is 6.7%.

Agreement

- The **EW** production of $Z(\rightarrow \nu\nu)\gamma jj$ with a γ ($E_T > 150$ GeV), is a **Probe** of the **EW** symmetry breaking mechanism in SM & is sensitive to **quartic gauge boson couplings (QGS)** via **vector-boson scattering (VBS)**

Z boson branching ratio $Z \rightarrow \nu\nu$ is larger than the branching ratio $Z \rightarrow ll$; the background is under better control than in the hadronic decay channel



Photon centrality relative to the **two jets** with the **highest p_T** values in the event is defined as

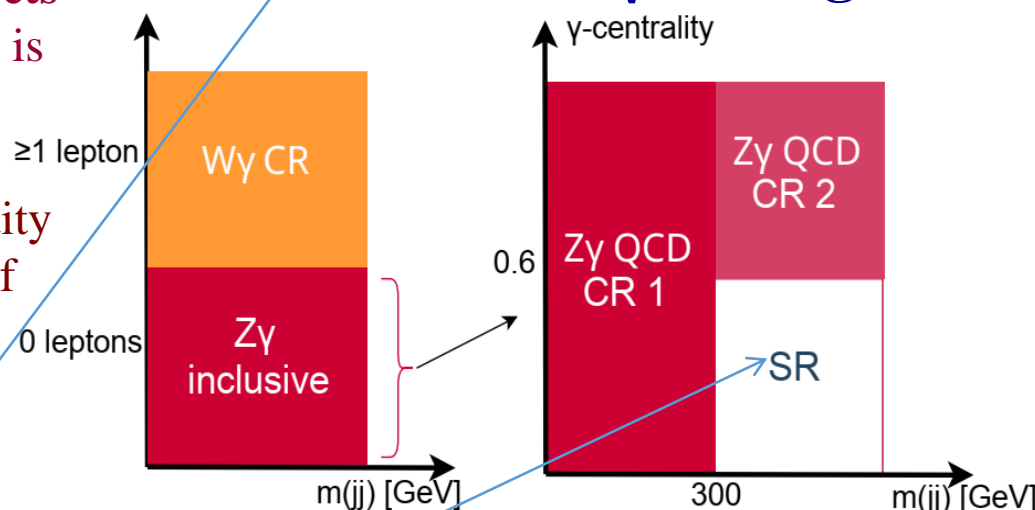
$$\gamma\text{-centrality} = \left| \frac{y(\gamma) - 0.5[y(j_1) + y(j_2)]}{y(j_1) - y(j_2)} \right|, \geq 1 \text{ lepton}$$

Where $y = 0.5 \times \ln[(E + p_z)/(E - p_z)]$ is the rapidity of the objects (p_z is the z-component of the momentum of a particle)

Fiducial region definition

Selections	Cut value
E_T^{miss}	> 120 GeV
E_T^γ	> 150 GeV
Number of isolated photons	$N_\gamma = 1$
Photon isolation	$E_T^{\text{cone40}} < 0.022 p_T + 2.45$ GeV, $p_T^{\text{cone20}}/p_T < 0.05$
Number of jets	$N_{\text{jets}} \geq 2$ with $p_T > 50$ GeV
Overlap removal	$\Delta R(\gamma, \text{jet}) > 0.3$
Lepton veto	$N_e = 0, N_\mu = 0$
$ \Delta\phi(\gamma, \vec{p}_T^{\text{miss}}) $	> 0.4
$ \Delta\phi(j_1, \vec{p}_T^{\text{miss}}) $	> 0.3
$ \Delta\phi(j_2, \vec{p}_T^{\text{miss}}) $	> 0.3
m_{jj}	> 300 GeV
γ -centrality	< 0.6

Definition of the $Z\gamma$ subregions

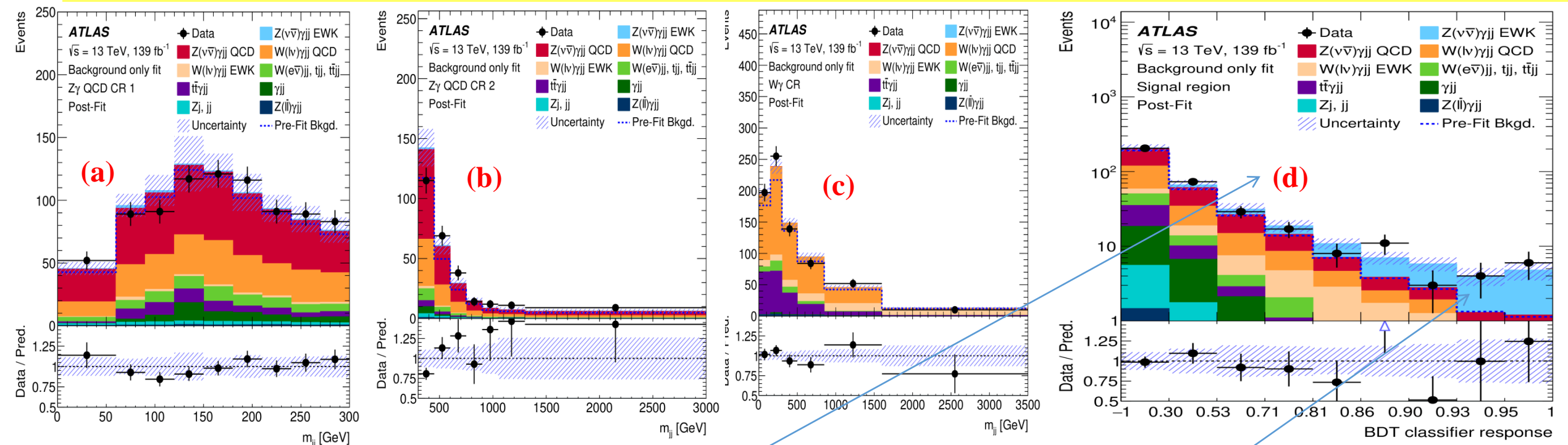


Feynman diagrams of electroweak

- ☐ $Z\gamma jj$ production involving the **VBS** subprocess (top left) or
- ☐ **non-VBS** subprocesses (top right)
- ☐ Of **QCD $Z\gamma jj$** production with gluon exchange (bottom left) or
- ☐ the **s-channel $gg-qq$** process (bottom right)

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- ☐ The **signal region (SR)** is required to have $m_{jj} > 300$ GeV and γ -centrality < 0.6 , where m_{jj} is defined as the invariant mass of the **2jets** with the highest values of p_T
- ☐ The **$Z\gamma$ QCD CR 1** requires events with $m_{jj} < 300$ GeV; it is used to estimate the $Z(\nu\nu)\gamma jj$ QCD background yield
- ☐ The **$Z\gamma$ QCD CR 2** has the same selection criteria as the **SR** but requires events with γ -centrality > 0.6 ; it is used to check for possible m_{jj} mismodelling. The values of the requirements are chosen to maximise the number of events and the purity of the targeted process in each region

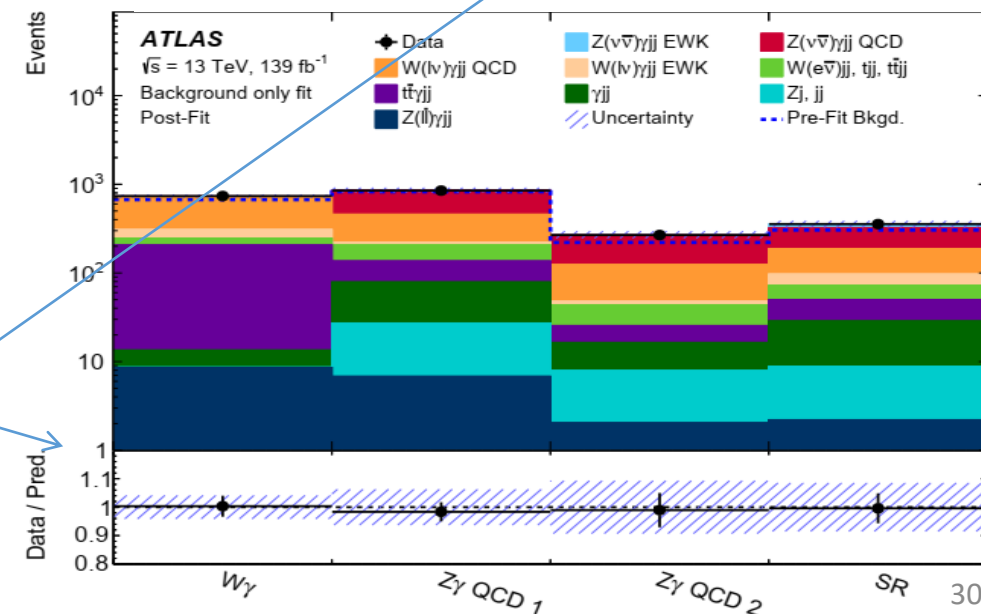


Summary of the yield for processes in all regions, after the fit. The dashed line shows the total background distribution before the fit.

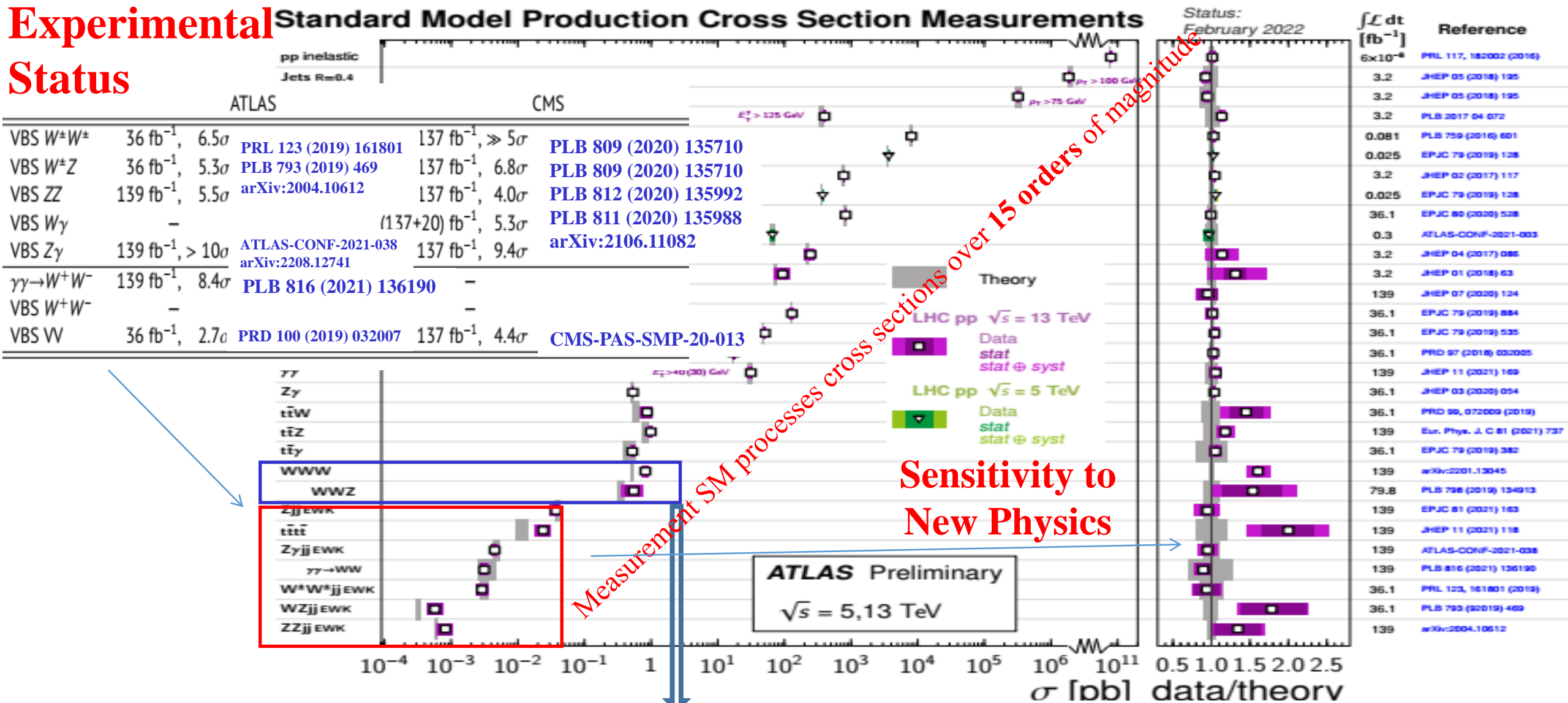
- The m_{jj} distributions for the
 - (a) CR 1: $Z\gamma$ QCD,
 - (b) CR 2: $Z\gamma$ QCD,
 - (c) CR: $W\gamma$, and the BDT classifier response distribution for the
 - (d) SR: after the fit in all regions
- The boosted decision tree (BDT) classifier response was remapped into equal-width bins for better representation. The dashed line shows the total background distribution before the fit.
- The signal selection uses a boosted decision tree (BDT)

➤ The observed fiducial cross section

$$\sigma_{Z\gamma\text{EWK}} = 0.77^{+0.34}_{-0.30} \text{ fb}$$



Experimental Status



Electroweak $VVjj$ production has been observed in all major channels

✓ They are amongst the **rarest** processes currently experimentally accessible



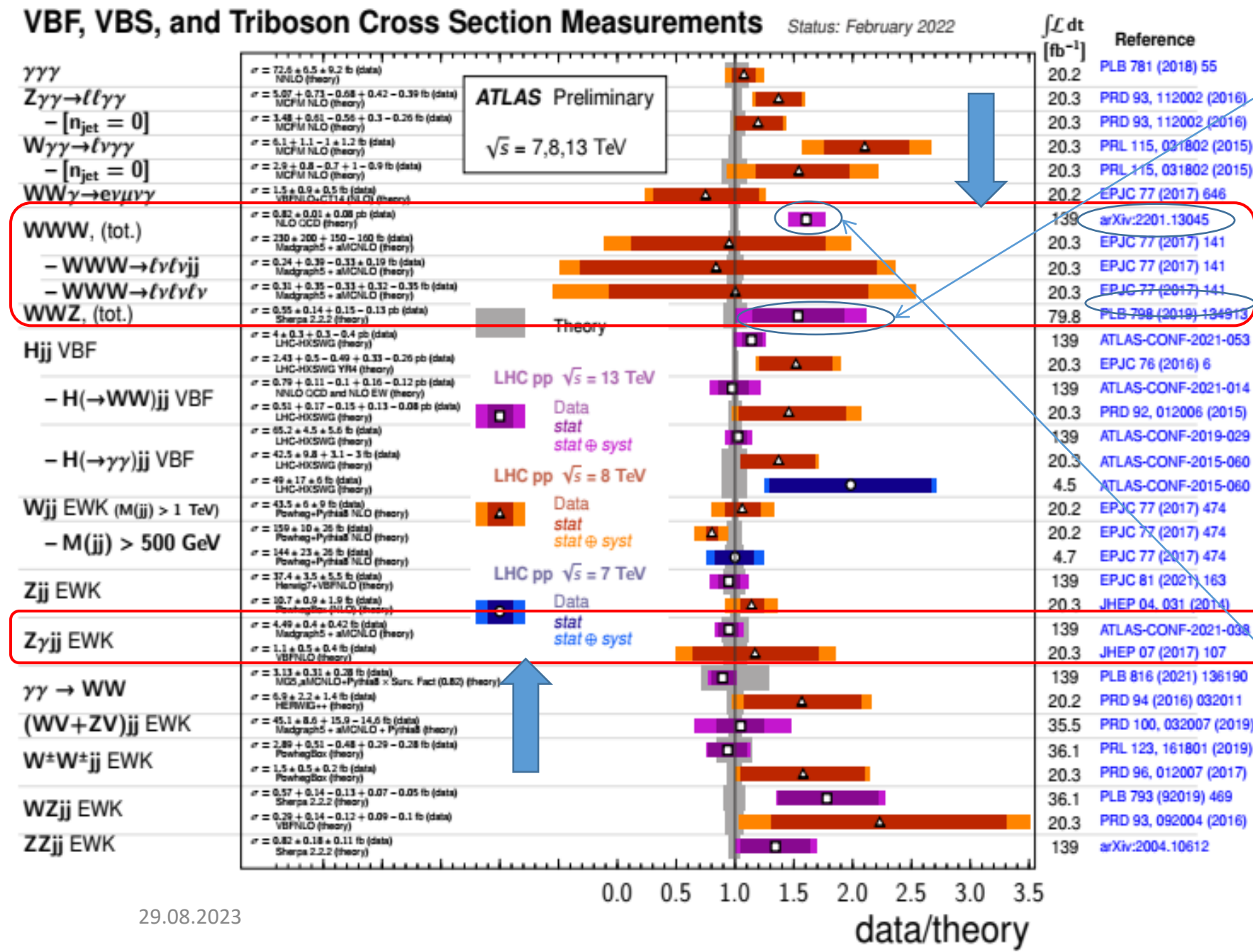
The background of the slide features a complex visualization of particle tracks within a detector. Two primary interaction vertices are highlighted with starburst patterns. The left vertex has tracks labeled e^+ (yellow) and MET (white). The right vertex has tracks labeled e^+ (yellow), μ^- (red), and MET (white). The text 'OBSERVATION OF WW PRODUCTION' is overlaid in a large, stylized font. The word 'OBSERVATION' is in yellow, 'OF' is in yellow, ' WW ' is in cyan, and 'PRODUCTION' is in yellow. The background is dark blue with various colored lines and shapes representing detector components and particle paths.

OBSERVATION OF WW PRODUCTION



Run: 349169
Event: 1043374730
2018-04-30 01:58:32 CEST

$WW \rightarrow e^+ \nu \ e^+ \nu \ \mu^- \nu$ candidate event
 $E_T^{\text{miss}} = 105 \text{ GeV}$



Experimental Status

1. Evidence for WVV from ATLAS

- ☐ $WWW \rightarrow \ell\nu\ell\nu qq$ ($\ell=e,\mu$)
- ☐ $WWW \rightarrow \ell\nu\ell\nu\ell\nu$
- ☐ $WWZ \rightarrow \ell\nu qq \ell\ell$
- ☐ $WWZ \rightarrow \ell\nu\ell\nu\ell\ell$
- ☐ $WZZ \rightarrow qq \ell\ell \ell\ell$

in 80 fb^{-1} (2015-2017 data)

- ☐ Evidence 4.1σ (WVV) and 3.2σ (WWW) PLB 798 (2019) 134913

New ATLAS result focuses on WWW in $\ell^+\nu\ell^+\nu jj$ and $\ell^+\nu\ell^+\nu\ell^+\nu$; Phys.Rev.Lett.129 (2022) 6, 061803; ArXiv 2201.13045

- ☐ $e^+e^+ + jj$,
- ☐ $e^+\mu^+ + jj$,
- ☐ $\mu^+\mu^+ + jj$
- ☐ $3l + E_T^{\text{miss}}$ with no same-flavour opposite-charge lepton pair

→ strong suppression of $VV + jets$

The $\sigma_{WWW+X} = 820 \pm 100$ (stat) ± 80 (syst) fb it is 2.6σ of predicted $\sigma_{WWW+X} = 511 \pm 18$ fb (NLO QCD and LO EW)

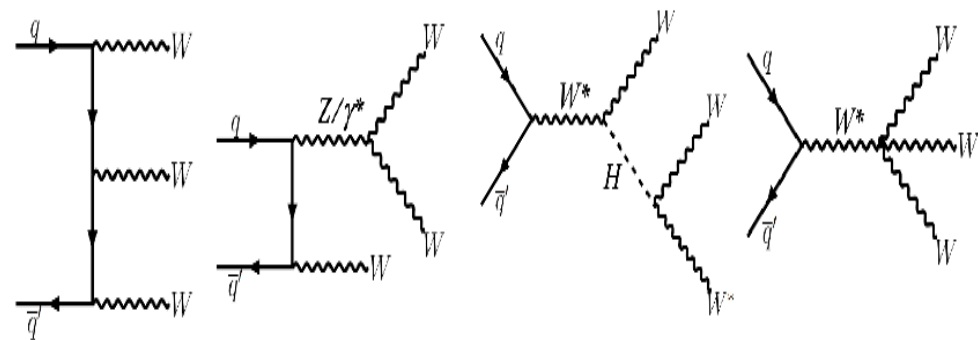
Observation of VVV from CMS

in 137 fb^{-1} (2015-2018 data)

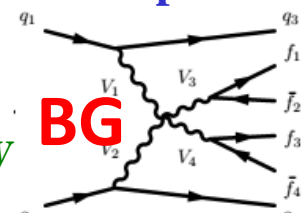
- ☐ 5.7σ (VVV), 3.3σ (WWW)

PRL 125 (2020) 151802

- The combined observed (expected) significance $pp \rightarrow WWW + X$ is found to be $8.0(5.4)\sigma$
- The inclusive cross section is 820 ± 100 (stat) ± 80 (syst) fb: 2.6σ from the predicted of 511 ± 18 fb (NLO QCD, LO EW accuracy)



Feynman diagrams at LO for the production of WWW bosons, including diagrams sensitive to triple & quartic gauge couplings



BG

WWW Simulation

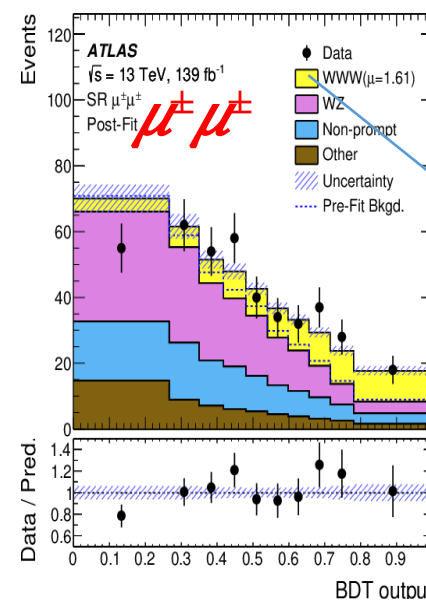
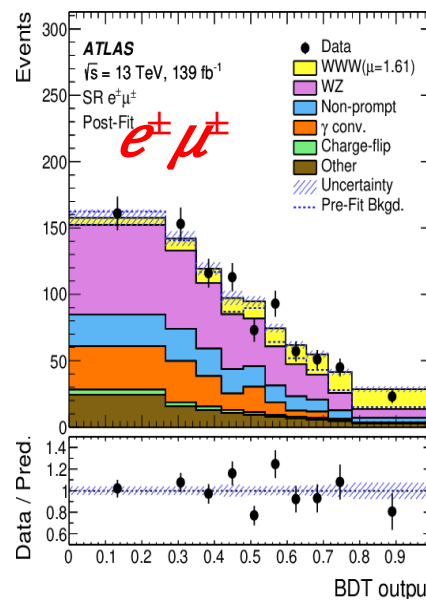
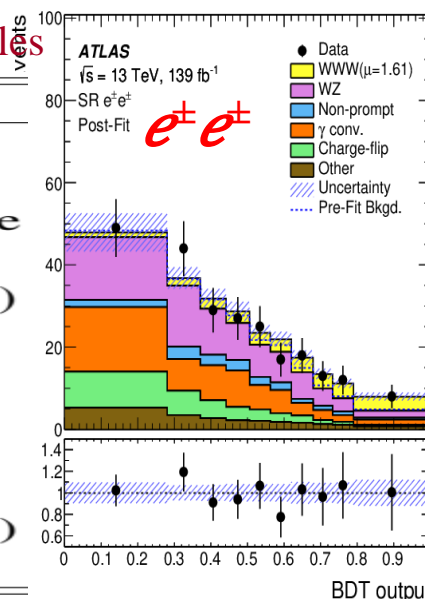
- on-shell WWW simulated with Sherpa 2.2.2 at NLO
- $WH \rightarrow WWW^*$ simulated with Powheg+Pythia8 at NLO
- spin correlations accounted for in W decays
- t - and u -channel production at $O(\alpha^6)$ as background

