

# Performance study of the anisotropic flow measurements of identified charged hadrons with fixed-target mode of the MPD experiment at NICA

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for the MPD Collaboration

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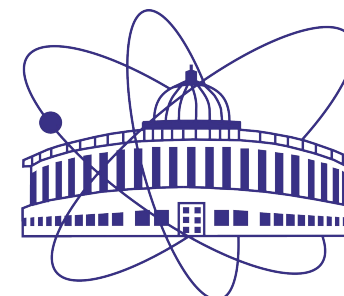
<sup>3</sup>BLTP JINR

XXV International Baldin Seminar on High Energy Physics Problems “Relativistic Nuclear Physics and Quantum Chromodynamics”

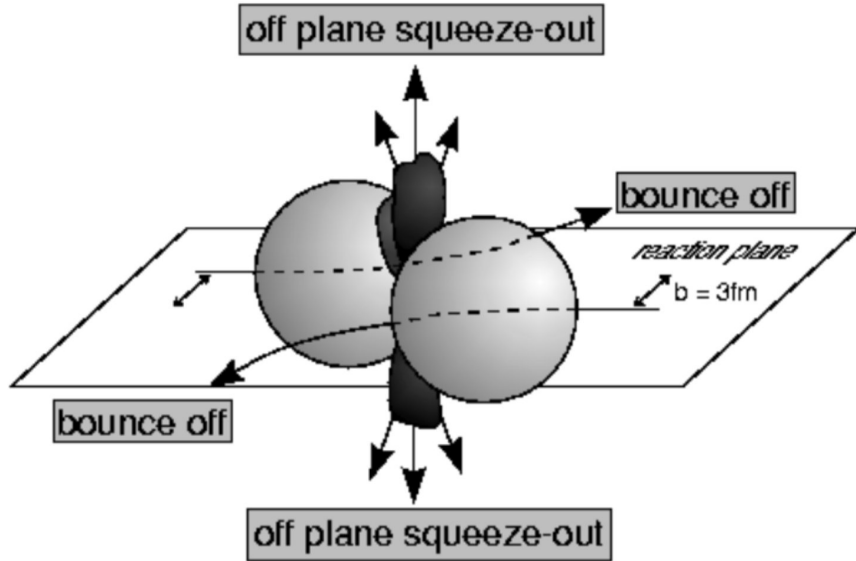
18-23 September 2023, Dubna, Russia

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# Sensitivity of the collective flow to the EOS



Azimuthal distribution of produced particles with respect to RP:

$$\rho(\varphi - \Psi_{RP}) = \frac{1}{2\pi} \left( 1 + 2 \sum_{n=1}^{\infty} v_n \cos n(\varphi - \Psi_{RP}) \right)$$

Coefficients of the decomposition are referred to as collective flow

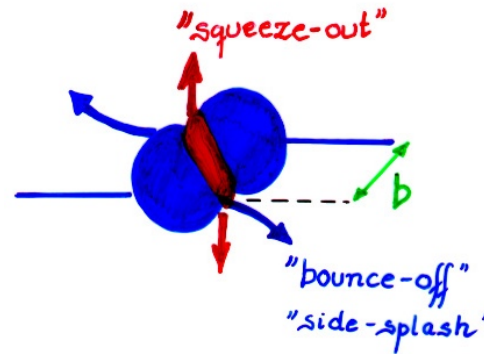
$$v_n = \langle \cos [n(\varphi - \Psi_{RP})] \rangle$$

$v_1$  is called directed and  $v_2$  is called elliptic flow

Collective flow is sensitive to:

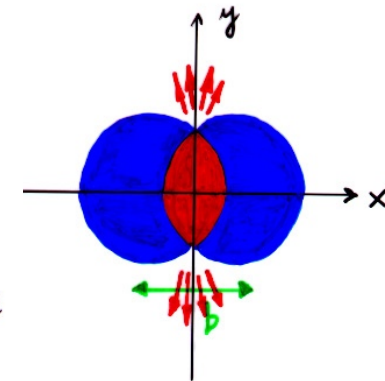
- Compressibility of the created in the collision matter ( $t_{exp} = R/c_s$ ,  $c_s = c \sqrt{dp/d\varepsilon}$ )
- Time of the interaction between the matter within the overlap region and spectators ( $t_{pass} = 2R/\gamma_{CM}\beta_{CM}$ )

Bounce-off



transverse directed flow

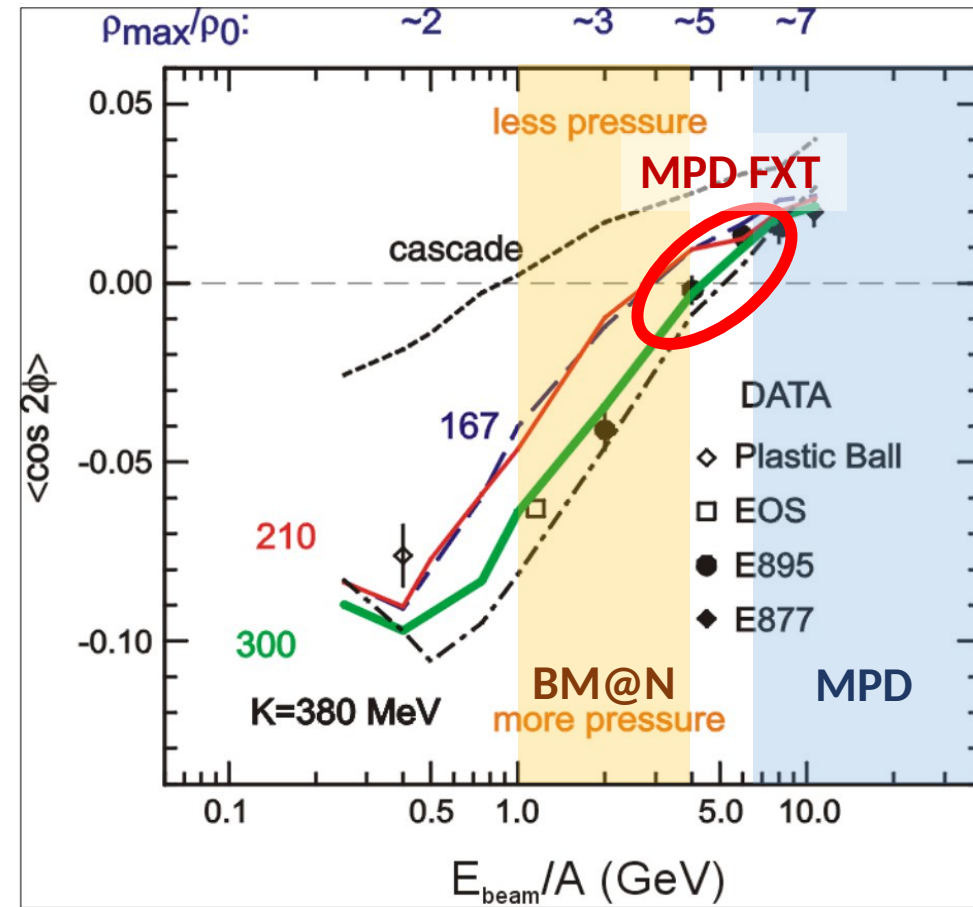
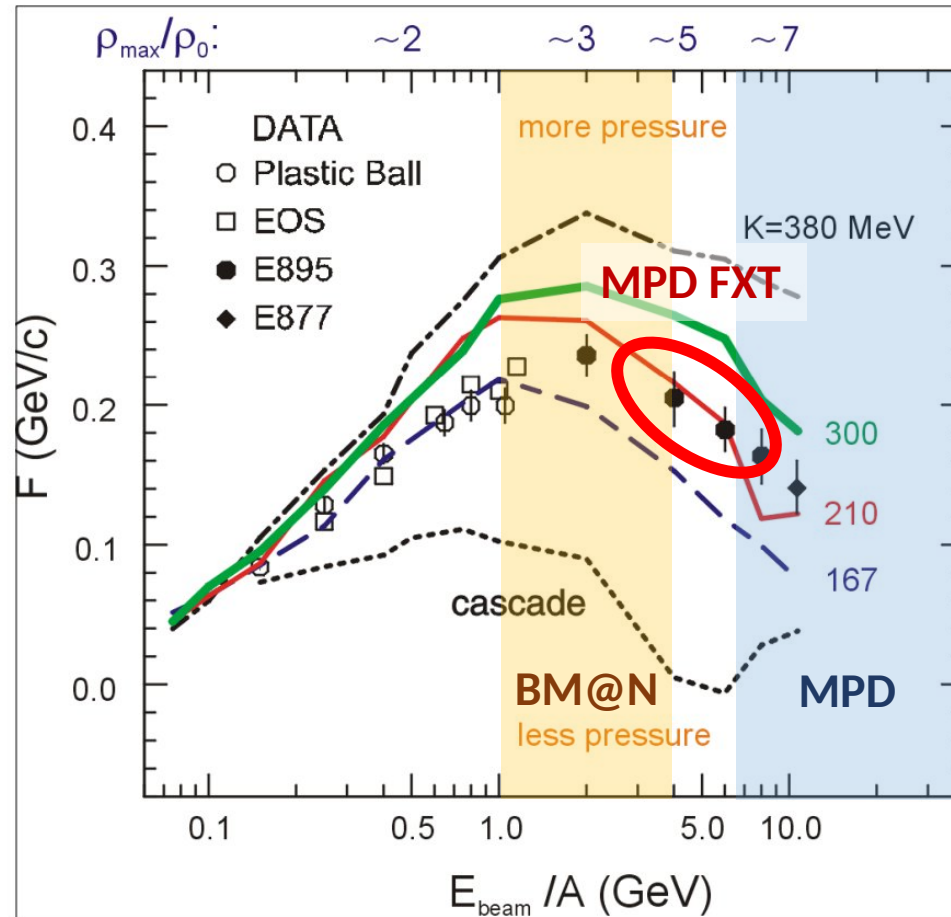
Squeeze-out



