

“Active Role of Gluons in Multiparticle Production”

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(JINR, RF и GSTU, Gomel, RB)

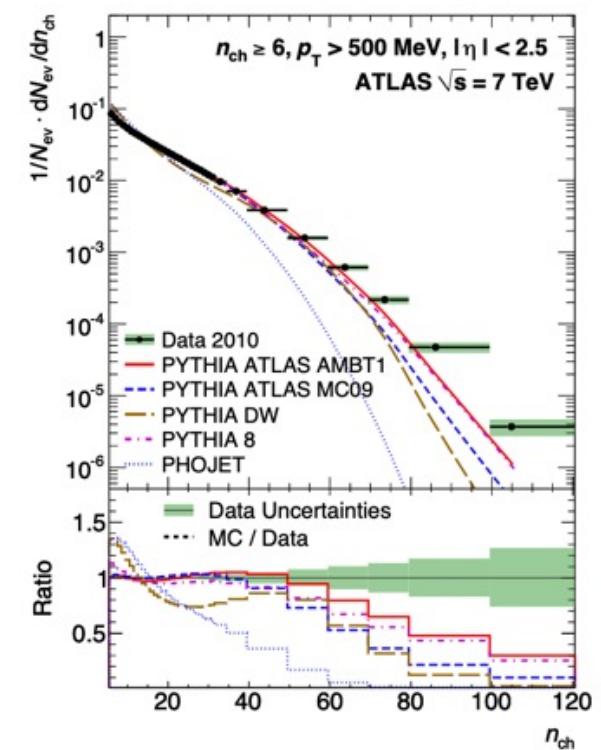
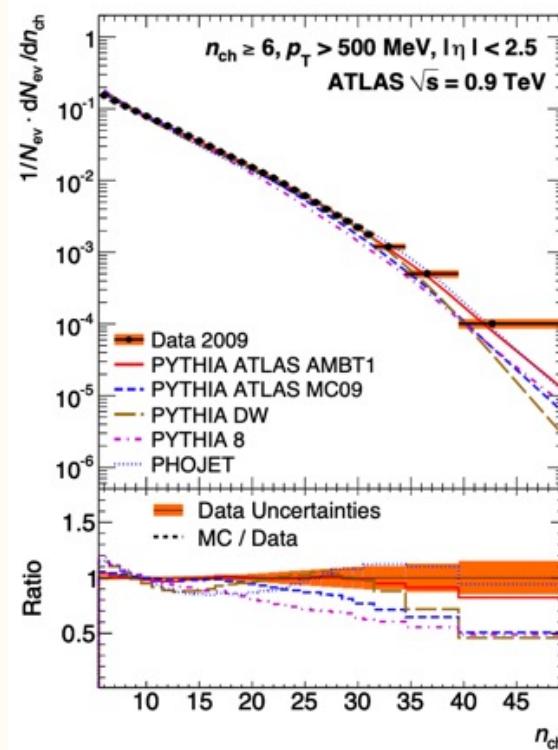
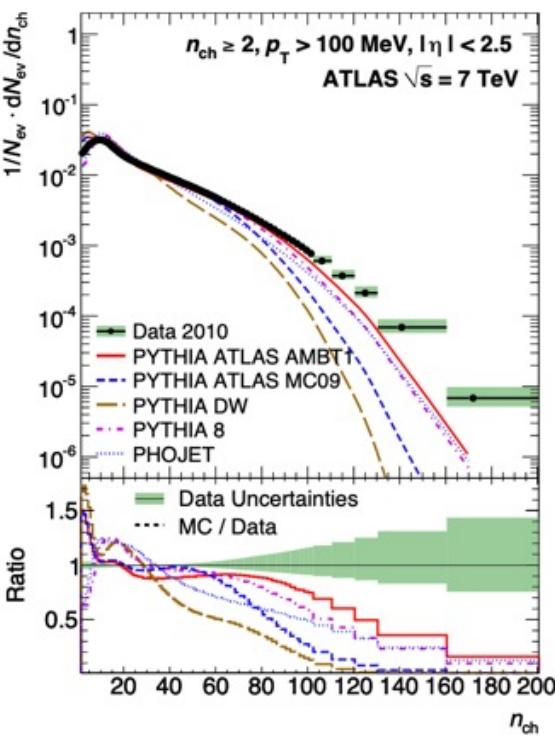
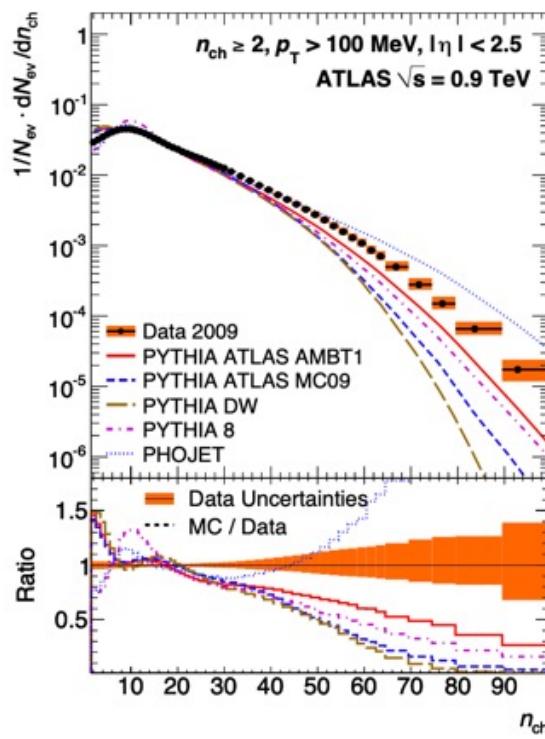
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Multiparticle processes studying in HEP

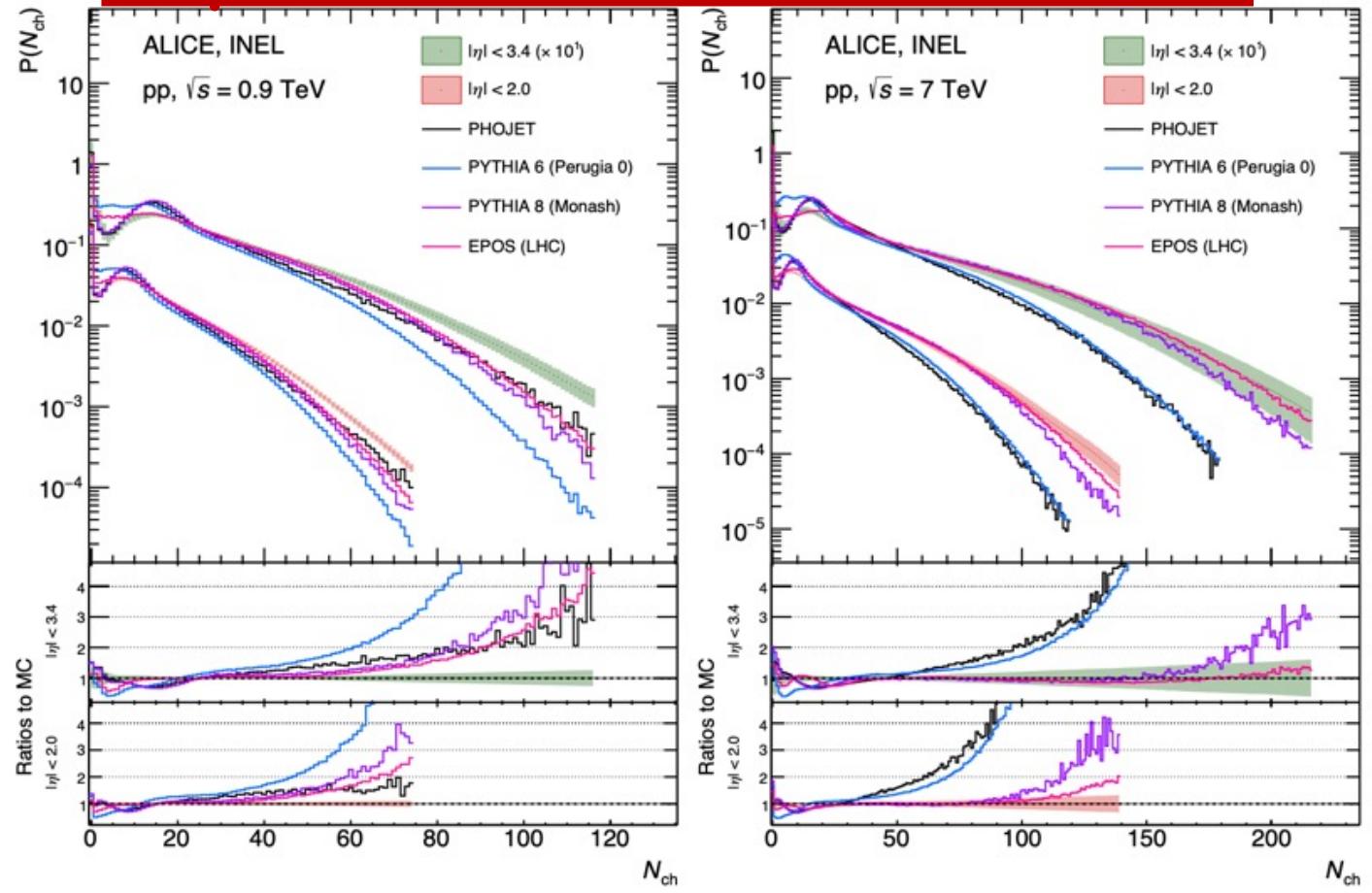
1. Electron-positron annihilation (e^+e^-).
2. Proton & proton-antiproton interactions.
Bottomonium decay.
3. "Thermalization" project (high multiplicity).
Collective phenomena.
4. Relativistic nuclear physics (AA).