$WWW \rightarrow l^\pm \nu l^\pm \nu jj$ is distinguished from VBS signatures by requiring at least 2 central jets with $m_{jj} < 160$ GeV & $\Delta\eta_{jj} < 1.5$

Signal is measured in a fit to a *boosted-decision-tree (BDT)* discriminant

12 BDT training variables

2ℓ
$ m_{jj} - m_W $
p_T (forward jet)
E_T^{miss} significance
$p_T(j_2)$
minimum $m(\ell, j)$
$m(\ell_2, j_1)$
$N(\text{jets})$
$p_T(\ell_2)$
$m_{\ell\ell}$
$ \eta(\ell_1) $
$N(\text{leptons in jets})$
$m(\ell_1, j_1)$



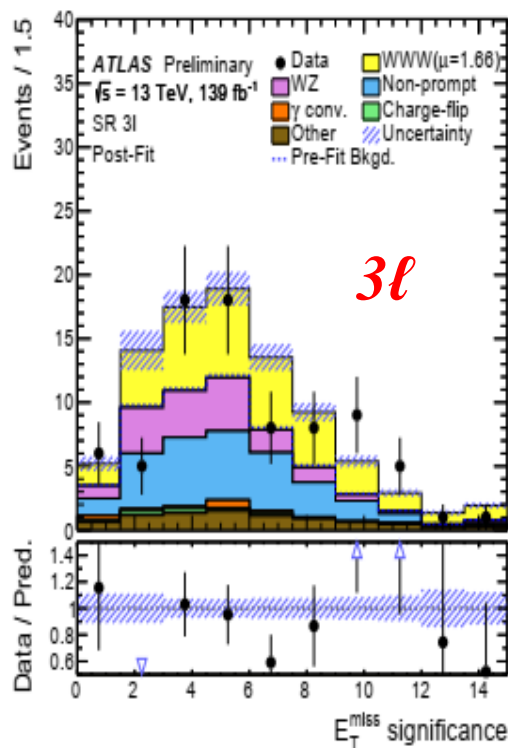
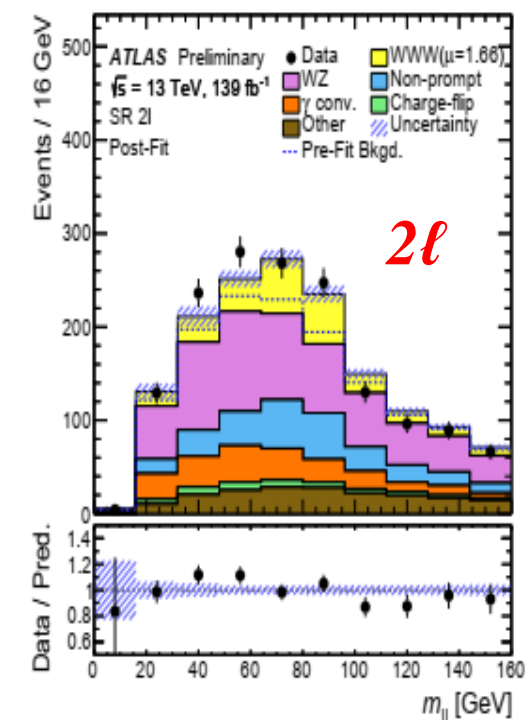
Number of events for postfit signal, BG & data observed in the 2ℓ & 3ℓ SRs

	e^+e^+	$e^+\mu^+$	$\mu^+\mu^+$	3ℓ
WWW signal	28.4 ± 4.3	124 ± 19	82 ± 12	34.8 ± 5.2
WZ	81.1 ± 5.7	346 ± 22	170 ± 10	16.4 ± 1.5
Charge-flip	31.1 ± 7.3	19 ± 5	...	1.7 ± 0.4
γ conversions	60.8 ± 8.5	139 ± 15	...	1.5 ± 0.1
Nonprompt	17.0 ± 4.0	145 ± 23	104 ± 21	26.6 ± 2.9
Other	22.3 ± 2.4	100 ± 10	58 ± 6	8.0 ± 0.9
Total predicted	241 ± 11	873 ± 22	415 ± 17	89.0 ± 5.4
Data	242	885	418	79

- Postfit BDT output distribution in the e^+e^+ , $e^+\mu^+$, $\mu^+\mu^+$ & 3ℓ channels for improving the separation between signal and background

Agreement

OBSERVATION OF $PP \rightarrow WW+X$ PRODUCTION #2



11 BDT training variables

3ℓ

E_T^{miss} significance $\times 10 / E_T^{\text{miss}}$

$p_T(\ell_2)$

$N(\text{jets})$

same flavor $m_{\ell\ell}$

$m_T(\ell\ell\ell, E_T^{\text{miss}})$

$m(\ell_2, \ell_3)$

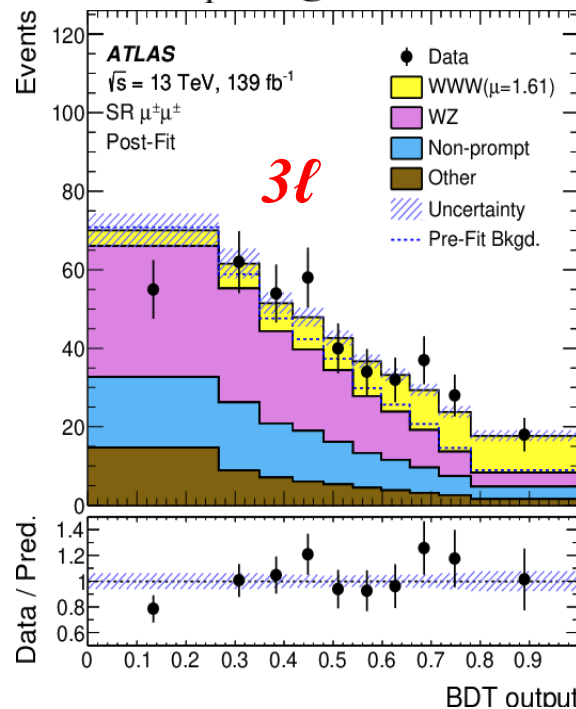
$\Delta\phi(\ell\ell\ell, E_T^{\text{miss}})$

minimum $\Delta R(\ell, \ell)$

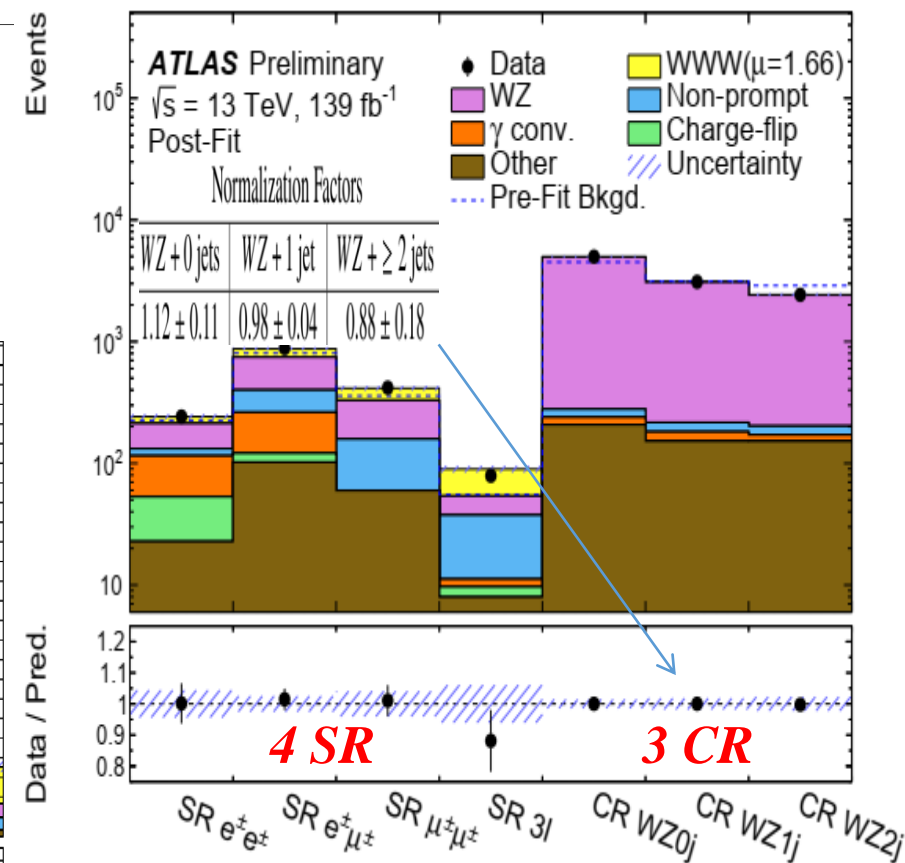
$p_T(\ell_3)$

$m_T(\ell_2, E_T^{\text{miss}})$

E_T^{miss} significance



Backgrounds and Background Estimation



- Higher signal purity in $3l$ events compared to l^+l^+jj
- ▣ Separate **boosted decision tree (BDT)** training to extract signal: Kinematic and angular variables, combining **leptons** and E_T^{miss}
- ▣ Similar validation of BDT modelling as for l^+l^+jj

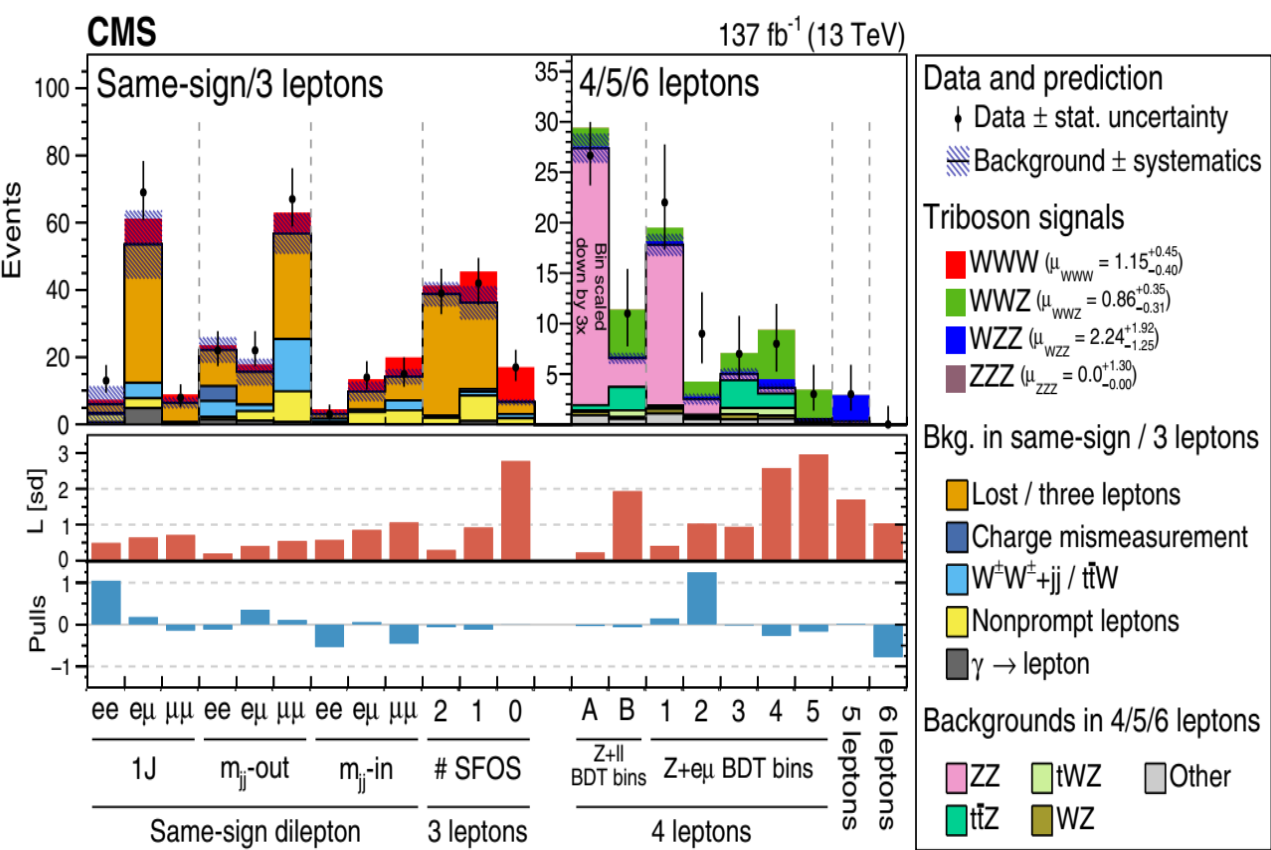
▣ The measured cross section, extrapolated to the total phase space:

✓ $\sigma_{WWW} = 820 \pm 100$ (stat) ± 80 (syst) fb; $\sigma_{WWW}^{\text{MC}} = 511 \pm 18$ fb
2.6 σ from the predicted of calculated at NLO QCD and LO EW accuracy

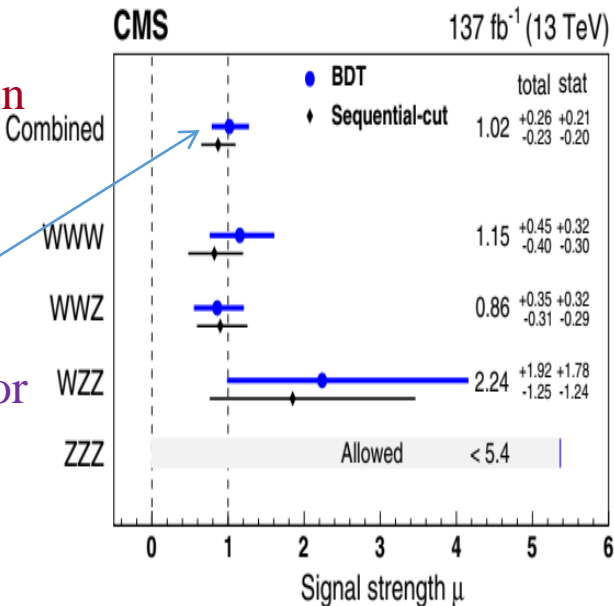
- ▣ Fit of BDT (SRs) and m_{3l} (WZ CRs)
- ▣ Largest source of background from WZ+jets production, constrained in control regions
- ▣ Other backgrounds are instrumental and estimated in data (misidentified leptons, γ -conversions, electron charge misreconstruction)



- The production of VVV ($V = W, Z$) bosons in pp collisions at $\sqrt{s}=13$ TeV
- 5 final states: $W^\pm W^\pm W^\mp \rightarrow l^\pm l^\pm 2\nu qq^-$, $W^\pm W^\pm W^\mp \rightarrow l^\pm l^\pm l^\mp 3\nu$, $W^\pm W^\mp Z \rightarrow l^\pm l^\mp 2\nu l^\pm l^\mp$, $W^\pm ZZ \rightarrow l^\pm \nu 2(l^\pm l^\mp)$, $ZZZ \rightarrow 3(l^\pm l^\mp)$



- ❑ The observed significance of the combined VVV production signal is **5.7 (5.9)**
- ❑ the corresponding measured cross section relative to SM prediction is **$1.02^{+0.26}_{-0.23}$**
- ❑ The observed significances for the individual triboson production processes:
 - 3.3 (3.1) WWW,
 - 3.4 (4.1) WWZ,
 - 1.7 (0.7) WZZ,
 - 0.0 (0.9) ZZZ.



Best fit values of the signal strengths for the BDT-based analyses and the sequential-cut analyses

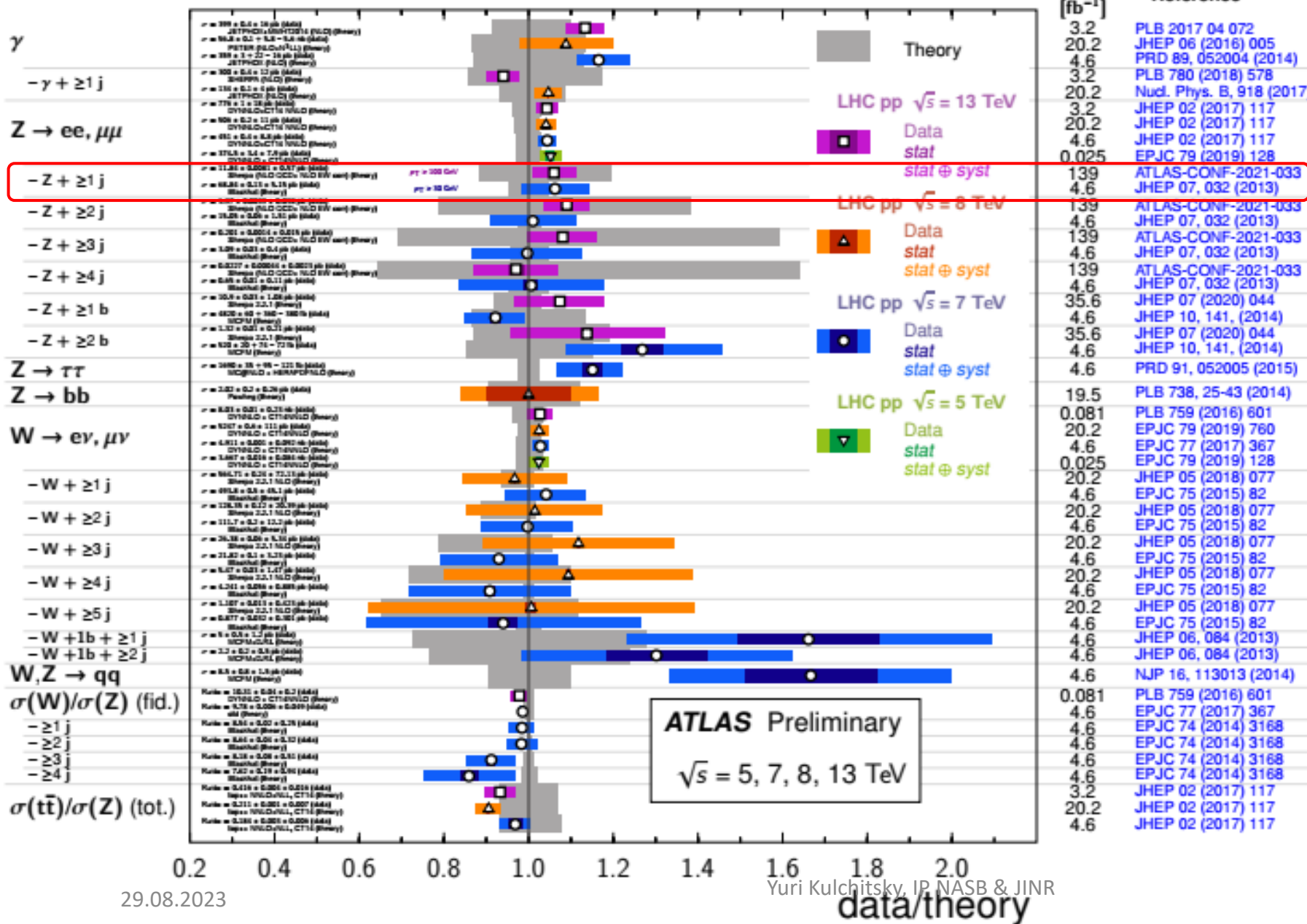
Process	Cross section (fb)	
	Treating Higgs boson contributions as	
	Signal	Background
VVV	$1010^{+210}_{-200} +150_{-120}$	$370^{+140}_{-130} +80_{-60}$
WWW	$590^{+160}_{-150} +160_{-130}$	$190^{+110}_{-100} +80_{-70}$
WWZ	$300^{+120}_{-100} +50_{-40}$	$100^{+80}_{-70} +30_{-30}$
WZZ	$200^{+160}_{-110} +70_{-20}$	$110^{+100}_{-70} +30_{-10}$
ZZZ	< 200	< 80

Comparison of the observed numbers of events to the predicted yields after fitting. For the WWW & WWZ channels, the results from the *boosted decision trees* (BDT) based selections are used. Events with two or more jets are categorized as “ m_{jj} -in” or “ m_{jj} -out”. The expected significance L in the middle panel represents the number of standard deviations with which the null hypothesis (no signal) is rejected; it is calculated for the fit for μ_{comb} . Pulls are the differences in the numbers of observed and predicted events normalized to the uncertainties in the numbers of predicted events.