out-of-plane elliptic flow

# Interpretation of the previous flow data

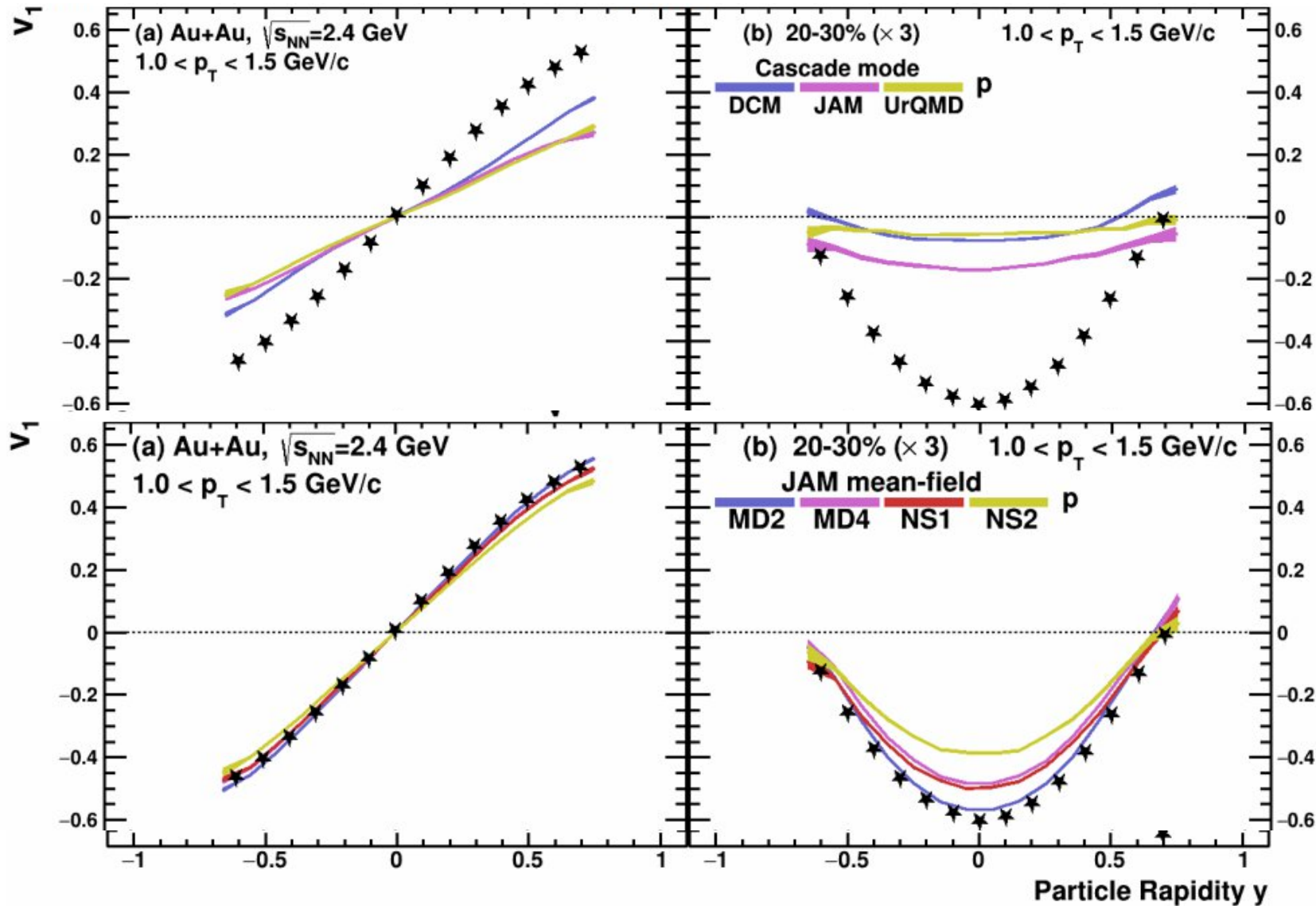
P. DANIELEWICZ, R. LACEY, W. LYNCH  
[10.1126/science.1078070](https://arxiv.org/abs/10.1126/science.1078070)



- The flow data from E895 experiment have ambiguous interpretation:  $v_1$  suggests soft EOS while  $v_2$  corresponds to hard EOS
- Additional measurements are essential to clarify the previous measurements

# Selecting the model

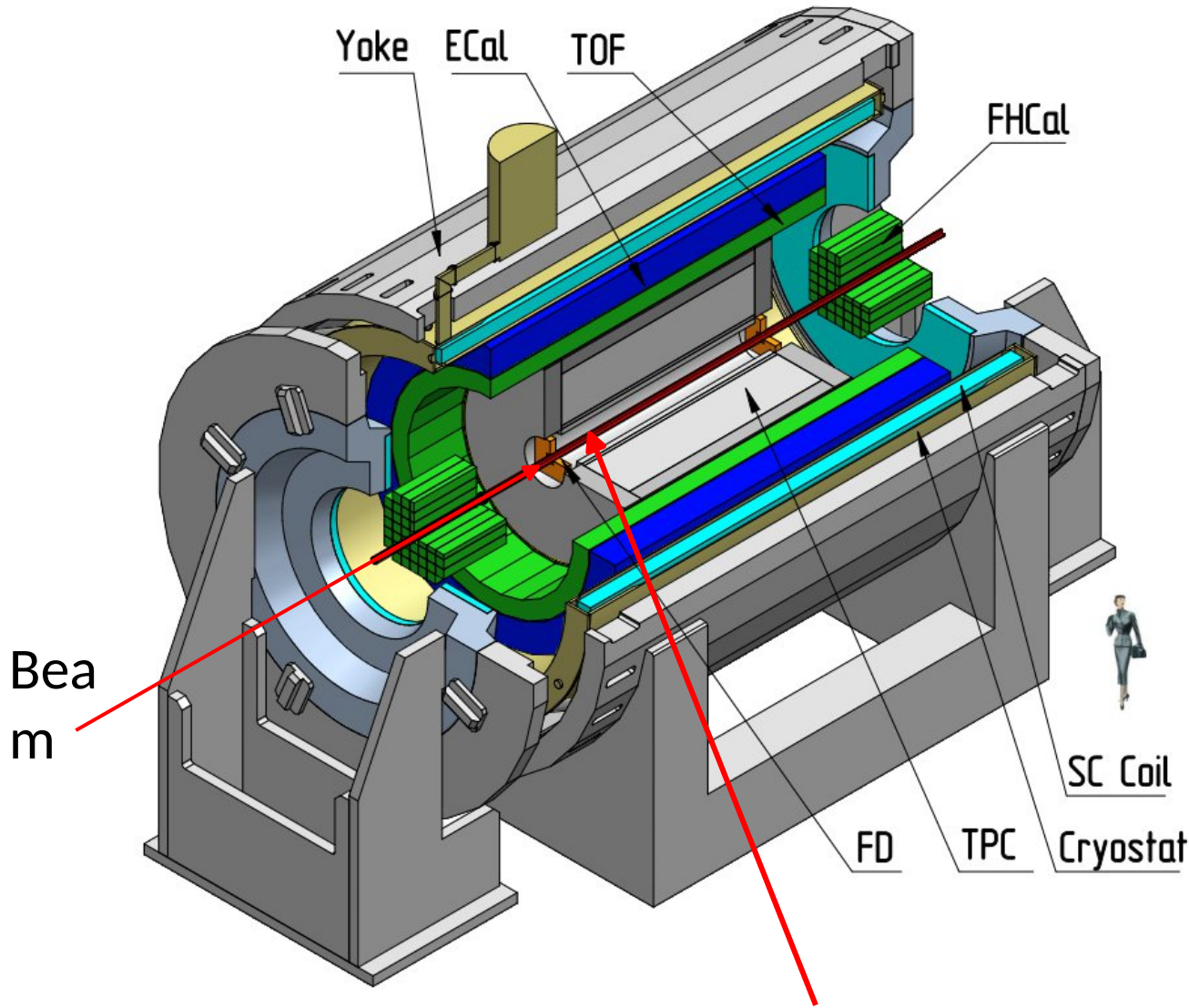
P.Parfenov Particles 5 (2022) 4, 561-579



Cascade models fail to reproduce  $v_n$  at low-energy heavy-ion collision

Mean field models reproduce the  $v_n$  rather well

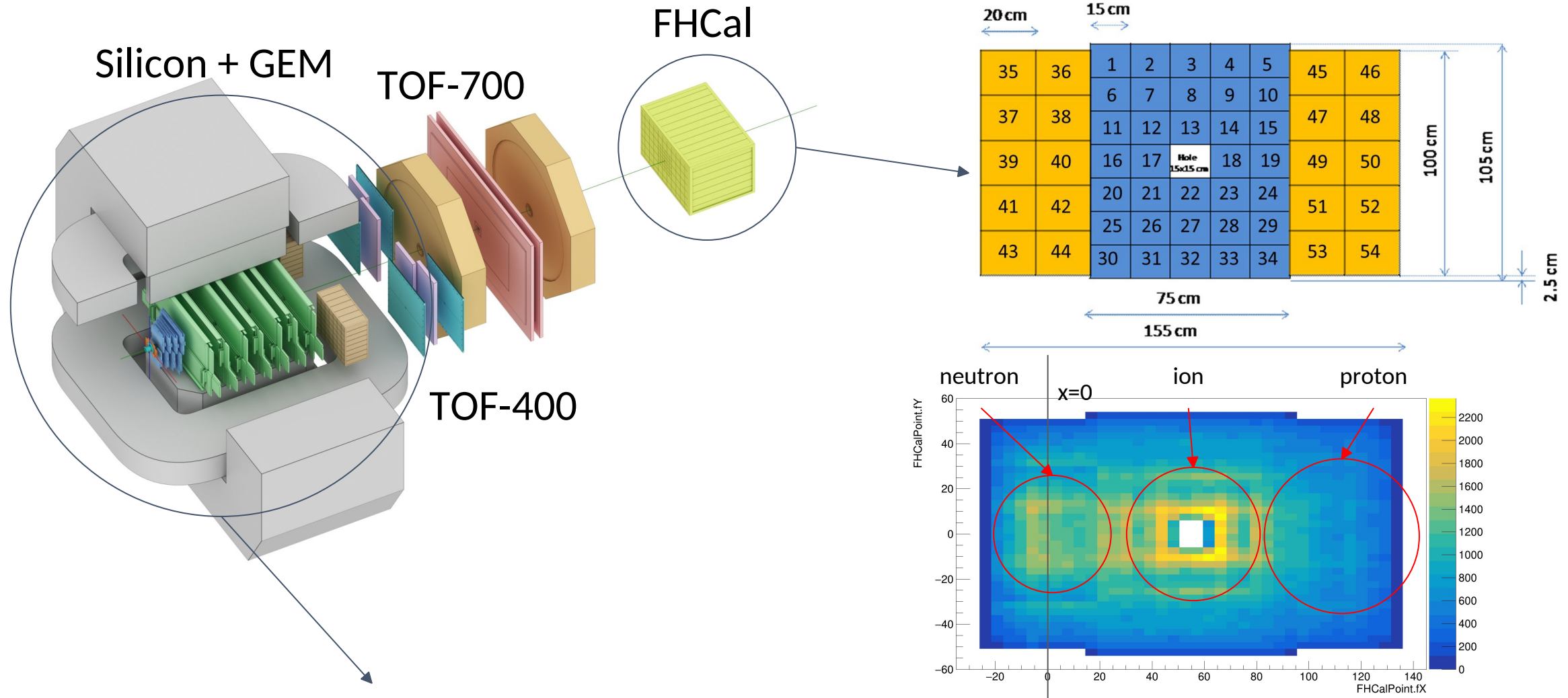
# MPD in Fixed-Target Mode (FXT)



