

Improved Muon-Capture Calculations in Light and Heavy Nuclei

Lotta Jokiniemi

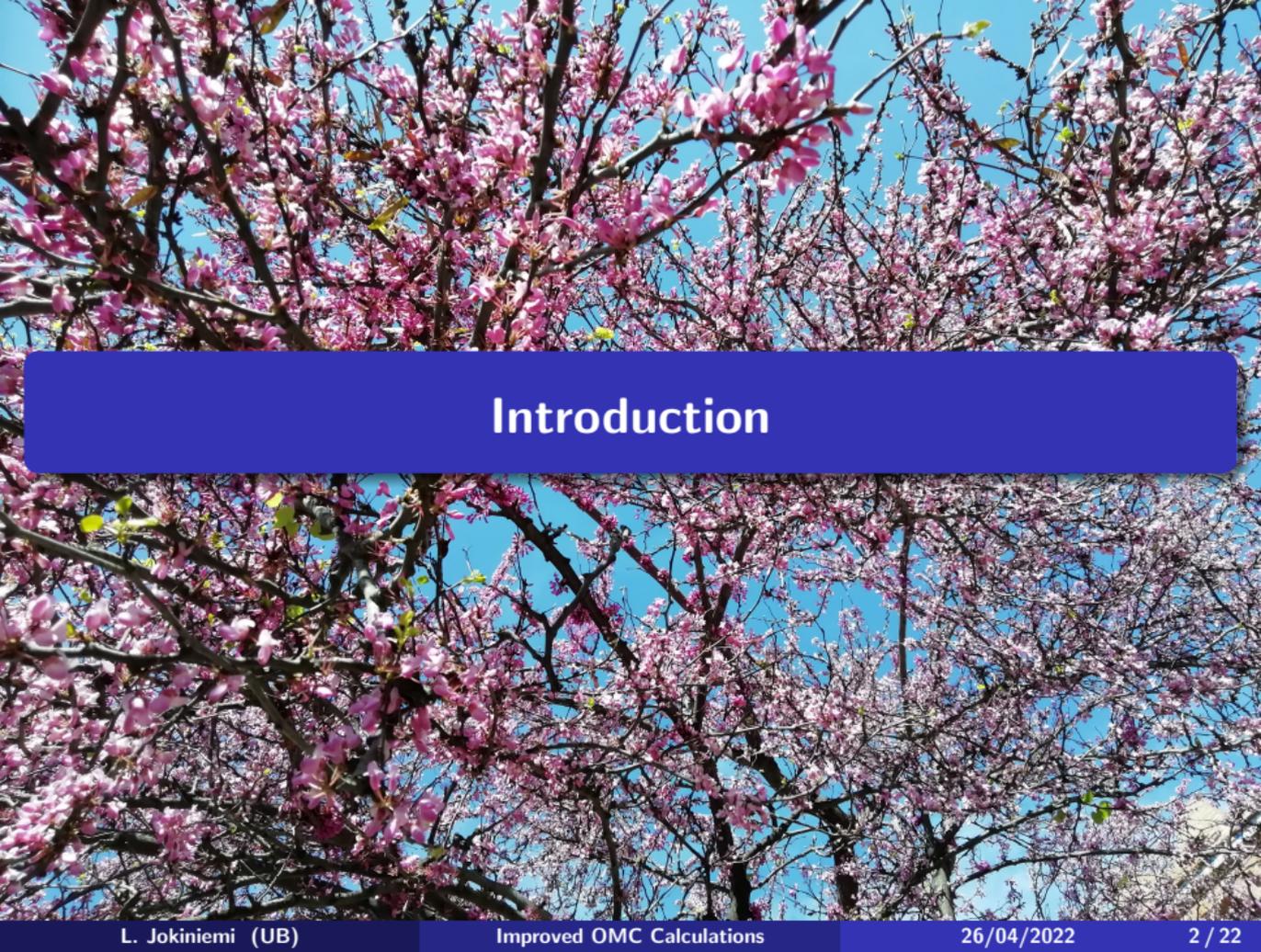
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University of Barcelona, Spain

OMC4DBD Collaboration Meeting, 26/04/2022



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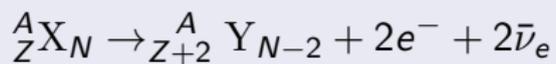
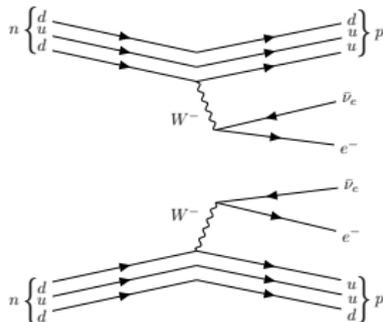
Introduction