VECTOR BOSON + X CROSS SECTION MEASUREMENTS

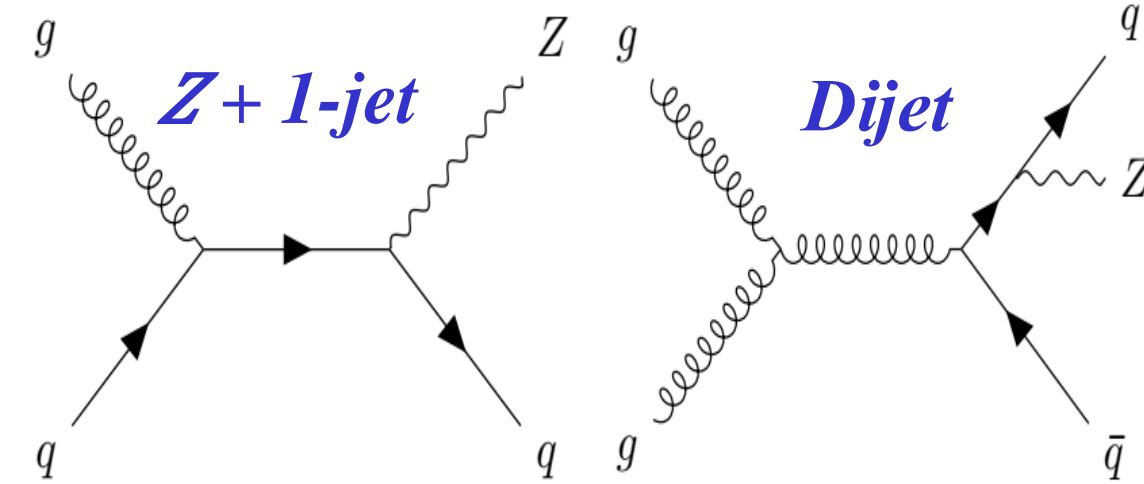
Vector Boson + X fid. Cross Section Measurements

Status: February 2022



- The **data/theory ratio** for several single-boson fiducial production cross section measurements, corrected for leptonic branching fractions
- All **theoretical expectations** were calculated at **NLO or higher**
- The dark-color error bar represents the statistical uncertainty
- The lighter-color error bar represents the full uncertainty, including systematics and luminosity uncertainties
- The luminosity used and reference for each measurement are also shown
- They were not always evaluated using the same prescriptions for PDFs and scales

- The $pp \rightarrow Z + \text{high-}p_T \text{ jets} + X$ provides a way to probe the *interplay between QCD & higher-order EW processes*
- The angular correlations between the Z boson and the *closest jet* are performed in events with at ≥ 1 jet with $p_T \geq 500 \text{ GeV}$

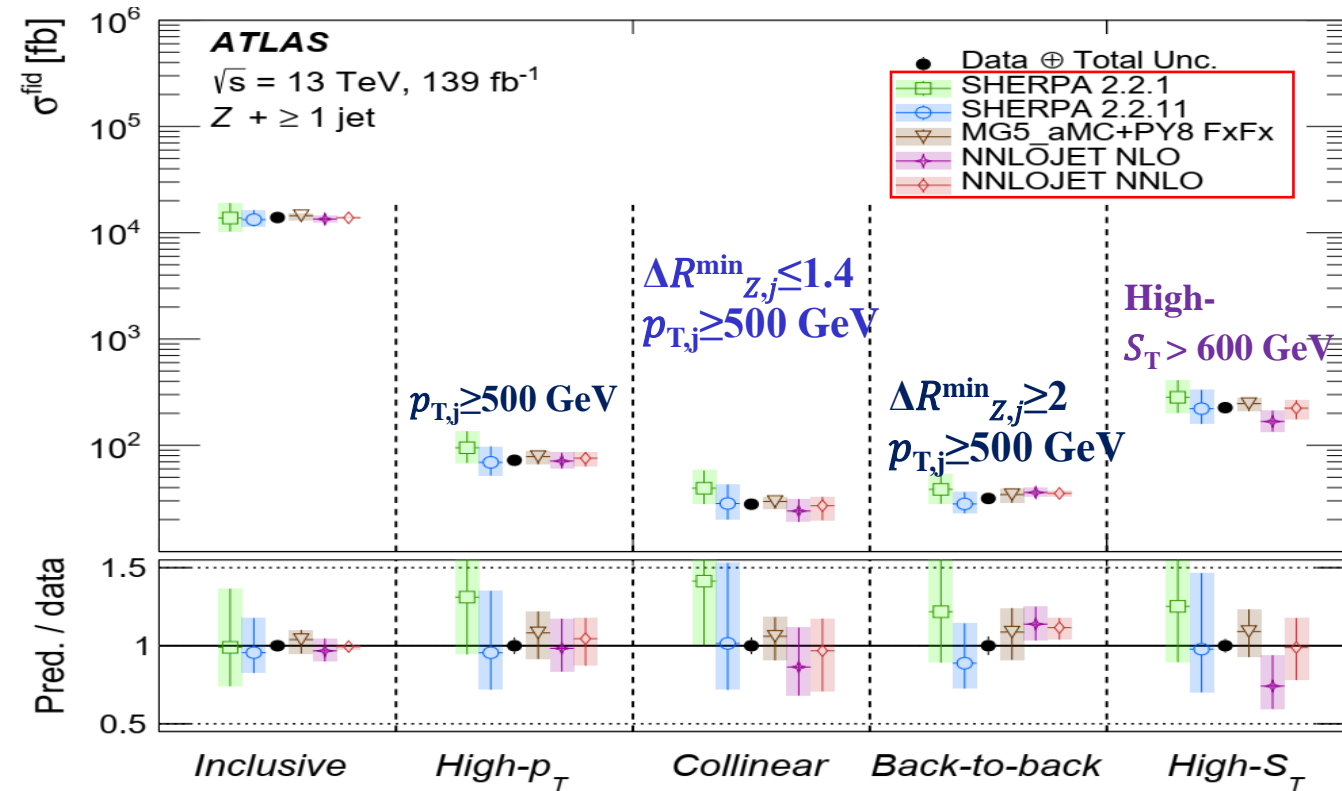


Feynman diagrams for the production of a Z boson in association with *high- p_T jets*

- The $Z + 1\text{-jet}$ events are expected to populate the **back-to-back** region where the Z boson is balanced against a *single high- p_T jet*
- In *dijet* events, the Z boson is expected to be **radiated** from the quark leg, with kinematics leading to *small values* of the **angular distance** between the Z boson and the closest *jet*, $\Delta R_{\min Z,j}$ and, therefore, populating the collinear region

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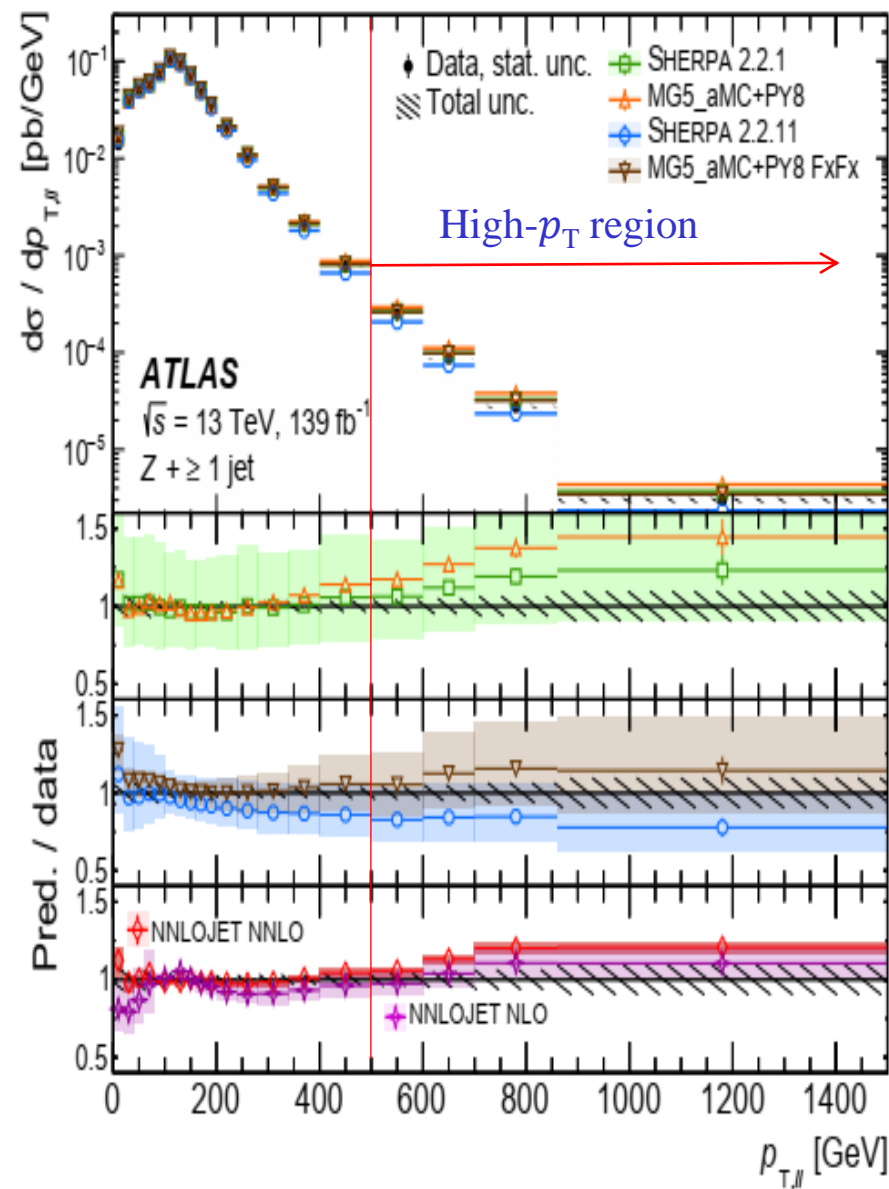
Yuri I



Summary of integrated fiducial cross-section results

- Data are compared with predictions from MC generators
- S_T is the scalar sum of the transverse momentum of all selected jets. *High- $S_T > 600 \text{ GeV}$*
- The *collinear* is for $\Delta R_{\min Z,j} \leq 1.4$
- The *back-to-back* regions is for $\Delta R_{\min Z,j} \geq 2.0$

$$S_T = \sum p_T^{\text{jet}} + \sum p_T^l$$



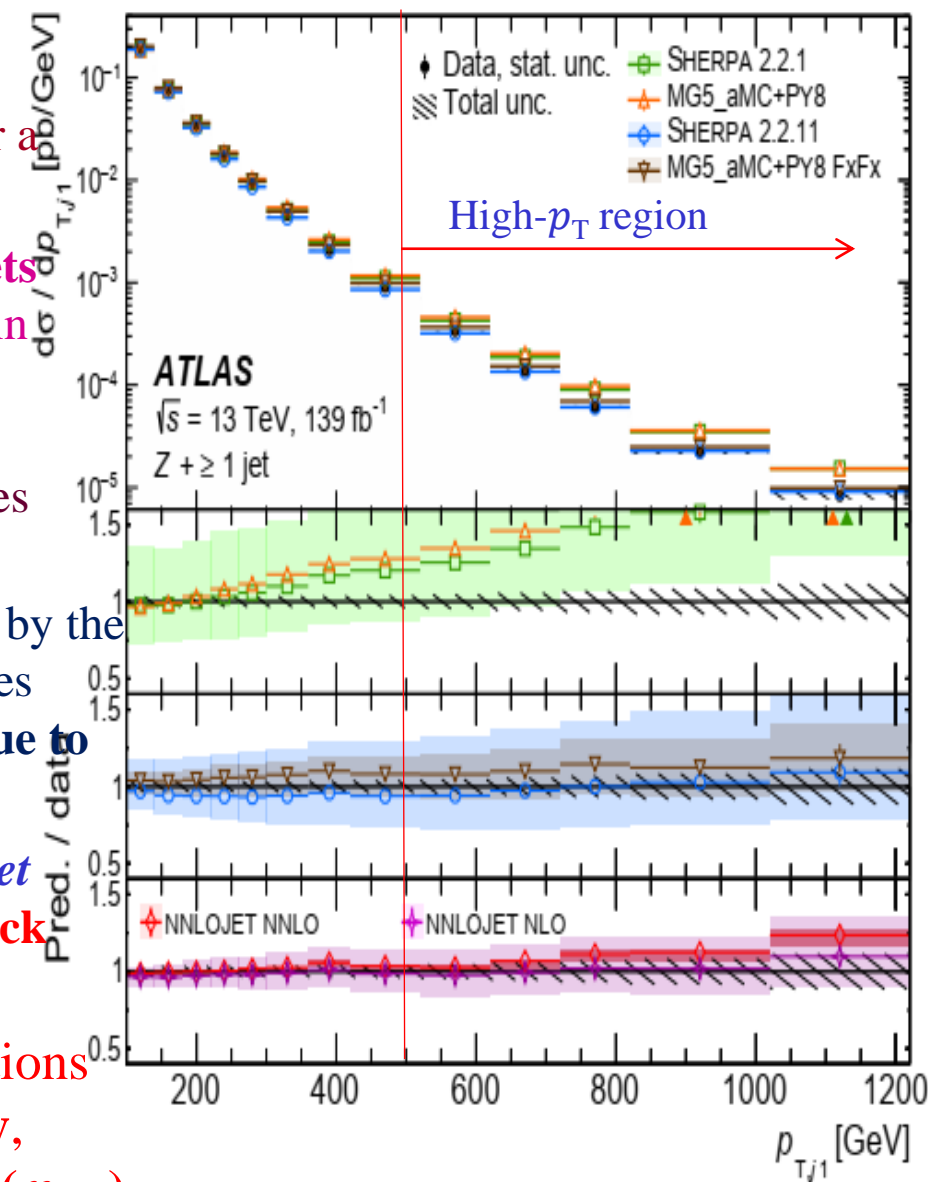
❖ The Z -boson p_T and jet p_T are observables of the $Z + \text{jets}$ process and probe perturbative QCD over a wide range of scales

❖ Understanding the kinematics of jets in events with V-bosons produced in association with several jets is essential for the modelling of backgrounds for other SM processes & searches beyond the SM

❑ The $\text{high-}p_{T,Z}$ region is dominated by the **back-to-back topology** and receives significant **negative corrections due to EW effects**

❑ In contrast, events with a $\text{high-}p_T \text{ jet}$ typically result in **both back-to-back and collinear topologies**

➤ The NNLOjet@NNLO predictions describe the data **very precisely**, except for **large values of $p_{T,\ell\ell}$ ($p_{T,j1}$)**



- **Heavy-flavor** partons in the initial state of a hard-scattering process are understood to **arise from gluon splittings** into $b\bar{b}, c\bar{c}$
- ❑ A **high-mass new particle decaying** into resonances naturally **generates high-momentum merged jets**, and the **high-momentum regime** is particularly **sensitive to modifications to SM** dynamics by **new physics**
- ❑ A large- R -jet is required as a proxy for a **high- p_T , hadronically decaying or splitting object** (high-energy gluon)

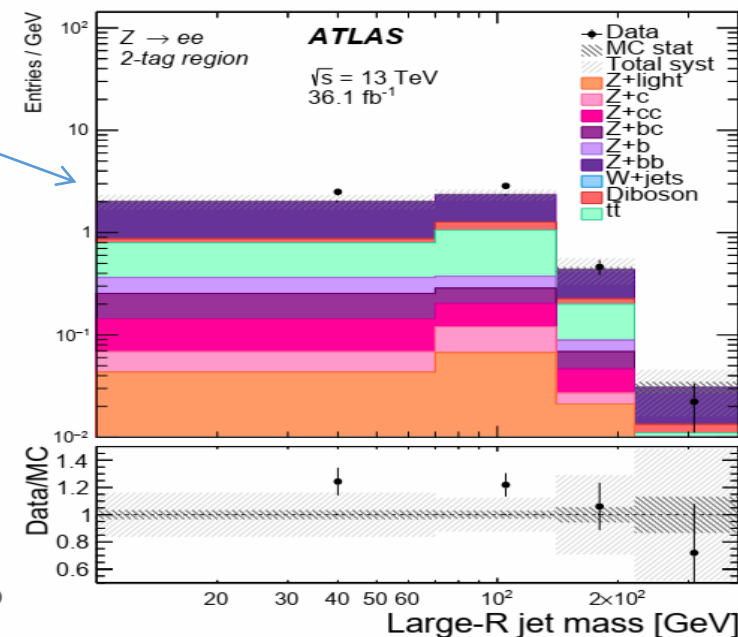
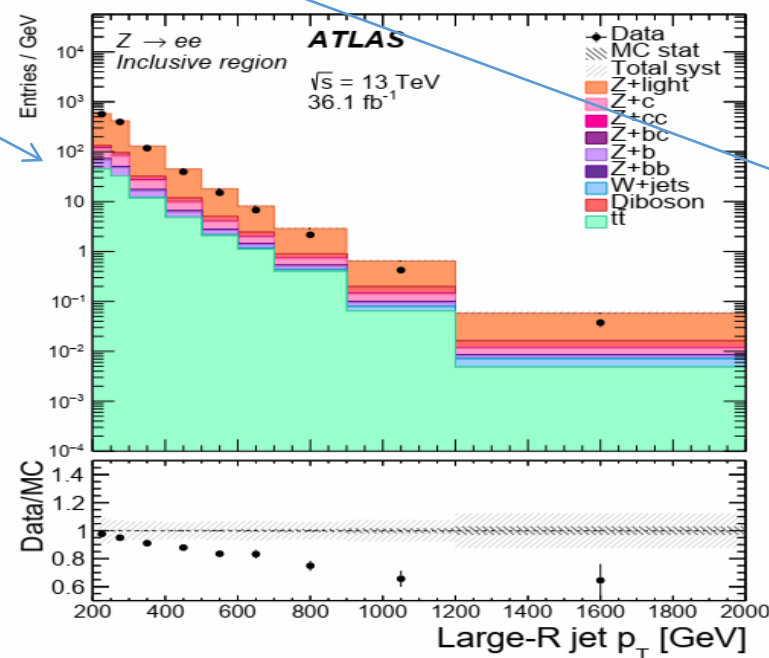
	Inclusive		2-tag	
	ee	$\mu\mu$	ee	$\mu\mu$
$Z+b\bar{b}$	324 ± 4	305 ± 4	163.8 ± 2.6	157.2 ± 2.5
$Z+c\bar{c}$	536 ± 10	530 ± 9	12.3 ± 1.8	19.3 ± 2.0
$Z+bc$	89 ± 2	81 ± 2	14.6 ± 1.2	12.1 ± 0.9
$Z+b$	2588 ± 13	2423 ± 12	14.8 ± 1.1	12.4 ± 1.3
$Z+c$	5073 ± 32	4862 ± 39	5.5 ± 1.3	6.9 ± 1.7
$Z+\text{light}$	$53\,808 \pm 164$	$51\,206 \pm 145$	9.4 ± 1.1	11.1 ± 1.5
$t\bar{t}$	5960 ± 46	5204 ± 43	82.7 ± 5.3	75.4 ± 5.6
$W+\text{jets}$	73 ± 4	7 ± 1	0.4 ± 0.1	< 0.1
Diboson	2042 ± 17	1834 ± 16	21.5 ± 1.4	20.7 ± 1.4
MC total	$70\,493 \pm 175$	$66\,452 \pm 158$	324.9 ± 6.8	315.1 ± 7.2
Data	66 481	65 034	391	384

- The **2-tag** region was defined as a subset of **inclusive** with requiring “**the large- R jet contains exactly two b -tagged subjects**” (**highest- p_T 2-tag large- R jet** in the 2-tag selection)

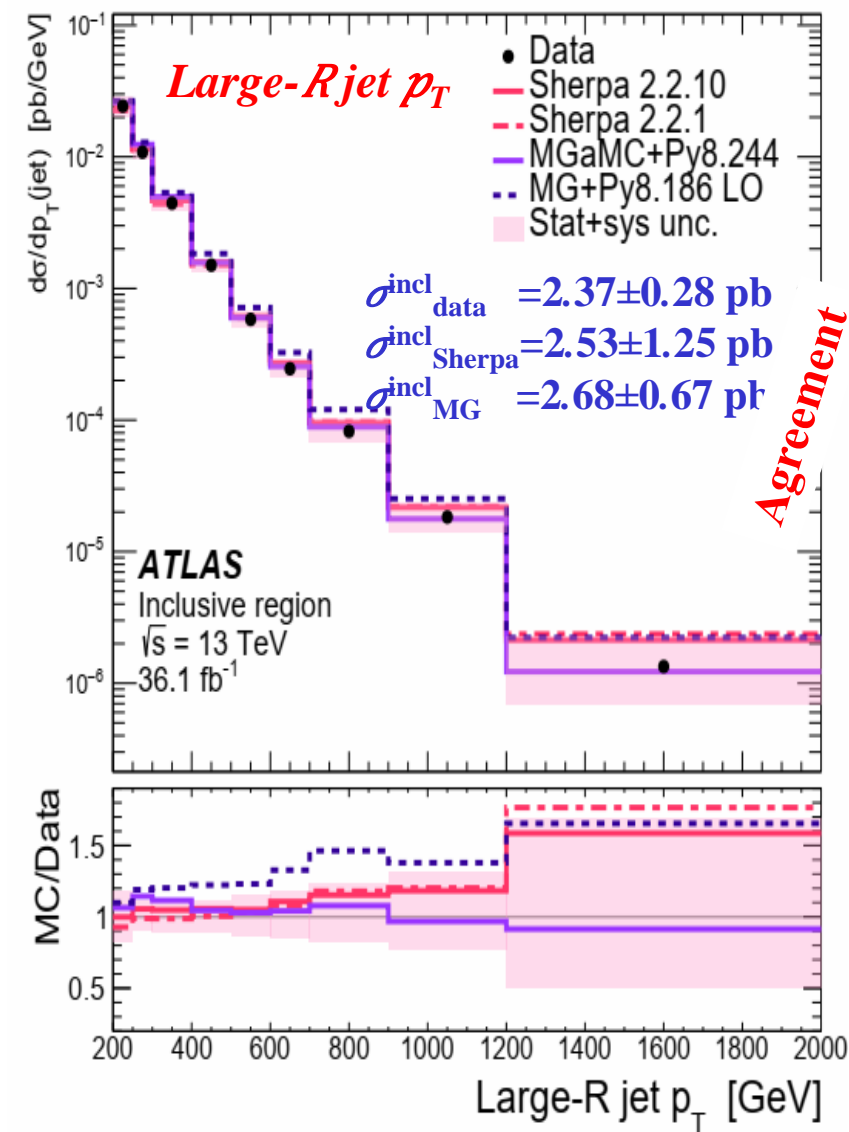
Subject separation: the angular separation, $\Delta R(b, \bar{b})$, between 2 b -tagged subjects

Reconstruction-level event-selection yields in the ee & $\mu\mu$ channels from each process's MC sample (with **Sherpa 2.2.1** used for the $Z+\text{jets}$ samples) and from collision data

- The **single-top process** was found to make a negligible contribution to all event selections and has been omitted
- **Multijet backgrounds** were estimated to be negligible by a data-driven method



Selected reconstruction-level observables, compared with **pre-fit MC simulation** with **Sherpa 2.2.1** used for the $Z+\text{jets}$ samples: the left shows the inclusive-selection ee large- R jet p_T and the 2-tag selection ee large- R jet mass distributions

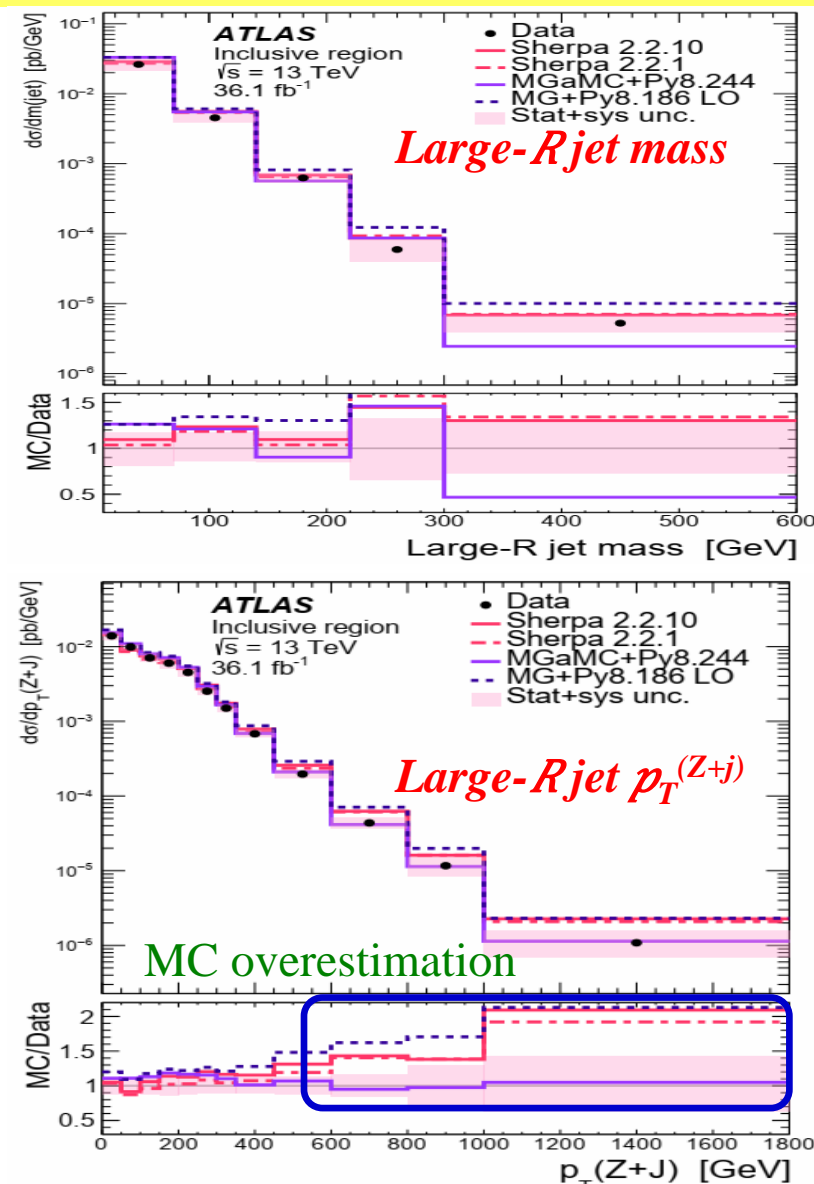


❑ The differential cross-sections indicate significant **mismodeling of QCD activity** in the **inclusive event selection** by MC models

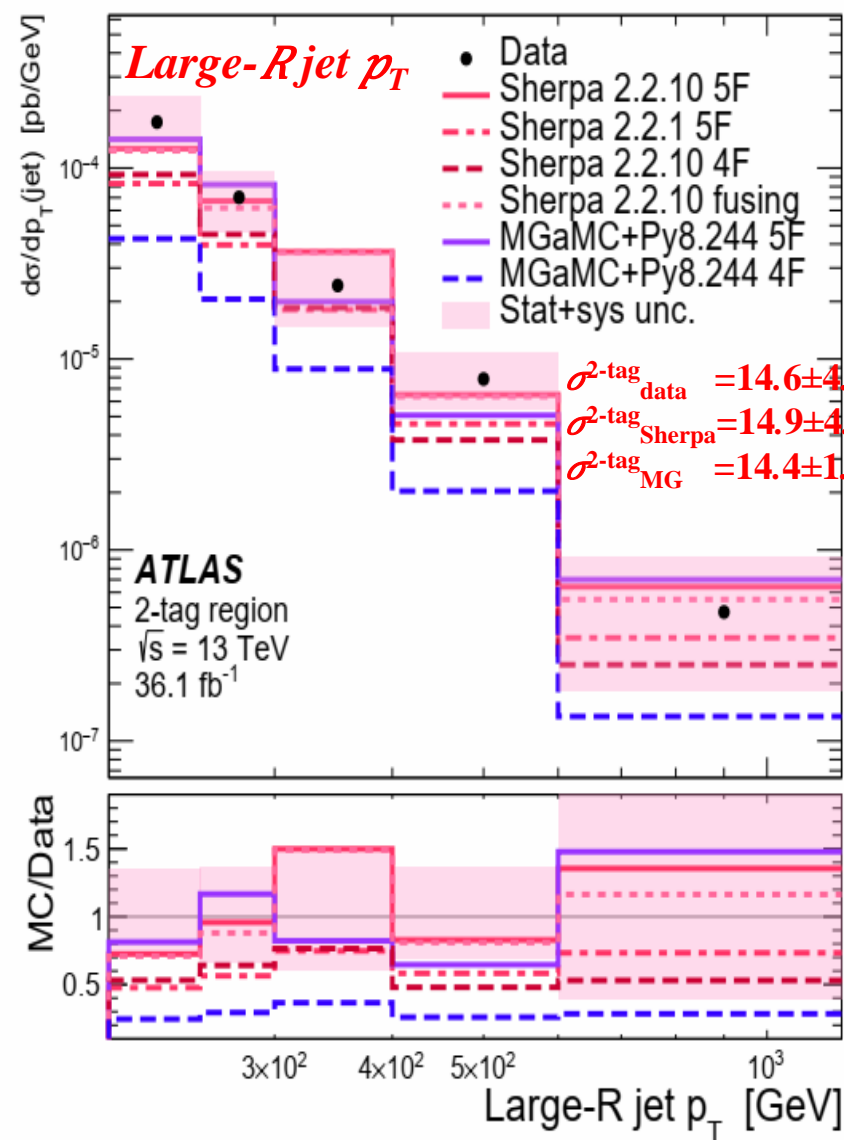
❑ The NLO Sherpa and LO MadGraph+Pythia8 event generators **predicting greater p_T and azimuthal decorrelation** in the **$Z+J$ system** than seen in the data. The **large- R jet** itself is consequently biased to higher p_T and mass values than in data, although to a lesser extent than the deviations in the **$Z+J$ -system** observables

❑ **The NLO MadGraph5_aMC@NLO + Pythia8 model describes all distribution shapes well**, with only a small overestimate of the inclusive fiducial cross-section.