Target ( $z=-115$  cm)

- Model used: UrQMD mean-field
  - Bi+Bi,  $E_{\text{kin}}=1.45$  AGeV
- Point-like target
- GEANT4 transport
- Particle species selection via true-PDG code of the associated MC particle

# The BM@N experiment (GEANT4 simulation for RUN8)



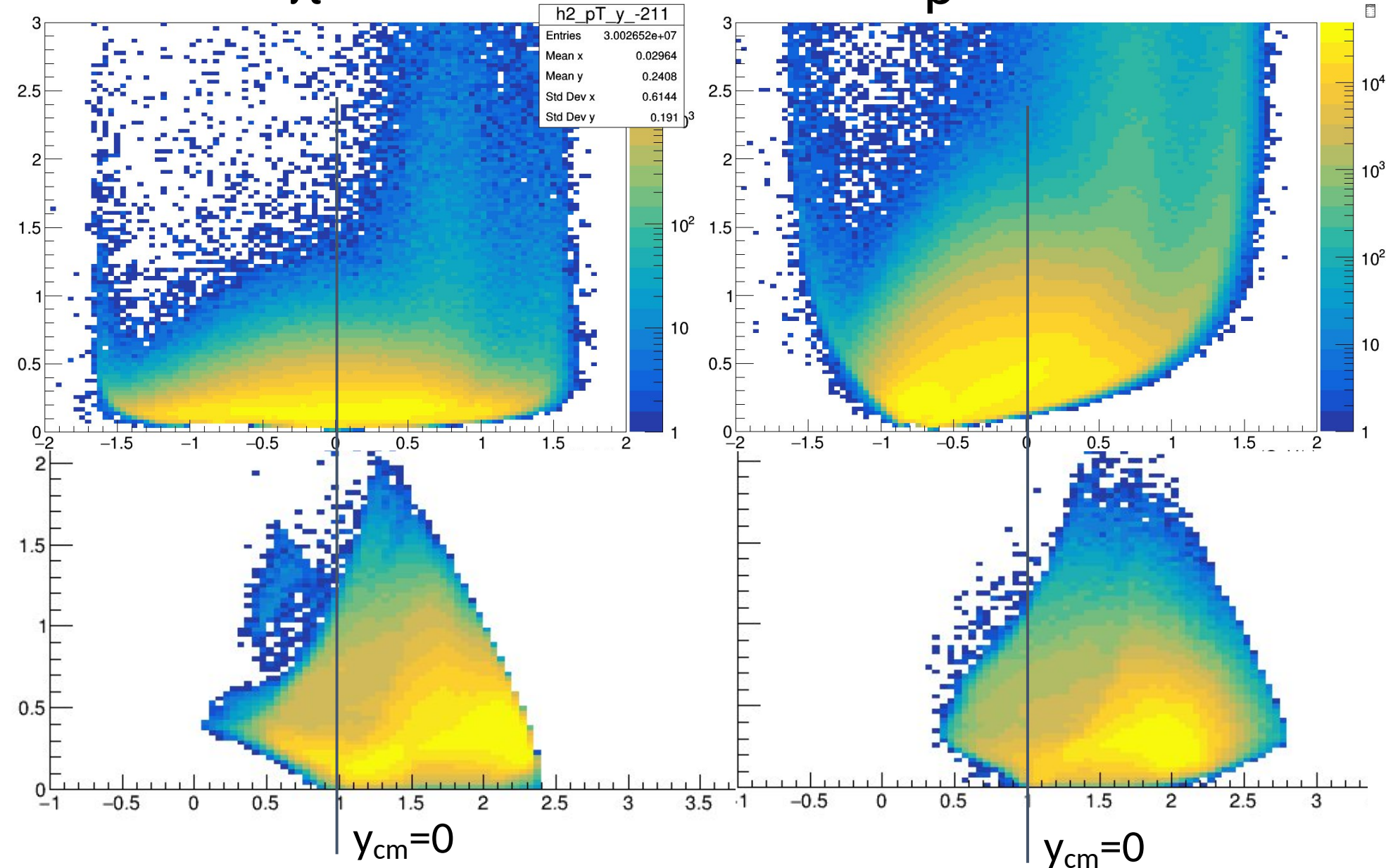
Square-like tracking system within the magnetic field deflecting particles along X-axis

Charge splitting on the surface of the FHCaI is observed due to magnetic field

# BM@N vs MPD: $p_T$ - $y$ acceptance

$\pi^-$

$p$

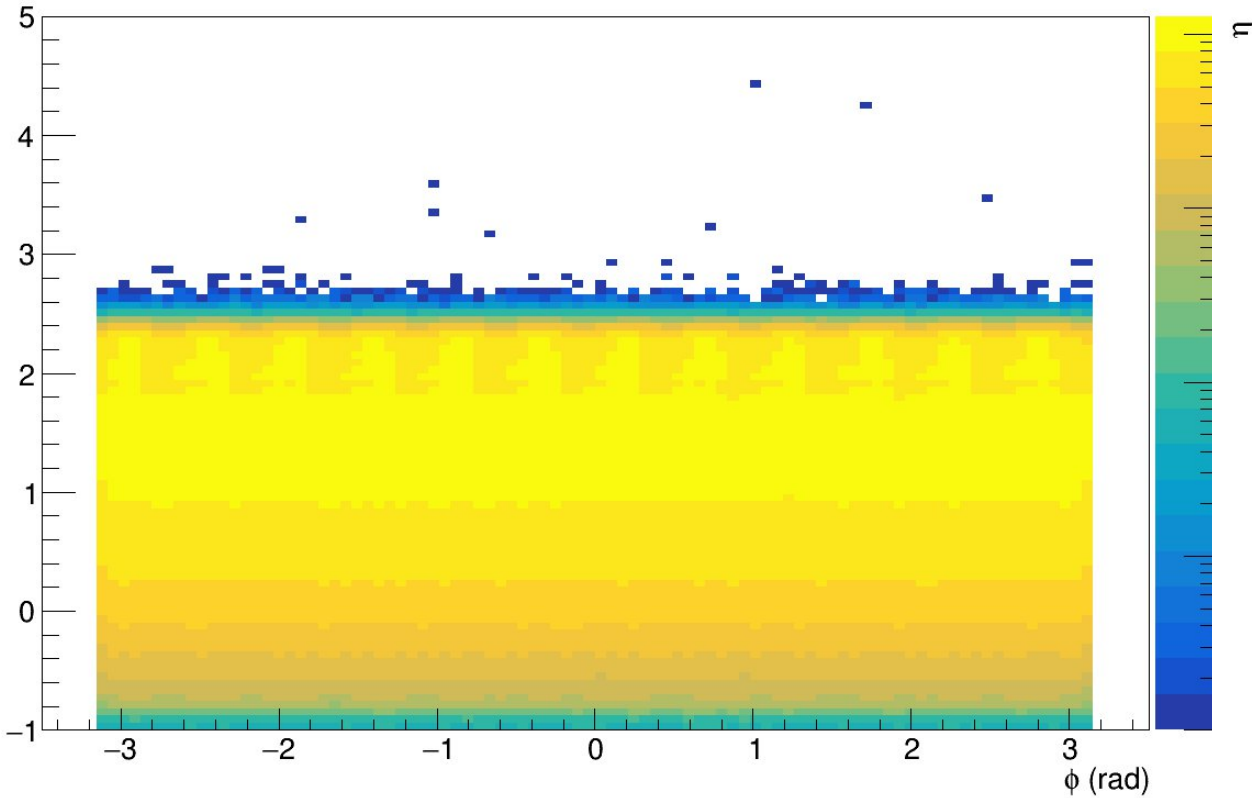


MPD has greater coverage of backward area (even covers projectile spectators) and MPD covers midrapidity region

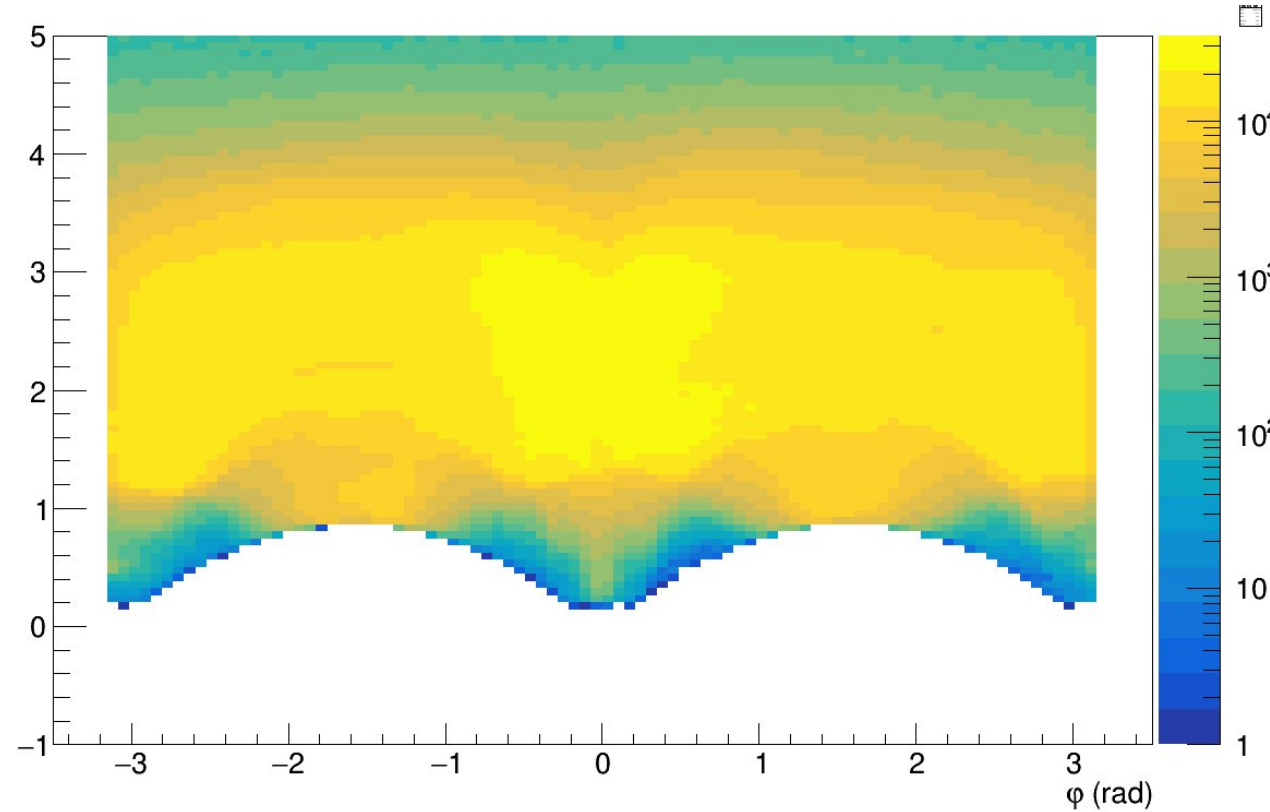
BM@N has greater coverage of forward area