Experiments at LHC (ATLAS)



ATLAS's data [2010] for MD in pp interactions at 0.9 and 7 TeV
Comparisons with Monte Carlo generators (PYTHIA, PHOJET).

Experiments at LHC (ALICE)



[arXiv:1708.01435, ALICE Coll.]: MD of charged particles in high pseudo-rapidity region (pp collisions at 0.9 и 7 ТэВ).

e^+e^- annihilation. I stage

Konishi et al. & Giovannini [NP, 1979] describe quark-gluon cascade in pQCD in view of elementary processes

1) Quark emission of gluon - $q \rightarrow q + g$, (\tilde{A})

2) Gluon fission - $g \rightarrow g + g$, (A)

3) Quark-antiquark pair creation from gluon - $g \rightarrow q + \bar{q}$.

$$\left[\begin{array}{l} \frac{\partial G}{\partial Y} = -AG + AG^2, \\ \frac{\partial Q}{\partial Y} = -\tilde{A}Q + \tilde{A}QG. \end{array} \right]$$

System of diff. eq. describing branching processes, leads to Polia (NBD) for q-jet and Farry MD for g-jet:

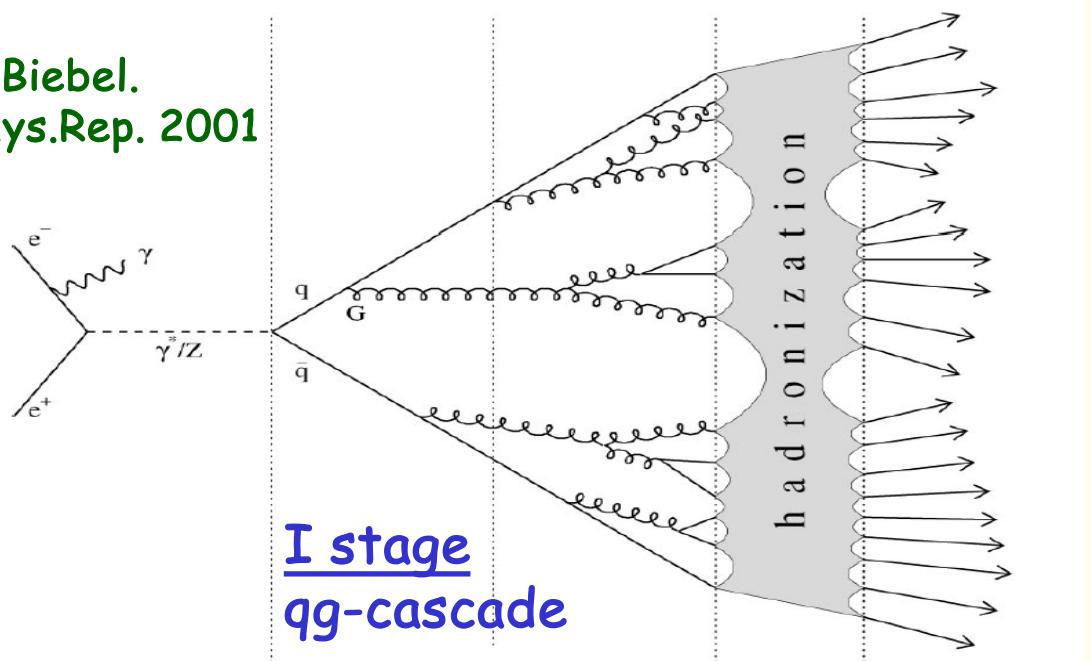
$$P_m^g = \frac{1}{\bar{m}} \left(1 - \frac{1}{\bar{m}}\right)^{m-1}, \quad P_m^q = \frac{k_p(k_p+1)\dots(k_p+m-1)}{m!} \left(\frac{\bar{m}}{\bar{m}+k_p}\right)^m \left(\frac{k_p}{\bar{m}+k_p}\right)^{k_p}.$$

$$Y = \frac{1}{2\pi b} \ln[1 + ab \ln(Q^2 / \mu^2)], \quad \tilde{A} \text{ и } A - \text{вероятности процессов 1) и 2), } k_p = \tilde{A}/A.$$

e^+e^- - annihilation

$$e^+e^- \rightarrow \gamma(Z^0) \rightarrow q\bar{q} \rightarrow (q,g) \rightarrow ? \rightarrow hadrons$$

O.Biebel.
Phys.Rep. 2001



II stage
(hadronization)

Multiplicity
Distribution (MD)

Generation
function (GF):

GF \leftrightarrow MD

Correlative moments, F_k :

$$P_n(s) = \frac{\sigma_n}{\sum_m \sigma_m}$$

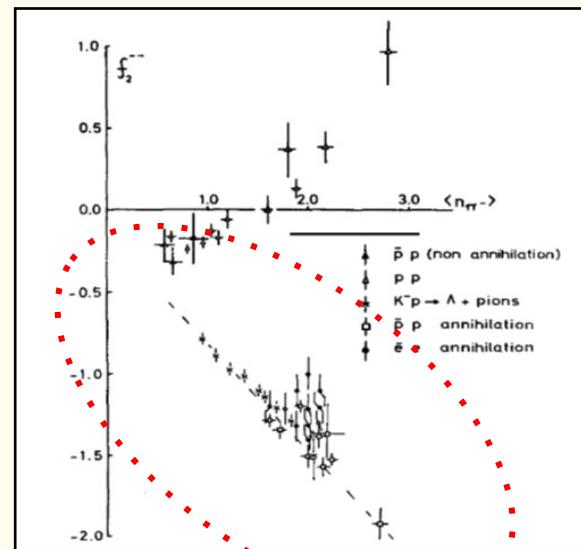
$$Q(s, z) = \sum_n P_n(s) z^n$$

$$P_n(s) = \frac{1}{n!} \left. \frac{\partial^n}{\partial z^n} Q(s, z) \right|_{z=0}$$

$$F_k(s) = \overline{n(n-1)...(n-k+1)} = \left. \frac{\partial^k}{\partial z^k} Q(s, z) \right|_{z=1}$$

e^+e^- annihilation - II stage

pQCD is unable to describe hadronization. The choice of MD at this stage is based on experimental behavior of the second correlative moment $f_2 = \langle n(n-1) \rangle - \langle n \rangle^2 = D_2 - \langle n \rangle$ (D_2 - variance).