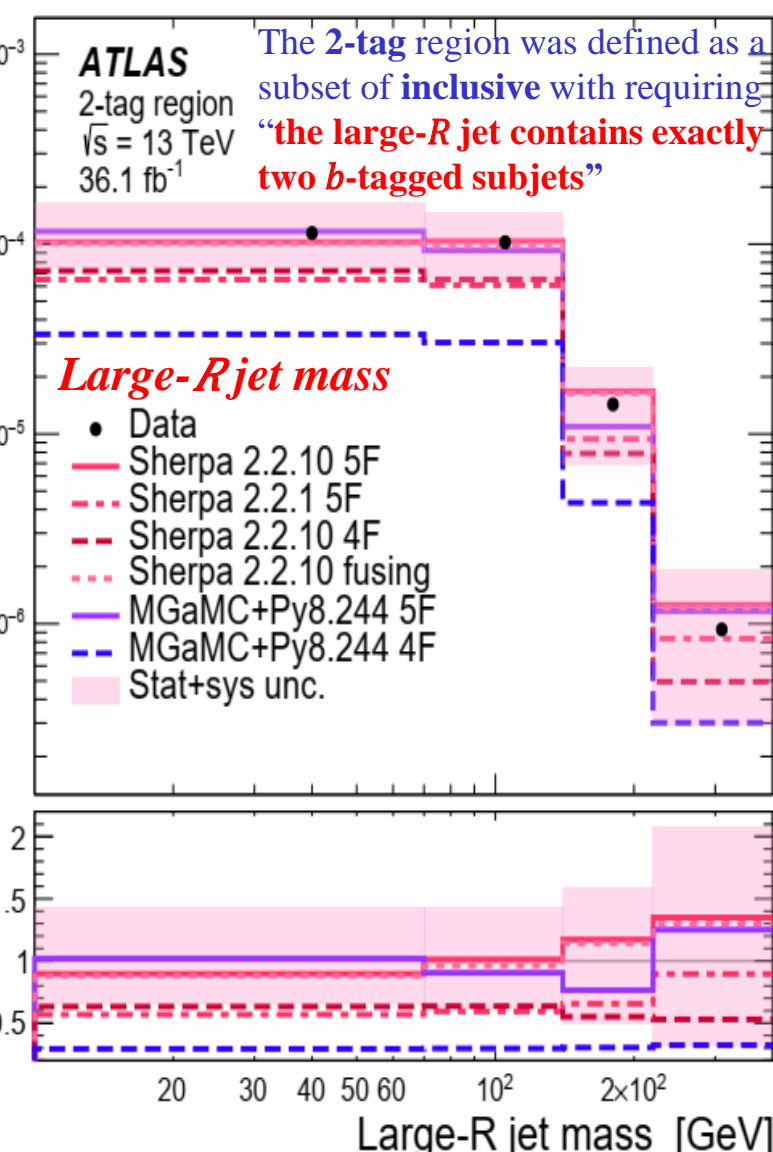
➤ All models somewhat overestimate this cross-section, with recent **Sherpa** versions providing the **best description**



Particle-level differential cross-sections in the inclusive event selection. In the legend, “MGaMC” refers to NLO configurations of the MadGraph5_aMC@NLO generator, and “MG” to LO MadGraph, both run in conjunction with Pythia 8. All models are using the “5FNS” refer to the flavor-number scheme used. These are a significant background to important Higgs-boson searches.



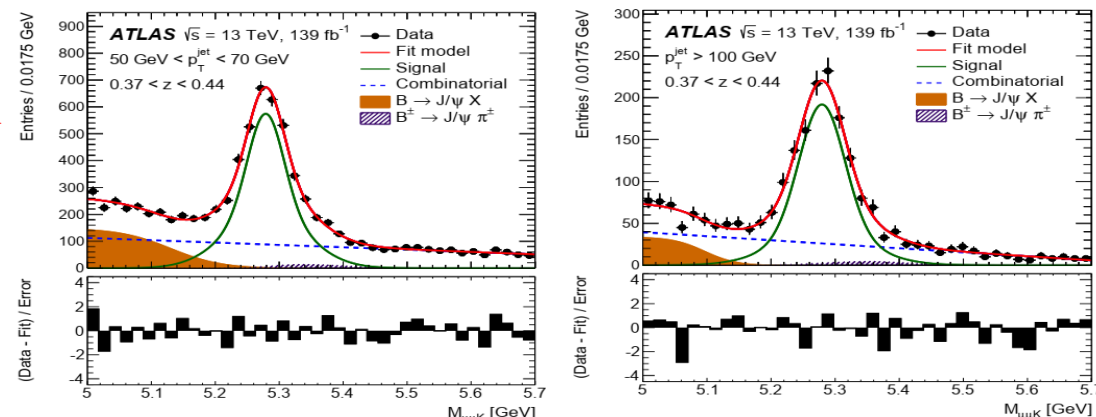
- $\sigma_{2\text{-tag}}^{\text{data}}/\sigma_{\text{incl}}^{\text{data}} = (0.62 \pm 0.12)\%$
- $\sigma_{2\text{-tag}}^{\text{Sher}}/\sigma_{\text{incl}}^{\text{Sher}} = (0.59 \pm 0.39)\%$
- $\sigma_{2\text{-tag}}^{\text{MG}}/\sigma_{\text{incl}}^{\text{MG}} = (0.54 \pm 0.21)\%$
- The 2-tag selection: there is **good shape agreement** between the data & MC models
- The strongest feature observed in this event-selection region is in normalization, with models using **four-flavor number scheme (4FNS, b-absent in PDF)** approach significantly **underestimating** the rate of $b\bar{b}$ boosted-jet production
- Five-four-flavor number scheme (5FNS, with b) approaches with modern tools **do much better** (Sherpa 2.2.10 & MadGraph5_aMC@NLO) providing accurate predictions for the 2-tag cross-section and the **ratio of 2-tag to inclusive**
- This results is important for future use of MC-derived **large- R jet** flavor composition in studies of the $VH(b\bar{b})$ process



The 2-tag region was defined as a subset of **inclusive** with requiring “the large- R jet contains exactly two b -tagged subjects”

Particle-level differential cross-sections in the 2-tag event selection. In the legend, “MGaMC” refers to NLO configurations of the MadGraph5_aMC@NLO generator, and “MG” to LO MadGraph, both run in conjunction with Pythia 8. All models are using the “4/5FNS” refer to the flavor-number scheme used.

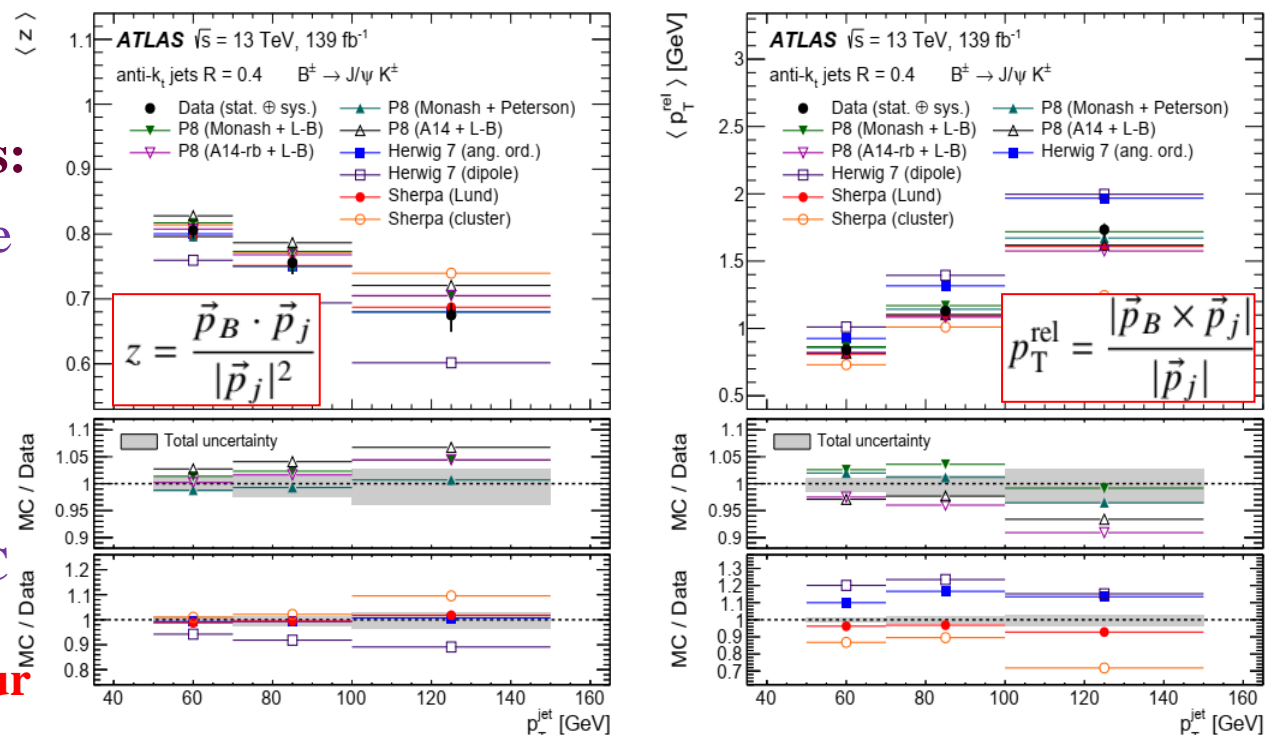
- The fragmentation of **heavy quarks** is a **crucial** aspect of QCD
- Detailed studies and precision measurements of the heavy-quark fragmentation properties allow a **deeper underst. of QCD**
- ❑ The fragmentation properties of jets containing b -hadrons are studied using $B^\pm \rightarrow J/\psi K^\pm$ in pp collisions, with $J/\psi \rightarrow \mu^+ \mu^-$
- ❑ The measurement determines the p_L and p_T profiles of the reconstructed B hadrons with respect to the axes of the *jets* to which they are geometrically associated



Fits to the invariant mass distributions of B^\pm candidates for $0.37 < z < 0.44$

The provides key measurements which help to **better understand the fragmentation functions of heavy quarks:**

- **Significant differences** among different MC models are **observed**, and also **between the models and the data**
- Some of the discrepancies are understood to arise from poor modelling of the $g \rightarrow b\bar{b}$ splittings, **to which the present analysis has substantial sensitivity**
- Including the present measurements in a **future tune** of the MC predictions **may help to improve the description and reduce the theoretical uncertainties** of processes where **heavy-flavour quarks are present in the final state**, such as *top quark pair production or Higgs boson decays into heavy quark pairs*

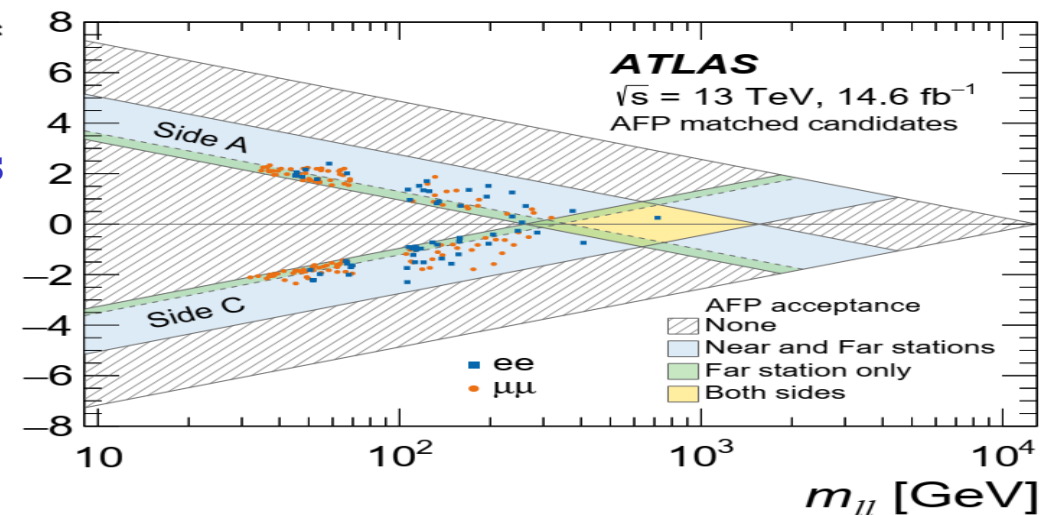


Average values of the longitudinal profile $\langle z \rangle$ and of the transverse profile $\langle p_T^{\text{rel}} \rangle$ as a function of the p_T^{jet}

- The observation of *forward p* scattering with 2 leptons ($e^+e^-/\mu^+\mu^-$) produced via **photon fusion**
- The scattered **proton** is detected by **Forward Proton Spectrometer** while the $\ell^+\ell^-$ are reconstructed by the **Central Detector**
- The 57 (123) ee ($\mu\mu$) data event candidates in the dilepton rapidity $y_{\ell\ell}$ vs $m_{\ell\ell}$ plane satisfying even selection and kinematic matching
- Proton-tagging techniques are introduced for X-section measurements

$$\sigma_{ee+p}^{\text{fid}} = 11.0 \pm 2.6 \text{ (stat)} \pm 1.2 \text{ (syst)} \pm 0.3 \text{ (lumi) fb}$$

$$\sigma_{\mu\mu+p}^{\text{fid}} = 7.2 \pm 1.6 \text{ (stat)} \pm 0.9 \text{ (syst)} \pm 0.2 \text{ (lumi) fb}$$

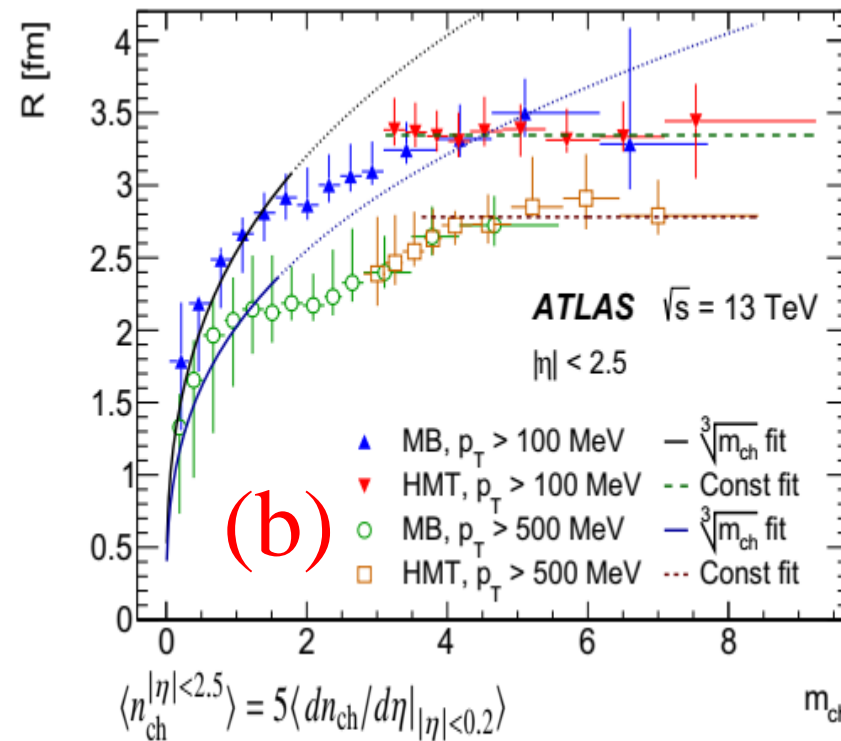
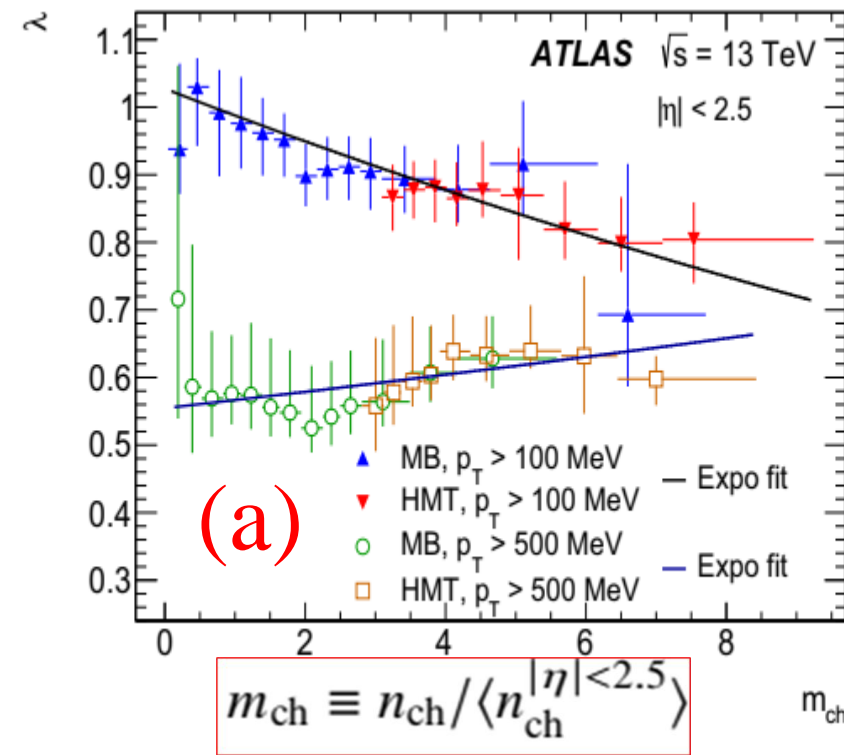


$\sigma_{\text{HERWIG+LPAIR}} \times S_{\text{SURV}}$	$\sigma_{ee+p}^{\text{fid}} \text{ (fb)}$	$\sigma_{\mu\mu+p}^{\text{fid}} \text{ (fb)}$
$S_{\text{SURV}} = 1$	15.5 ± 1.2	13.5 ± 1.1
S_{SURV} using Refs. [33, 34]	10.9 ± 0.8	9.4 ± 0.7
SUPERCHIC 4 [97]	12.2 ± 0.9	10.4 ± 0.7
Measurement	11.0 ± 2.9	7.2 ± 1.8

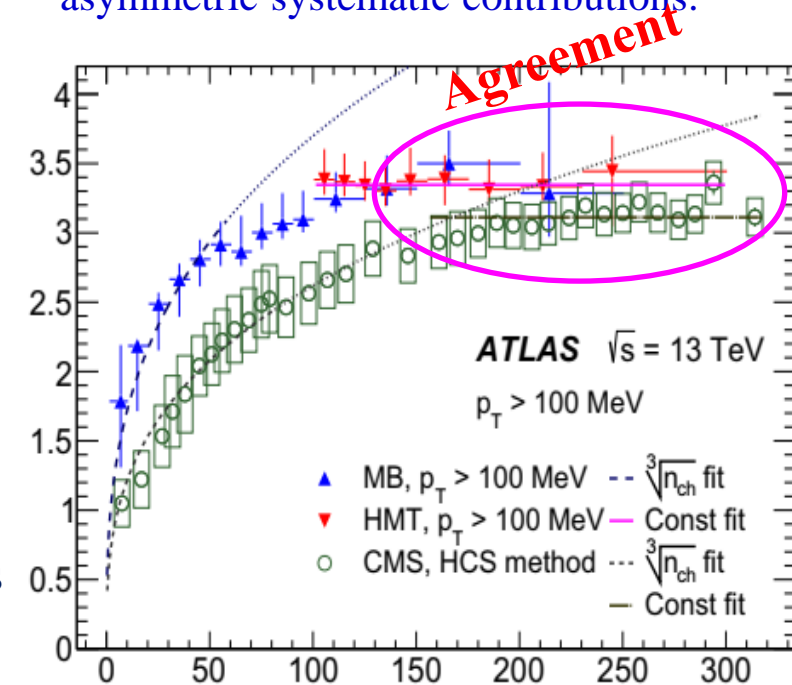
Agreement

- $pp \rightarrow p(\gamma\gamma \rightarrow \ell^+\ell^-)p^{(*)}$ is observed with a significance exceeding **5 σ** in both the e^+e^-+p & $\mu^+\mu^-+p$ final states
- These results demonstrate that the **ATLAS Forward Proton** spectrometer performs well in high-luminosity data taking
- Proton tagging is introduced for cross-section measurements of **photon fusion processes** at the electroweak scale

- Compares X-sections with the combined **HERWIG+LPAIR** predictions assuming unit **soft-survival factors** $S_{\text{SURV}}=1$
- Soft-survival effects are included using an $m_{\ell\ell}$ -**dependent** reweighting of these predictions to S_{SURV} calculated for **exclusive** processes; **LPAIR** predictions are additionally **scaled down by 15%** to account for S_{SURV} being lower for single-dissociative processes. **SUPERCHIC4** predictions include **full kinematic dependence** on S_{SURV} for exclusive, single-, and double-dissociative processes.