# BM@N vs MPD: $\eta$ - $\phi$ acceptance

MPD



BM@N



- MPD has more uniform acceptance along  $\phi$ -axis
- BM@N has non-uniform acceptance due to square-like shape of the tracking system



# Flow vectors

From momentum of each measured particle define a  $u_n$ -vector in transverse plane:

$$u_n = e^{in\phi}$$

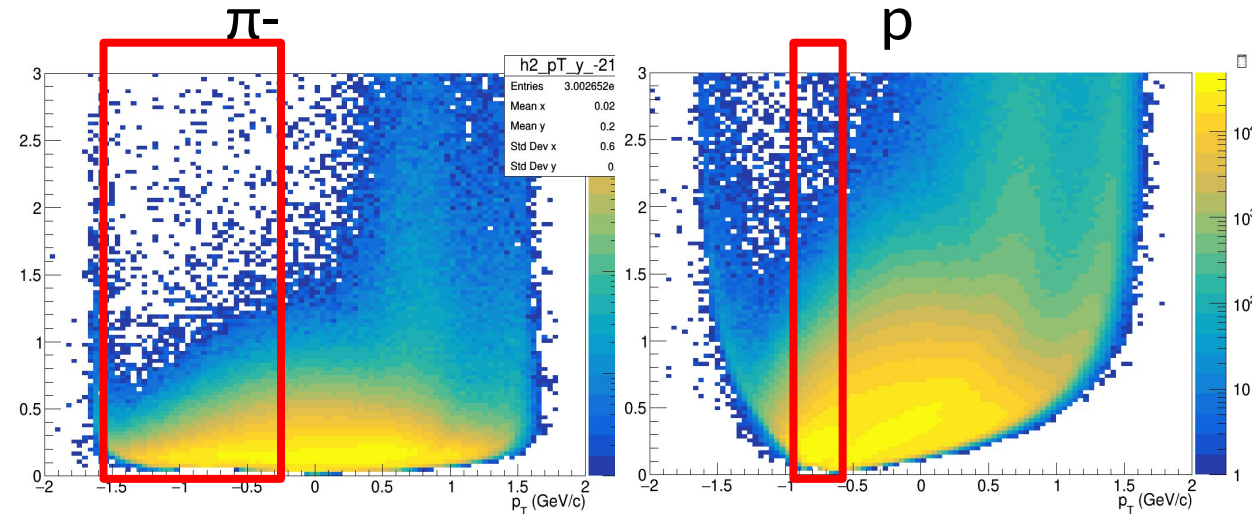
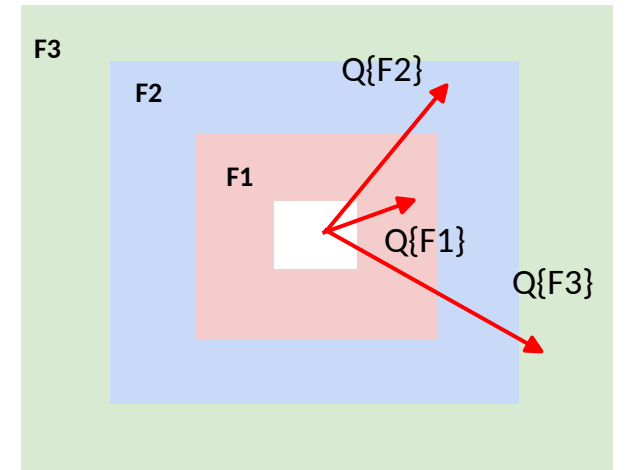
where  $\phi$  is the azimuthal angle

Sum over a group of  $u_n$ -vectors in one event forms  $Q_n$ -vector:

$$Q_n = \frac{\sum_{k=1}^N w_n^k u_n^k}{\sum_{k=1}^N w_n^k} = |Q_n| e^{in\Psi_n^{EP}}$$

$\Psi_n^{EP}$  is the event plane angle

Modules of FHCaI divided into 3 groups



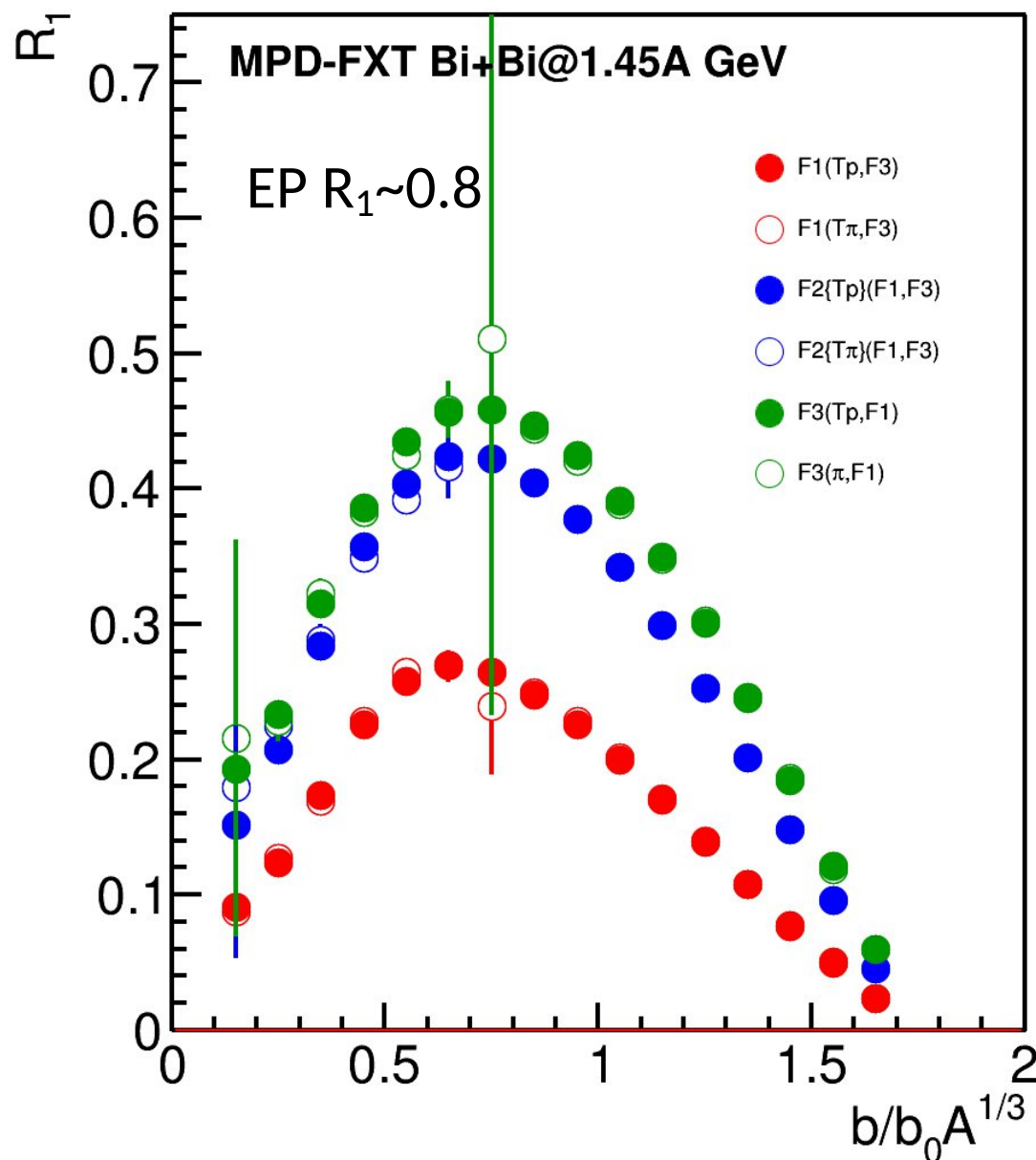
**Additional subevents from tracks not pointing at FHCaI:**

**Tp:** p;  $-1.0 < y < -0.6$ ;

**Tπ:** π-;  $-1.5 < y < -0.2$ ;

# Scalar Product method using FHCAL symmetry plane

Tested in HADES: M Mamaev et al 2020 PPNuclei 53, 277-281  
M Mamaev et al 2020 J. Phys.: Conf. Ser. 1690 012122



Scalar product (SP) method:

$$v_1 = \frac{\langle u_1 Q_1^{F1} \rangle}{R_1^{F1}} \quad v_2 = \frac{\langle u_2 Q_1^{F1} Q_1^{F3} \rangle}{R_1^{F1} R_1^{F3}}$$

Where  $R_1$  is the resolution correction factor:

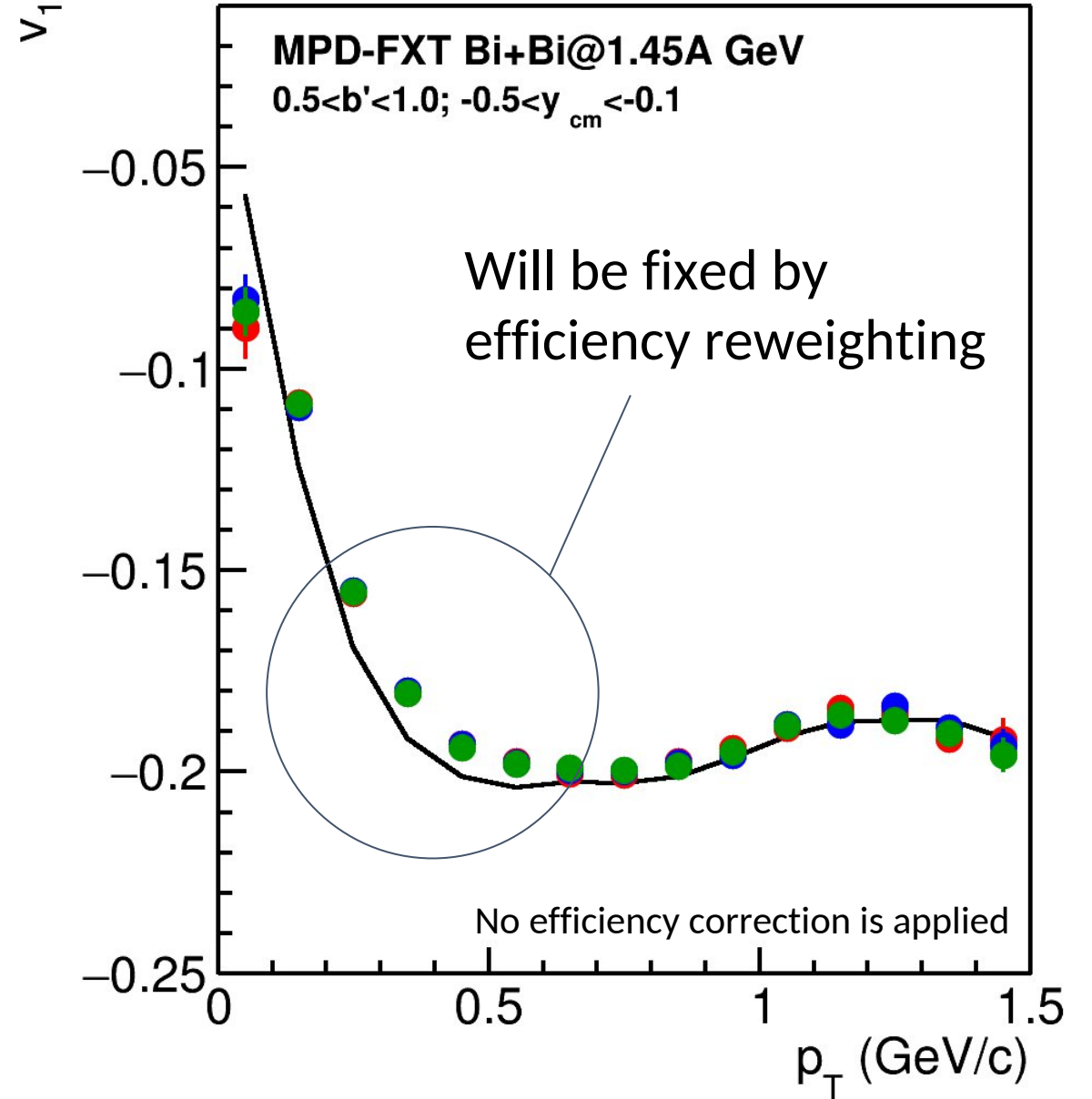
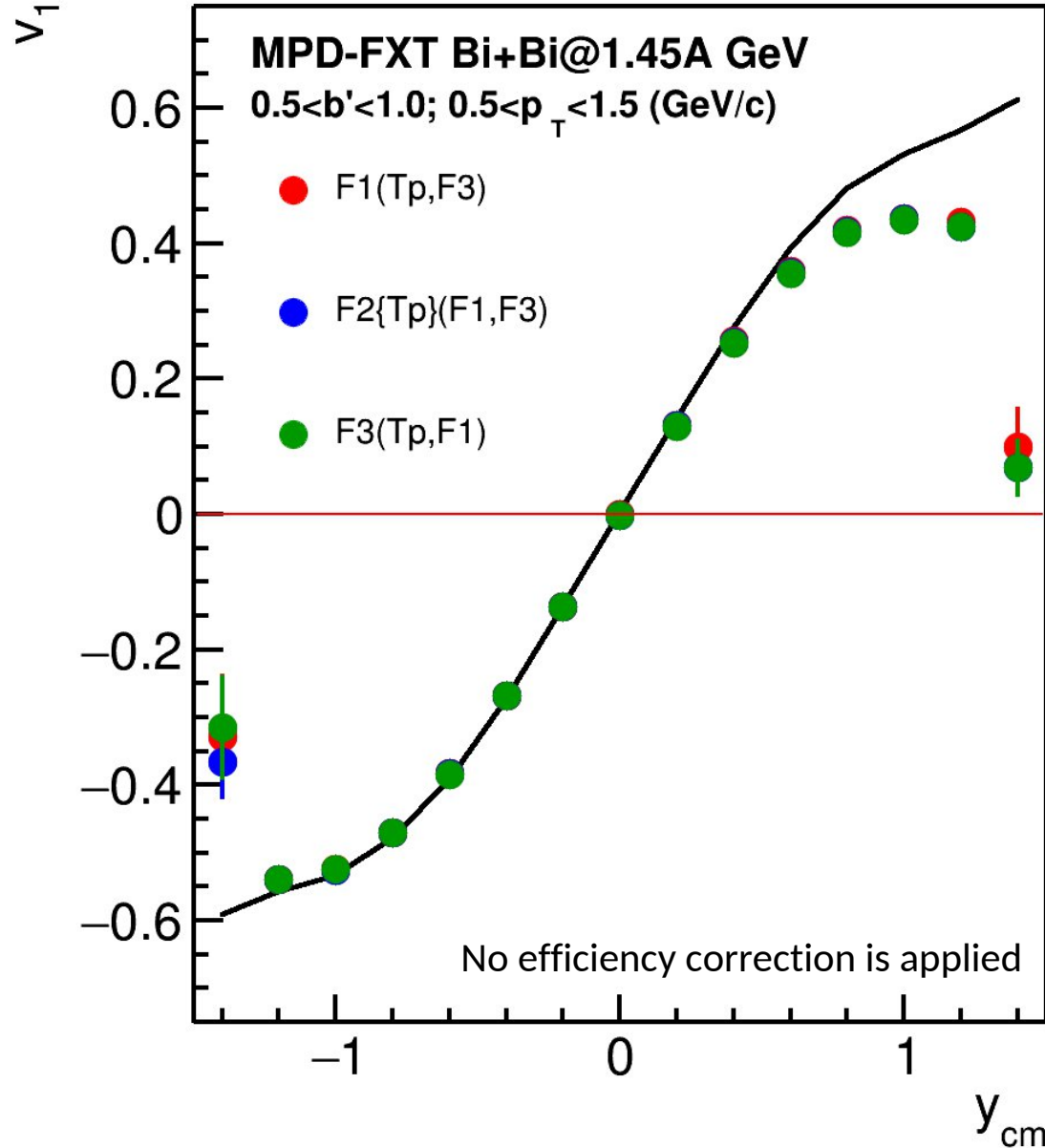
$$R_1^{F1} = \langle \cos(\Psi_1^{F1} - \Psi_1^{RP}) \rangle$$

Symbol "F2{Tp}(F1,F3)" means  $R_1$  calculated via (4S resolution):

$$R_1^{F2\{Tp\}(F1,F3)} = \langle Q_1^{F2} Q_1^{Tp} \rangle \frac{\sqrt{\langle Q_1^{F1} Q_1^{F3} \rangle}}{\sqrt{\langle Q_1^{Tp} Q_1^{F1} \rangle \langle Q_1^{Tp} Q_1^{F3} \rangle}}$$