Motivation

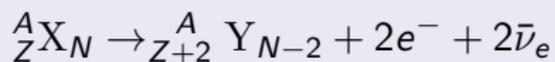
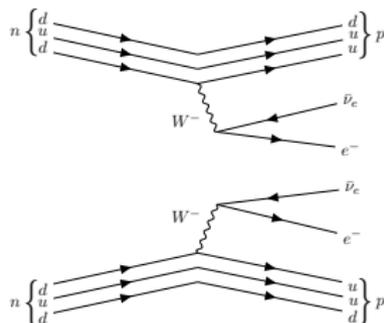
- Current knowledge on particles and interactions between them is based on the Standard Model (SM) of particle physics
- According to the SM, **neutrinos** are extremely weakly interacting, **massless** fermions
- Yet we know from solar neutrino experiments that neutrinos must have a **non-zero mass**
 - But what is the absolute mass scale?
 - What else is there beyond the SM?
- *Observing neutrinoless double-beta decay would provide answers!*



Two-Neutrino Double-Beta ($2\nu\beta\beta$) Decay

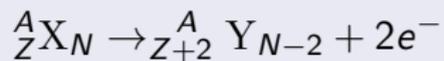
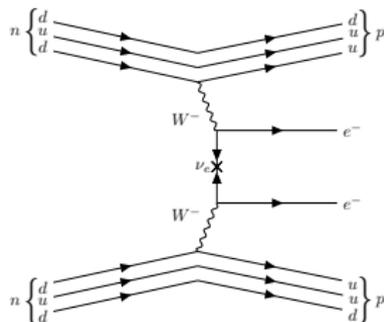


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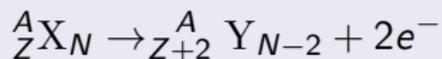
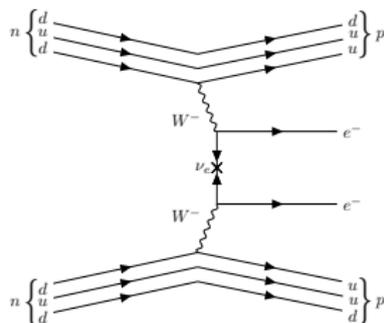


- May happen, when β -decay is not energetically allowed
- Allowed by the Standard Model
- Measured in ≈ 10 isotopes
 - Half-lives of the order 10^{20} years or longer

Neutrinoless Double-Beta ($0\nu\beta\beta$) Decay



Neutrinoless Double-Beta ($0\nu\beta\beta$) Decay



- Requires that the neutrino is a Majorana particle
- Violates the lepton-number conservation law by two
- $\frac{1}{t_{1/2}^{(0\nu)}} \propto |\langle m_\nu \rangle|^2$

Half-life of $0\nu\beta\beta$ Decay

$$\frac{1}{t_{1/2}^{0\nu}} = g_A^4 G_{0\nu} |M^{0\nu}|^2 \left(\frac{m_{\beta\beta}}{m_e} \right)^2$$

¹Agostini *et al.*, arXiv:2202.01787 (2022)

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New physics

- Axial-vector coupling
($g_A^{\text{free}} \approx 1.27$)

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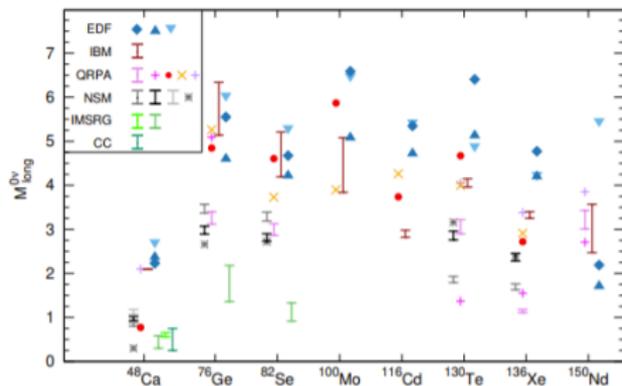
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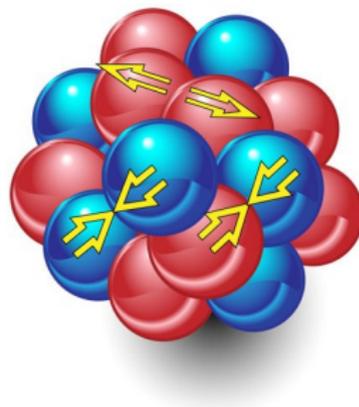


Matrix elements of $0\nu\beta\beta$ decays¹

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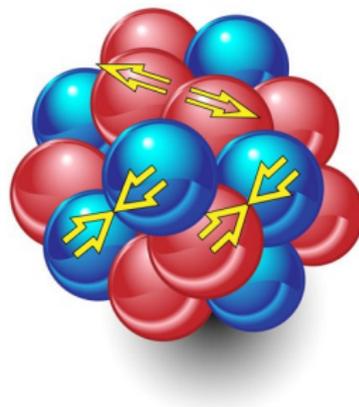
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- Ab initio methods (CC, **IMSRG**, ...)