The independent formation of particles is described by the Poisson distribution with $f_2 = 0$. Polya MD at the first stage corresponds to **negative binomial distribution (NBD)**:

$$f_2 = \overline{n}(n-1) - \overline{n}^2 \rightarrow \frac{\overline{m}^2}{k_p} > 0$$

We chose **binomial MD (Bernoulli)** for II-stage:

$$P_P^H(n) = C_{N_p}^n \left(\frac{\overline{n}_p^h}{N_p} \right)^n \left(1 - \frac{\overline{n}_p^h}{N_p} \right)^{N_p-n}, P = q, g.$$

Convolution of two stages

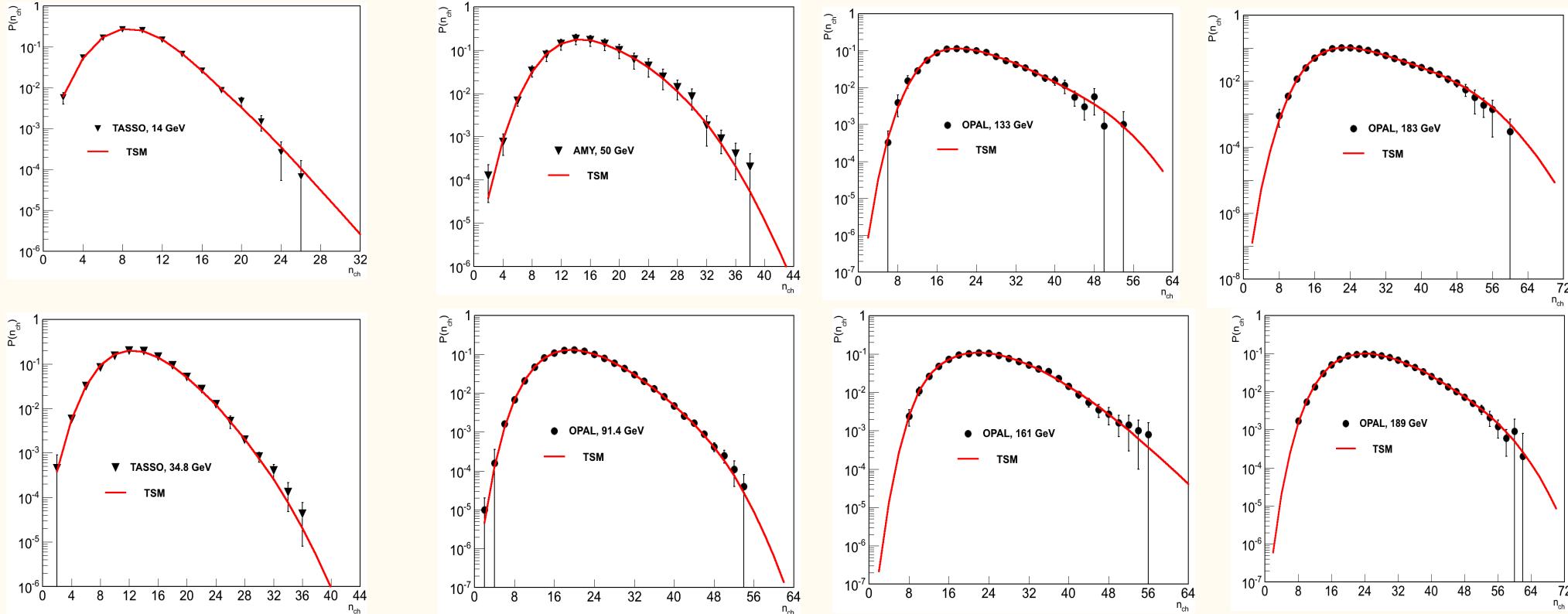
$$Q(s, z) = \sum_m P_m^P Q^H(m, s, z) \quad (\text{soft decoloration}).$$

$$P_n(s) = \Omega \sum_{m=0}^{M_g} P_m^P C_{(2+\alpha m)N}^n \left(\frac{\bar{n}^h}{N} \right)^n \left(1 - \frac{\bar{n}^h}{N} \right)^{(2+\alpha m)N-n}$$

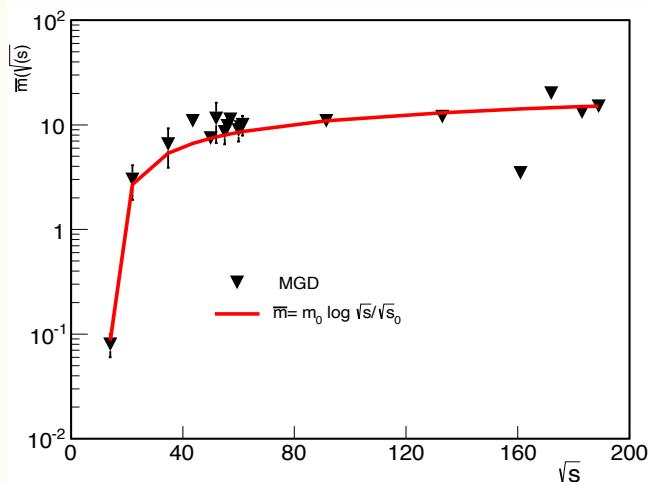
$$Q_p^H = \left[1 + \frac{\bar{n}_p^h}{N_p} (z - 1) \right]^{N_p}, \quad p = q, g, \quad f_2 = -\frac{(\bar{n}_p^h)^2}{N_p} < 0.$$

Model parameters: $N_q = N$, $\alpha = N_g/N$, \bar{n}^h .

MD in e^+e^- annihilation (14 - 189 Γ_{eB})

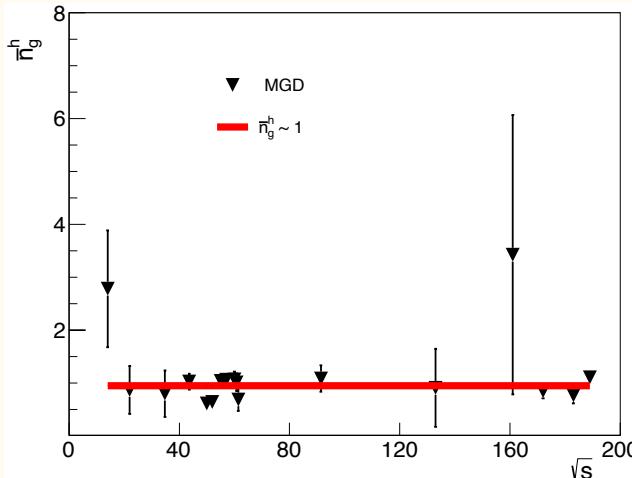


Parameters of Model

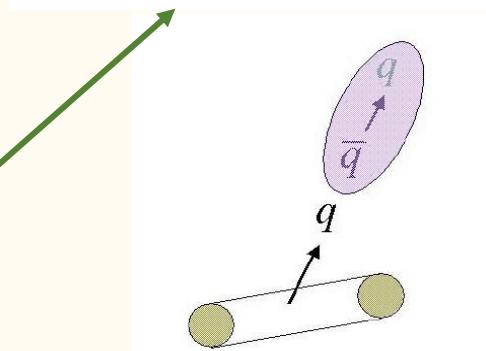


$$\bar{m} \sim \log s.$$

Hypothesis of parton-hadronic duality (LoPHD)
 $\langle m \rangle \sim \langle n \rangle$

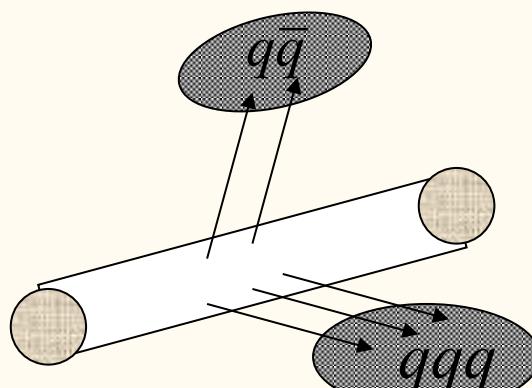


Fragmentation mechanism



Recombination mechanism

[B. Muller. 2004]



Average number of hadrons, \bar{n}_g^h , formed from gluon, is close to 1, which testifies the fragmentation mechanism of hadronization.

Parameter of model k_p

k_p is determined by the ratio of the bremsstrahlung contributions of active gluons ($q \rightarrow q + g$) to their division ($g \rightarrow g + g$), \tilde{A}/A , at the cascade stage. It takes on values more than 1, which indicates the predominance of bremsstrahlung over fission. It tends to decrease with increasing energy. Within the statistical bootstrap model, it was shown, k_p^{-1} can be interpreted as the "temperature" at that stage