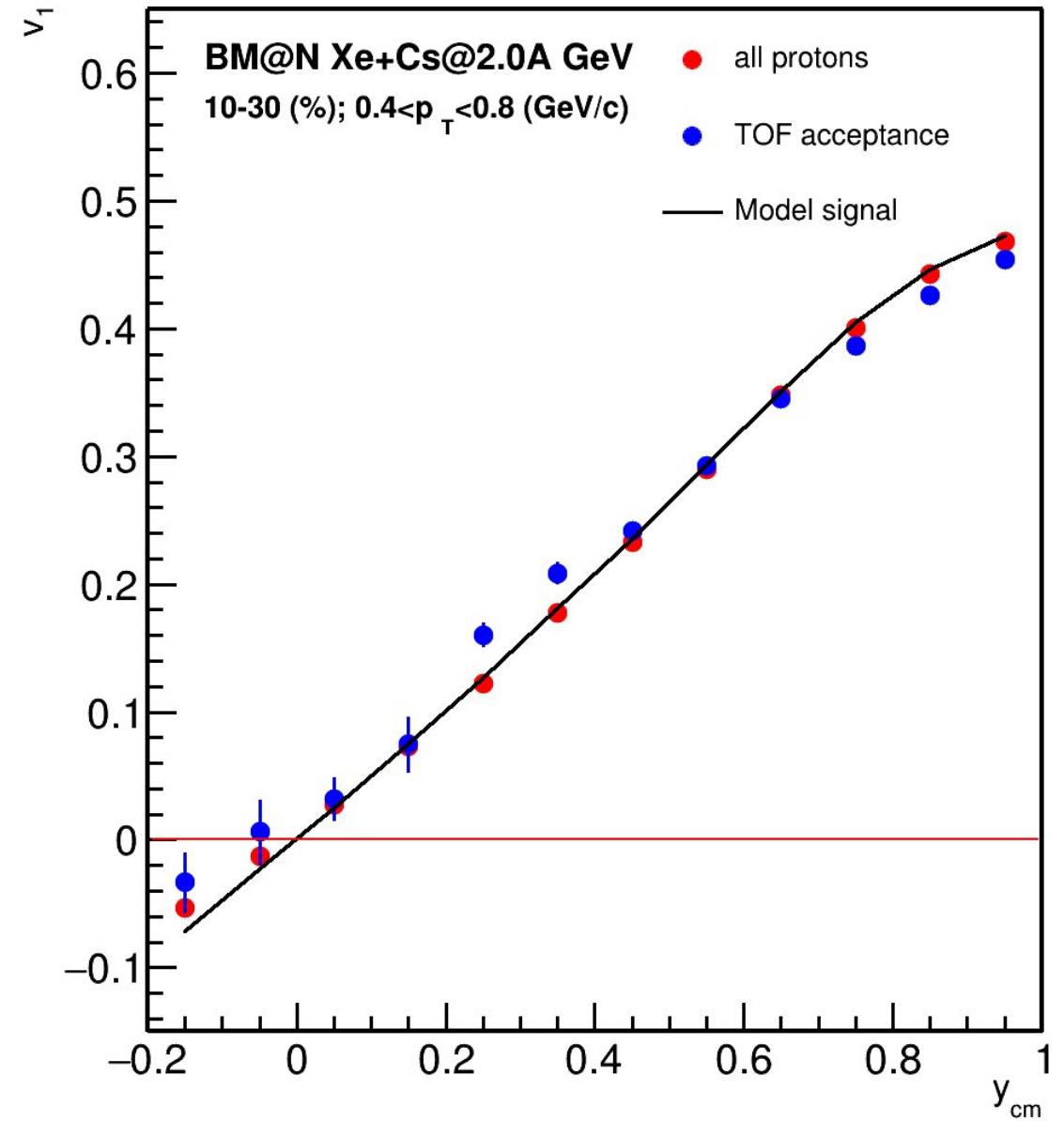
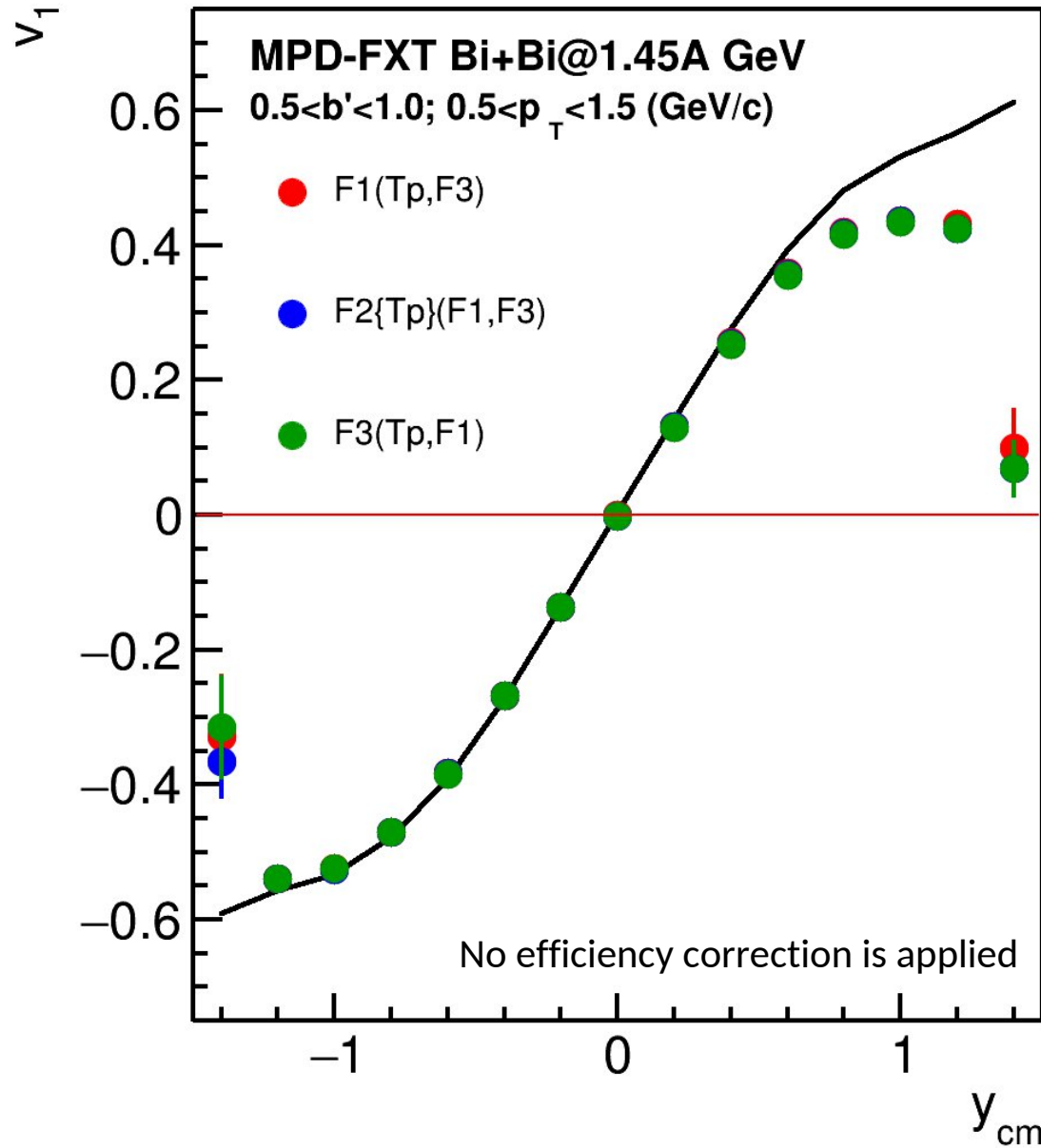
Good agreement between  $R_1$  calculated using different combinations of Q-vectors with significant rapidity separation

# MPD-FXT: $v_1$ for protons



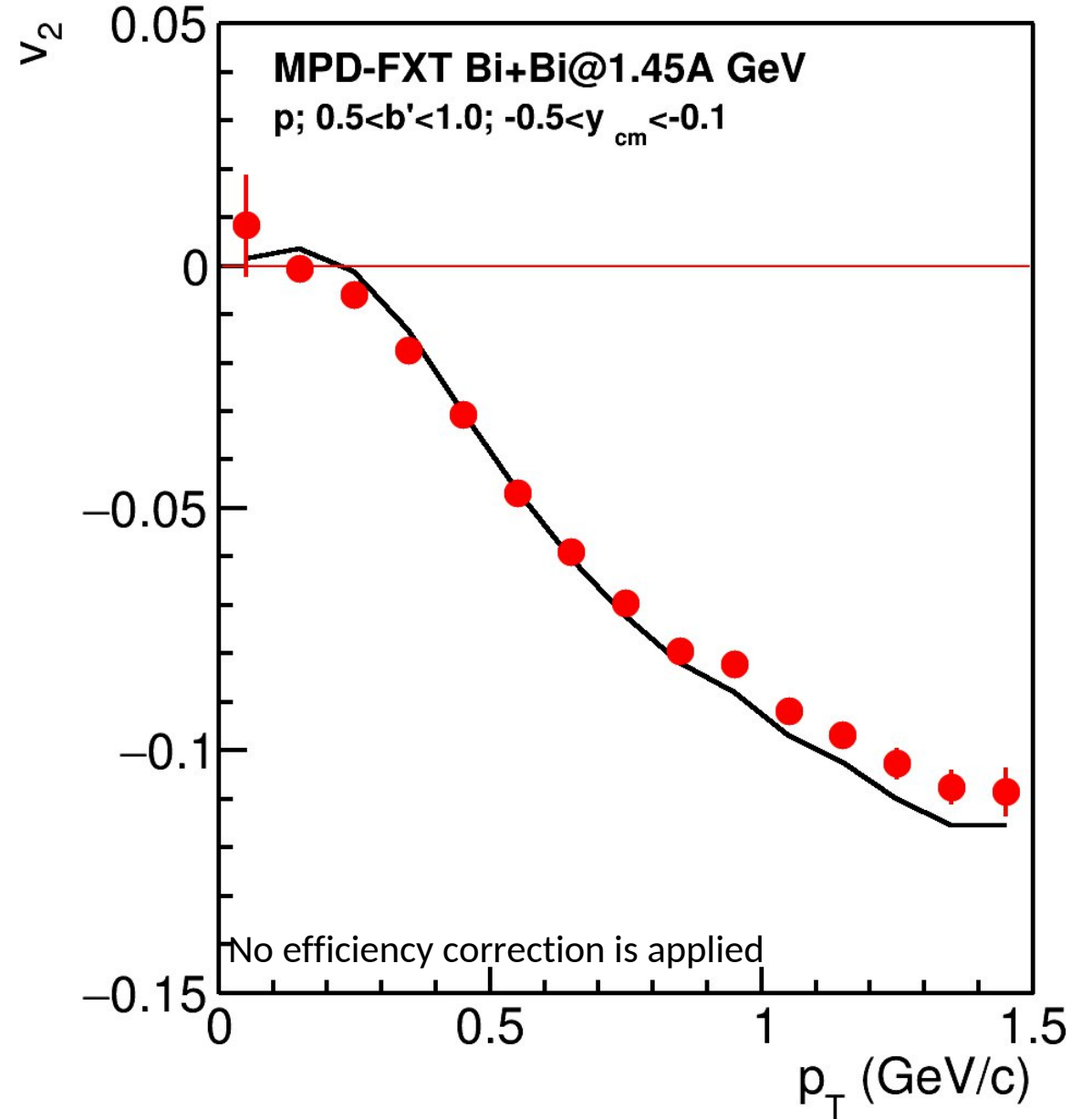
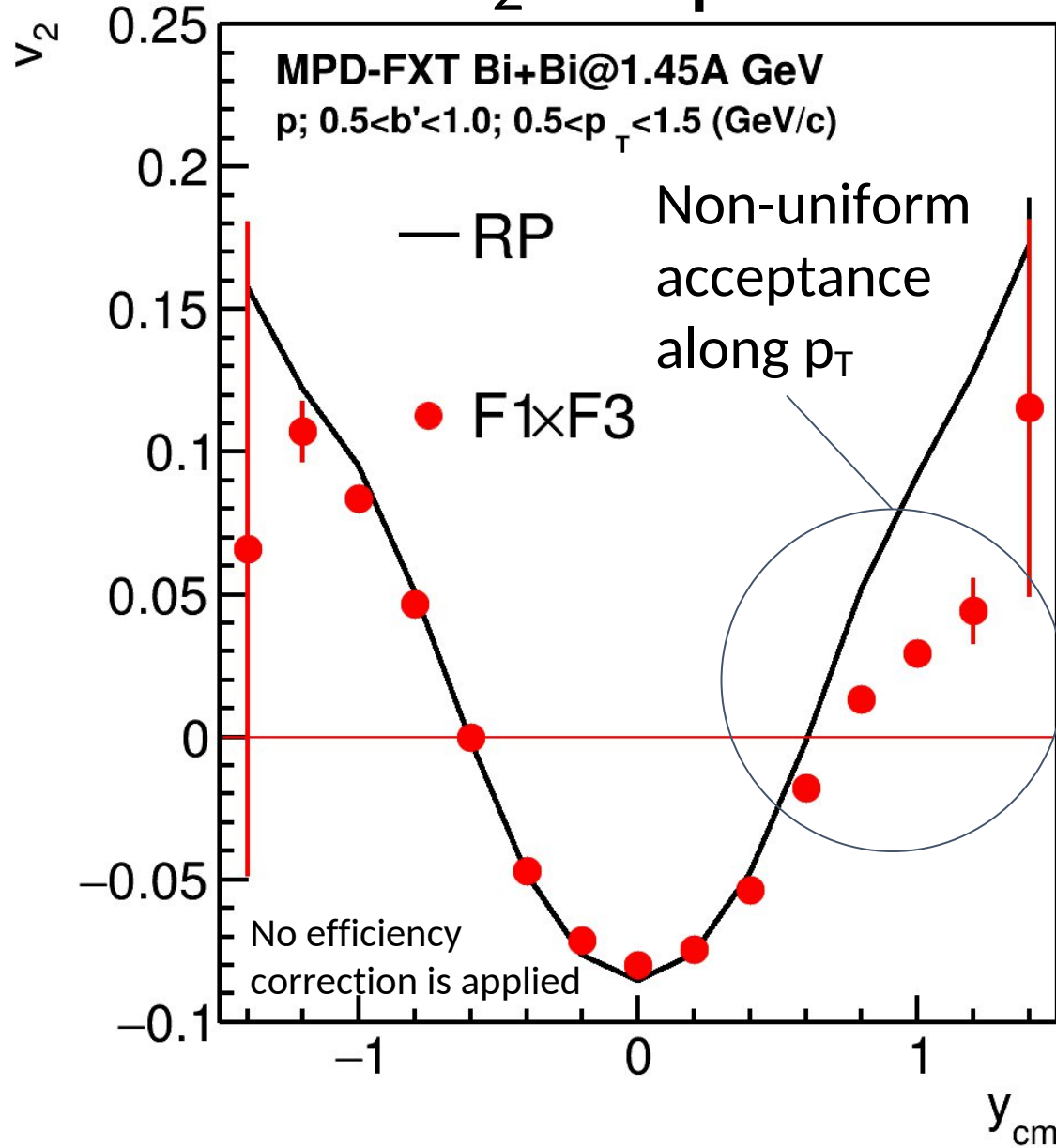
$v_1$  is consistent with model signal for  $y < 0.5$