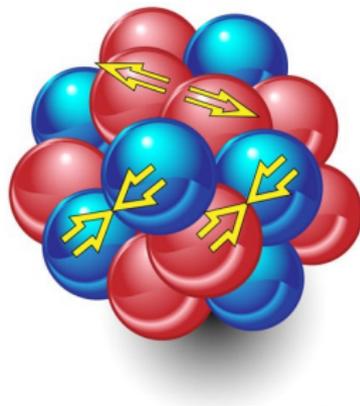
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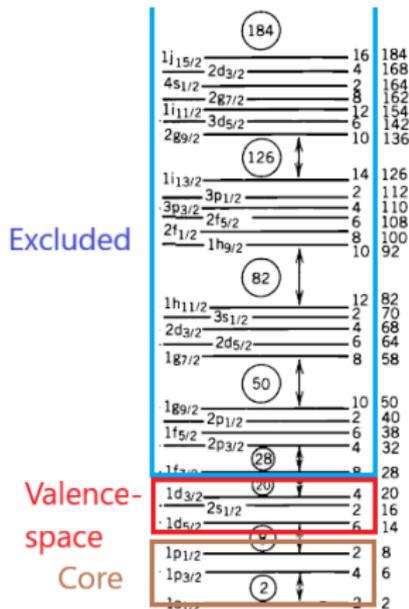
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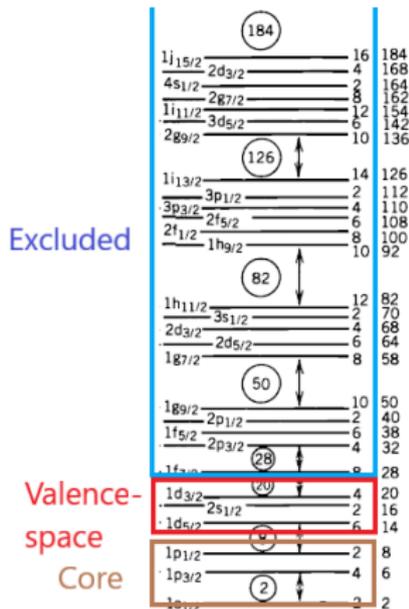
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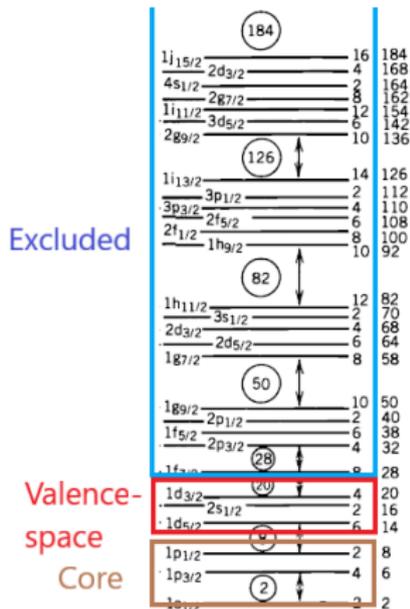
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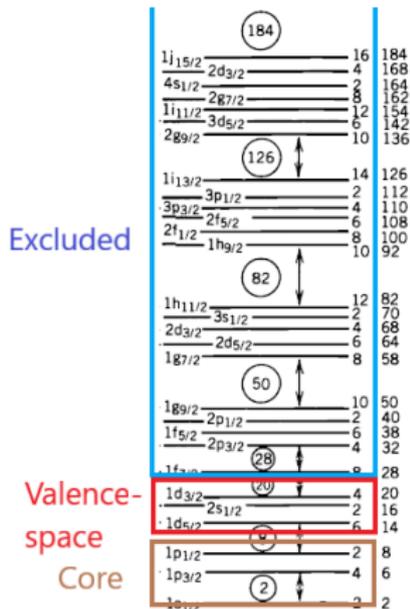
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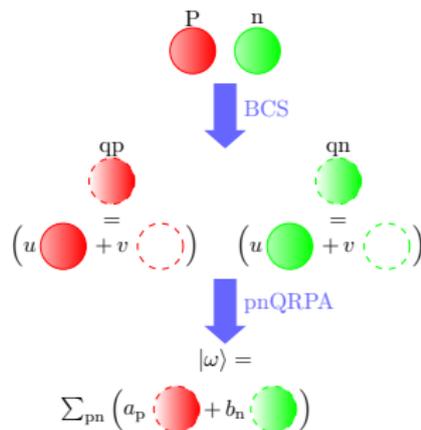
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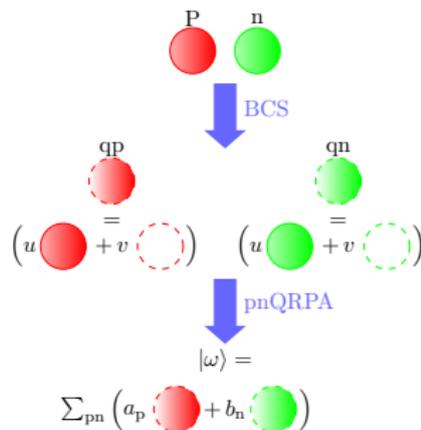
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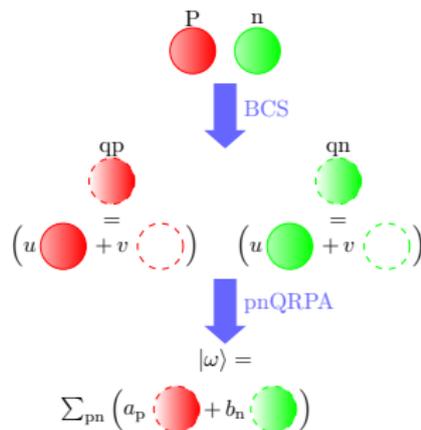
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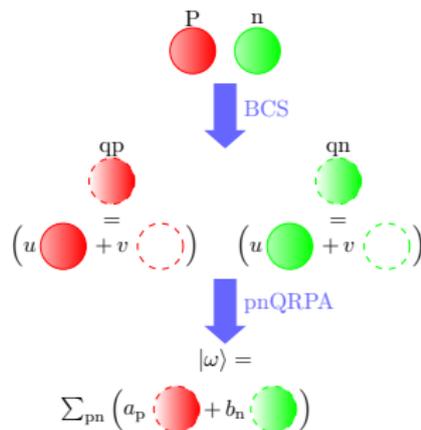
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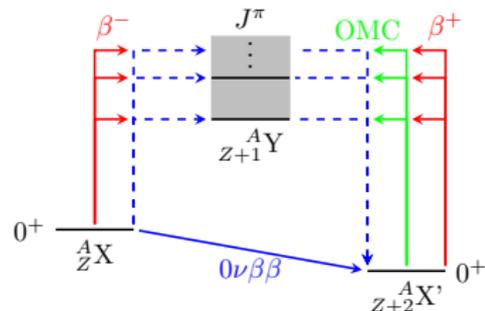
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Ordinary Muon Capture as a Probe of $0\nu\beta\beta$
Decay