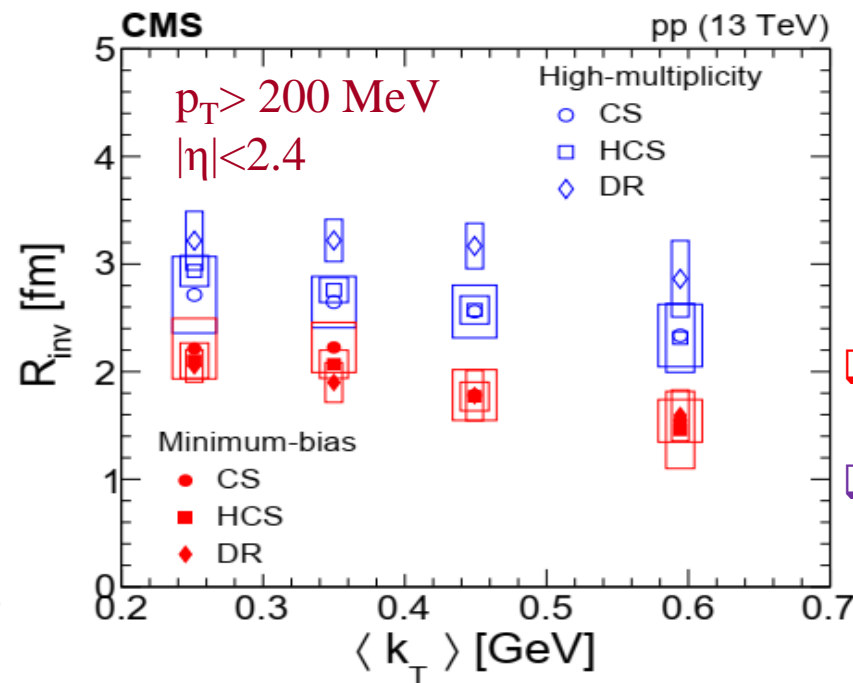
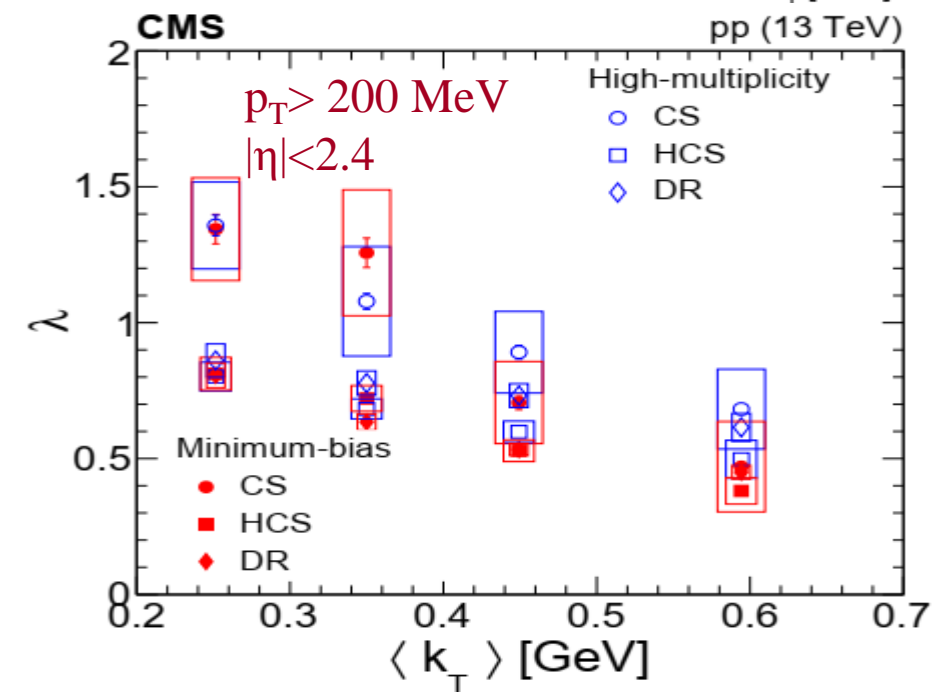
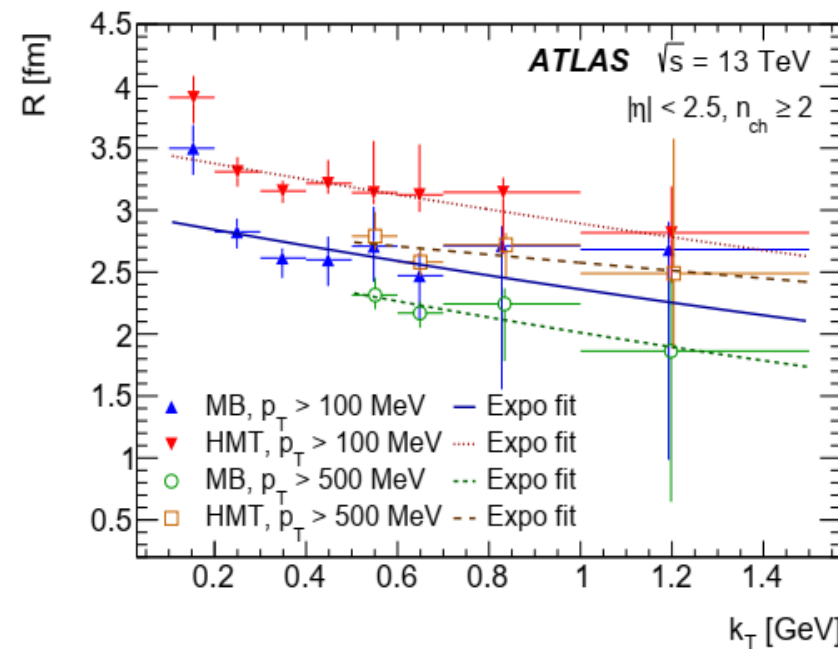
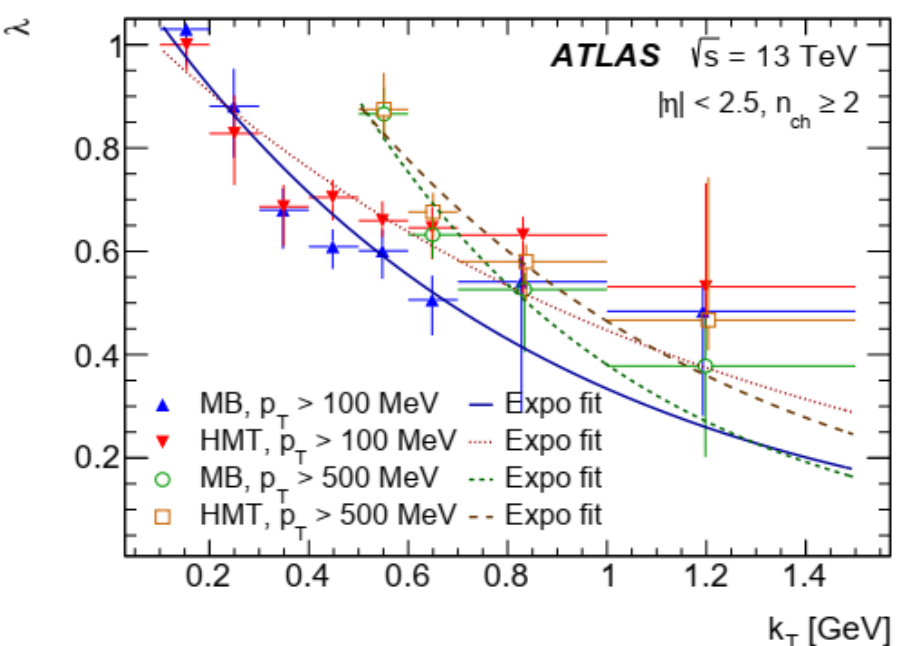
- The dependence of the $\lambda(m_{ch})$ on rescaled multiplicity obtained from the exponential fit of the $R_2(Q)$ correlation functions for tracks with $p_T > 100$ MeV and $p_T > 500$ MeV at $\sqrt{s} = 13$ TeV for MB & HMT data.
- The dependence of the $R(m_{ch})$ on m_{ch} . The uncertainties represent the sum in quadrature of the statistical and asymmetric systematic contributions.



- The parameter $R(m_{ch})$ is found to increase as $\alpha m_{ch}^{0.33}$ for multiplicity up to $m_{ch} \approx 2$: $p_T > 100$ MeV, $\alpha = 2.54^{+0.12}_{-0.22}$ fm
- For $m_{ch} \geq 3$, the source radius R saturates at a value $R = 3.35^{+0.20}_{-0.09}$ fm
- Results indicate a direct dependence of R on the p_T threshold
- The parameter $\lambda(m_{ch})$ decreases with multiplicity for the lower p_T threshold and is lower for the higher p_T threshold but increases slightly with multiplicity

- The behaviour of R is qualitatively similar.
- There is a clear difference in the results at low n_{ch} .

Comparison with CMS for $p_T > 100$ MeV and $|\eta| < 2.4$.



□ **ATLAS:** The k_T dependence of the correlation strength, $\lambda(k_T)$, and the source radius, $R(k_T)$, obtained from the exponential fit to the $R_2(Q)$ correlation functions for events with multiplicity $n_{ch} \geq 2$ and transfer momentum of tracks with $p_T > 100$ MeV and $p_T > 500$ MeV at 13 TeV for the minimum-bias (MB) and high-multiplicity track (HMT) events.

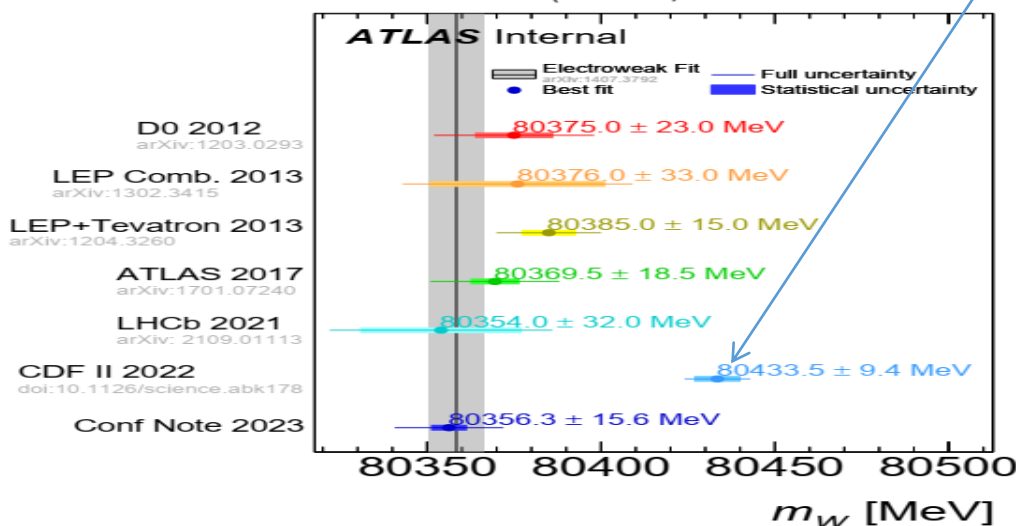
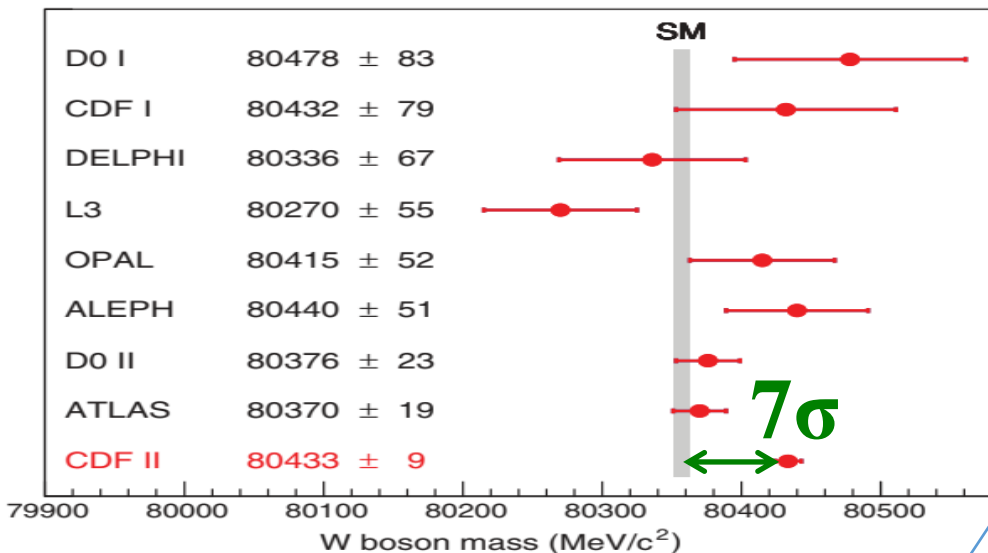
□ The uncertainties represent the sum in quadrature of the statistical & systematic contributions.

□ The curves represent the exponential fits to $\lambda(k_T)$ and $R(k_T)$.

□ **CMS:** Results for R_{inv} and λ from the three methods as a function of k_T

□ In the lower plots, statistical & systematic uncertainties are shown as error bars and open boxes, respectively

CDF II Collaboration, High-precision measurement of the W boson mass with the CDF detector. *PoS ICHEP2022 (2022) 898*



The ATLAS Collaboration has reported a measurement
 $M_W = 80370 \pm 7$ (stat) ± 11 (syst) ± 1 (mod. syst) = 80370 ± 19 MeV
 $M_W = 80356 \pm 5$ (stat) ± 15 (syst) = 80356 ± 16 MeV

□ **CDF II** measure the W boson mass, M_W , using data corresponding to **8.8 fb⁻¹** collected in proton-antiproton collisions at a **1.96 TeV** with the CDF II detector at the Fermilab Tevatron collider. A sample of approximately **4x10⁶ W boson** is used to obtain **$M_W = 80433.5 \pm 6.4$ (stat) ± 6.9 (syst) = 80433.5 ± 9.4 MeV**. The W bosons are identified using their decays to **$e\nu$ & $\mu\nu$** and the mass is measured by fitting template distributions of transverse momentum and mass: $m_T = \sqrt{2p_T^\ell p_T^\nu (1 - \cos \Delta\phi)}$

□ A comparison with the **SM expectation** of **$M_W = 80357 \pm 6$ MeV**, treating the quoted uncertainties as independent, yields a difference with a significance of **7 σ** . The suggests are:

- the improvements to the SM calculation or
- of extensions to the SM

□ SM result includes the published estimates of the uncertainty (4 MeV) due to missing higher-order quantum corrections and the uncertainty (4 MeV) from other global measurements used as input to the calculation

➤ **ATLAS Collaboration**, Measurement of the WW-boson mass in pp collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector, **EPJ C 78, 110 (2018), EPJ C 78, 898 (2018), arXiv:1701.07240**; A measurement of the mass of the W boson is presented based on proton-proton collision data recorded in 2011 at a **7 TeV** with the **ATLAS** detector at the LHC, and corresponding to **4.6 fb⁻¹** of integrated luminosity. The selected data sample consists of **7.8x10⁶ candidates** in the **$W \rightarrow \mu\nu$** channel and **5.9x10⁶ candidates** in the **$W \rightarrow e\nu$** channel.

ATLAS/CMS: Run-2 at 13 TeV is 139/137 fb⁻¹

1. Many new SM results with the latest Run 2 dataset by the ATLAS and the CMS Collaborations were published
2. Measurement SM processes cross sections over 15 orders of magnitude
3. Comparison of the measurement SM processes to theory expectations
4. Evidence or Observation of rare processes
5. The most precise measurement of the total cross section pp-interactions
6. Observation of the BEC radius saturation for very high multiplicity

List of the SM papers:

CMS Collaboration: <https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsSMP>

ATLAS Collaboration: <https://twiki.cern.ch/twiki/bin/view/AtlasPublic/StandardModelPublicResults>



THANK YOU
VERY MUCH
FOR ATTENTION!



CMS Experiment at the LHC, CERN

Data recorded: 2015-May-20 20:51:09.209152 GMT

Run / Event / LS: 245155 / 123253393 / 363

BACKUP SLIDES

STANDARD MODEL: LAGRANGIAN

- **The Quantum chromodynamics sector (QCD)** sector defines the interactions between quarks and gluons, which is a Yang–Mills gauge theory with SU(3) symmetry. Since **leptons do not interact with gluons**, they are not affected by this sector. The Dirac Lagrangian of the quarks coupled to the gluon fields is given by

$$\mathcal{L}_{\text{QCD}} = \sum_{\psi} \bar{\psi}_i \left(i\gamma^{\mu} (\partial_{\mu} \delta_{ij} - ig_s G_{\mu}^a T_{ij}^a) \right) \psi_j - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu},$$

ψ_i is the Dirac spinor of the quark field, where $i=\{r, g, b\}$ represents color,
 γ_{μ} are the Dirac matrices,
 G_{μ}^a is the 8-component ($a=1,2,\dots,8$) SU(3) gauge field,
 T_{ij}^a are the 3×3 Gell-Mann matrices, generators of the SU(3) color group,
 $G_{\mu\nu}^a$ represents the gluon field strength tensor,
 g_s is the strong coupling constant.

- **The Electroweak sector** is a Yang–Mills gauge theory with the symmetry group $U(1) \times SU(2)_L$,

$$\mathcal{L}_{\text{EW}} = \sum_{\psi} \bar{\psi} \gamma^{\mu} \left(i\partial_{\mu} - g' \frac{1}{2} Y_W B_{\mu} - g \frac{1}{2} \vec{\tau}_L \vec{W}_{\mu} \right) \psi - \frac{1}{4} W_{\mu\nu}^a W_{\mu\nu}^a - \frac{1}{4} B_{\mu\nu} B_{\mu\nu},$$

B_{μ} is the U(1) gauge field,
 Y_W is the weak hypercharge – the generator of the U(1) group,
 W_{μ} is the 3-component SU(2) gauge field,
 τ_L are the Pauli matrices – infinitesimal generators of the SU(2) group – with subscript L to indicate that they only act on left-chiral fermions,
 g' and g are the U(1) and SU(2) coupling constants respectively,
 $W_{\mu\nu}^a$ ($a=1,2,3$) and $B_{\mu\nu}$ are the field strength tensors for the weak isospin and weak hypercharge fields.

- **Higgs sector:** In the Standard Model, the Higgs field is a complex scalar of the group SU(2)_L:

$$\mathcal{L}_{\text{H}} = \left| \left(\partial_{\mu} + \frac{i}{2} (g' Y_W B_{\mu} + g \vec{\tau} \vec{W}_{\mu}) \right) \varphi \right|^2 - \frac{\lambda^2}{4} (\varphi^{\dagger} \varphi - v^2)^2.$$

- **Yukawa sector:** The Yukawa interaction terms are

$$\mathcal{L}_{\text{Yukawa}} = \bar{U}_L G_u U_R \varphi^0 - \bar{D}_L G_u U_R \varphi^{-} + \bar{U}_L G_d D_R \varphi^{+} + \bar{D}_L G_d D_R \varphi^0 + \text{h. c.},$$

JET PHYSICS IN PP-COLLISIONS MC MODELS

Jets are crucial for understanding of the SM. Probing of the QCD \rightarrow Jets are the result of fragmentation of partons produced in a scattering process

In High-Energy Particle collisions – two main phases:

- ❑ **Perturbative phase:** partons with *high-transverse momentum* are produced in a hard-scattering process at a scale Q
- ❑ **Non-perturbative phase:** partons convert in hadrons *emitting gluons* and $q\bar{q}$ -pairs an interplay between *Hadronization Process (HP)* and *Underlying Event (UE)*:

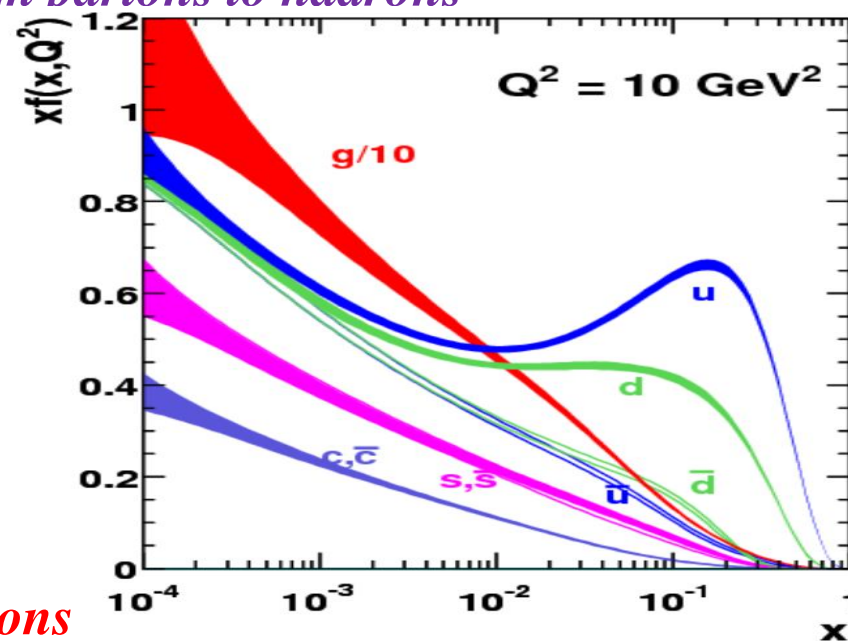
➤ **Hadronization Process:** *transition from partons to hadrons*

➤ **Underlying Event:**

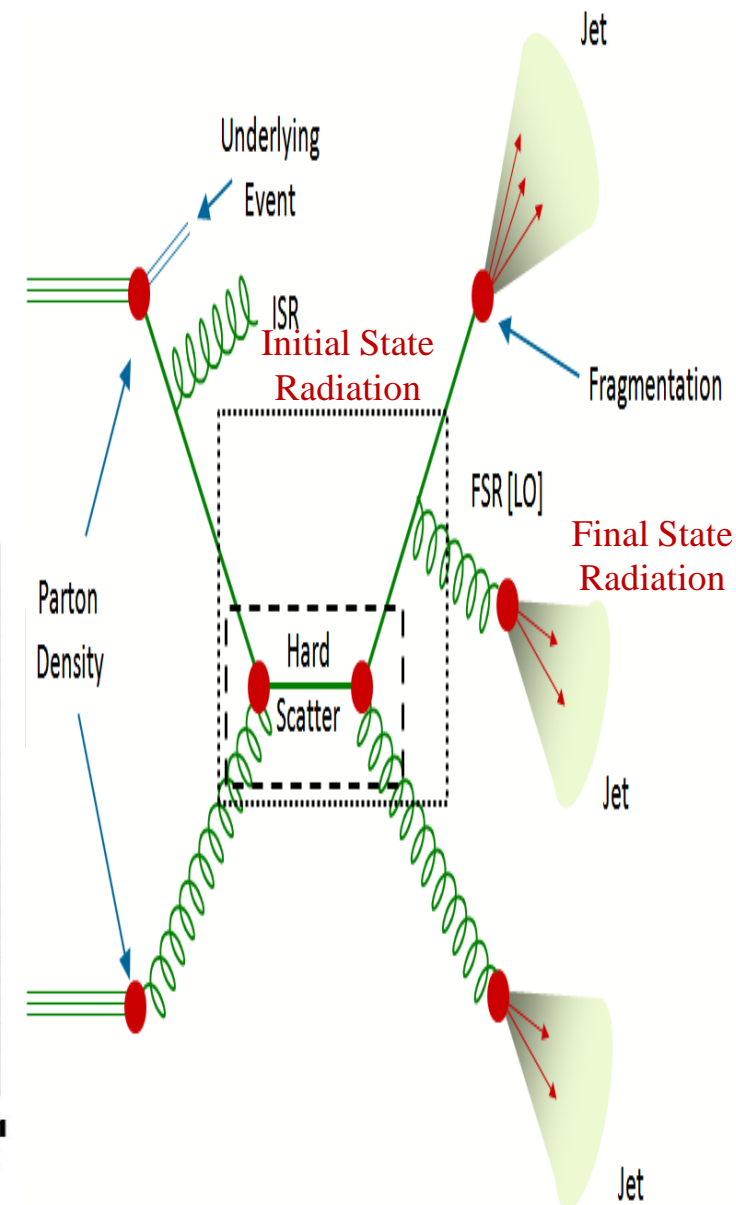
- ❑ *initial-state radiation (ISR);*
- ❑ *final-state radiation (FSR);*
- ❑ *multiple-parton interactions;*
- ❑ *colour-reconnection effects*