$$k_p^{-1} = T_0 + \frac{1}{c} E$$

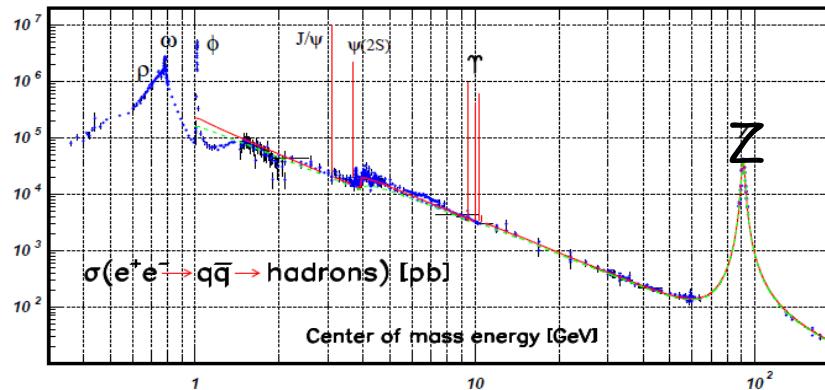
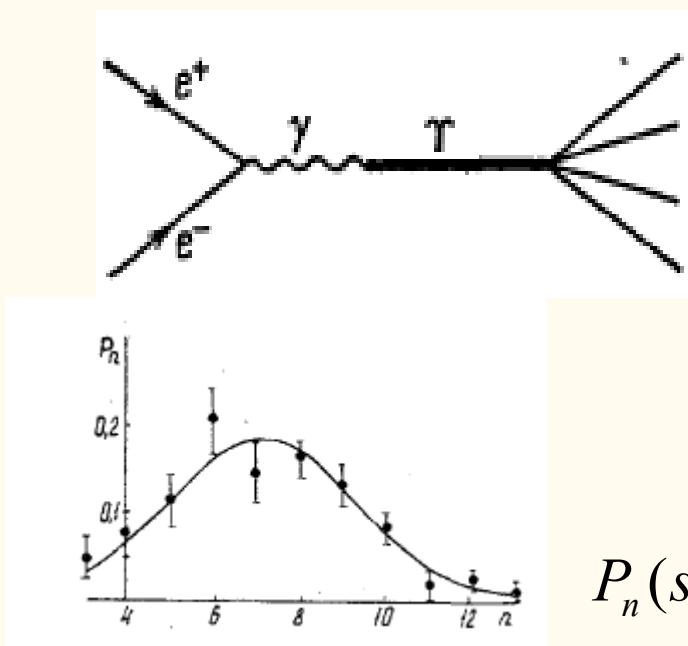
it rises with energy.

f_2 Sign changing with energy growth from “-” to “+”

$$\begin{aligned} f_2 = F_2 - F_1^2 &= \sum_{m=0} (2 + \alpha m) \left(2 + \alpha m - \frac{1}{N} \right) P_m^q \cdot (\bar{n}^h)^2 - [(2 + \alpha \bar{m}) \bar{n}^h]^2 = \\ &= \left[\alpha^2 \frac{\bar{m}^2}{k_p} + \alpha^2 \bar{m} - \frac{2 + \alpha \bar{m}}{N} \right] (\bar{n}^h)^2. \end{aligned}$$

Parameters: $\alpha \sim 1$, $N \sim 6$. At $\sqrt{s} < 5$ $\bar{m} \ll 1$ and $f_2 < 0$.
At $\sqrt{s} > \gtrsim 10$, $\bar{m} > 10$ and the sign of the second
correlative moment is changed : $f_2 > 0$.

Three-gluon decay of quarkoniums $\Upsilon(9.46)$, $\Upsilon(10.02)$



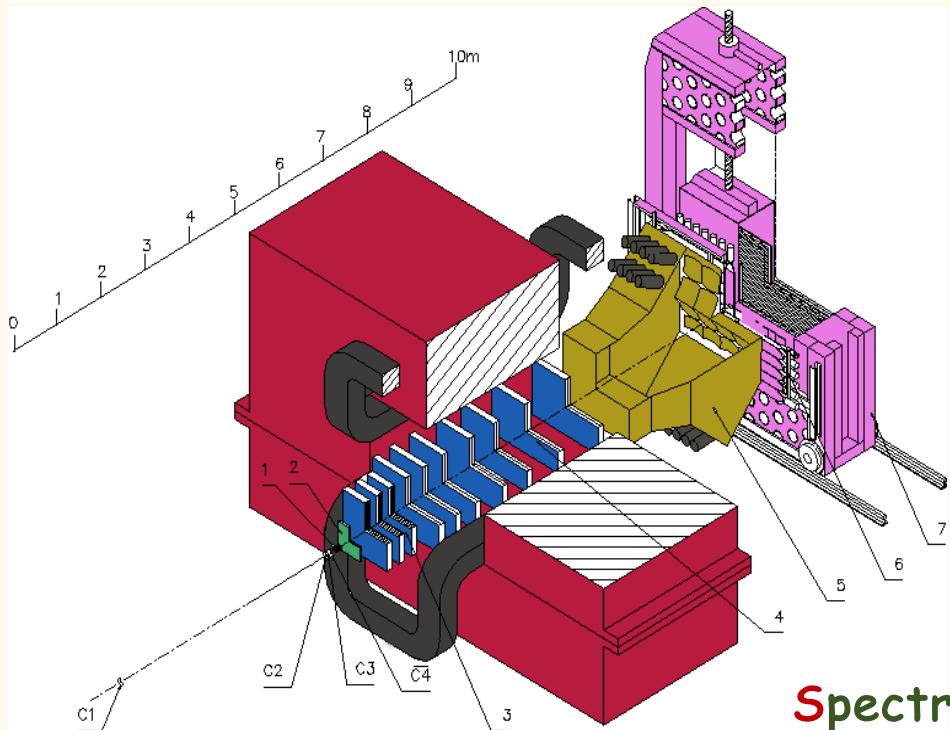
$$P_n(s) = \sum_{m'=0} \frac{(m'-1)(m'-2)}{2(\bar{m}/3)^2} \left(1 - \frac{1}{\bar{m}/3}\right)^{m'} C_{(3+m')N_g}^n \left(\frac{\bar{n}_g^h}{N_g}\right)^n \left(1 - \frac{\bar{n}_g^h}{N_g}\right)^{(3+m')N_g - n}$$

$$\Delta \bar{n}_{TSM}(s) = [\alpha(\bar{m}' - \bar{m}_{(q)}) - 3(\alpha - 2/3)] \bar{n}_q^h \quad \Delta \bar{n}_{\text{exp}}(s) \approx \Delta \bar{n}_{TSM}(s) \approx 0.8$$

Hadron Interactions with High Multiplicity (HM).

"Thermalization" project (IHEP+SINP MSU+JINR)

$$p + p \rightarrow 2N + \pi_1 + \pi_2 + \cdots + \pi_n$$



Spectrometer with Vertex Detector
(SVD-2)

Project started in 2005, SVD-2 setup at U-70 accelerator, 50-GeV p-beam, H₂ target, high multiplicity trigger. Kinematical limit (total energy spent on secondary particle production at rest) is close to 59 pions.

Gluon Dominance Model (GDM)

Modification of the two stage model has been carried out: at the initial stage valence quarks and nascent gluons in accordance with pQCD branch due elementary fission processes (1-3). Then they fragment into real hadrons in accordance with the binomial distribution confirmed in e^+e^- annihilation. The scheme of convolution combines both stages.

A comparison of this model with data (Mirabelle, U-70) showed that the hadronization parameters differ considerably from the values obtained in e^+e^- annihilation, they were much less than 1. It meant that the number of valence quarks involved in branching should be reduced. Participation of all three pairs of quarks is unlikely.

Gluon Dominance Model (GDM)

Our study showed: hadronization parameters grow with decreasing of valence quark pairs. Only their complete exclusion improved considerably χ^2 for MD description, but hadronization parameter, the average number of hadrons produced from gluon at the second stage increases with energy, exceeding 1.

Valence quarks are staying in leading particles, secondaries are created by active gluons, which decay into quark-antiquark pairs, and form colorless hadrons.

Soft gluons (~50%) remain in the qg-system and partly re-emit soft photons leading to their excess yield.

GDM, g-fission:



$$P_k = \frac{e^{\bar{k}^k} \bar{k}^k}{k!}; \quad k > 1, P_m^P = \frac{1}{m} \left(1 - \frac{1}{m}\right)^{m-1},$$

$$P_n^h = C_{\delta m N_g}^{n-2} \left(\frac{\bar{n}_g^h}{N_g}\right)^{n-2} \left(1 - \frac{\bar{n}_g^h}{N_g}\right)^{\delta m N_g - (n-2)};$$

Almost a half of gluons ($\sim 47\%$) remain in qg-system, without being fragmented into hadrons.