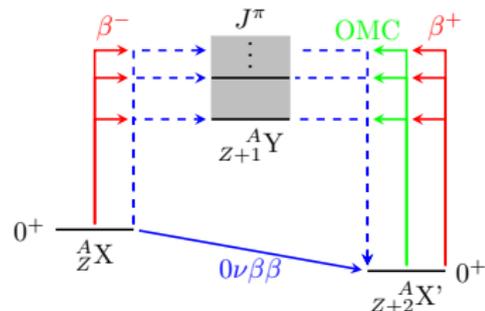
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- OMC leads to transitions to all J^π states up to high energies



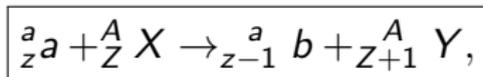
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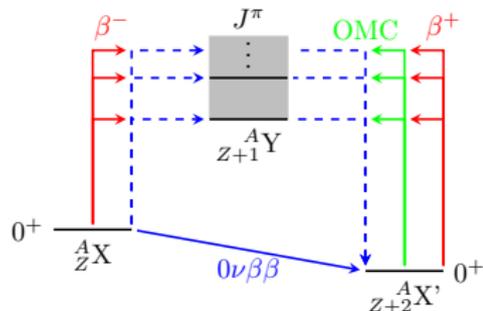


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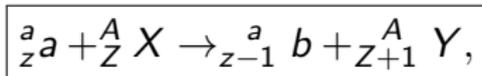


where (a, b) can be (p, n) , $({}^3\text{He}, t)$, ...



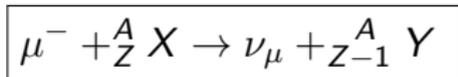
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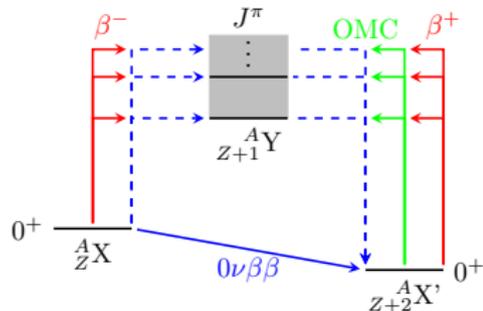


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- Ordinary muon capture (OMC)



can probe the right-hand side



Advantages of OMC as a Probe of $0\nu\beta\beta$ Decay

- Both OMC and $0\nu\beta\beta$ decay involve couplings g_A and g_P :