❖ *Effects of HP and UE depend on Jet radius parameter and are most pronounced at low p_T*

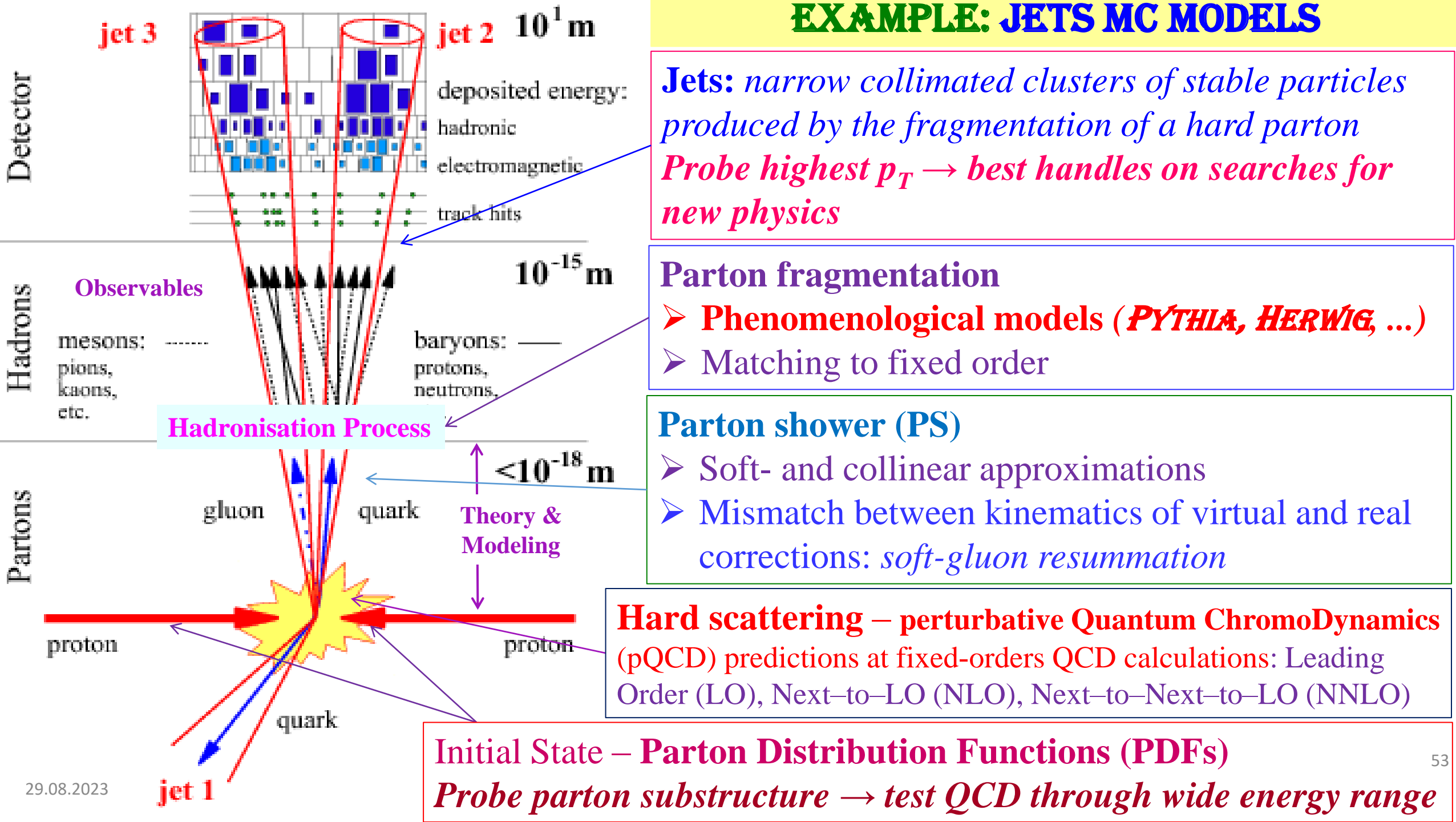
❖ *All these aspects of high energy collisions can be Probed in the Jet Physics*



Parton density functions (PDF)

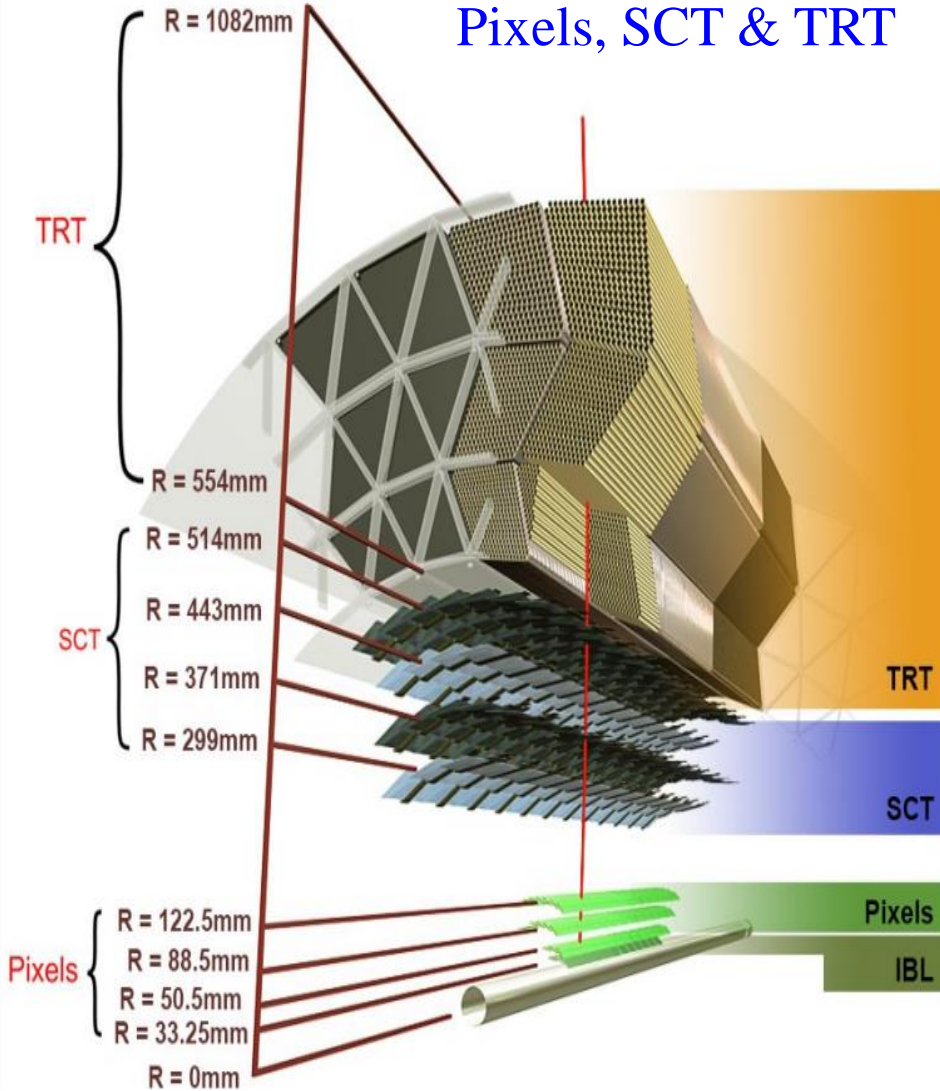


EXAMPLE: JETS MC MODELS

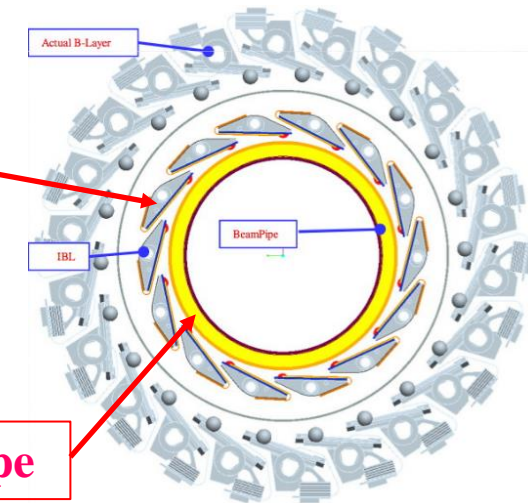


INNER DETECTORS (ID)

ATLAS tracking detectors: Pixels, SCT & TRT



- ☐ **New innermost 4-th layer** for the Pixel detector
[IBL = Insertable B-Layer]
- ☐ Required complete removal of the ATLAS Pixel volume
- ☐ IBL fully operational

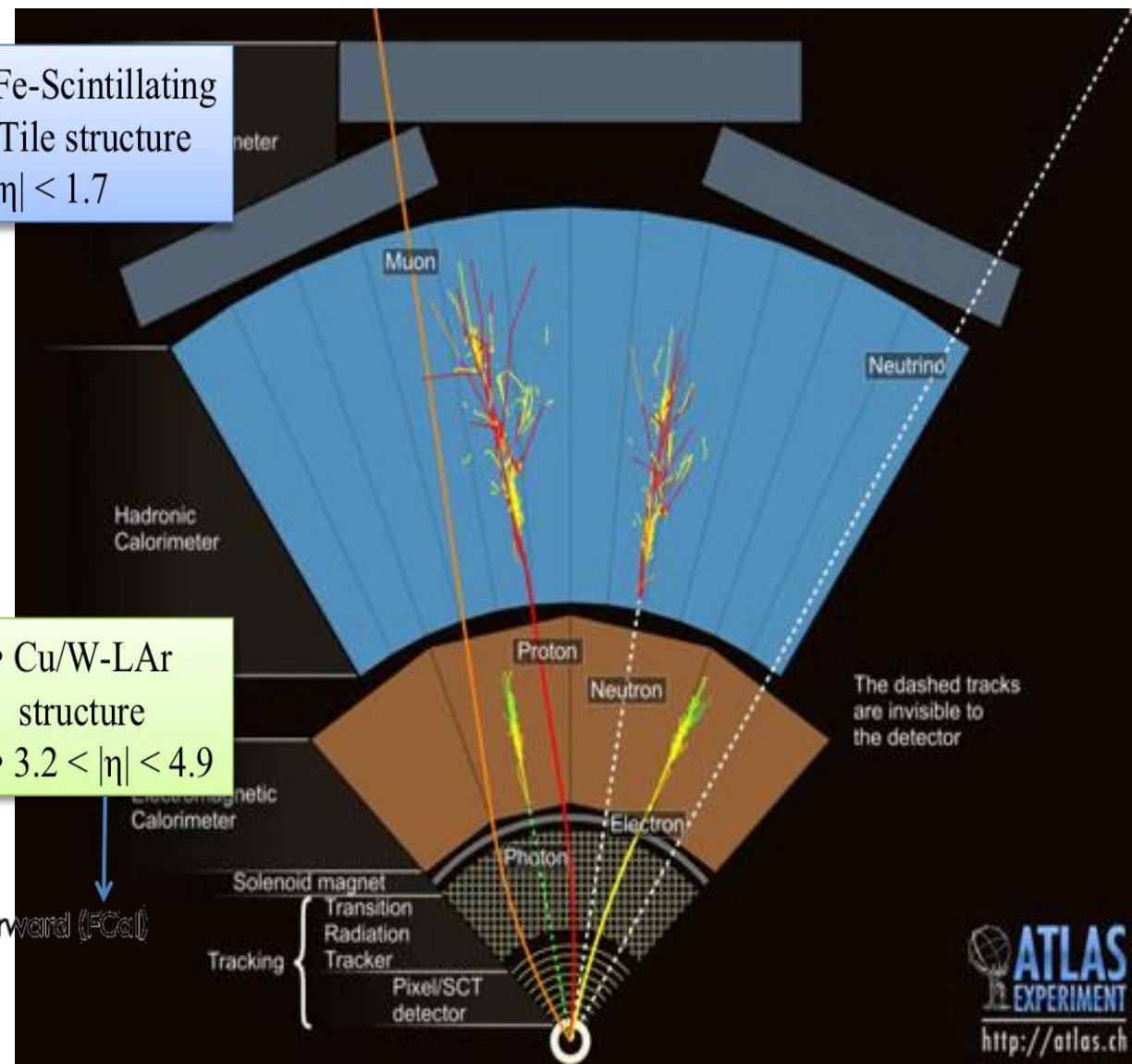
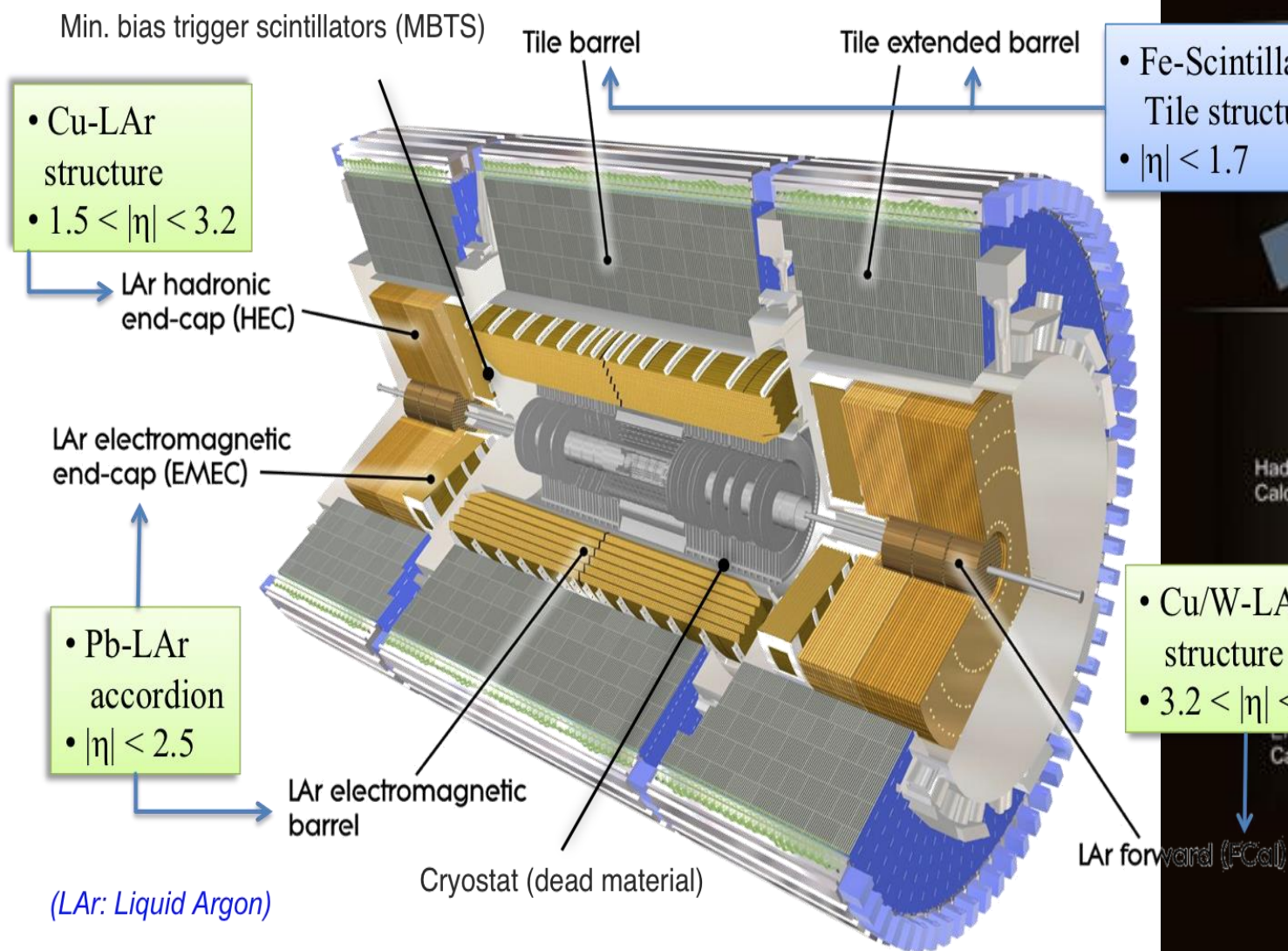


Two times better tracks impact parameters resolution at 13 TeV!

ATLAS CALORIMETERS

- **LAr & TileCal >> Very stable performance**
- **Improved stability of new Tile power supplies**

- **Good operation efficiency: ~100% for LAr & Tile**
- LAr using 4 sample readout to achieve 100 kHz



❑ Search for **seed energy clusters** in the EM calorimeter with significant energy

➤ For $|\eta| < 2.5$ EM LAr calorimeter is divided into 3 layers in depth, which are finely segmented in η and ϕ ; a thin LAr presampler layer covering $|\eta| < 1.8$;

➤ Form a cluster from cells in a rectangular region $\Delta\eta \times \Delta\phi = 0.125 \times 0.1715$ around seed

❖ Selected in **barrel** $|\eta^{\gamma}| < 1.475$ & **two end-cap** $1.375 < |\eta^{\gamma}| < 3.2$;

❖ **Photon identification**: classify as electron, photon, or converted photon matching cluster with tracks; use lateral and longitudinal energy profiles of the *photon/electron* electromagnetic shower

❑ **Calorimeter isolation** in region $\Delta R_{\gamma\gamma} > 0.4$ around photon with requirement $E_T^{iso} < 0.05 \times p_T^{\gamma}$

➤ Converted and unconverted γ -s are calibrated separately use the tracking information to correct the Calorimeter response for upstream energy losses and leakage

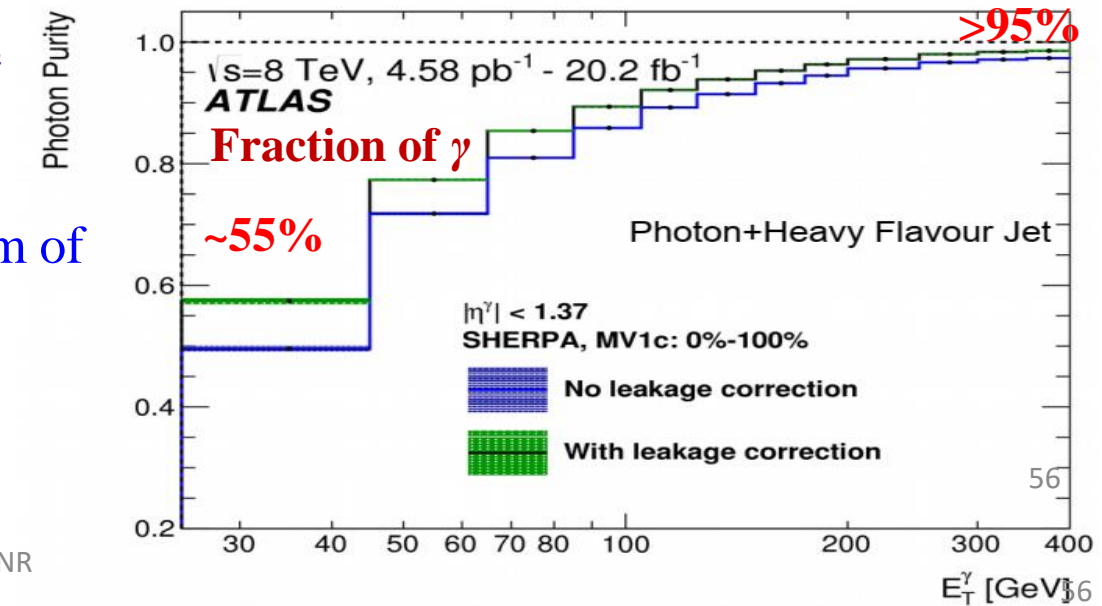
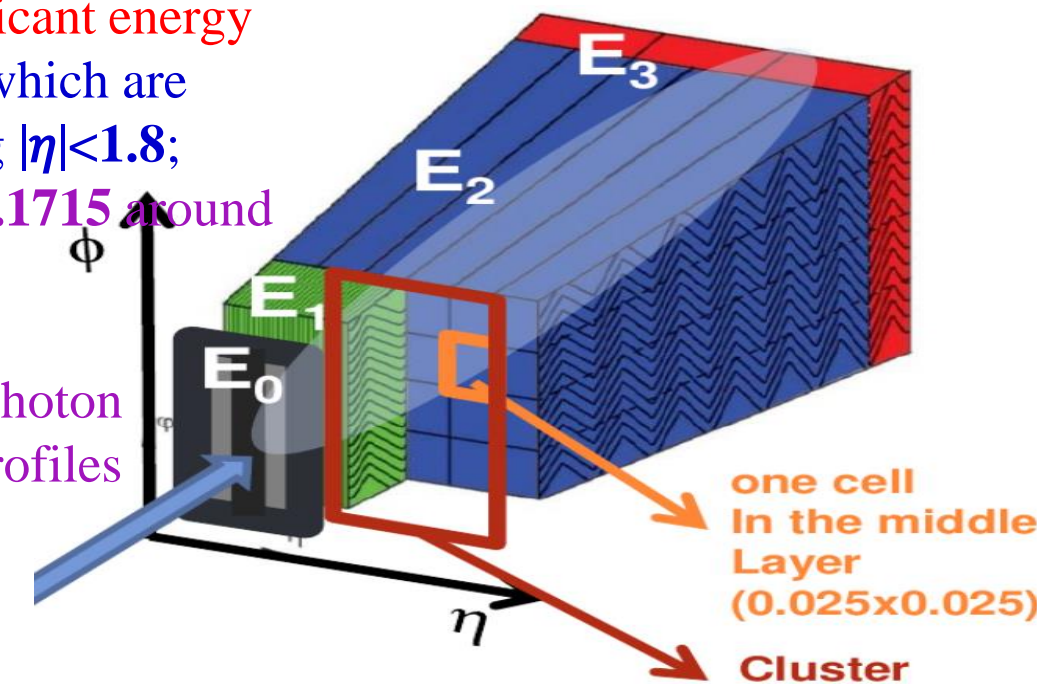
❖ **Calculate energy and direction**: photon energy a weighted sum of layer energies, with corrections for detector effects