GDM, scheme without gluon fission

$$P_n(s) = \Omega \sum_{m=1}^{ME} \frac{\bar{m}^m e^{-\bar{m}}}{m!} \cdot C_{mN}^{n-2} \left(\frac{\bar{n}^h}{N} \right)^{n-2} \left(1 - \frac{\bar{n}^h}{N} \right)^{mN-(n-2)}.$$

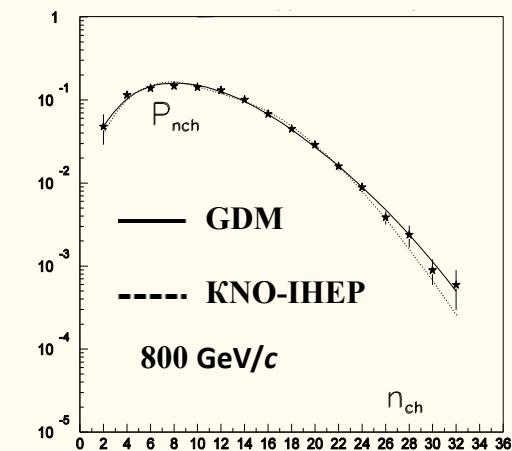
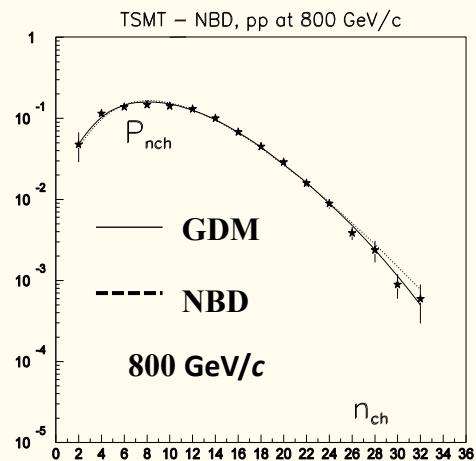
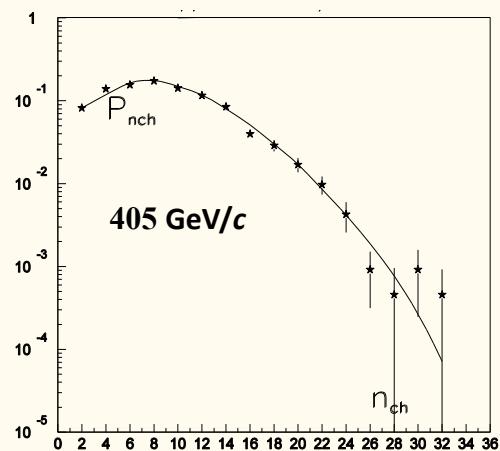
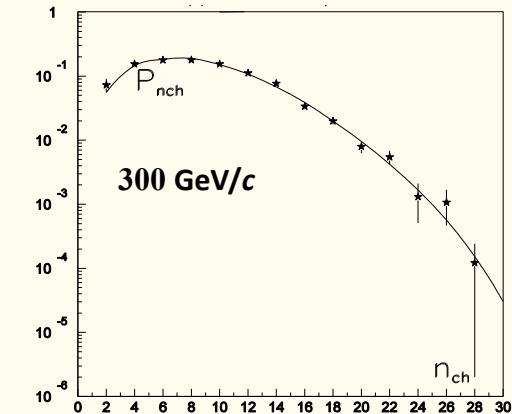
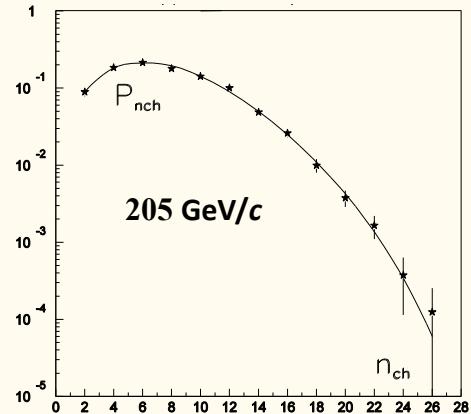
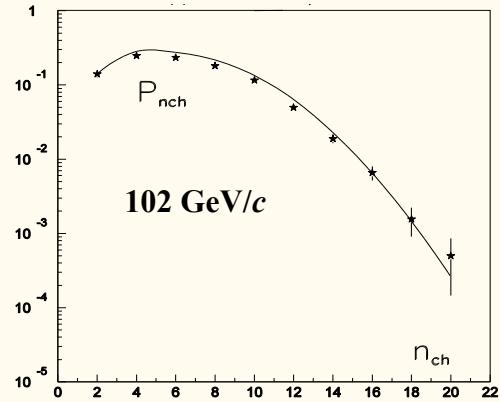
Parameters: $N = 4.24 \pm .13$, $\bar{m} = \bar{m}(s) = 2.48 \pm .20$, $\bar{n}^h = 1.63 \pm .12$, $\chi^2 = 2.0$.

$p_{\Gamma \Theta B/c}$	\bar{m}	M_g	N	\bar{n}_g^h	Ω	χ^2/ndf
102	2.75 ± 0.08	8	3.13 ± 0.56	1.64 ± 0.04	1.92 ± 0.08	2.2/5
205	2.82 ± 0.20	8	4.50 ± 0.10	2.02 ± 0.12	2.00 ± 0.07	2.0/8
300	2.94 ± 0.34	10	4.07 ± 0.86	2.22 ± 0.23	1.97 ± 0.05	9.8/9
405	2.70 ± 0.30	9	4.60 ± 0.24	2.66 ± 0.22	1.98 ± 0.07	16.4/12
800	3.41 ± 2.55	10	20.30 ± 10.40	2.41 ± 1.69	2.01 ± 0.08	10.8/12

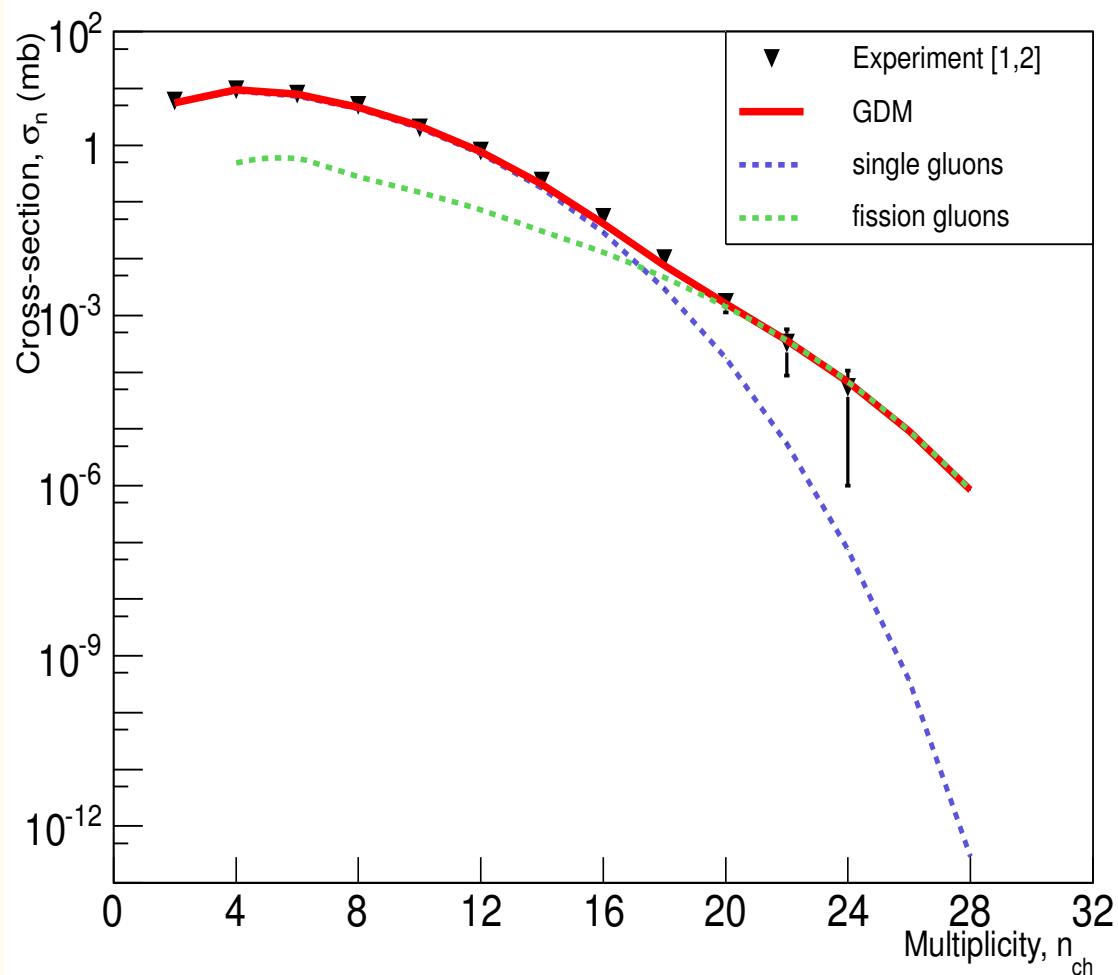
ISR energy 60 GeV: $\bar{n}_g^h \approx 3.3$