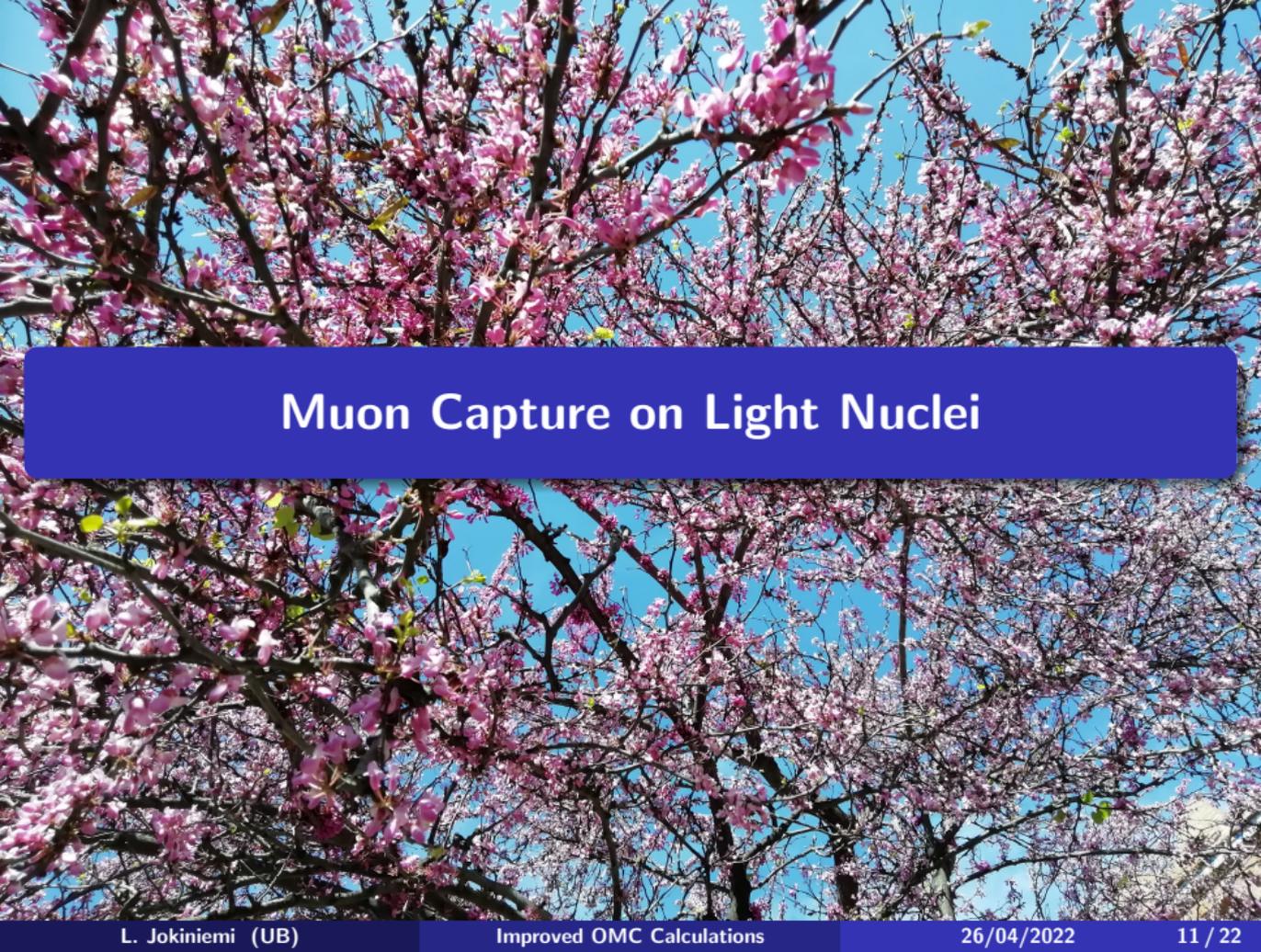
$$W^{(OMC)} \propto |g_A M_A + g_V M_V + g_P M_P|^2$$

$$M^{0\nu} = M_{GT}^{0\nu}(g_A, g_P, g_M) - \left(\frac{g_V}{g_A}\right)^2 M_F^{0\nu}(g_V) + M_T^{0\nu}(g_A, g_P, g_M),$$

$$[t_{1/2}^{0\nu}]^{-1} = g_A^4 G_{0\nu} |M^{0\nu}|^2 \left(\frac{m_{\beta\beta}}{m_e}\right)^2$$

- ...so if
 - we know the involved nuclear structure precisely enough, and
 - OMC rates to individual nuclear states can be measured