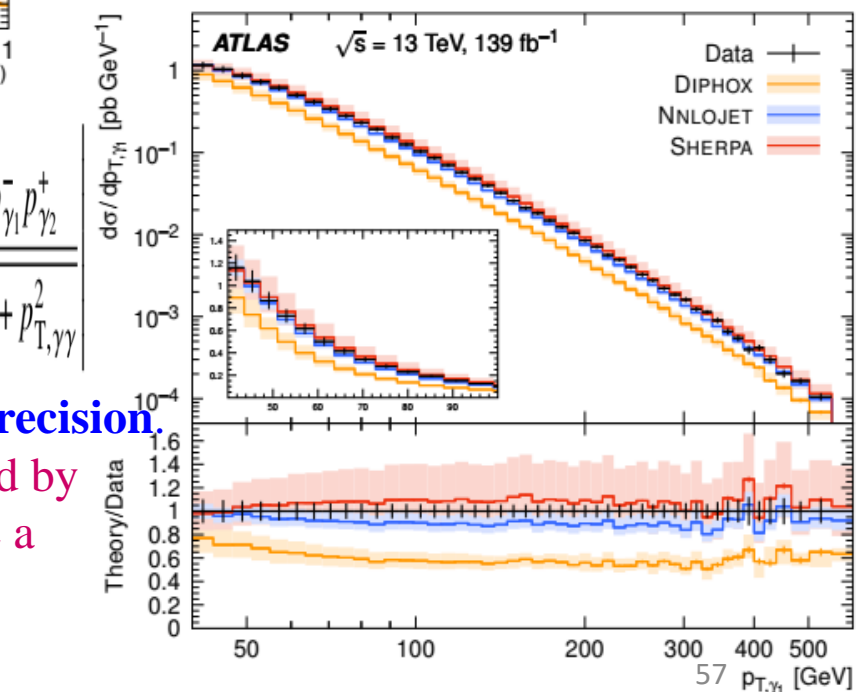
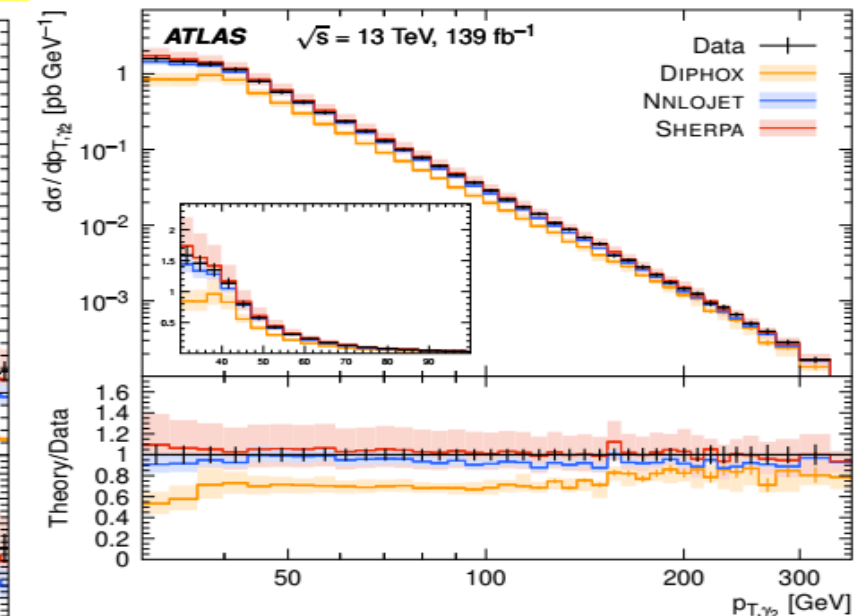
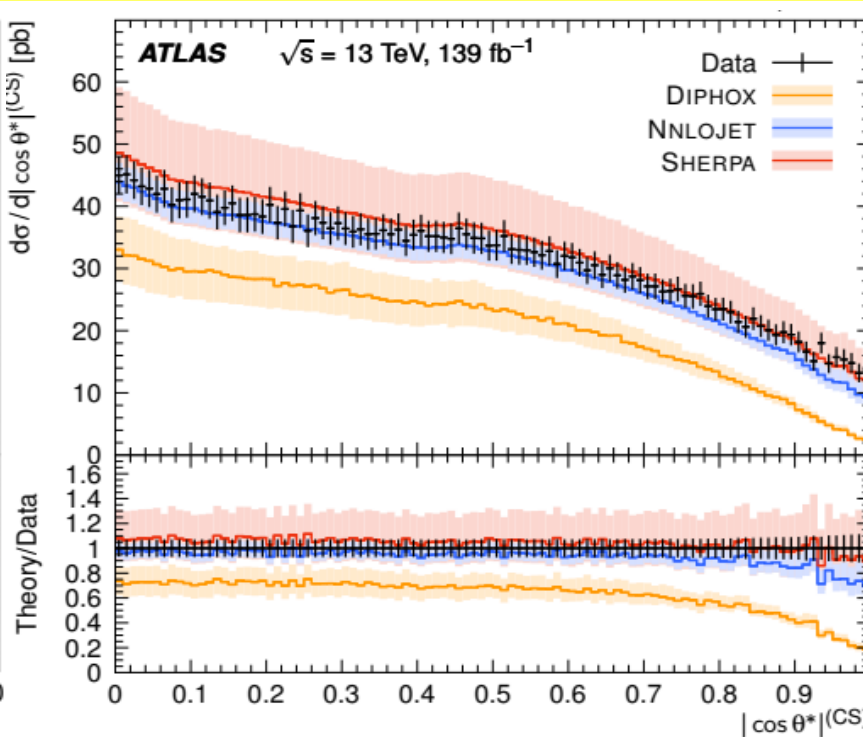
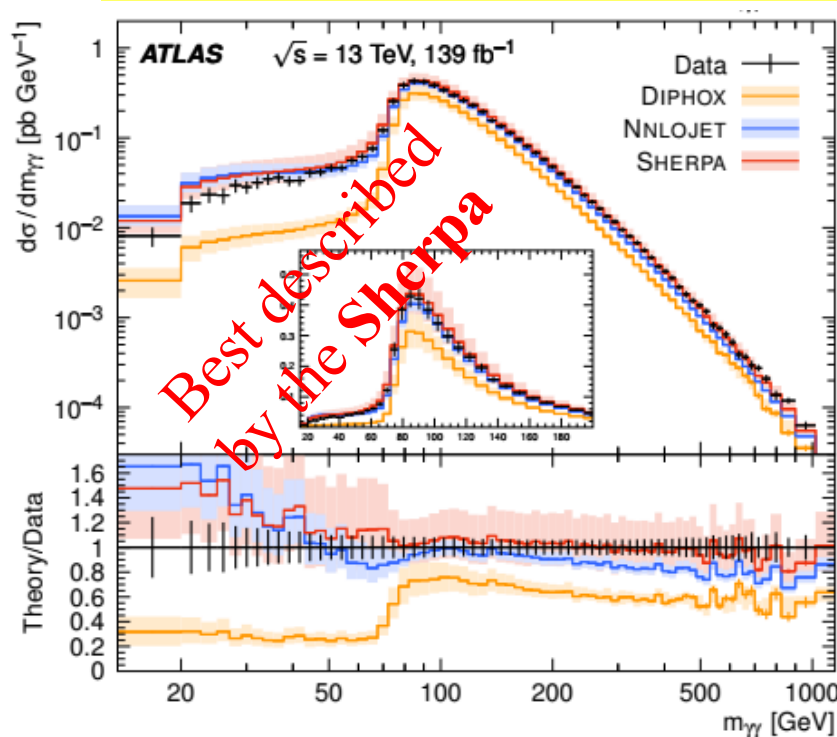
✓ Corrected for **pileup** using jet area method

➤ Use 2D-sidebands for remaining background

❖ Remove hadron and τ background

✓ Small electron background removed using MC



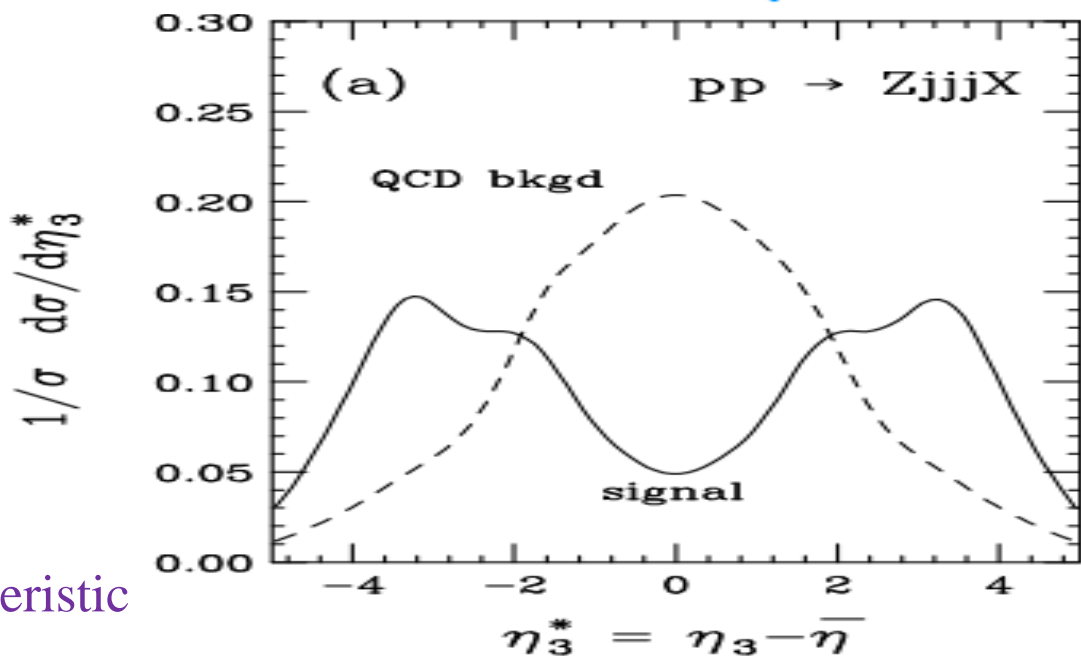
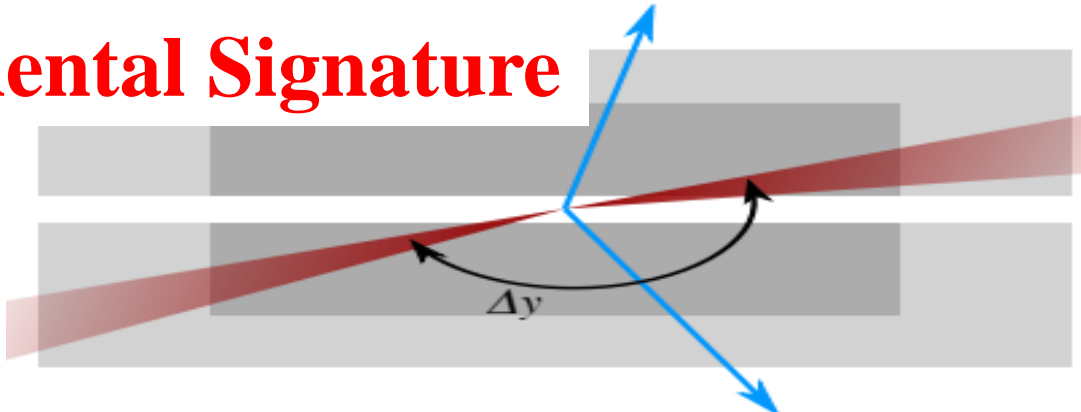
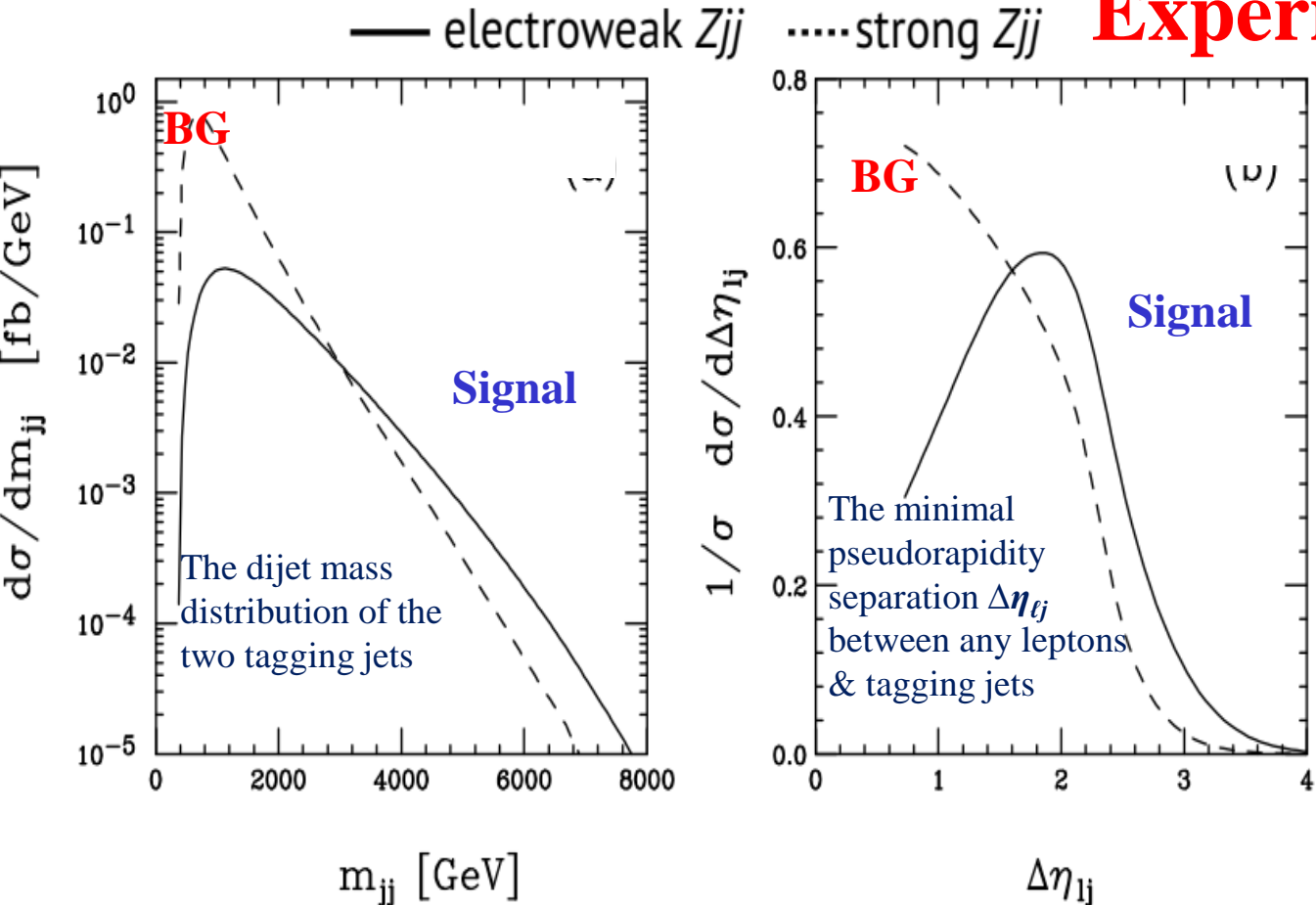


Differential cross sections measured as functions of $m_{\gamma\gamma}$, $p_{T,\gamma 1}$, $p_{T,\gamma 2}$, $|\cos \theta^*|^{(CS)}$ compared with the predictions from **Diphox NLO**, **Nnlojet NNLO**, **Sherpa MEPS@NLO**.

$$|\cos \theta^*|^{(CS)} = \frac{\sinh(\Delta\eta_{\gamma\gamma})}{\sqrt{1 + (p_{T,\gamma\gamma}/m_{\gamma\gamma})^2}} \cdot \frac{2p_{T,\gamma 1}p_{T,\gamma 2}}{m_{\gamma\gamma}^2} = \frac{p_{\gamma 1}^+ p_{\gamma 2}^- - p_{\gamma 1}^- p_{\gamma 2}^+}{m_{\gamma\gamma} \cdot \sqrt{m_{\gamma\gamma}^2 + p_{T,\gamma\gamma}^2}}$$

- Good agreement is generally found with the predictions at the **highest theoretical precision**.
- Only the merged approach with **multi-leg matrix elements at NLO**, as implemented by **Sherpa**, and a **fixed-order NNLO** calculation, as implemented by **NNLOJET**, give a **satisfactory description of the data**.

Experimental Signature



- ❑ No colour connection between scattering quarks leads to characteristic signature
- ❑ Additional activity in the event measured relative to centre of “tagging jets”:

29.08.2023 $\zeta_X = \left| \frac{y_X - (y_{j1} + y_{j2})/2}{y_{j1} - y_{j2}} \right|, \quad C_X = \exp \left[-4 \left(\frac{\eta_X - (\eta_{j1} + \eta_{j2})/2}{\eta_{j1} - \eta_{j2}} \right)^2 \right]$ Yuri Kulchitsky, P NASB & JINR

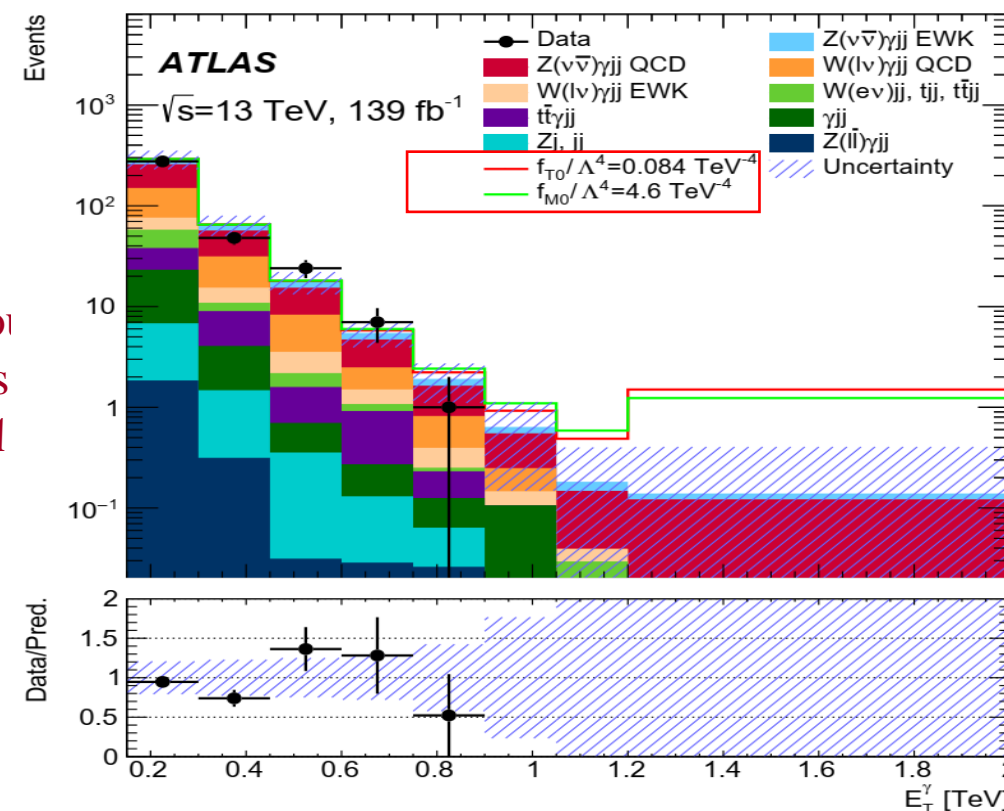
Characteristics of the third - soft jet in Z_{jjj} signal (solid) and background (dashed) events at the LHC. The pseudorapidity η_3^* is measured with respect to the center of the two tagging jets, $\langle \eta \rangle = (\eta_j^{\text{tag1}} + \eta_j^{\text{tag2}})/2$, and the distributions are normalized to unit area.

- The combination of **EW** $Z(\rightarrow \nu\nu)\gamma jj$ ($E_T^\gamma > 150 \text{ GeV}$) & ($15 < E_T^\gamma < 110 \text{ GeV}$) production yields an observed (expected) **signal significance of 6.3σ (6.6σ)**
- Limits on **anomalous Quartic Gauge Couplings (aQGC)** are obtained in the framework of **Effective Field Theory (EFT)** with **dimension-8** operators

The effect of new physics introduced by **aQGCs** can be realised using an **EFT** linearly parameterised by an **effective Lagrangian**:

$\mathcal{L} = \mathcal{L}^{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \sum_j \frac{f_j}{\Lambda^4} \mathcal{O}_j$, where \mathcal{O}_i & \mathcal{O}_j are **dimension-6** or **dimension-8** operators induced by integrating out the new degrees of freedom, while c_i & f_j represent the numerical coefficients that are meant to be derivable from a more complete high-energy theory. The Λ term is a **mass-dimension parameter** associated with the **energy scale** of the new degrees of freedom that have been integrated out.

- Having found no significant deviations from SM predictions**, the data are used **to set limits on anomalous Quartic Gauge Couplings**
- The limits are set on **Effective Field Theory** dimension-8 operators (f_{T0}/Λ^4 , f_{T5}/Λ^4 , f_{T8}/Λ^4 , f_{T9}/Λ^4 , f_{M0}/Λ^4 , f_{M1}/Λ^4 & f_{M2}/Λ^4)
- These constraints are either competitive with or **more stringent** than those previously published by **CMS**



The E_T^γ distribution in the **SR** after the fit in the **CRs**. The **red** (**green**) line shows the expected number of events in the case of non-zero **EFT** coefficient f_{T0}/Λ^4 (f_{M0}/Λ^4), values shown in the legend.

Observed (expected) significances and measured signal strengths for the individual and combined channels

Fit	Observed (expected) significances [σ]	$\mu(WWW)$
$e^\pm e^\pm$	2.3 (1.4)	1.69 ± 0.79
$e^\pm \mu^\pm$	4.6 (3.1)	1.57 ± 0.40
$\mu^\pm \mu^\pm$	5.6 (2.8)	2.13 ± 0.47
2ℓ	6.9 (4.1)	1.80 ± 0.33
3ℓ	4.8 (3.7)	1.33 ± 0.39
Combined	8.0(5.4)	1.66 ± 0.28

Uncertainty source	$\Delta\sigma/\sigma$ [%]
Data-driven background	5.3
Prompt-lepton-background modeling	3.3
Jets and E_T^{miss}	2.8
MC statistics	2.8
Lepton	2.1
Luminosity	1.9
Signal modeling	1.5
Pile-up modeling	0.9
Total systematic uncertainty	9.5
Data statistics	11.2
WZ normalizations	3.3
Total statistical uncertainty	11.6

Measured signal strength,
overall & in individual channels

Background-only hypothesis
rejected with 8.0σ , where 5.4σ
are expected

➤ The measured cross section, extrapolated to the total phase space, is:

$$\sigma_{WWW}^{\text{data}} = 820 \pm 100 \text{ (stat)} \pm 80 \text{ (syst) fb}$$

approximately 2.6σ from the predicted cross section of

$$\sigma_{WWW}^{\text{MC}} = 511 \pm 18 \text{ fb}$$

calculated at NLO QCD and LO electroweak accuracy

Wilson coefficients of **the Standard Model Effective Field Theory (SM EFT)** are constrained in a combined fit of differential cross-section measurements of the productions: WW & WZ in leptonic final states, **4 charged leptons**, a leptonically decaying **Z boson** in vector-boson-fusion topology. No significant deviations from the SM expectation are found .