Observed growth of \bar{n}_g^h evidences a change of mechanism from fragmentation to recombination

GDM describes MD in pp at 100-800 GeV/c



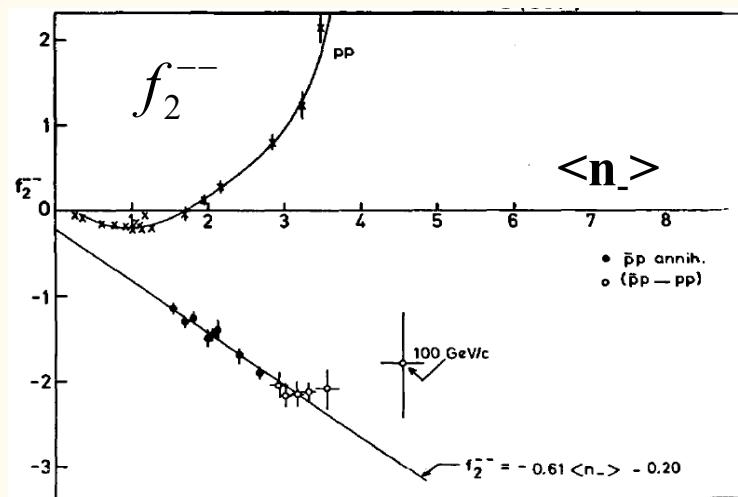
GDM with gluon fission



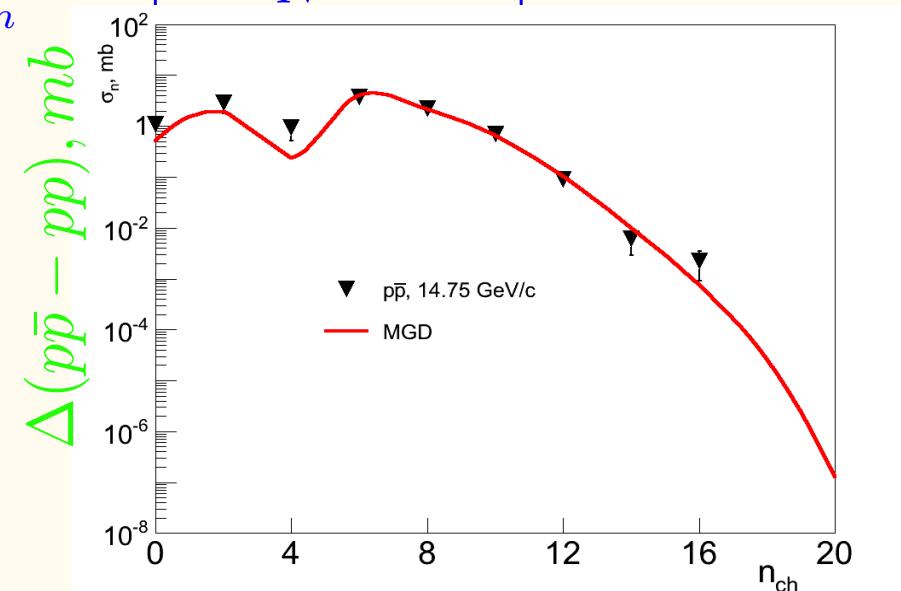
Data [Mirabelle & SVD-2] for pp at 50 GeV/c stitching along topological cross sections, σ_n . At calculation MD, GDM takes into account two types contributions: gluon hadronization without their fission (green line) and with fission (blue line). Superposition of their contributions presents red line. High multiplicity stipulates namely by gluon fission.

Proton-antiproton annihilation in GDM

$$Q(z) = c_0 \sum_m P_m^G \left[1 + \frac{\bar{n}^h}{N} (z-1) \right]^{mN} + c_2 z^2 \sum_m P_m^G \left[1 + \frac{\bar{n}^h}{N} (z-1) \right]^{mN} + \\ + c_4 z^4 \sum_m P_m^G \left[1 + \frac{\bar{n}^h}{N} (z-1) \right]^{mN},$$



J.G. Rushbrooke, B.R. Webber.
Phys.Rep. 44 (1978) 1



$$\Delta(p\bar{p} - pp) = \sigma_n(p\bar{p}) - \sigma_n(pp)$$

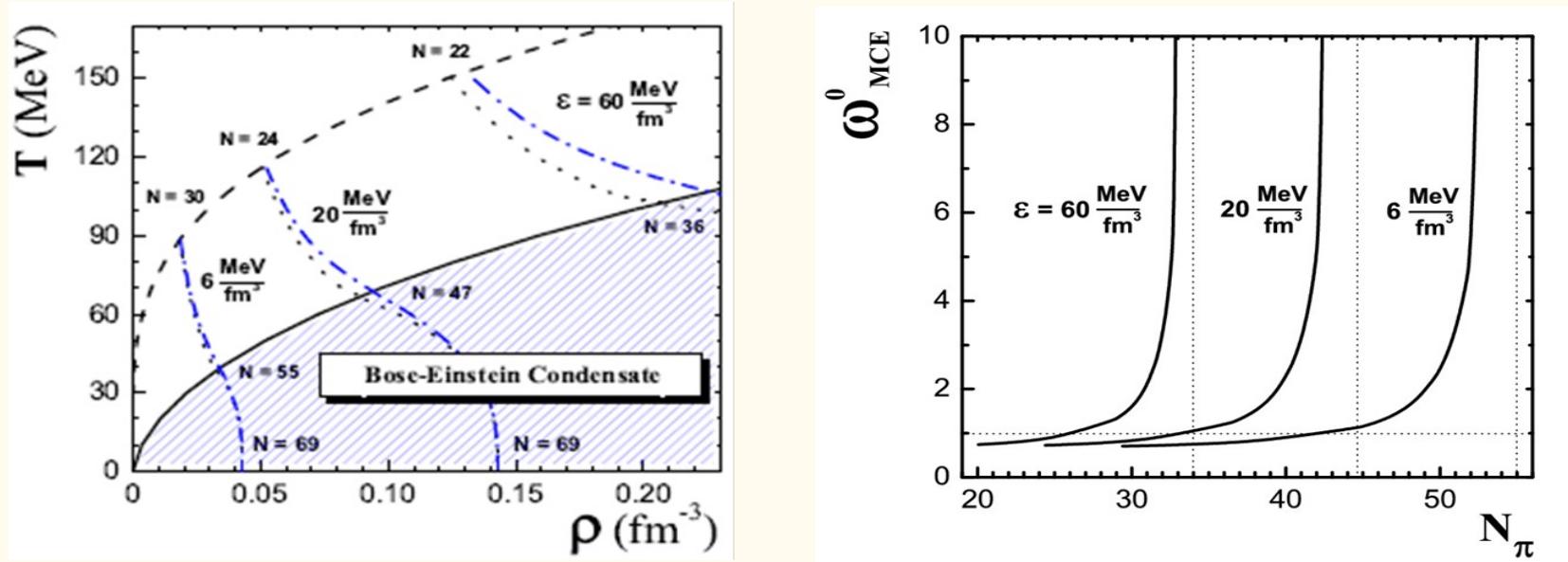
Neutral pion fluctuations at high total multiplicity

V. Begun & M. Gorenstein put us the task on searching for pionic (Bose-Einstein, BEC) condensate [Phys.Lett., 2007, Phys.Rev. 2008] in pp interactions at U-70 for HM. For this purpose, we only had to measure the scaled variance

$$\omega^0 = D / \langle N_0(N_{\text{tot}}) \rangle, \quad D = \langle N_0^2 \rangle - \langle N_0 \rangle^2,$$

of π^0 -meson number with growth of total multiplicity ($n_{\text{tot}} = n_{\text{ch}} + n_0$). Abrupt growth of ω^0 would be signal of BEC formation.