# BM@N vs MPD: $v_1$ vs $y$ for protons



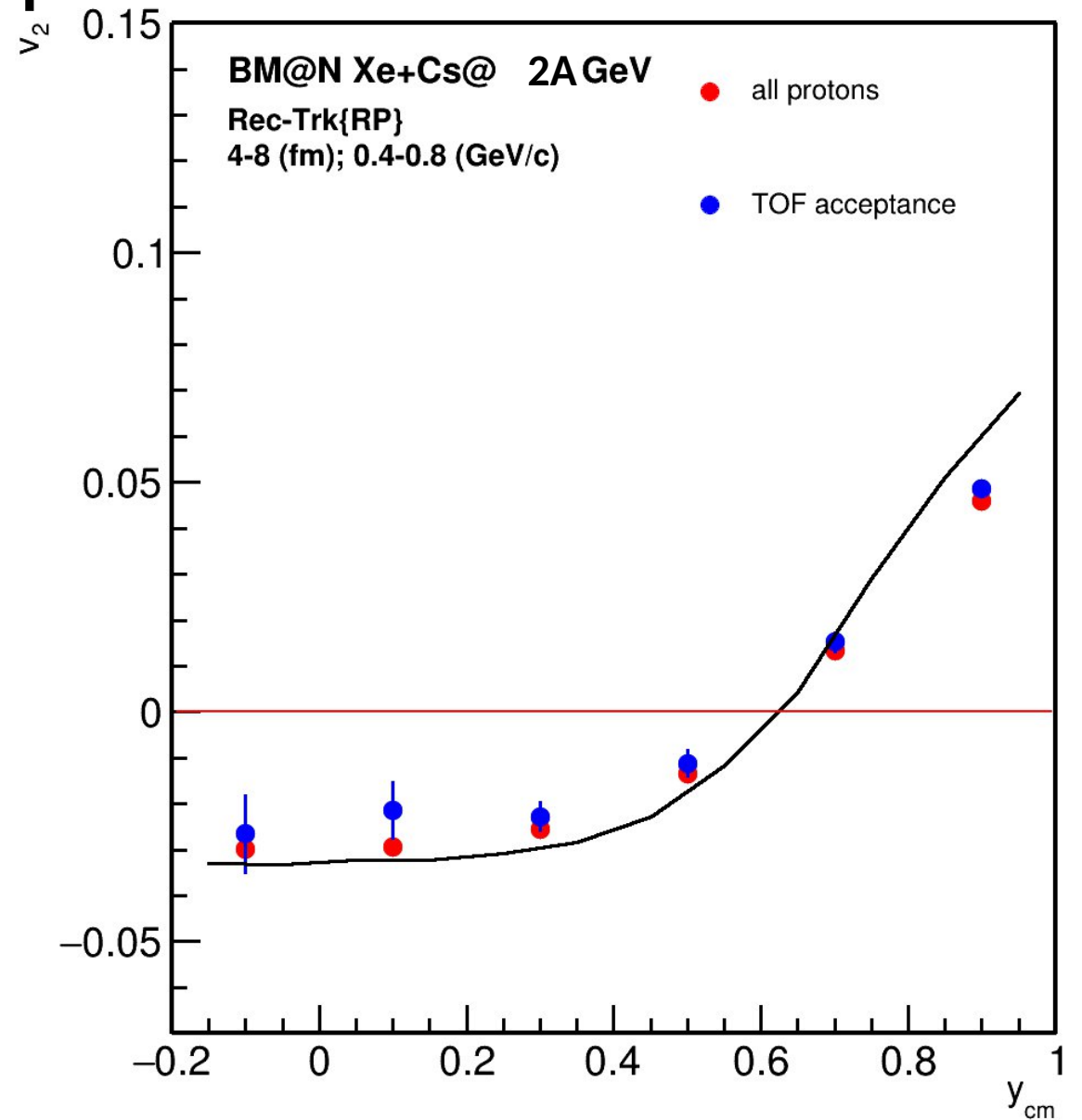
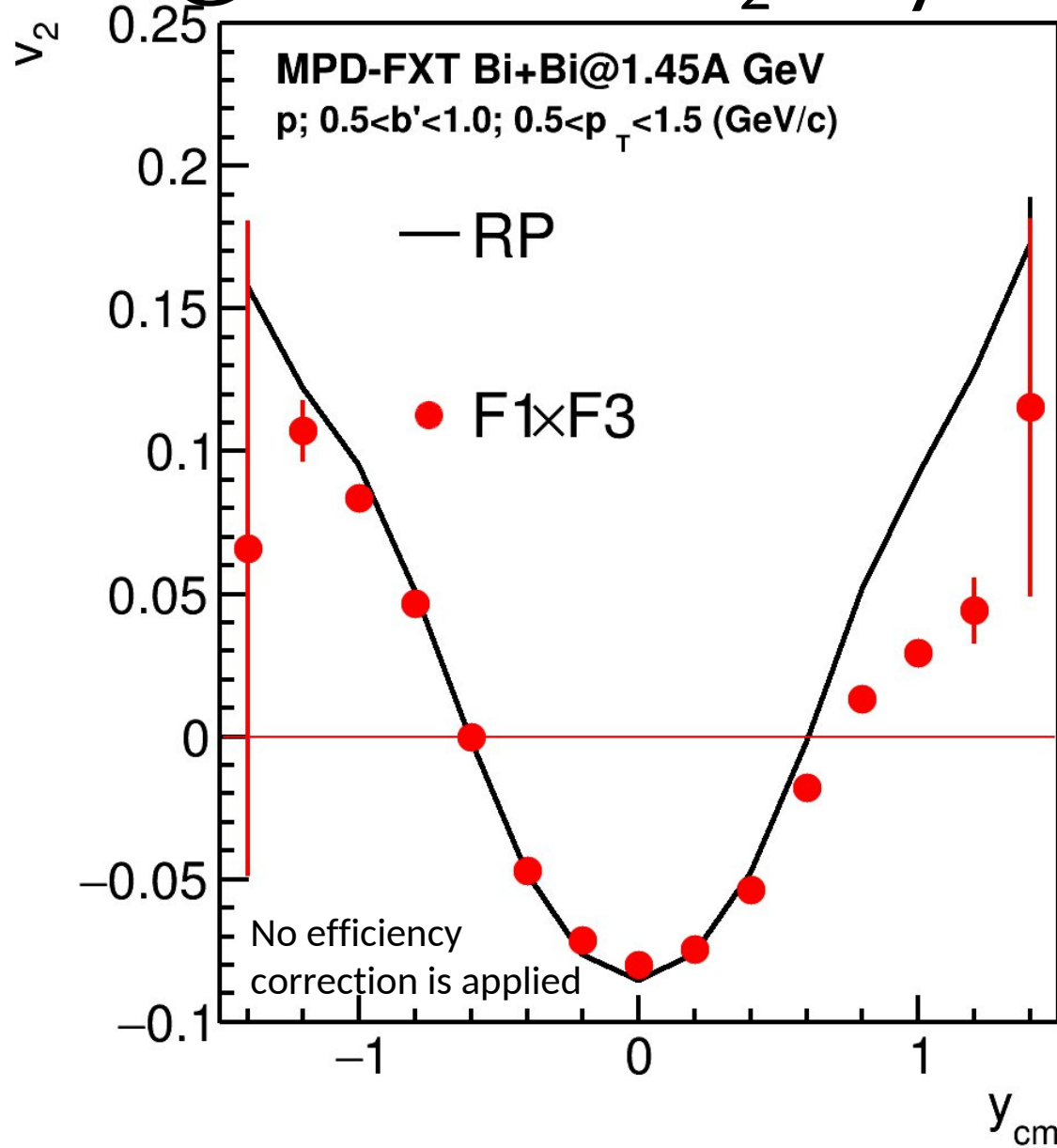
BM@N has better coverage in positive rapidities close to  $y_{beam}$

# MPD-FXT: $v_2$ for protons



$v_2$  is consistent with model signal for  $y < 0.5$

# BM@N vs MPD: $v_2$ vs $y$ for protons



BM@N has better coverage in positive rapidities close to  $y_{beam}$

# Summary

- The feasibility study for the flow measurements in the MPD experiment in a fixed-target mode was carried out with GEANT4 detector simulation and UrQMD Bi+Bi@1.45A GeV events as an input
- Acceptances of the BM@N and MPD facilities were compared:
  - MPD has greater coverage of the backward rapidities and midrapidity region
  - MPD has more uniform coverage for the azimuthal angle
- The procedure for the resolution correction factor  $R_1$  with 3 sub-event method and rapidity-separated combinations of Q-vectors was employed
  - Two separate estimations for the  $R_1$  for each symmetry plane were found in a good agreement
- Directed and elliptic flow for protons and light mesons were measured
  - For each particle species  $v_1$  and  $v_2$  are consistent with the model signal mostly in backward rapidities
- Perform flow measurements for MPD FXT for different energies:  $\sqrt{s_{NN}}=2.5, 3.0, 3.5$  GeV