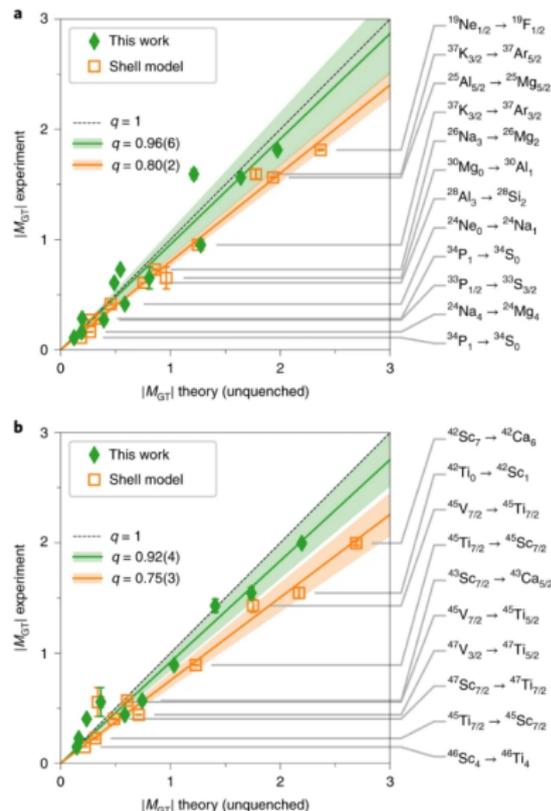
...we can probe g_A and g_P on the relevant momentum-exchange regime for $0\nu\beta\beta$ decay



Muon Capture on Light Nuclei

g_A Quenching at High Momentum Exchange?

- Recently, first *ab initio* solution to g_A quenching puzzle was proposed for β -decay²
- How about g_A quenching at high momentum transfer $q \approx 100$ MeV/c?
 - OMC could provide a hint!



²P. Gysbers *et al.*, *Nature Phys.* **15**, 428 (2019)

Ingredients 1: Particle Physics

Muon-Capture Theory

- Interaction Hamiltonian \rightarrow capture rate:

$$W(J_i \rightarrow J_f) = \frac{2J_f + 1}{2J_i + 1} \left(1 - \frac{q}{m_\mu + AM} \right) q^2 \sum_{\kappa U} |g_V M_V + g_A M_A + g_P M_P|^2$$

PHYSICAL REVIEW

VOLUME 118, NUMBER 2

APRIL 15, 1960

Theory of Allowed and Forbidden Transitions in Muon Capture Reactions*

MASATO MORITA

Columbia University, New York, New York

AND

AKIHIKO FUJII†

Brookhaven National Laboratory, Upton, Long Island, New York

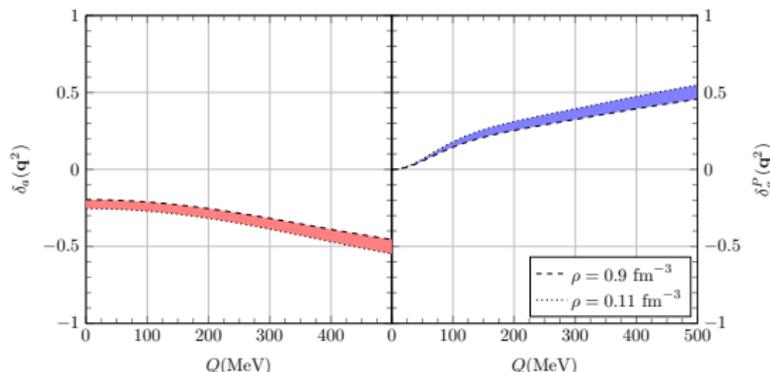
(Received November 9, 1959)

Ingredients 1: Particle Physics

NEW!! Hadronic Two-Body Currents

- Correct the couplings by effective one-body currents ³

$$g_A \rightarrow (1 + \delta_a(\mathbf{q}^2))g_A \text{ and } g_P \rightarrow \left(1 - \frac{q^2 + m_\pi^2}{q^2} \delta_a^P(\mathbf{q}^2)\right) g_P$$



³Hoferichter *et al.*, *Phys. Rev. C* **102**, 074018 (2020)

Ingredients 3: Atomic Physics

Bound-Muon Wave Function

- Solve the Dirac equations for the muon:

$$\begin{cases} \frac{d}{dr} G_{-1} + \frac{1}{r} G_{-1} = \frac{1}{\hbar c} (mc^2 - E + V(r)) F_{-1} \\ \frac{d}{dr} F_{-1} - \frac{1}{r} F_{-1} = \frac{1}{\hbar c} (mc^2 + E - V(r)) G_{-1} \end{cases}$$

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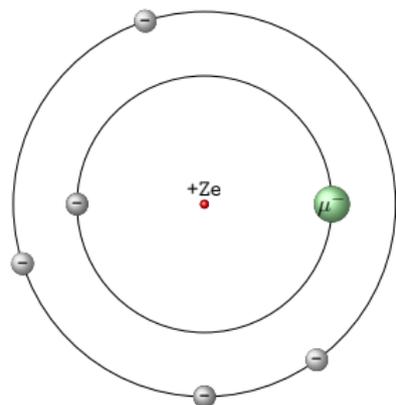
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- For a **point-like** or **finite-size** nucleus