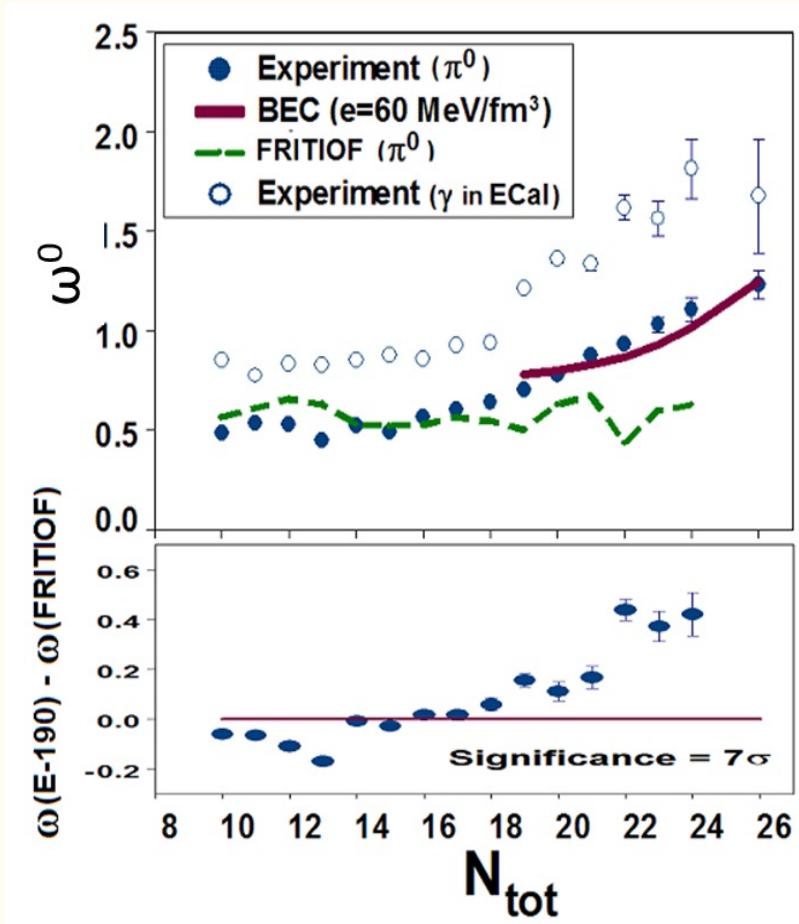
Fluctuations of π^0 -mesons at High multiplicity



Phase diagram of pionic gas at $\mu_Q = 0$. Dash line corresponds to $\rho_\pi(T, \mu_\pi = 0)$, solid - BEC. Energy densities 6, 20 и 60 Mev/fm 3 .

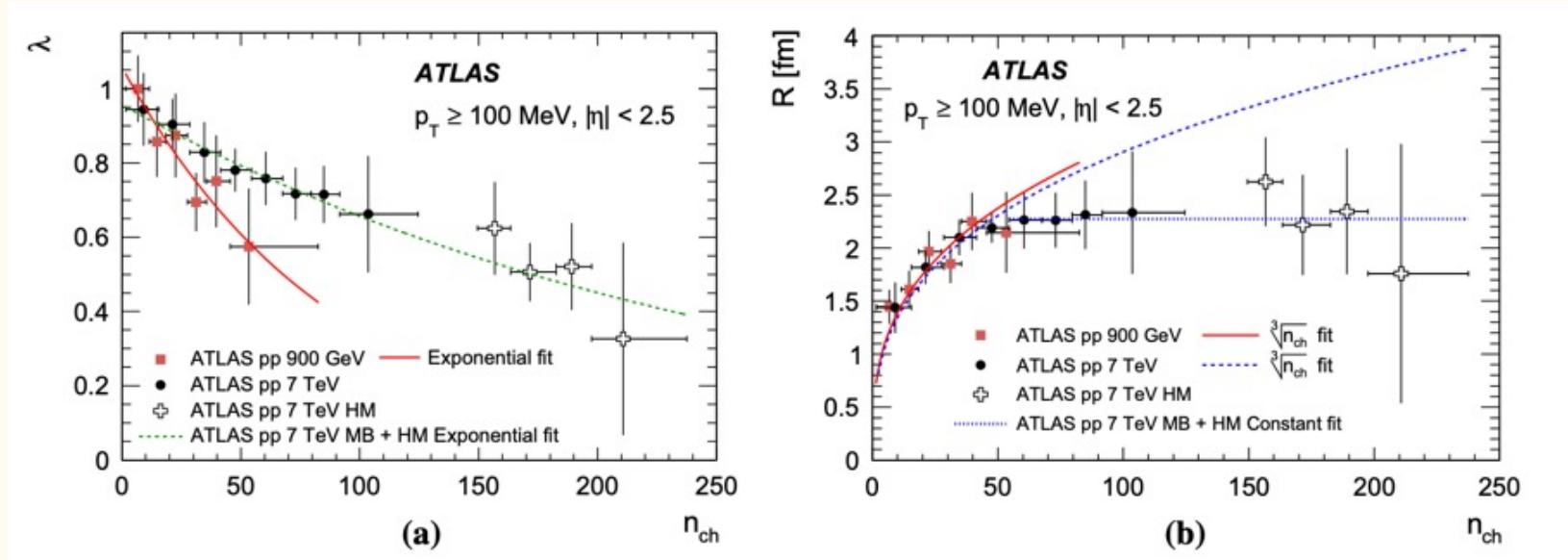
$$\frac{T_C(\pi)}{T_C(A)} \approx \frac{m_A}{m} \left(\frac{r_A}{r_\pi} \right)^2 \cong \frac{m_A}{m} 10^{10} \rightarrow T_C(\pi) \gg T_C(A).$$

Fluctuations of π^0 -mesons at High multiplicity



The deviation of the scaled variance, ω^0 measured on the SVD-2 from the Monte Carlo predictions in the HM region is 7σ at $N_{\text{tot}} \sim 25$ [EPJ, 2012, ICHEP 2012].

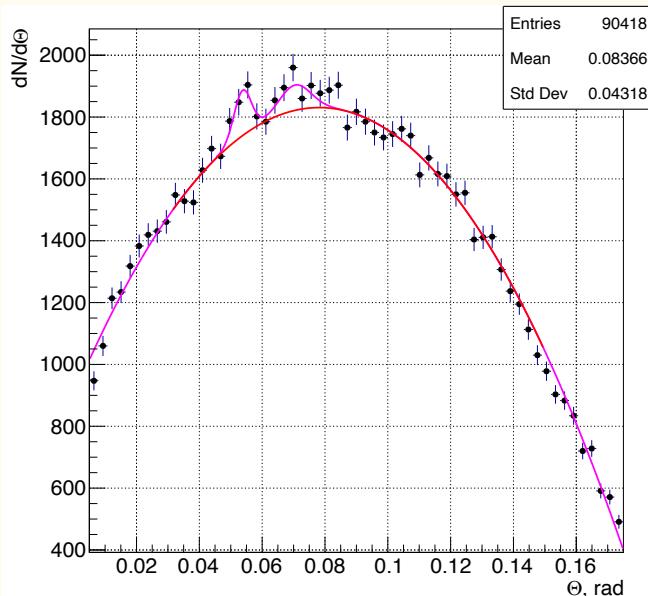
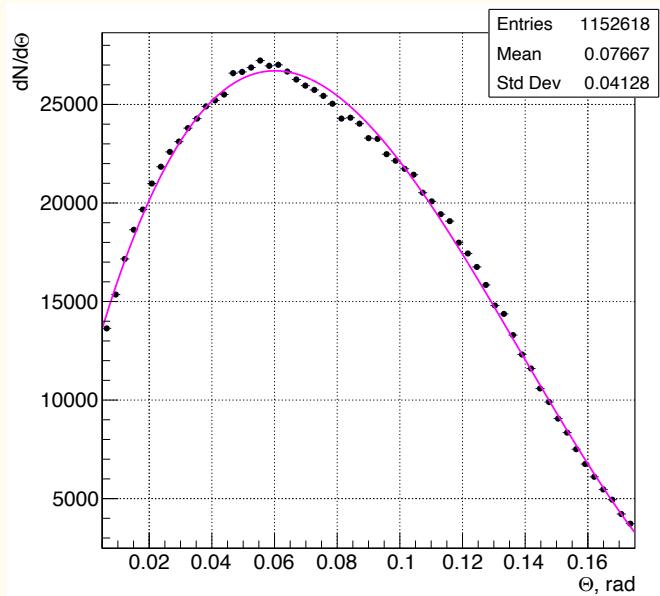
Bose-Einstein correlations in pp at HM (ATLAS)



Two-particle Bose-Einstein correlations in pp at 0.9 и 7 TeV. ATLAS Collab. EPJ, 75 (2015). $C_2(Q) = \rho(Q)/\rho_0(Q) = C_0 [1 + \Omega(\lambda, QR)](1 + \varepsilon Q)$, $Q^2 = (p_1 - p_2)^2$.