- Interpretation of multiboson measurements in the SMEFT
- Expansion of SM Lagrangian in increasing powers of inverse scale of new physics, $1/\Lambda$

$$\mathcal{L}_{\text{SMEFT}} \approx \mathcal{L}_{\text{SM}}^{(4)} + \sum_i \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \sum_j \frac{c_j^{(8)}}{\Lambda^4} \mathcal{O}_j^{(8)}$$

- Leading SMEFT effect expected from interference of dim-6 operators with SM:

$$\sigma \propto |\mathcal{M}_{\text{SMEFT}}|^2 = |\mathcal{M}_{\text{SM}}|^2 + \underbrace{\sum_i \frac{c_i^{(6)}}{\Lambda^2} 2\text{Re} \left(\mathcal{M}_i^{(6)} \mathcal{M}_{\text{SM}}^* \right)}_{\text{Linear model}} + \underbrace{\sum_i \frac{\left(c_i^{(6)} \right)^2}{\Lambda^4} \left| \mathcal{M}_i^{(6)} \right|^2}_{\text{Quadratic terms}} + \underbrace{\sum_{i < j} \frac{c_i^{(6)} c_j^{(6)}}{\Lambda^4} 2\text{Re} \left(\mathcal{M}_i^{(6)} \mathcal{M}_j^{(6)*} \right)}_{\text{Cross terms}}$$

- Quadratic term at the same order, $\mathcal{O}(\Lambda^{-4})$, as SM + dim-8 interference
- Focus on operators at dim-6: 33 CP-even operators studied, assuming flavour symmetry and neglecting Higgs

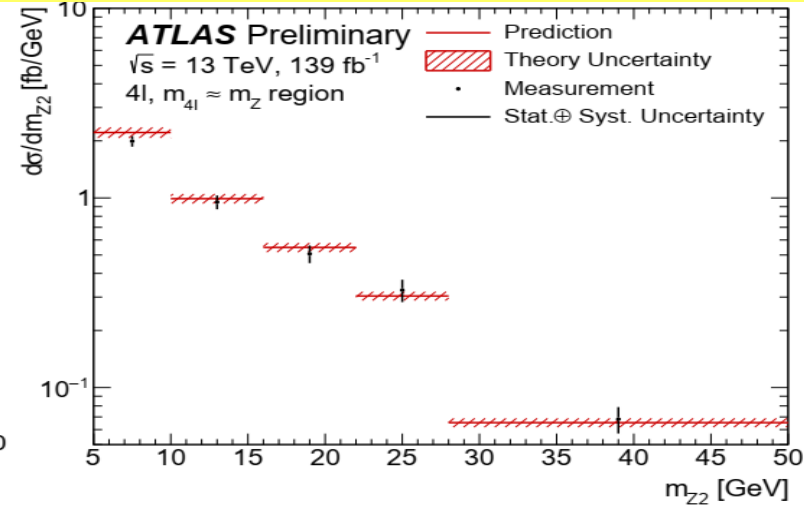
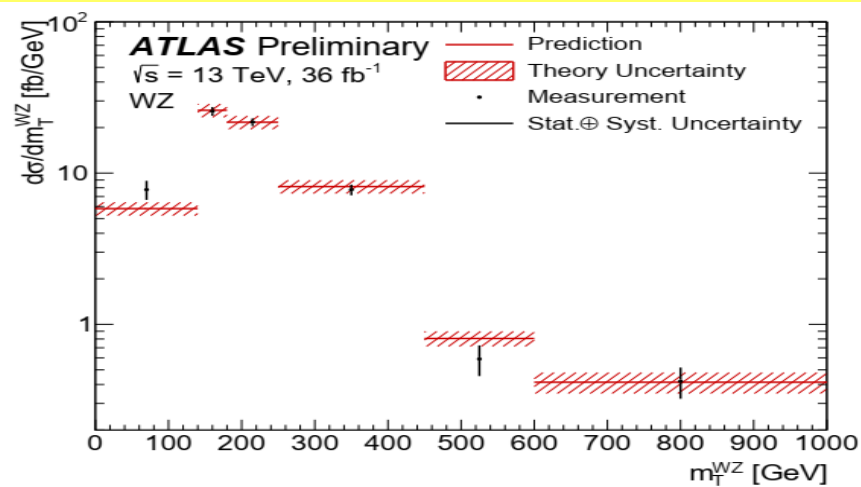
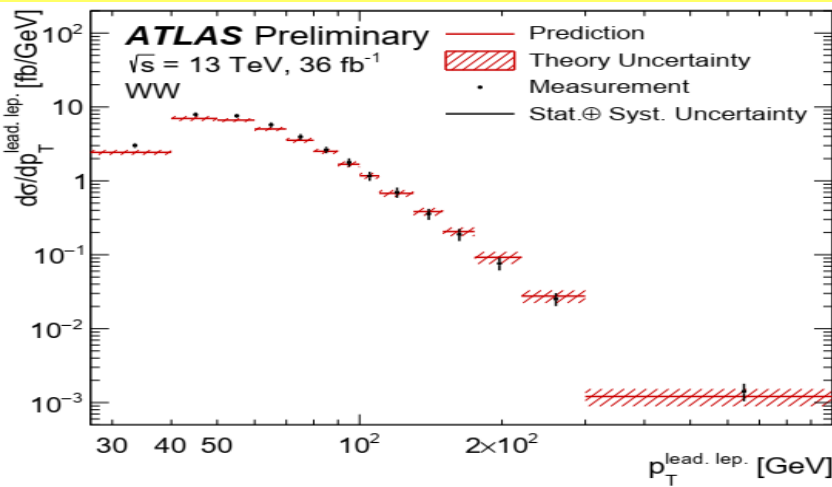
Combination of several multiboson measurements

- $pp \rightarrow WW \rightarrow e\mu\nu\nu$: Eur. Phys. J. C 79 (2019) 884 using 36 fb^{-1}
- $pp \rightarrow WZ \rightarrow lll'\nu$: Eur. Phys. J. C 79 (2019) 535 using 36 fb^{-1}
- $pp \rightarrow 4l \rightarrow 'lll'l'$: JHEP 07 (2021) 005 using 139 fb^{-1}
- $pp \rightarrow Zjj \rightarrow lljj$: Eur. Phys. J. C 81 (2021) 163 using 139 fb^{-1}

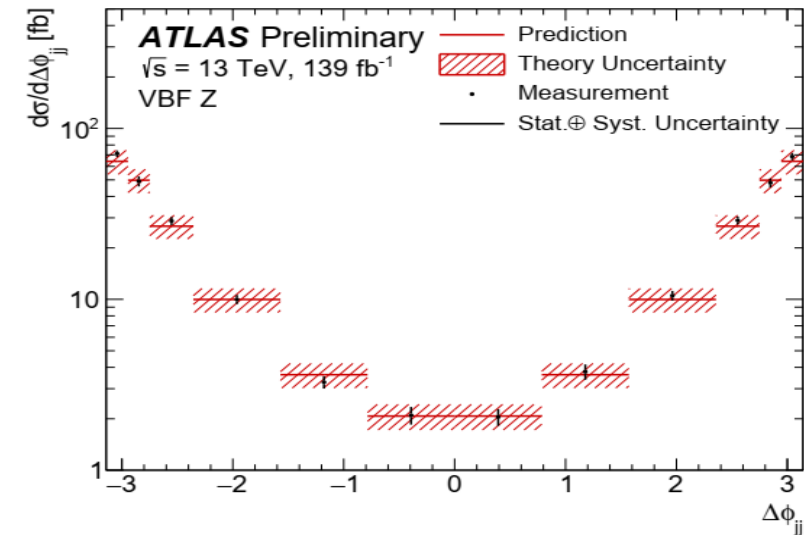
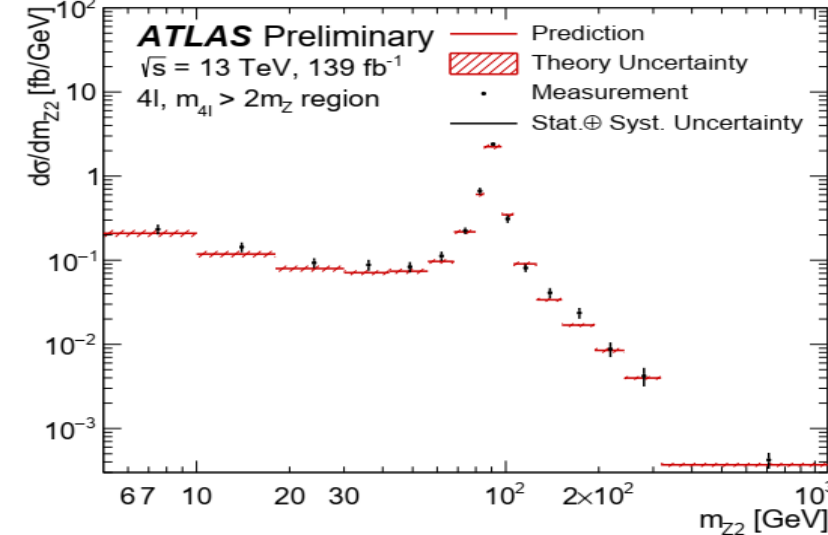
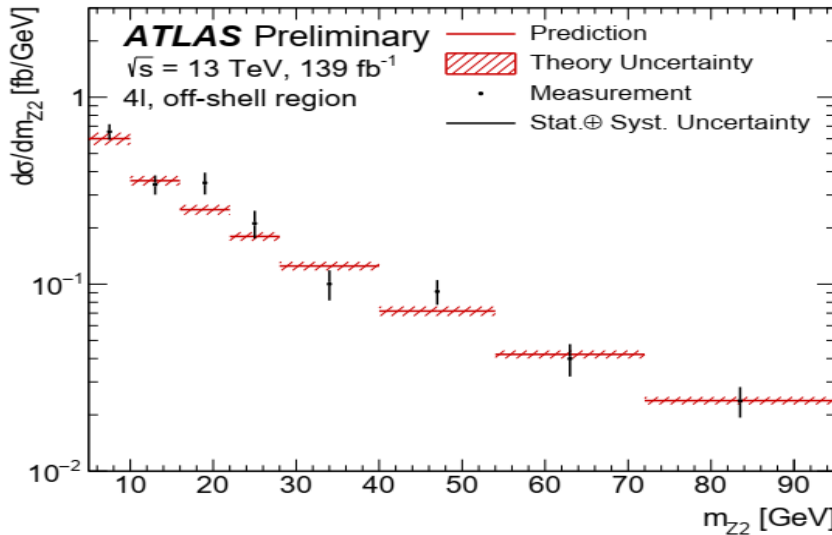
- Higgs-boson production kinematically suppressed:

- see ATLAS-CONF-2020-053 for dedicated EFT study
- see ATL-PHYS-PUB-2021-010 for a $H!WW^*$ and WW combination

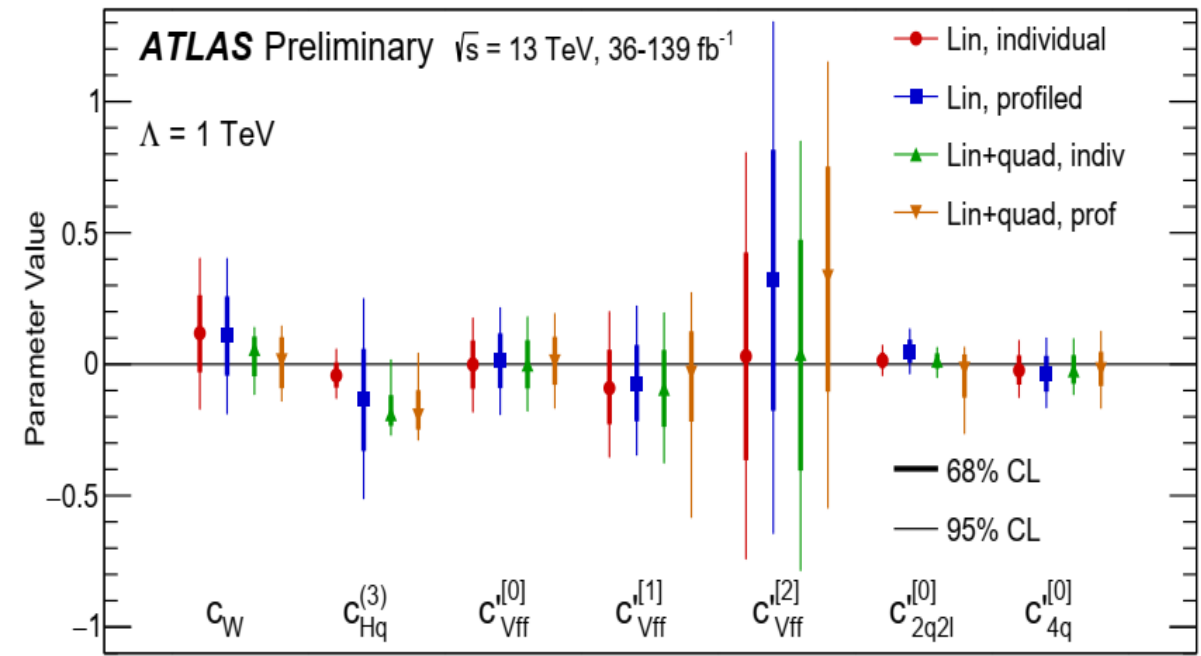
- Measurements with high precision and small background contributions
- Sensitive to a large number of dim-6 operators affecting
 - gauge-boson self-couplings
 - couplings of gauge bosons and fermions
 - four-fermion couplings



Assuming smooth EFT effects have the folding matrix of the SM. Injection of BSM physics has been explicitly tested in m_{4l}



Measured differential cross-sections compared to the SM theory predictions used in this analysis. The $p_T^{\text{lead.lep.}}$ in WW production, the WZ transverse mass in WZ production, the invariant mass of the secondary lepton pair in the three $4l$ analysis regions, the signed azimuthal angle between the two jets in $Z+\text{jets}$ production with a VBF topology.



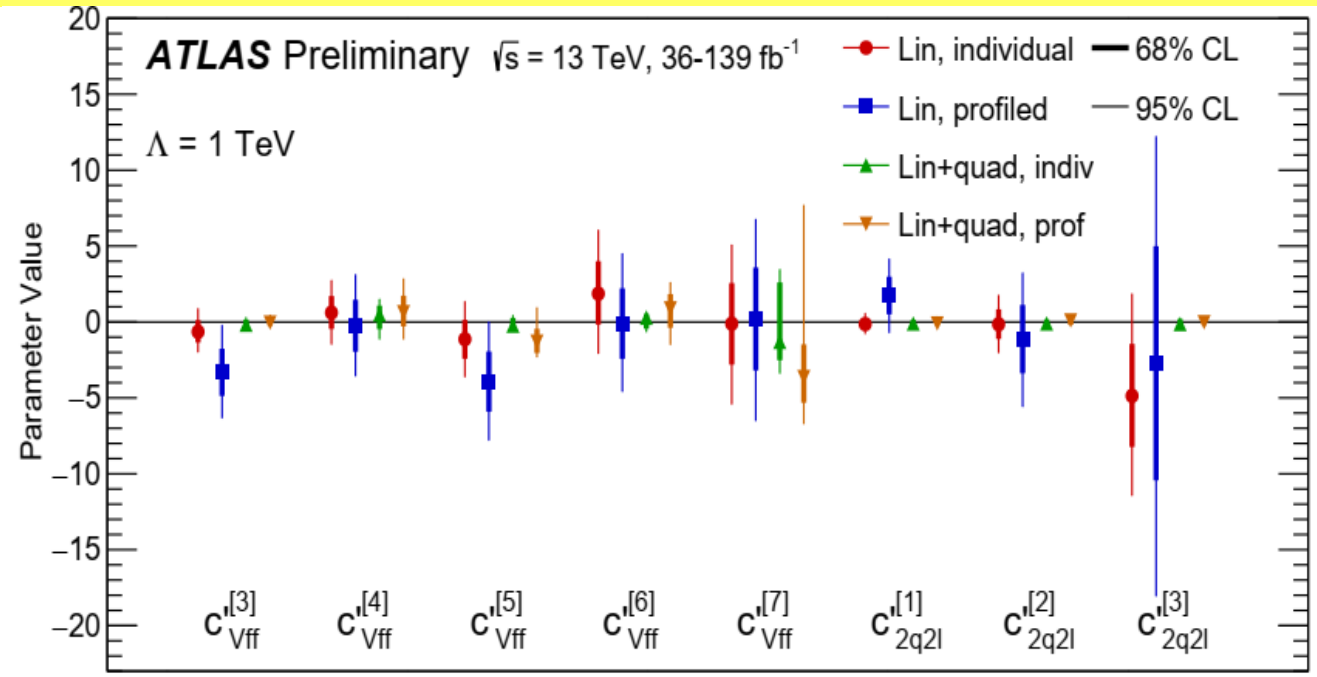
$$C_{Vff}^{[0]} \approx 0.81c_{HWB} + 0.38c_{HD} + 0.13c_{Hl}^{(1)} + 0.37c_{Hl}^{(3)} - 0.14c_{ll}^{(1)} + 0.12c_{Hq}^{(1)}$$

$$C_{2q2l}^{[0]} \approx -0.37c_{lq}^{(1)} + 0.89c_{lq}^{(3)} - 0.11c_{lu} - 0.21c_{eu} - 0.13c_{qe}$$

$$C_{Vff}^{[1]} \approx 0.73c_{Hl}^{(1)} - 0.28c_{Hl}^{(3)} - 0.48c_{He} + 0.38c_{ll}^{(1)} + 0.13c_{Hq}^{(1)}$$

$$C_{4q}^{[0]} \approx 0.11c_{qq}^{(11)} + 0.22c_{qq}^{(18)} + 0.95c_{qq}^{(31)} - 0.2c_{qq}^{(38)}$$

$$C_{Vff}^{[2]} \approx 0.37c_{HWB} + 0.17c_{HD} - 0.31c_{Hl}^{(1)} - 0.53c_{Hl}^{(3)} + 0.25c_{He} + 0.59c_{ll}^{(1)} - 0.21c_{Hq}^{(1)}$$



$$C_{Vff}^{[3]} \approx -0.19c_{Hl}^{(1)} - 0.14c_{Hl}^{(3)} + 0.86c_{Hq}^{(1)} + 0.41c_{Hu} - 0.17c_{Hd}$$

$$C_{Vff}^{[4]} \approx -0.35c_{HWB} + 0.49c_{HD} + 0.26c_{Hl}^{(1)} + 0.35c_{Hl}^{(3)} + 0.51c_{He} + 0.38c_{ll}^{(1)} + 0.18c_{Hq}^{(1)}$$

$$C_{Vff}^{[5]} \approx 0.25c_{HD} + 0.33c_{Hl}^{(1)} - 0.22c_{Hl}^{(3)} + 0.18c_{He} - 0.35c_{ll}^{(1)} - 0.3c_{Hq}^{(1)} + 0.71c_{Hu} - 0.16c_{Hd}$$

$$C_{Vff}^{[6]} \approx -0.22c_{Hl}^{(1)} + 0.52c_{Hl}^{(3)} - 0.39c_{He} + 0.44c_{ll}^{(1)} - 0.22c_{Hq}^{(1)} + 0.52c_{Hu}$$

$$C_{Vff}^{[7]} \approx -0.28c_{HWB} + 0.71c_{HD} - 0.31c_{Hl}^{(1)} - 0.21c_{Hl}^{(3)} - 0.5c_{He} - 0.14c_{ll}^{(1)}$$

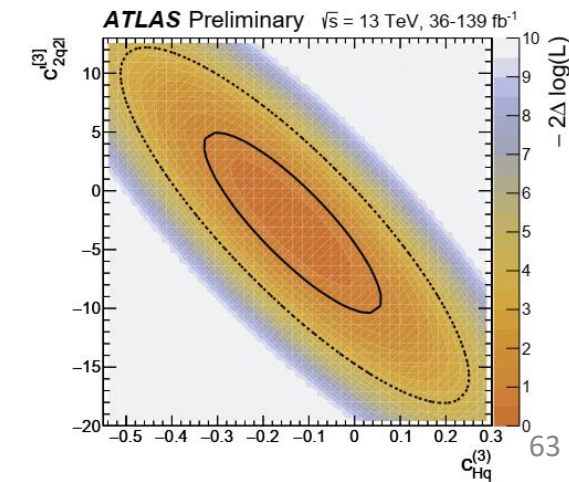
$$C_{2q2l}^{[1]} \approx 0.56c_{lq}^{(1)} + 0.44c_{lq}^{(3)} + 0.61c_{eu} - 0.1c_{ed} + 0.34c_{qe}$$

$$C_{2q2l}^{[2]} \approx 0.68c_{lq}^{(1)} + 0.15c_{lq}^{(3)} + 0.33c_{lu} - 0.51c_{eu} + 0.13c_{ed} - 0.37c_{qe}$$

$$C_{2q2l}^{[3]} \approx -0.27c_{lq}^{(1)} + 0.79c_{lu} - 0.39c_{ld} + 0.26c_{eu} - 0.22c_{ed} - 0.16c_{qe}$$

Confidence intervals for the 15 parameters included in the combined maximum likelihood fit. Results are quoted both for fits linear in the parameters and for fits that take into account also quadratic contributions. The first case corresponds to a model in which only the $\mathcal{O}(\Lambda^{-2})$ contributions to the cross-section prediction, the interference between SM and dimension-six operators, is included. The latter case also includes quadratic dimension-six contributions, which are part of the $\mathcal{O}(\Lambda^{-4})$ contributions. Comparisons of the two results can be used to estimate uncertainties due to the truncation of the EFT expansion.

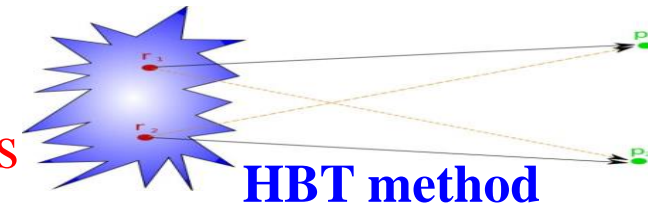
- ❑ Limits set at 95% confidence level, both for the “linear” and “linear plus quadratic” models
- ❑ Fits of individual coefficients, as well as combined fit



- ❑ An analysis at NNLO order in the theory of QCD for the determination of a new set of parton distribution functions (PDF) using diverse measurements in pp collisions at $\sqrt{s} = 7, 8$ and 13 TeV, performed by the ATLAS at the LHC, together with deep inelastic scattering data from ep collisions at the HERA.
- ❑ The ATLAS data sets considered are differential cross-section measurements of inclusive W^\pm and Z/γ^* boson production, W^\pm and Z boson production in association with jets, $t\bar{t}$ production, inclusive jet production and direct photon production. The resulting set of PDF is called **ATLASpdf21**.
- It is observed that the addition of the **ATLAS** data sets to the **HERA** data brings the PDFs much closer to the **global PDFs** of MSHT, CT and NNPDF than to HERAPDF2.0.
- The **ATLASpdf21** PDFs agree with these global fits as well as they agree with each other.
- Thus, ATLAS data seem to be able to replicate many of the features that the fixed-target deep inelastic scattering and Drell–Yan data plus the Tevatron Drell–Yan data bring to the global PDFs.
- Using only the HERA and ATLAS data allows a detailed treatment of correlated systematic uncertainties.

Correlations in phase space between two identical bosons from symmetry of wave functions.

- ▶ Enhances likelihood of *two particles close* in phase space
- ▶ Allows one to ‘probe’ the source of the bosons *in size and shape*
- ▶ Dependence on particle multiplicity and transverse momentum probes *the production mechanism*



Correlation function $C_2(Q)$ a ratio of probabilities:

$$C_2(Q) = \frac{\rho(p_1, p_2)}{\rho_0(p_1, p_2)} = C_0(1 + \Omega(\lambda, RQ)) \cdot (1 + Q\varepsilon),$$

$$\Omega^E(\lambda, RQ) = \lambda e^{-RQ}$$

$$\Omega^G(\lambda, RQ) = \lambda e^{-R^2 Q^2}$$

$$Q^2 = -(p_1 - p_2)^2$$

C_0 is a normalisation, ε accounts for long range effects, R is the effective radius parameter of the source, λ is the strength of the effect parameter, 0/1 for coherent/chaotic source.

Two possible parameterisation: **Gaussian and Exponential.**

$$C_2(Q) = \frac{N^{++,--}(Q)}{N^{ref}(Q)}$$

N_{ref} without BEC effect from: unlike-charge particles (UCP), opposite hemispheres, event mixing.
Basic Reference: distribution of UCP pairs of non-identical particle taken from the same event.

$$R_2(Q) = \frac{C_2^{Data}(Q)}{C_2^{MC}(Q)} = \frac{\rho^{++,-,-} / \rho^{+-}}{\rho^{MC}^{++,-,-} / \rho^{MC}^{+-}}$$

29.08.2023

The studies are carried out using the **double ratio correlation function**. The $R_2(Q)$ eliminates problems with energy-momentum conservation, topology, resonances, hadronic jets, mini-jets etc. **MC without BEC.**