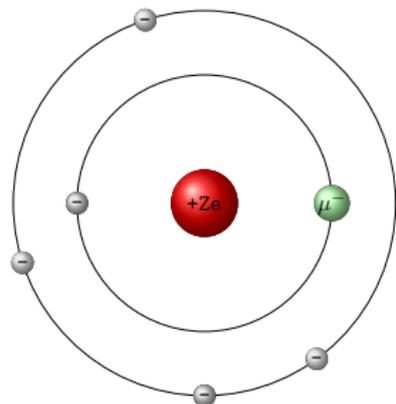
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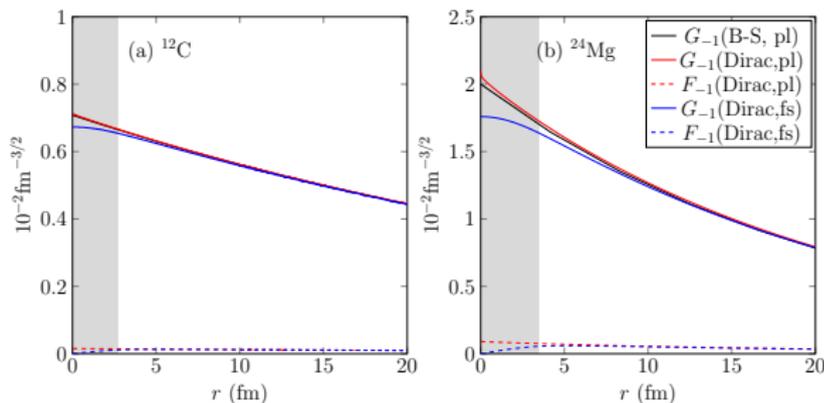
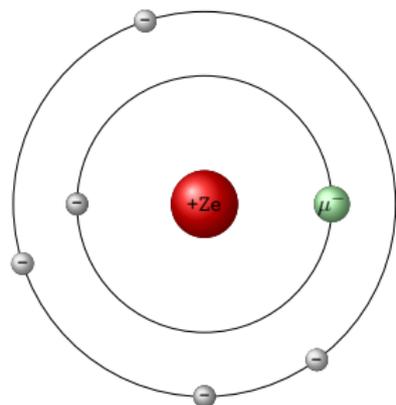
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Test case: Capture Rates on Low-Lying States in ^{12}C

Nuclear Shell Model + Two-Body Currents + Realistic Muon Wave Function

- Nuclear shell-model calculation in p -shell with chiral two-body currents and realistic bound-muon wave functions
- Quite good agreement with experiment

J_i^π	E_{exp} (MeV)	Rate (10^3 1/s)			
		Exp.		NSM	
		Measday ⁴	Double Chooz ⁵	1bc	1bc+2bc
1_{gs}^+	0	6.04 ± 0.35	$5.68^{+0.14}_{-0.23}$	6.48	4.56–4.86
2_1^+	0.953	0.21 ± 0.1	$0.31^{+0.09}_{-0.07}$	0.42	0.32–0.34
2_2^+	3.759	-	$0.026^{+0.015}_{-0.011}$	0.011	0.009–0.009

⁴Measday, *Phys. Rep.* **354**,243 (2001)

⁵Abe *et al.*, *Phys. Rev. C* **93**,054608 (2016)

Results: Capture Rates on Low-Lying States in ^{24}Na

J_i^π	E_{exp} (MeV)	Rate (10^3 1/s)				
		Exp. ⁶	NSM		IMSRG	
			1bc	1bc+2bc	1bc	1bc+2bc
1_1^+	0.472	(21.0 ± 6.6)	4.0	3.0	22.3	15.2
1_2^+	1.347	17.5 ± 2.3	32.7	21.3	7.7	4.9
Sum(1^+)		38.5 ± 8.9	36.7	24.5	30.0	20.0
2_1^+	0.563	17.5 ± 2.1	1.0	0.7	0.5	0.3
2_2^+	1.341	3.4 ± 0.5	3.1	2.5	1.0	0.9
Sum(2^+)		20.9 ± 2.6	4.1	3.2	1.5	1.2

[LJ, T. Miyagi, S.R. Stroberg, J.D. Holt, J. Kotila and J. Suhonen, arXiv:2111.12992]

⁶P. Gorringer *et al.*, Phys. Rev. C **60**, 055501 (1999)