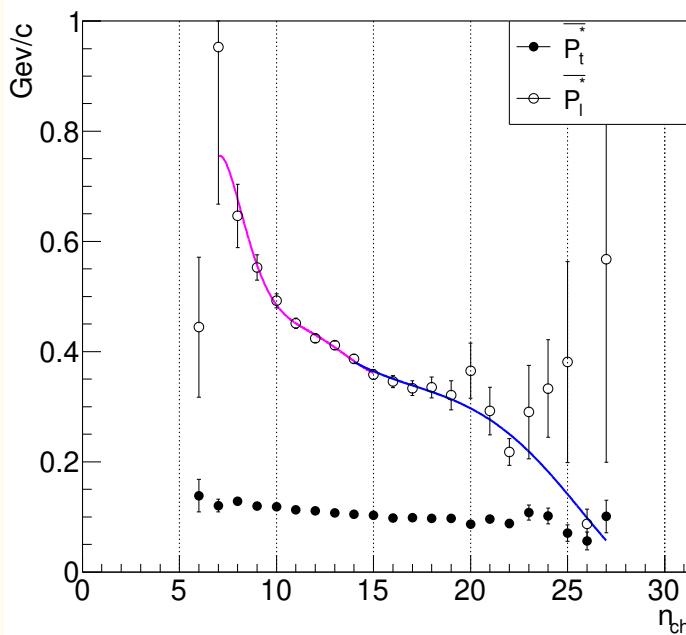
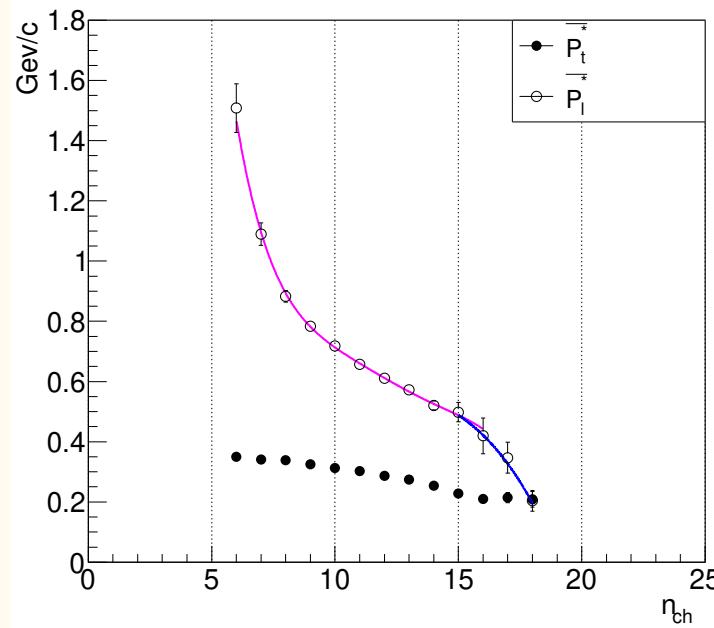
λ close to 1 characterizes chaotic emission of particles, λ close to 0 - coherent emission (is characteristic for BEC), R defines the size of emission region, for BEC it's hard time [hep-ph 1501.04530]. Their wave functions entangle . LHCb confirms 1709.01769 [hep-ex].

Polar angle (θ) distributions. All and HM.



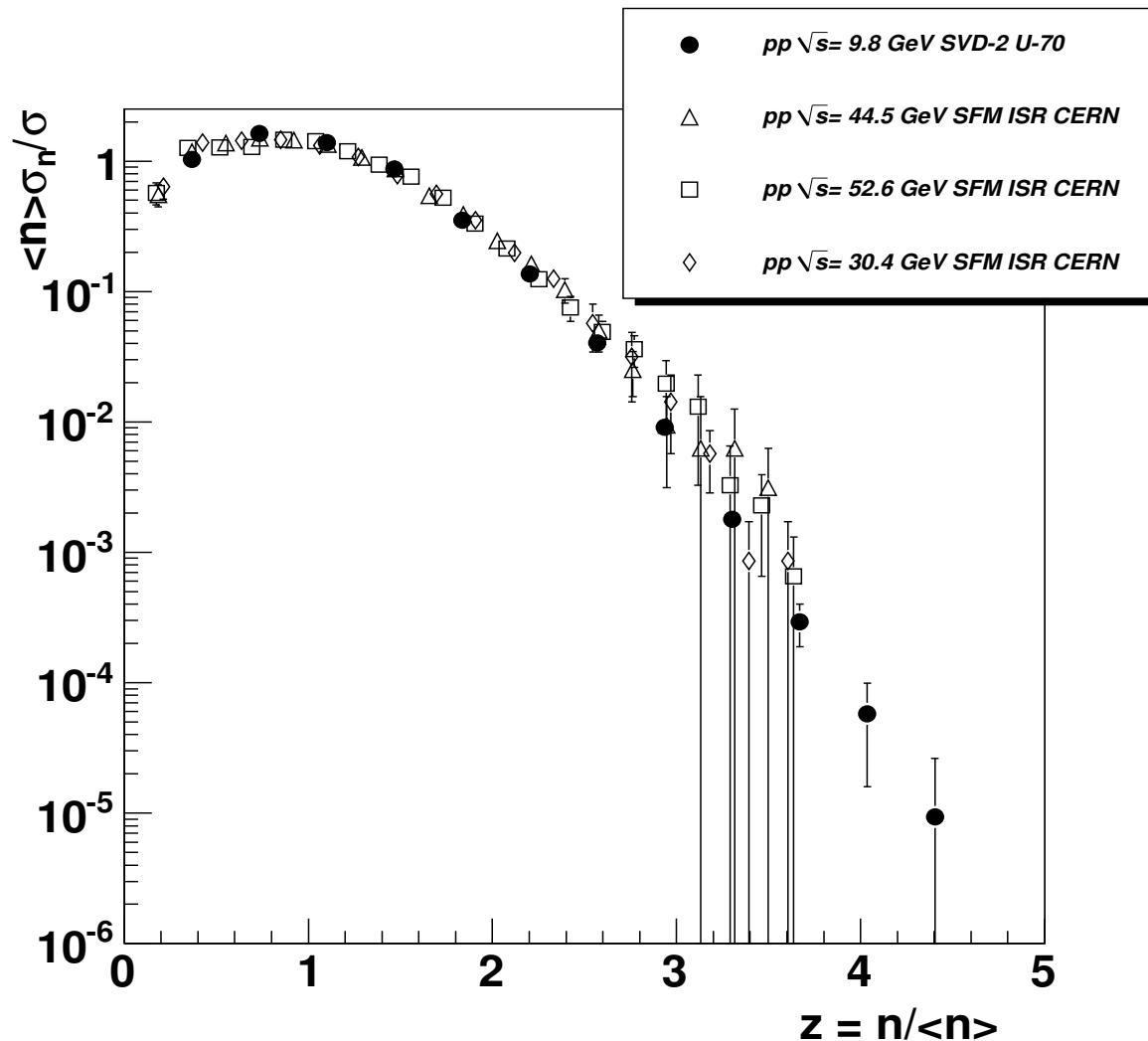
Angle distributions on the polar angle Θ . In HM region we observe two-humped structure, which it's interpreted as Cherenkov radiation gluon by quark. $\Theta_{\text{Cher}} = 0.05377 \pm 0.00273$ rad with $CL3.1\sigma$. For gluon rings $\cos \Theta = 1/\beta n_r$, where n_r refraction coefficient $n_r = 1.0016 \pm 0.0001(4)$, close to 1. It testifies about rarity of qg-medium.

Продольная и поперечная компонента при БМ



Средние значения поперечной и продольной компонент импульса заряженных частиц. Слева: М.К.-моделирование, справа: экспериментальные данные.

В области $n_{ch} \sim 16$ наблюдается образование БЭК. При $n_{tot} > 18$ растет ω^0 и исчезают лидирующие частицы, система становится изотропной по всем направлениям.



KNO scaling

The world KNO distribution
with addition of the SVD-2
points ($\sqrt{s} = 9.8 \text{ GeV}$).

RESULTS

The study of events with HM allow us to develop and supplement the mechanism of multiple production with a description of the hadronization stage by the Bernoulli distribution for various processes: e^+e^- annihilation, bottomonium decays, $p\bar{p}$ interactions, and proton and antiproton annihilation.

In the HM region, the collective behaviour of secondary particles has been discovered and confirmed, which gives us new ideas about the mechanism of multiple production, in particular, the active role of the gluon component in this process.

Afterwords

"Perhaps there are no discoveries in elementary or higher mathematics, or even, perhaps, in any other field that could be made ... without analogy."

George Poiya.

DNA Replication ~ gluons \rightarrow hadrons