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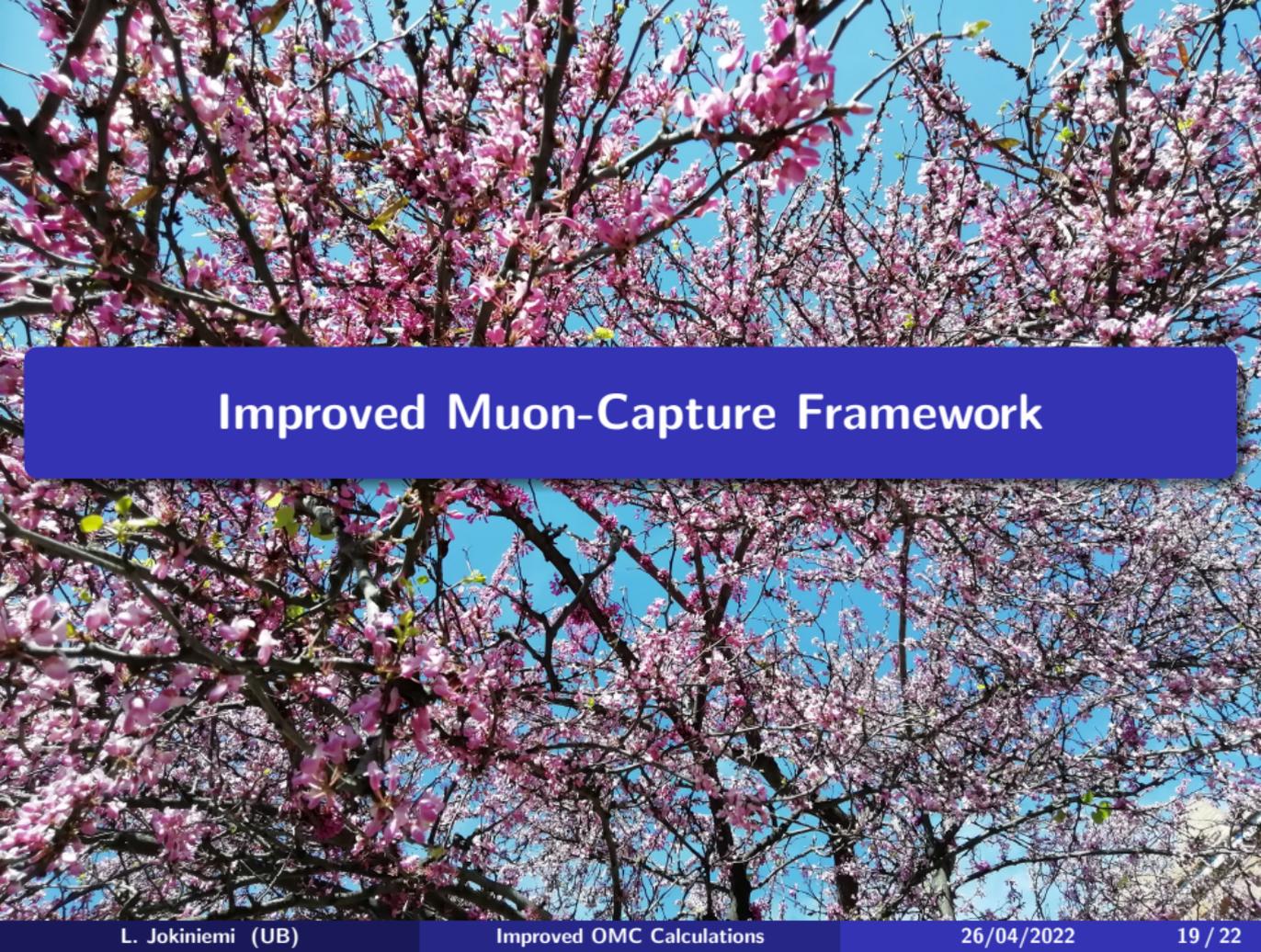
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- Generally, IMSRG gives smaller capture rates
- 1^+ states mixed
- Agreement with experiment could be better

⁶P. Gorringer *et al.*, Phys. Rev. C **60**, 055501 (1999)



Improved Muon-Capture Framework

Improvements to Morita-Fujii Formalism

- Morita-Fujii⁷ (MF) and Walecka⁸ formalisms for OMC combined with pnQRPA tend to give different capture rates

$$g_A^{\text{eff}}(\text{Jok.})^9 \approx 0.6$$

$$g_A^{\text{eff}}(\check{\text{Sim.}})^{10} \approx 1.27$$

$$g_A^{\text{eff}}(\text{Cic.})^{11} \approx 1.0$$

- Not straightforward to compare the two formalisms
- Something has to be done (ongoing work w/ E. Ydrefors and J. Suhonen)
 - 1 Check the assumptions made in the MF formalism
 - 2 Introduce 'Walecka-like' multipole operators into MF formalism → compare (first test case: ¹⁰⁰Mo)

⁷Morita and Fujii, *Phys. Rev.* **118**, 606 (1960)

⁸Walecka, *Muon Physics II* p.113 (Academic Press, New York) (1975)

⁹LJ, Suhonen, *Phys. Rev. C* **100**, 014619 (2019)

¹⁰ $\check{\text{S}}$ imkovic *et al.*, *Phys. Rev. C* **102**, 034301 (2020)

¹¹Ciccarelli *et al.*, *Phys. Rev. C* **102**, 034306 (2020)

- By studying OMC we can shed light on the unknown effective value of g_A
- In order to probe the effective value of g_P we would need to have data on capture rates to individual states
- First *ab initio* muon-capture studies performed
- Next: improved muon-capture theory

A photograph of a tree in full bloom with numerous small, light pink flowers. A green parrot with a red face is perched on a dark branch in the lower-left quadrant. The background shows a clear blue sky and a yellow building partially obscured by the branches.

Thank you!