**Thank you for your attention!**

# Backup



# Flow methods for $v_n$ calculation

Tested in HADES: M Mamaev et al 2020 PPNuclei 53, 277–281  
 M Mamaev et al 2020 J. Phys.: Conf. Ser. 1690 012122

Scalar product (SP) method:

$$v_1 = \frac{\langle u_1 Q_1^{F1} \rangle}{R_1^{F1}} \quad v_2 = \frac{\langle u_2 Q_1^{F1} Q_1^{F3} \rangle}{R_1^{F1} R_1^{F3}}$$

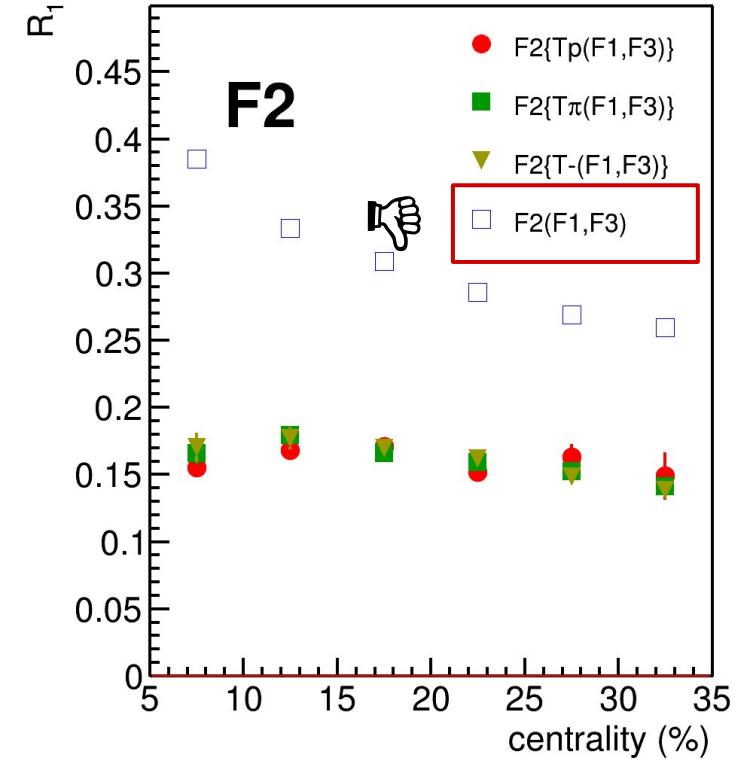
Where  $R_1$  is the resolution correction factor

$$R_1^{F1} = \langle \cos(\Psi_1^{F1} - \Psi_1^{RP}) \rangle$$

Symbol “F2(F1,F3)” means  $R_1$  calculated via  
 (3S resolution):

$$R_1^{F2(F1,F3)} = \frac{\sqrt{\langle Q_1^{F2} Q_1^{F1} \rangle \langle Q_1^{F2} Q_1^{F3} \rangle}}{\sqrt{\langle Q_1^{F1} Q_1^{F3} \rangle}}$$

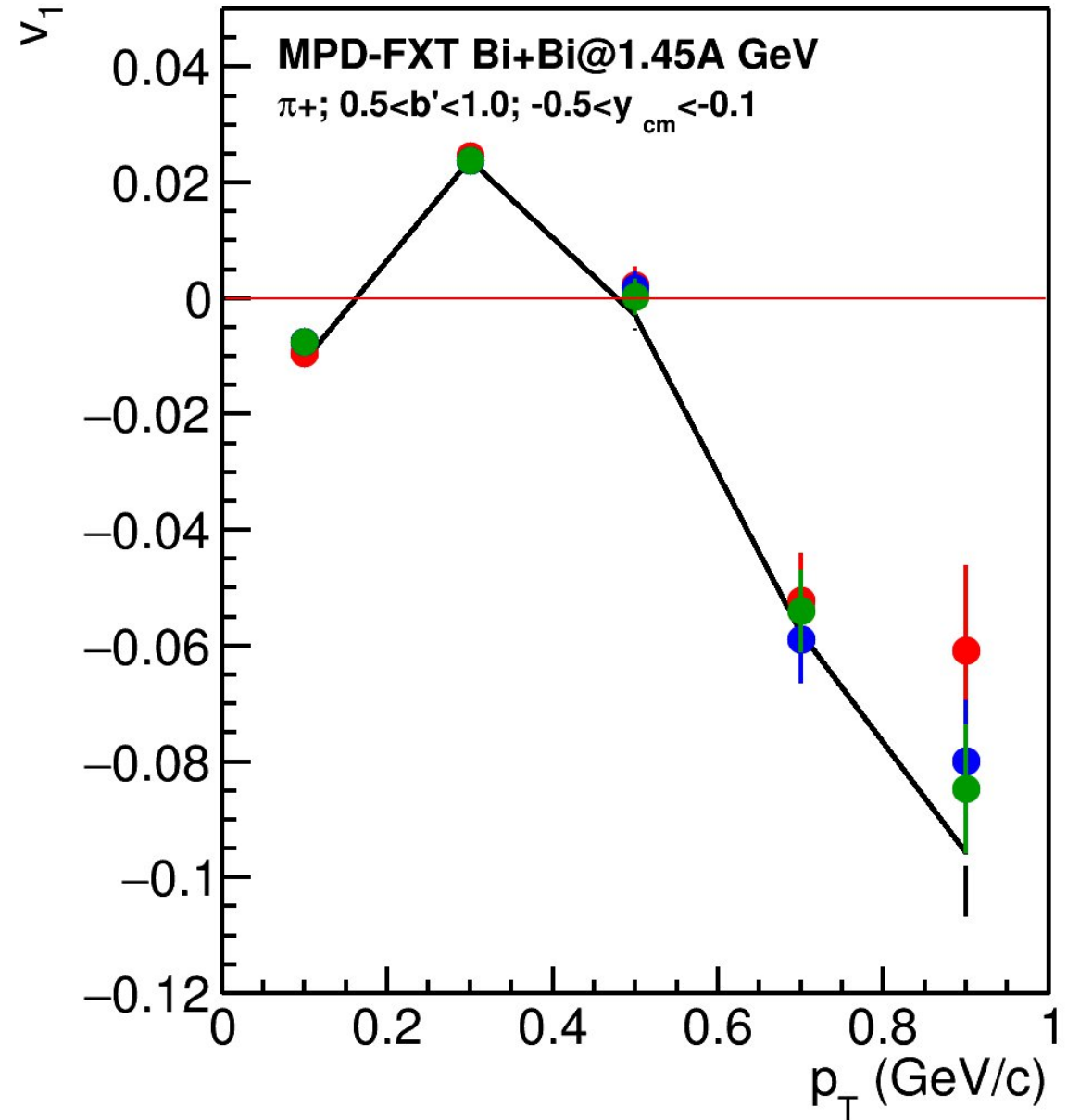
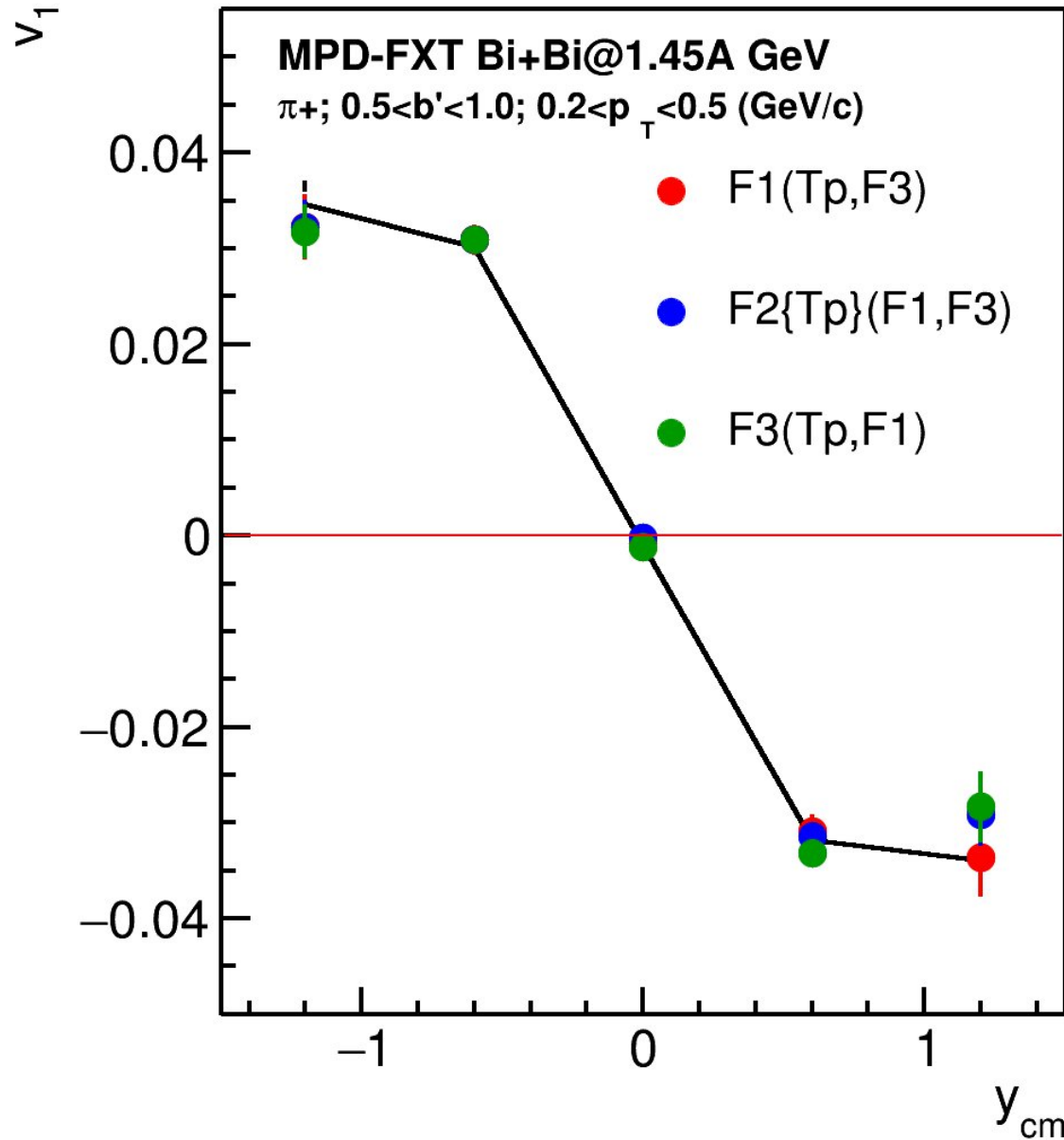
Method helps to eliminate non-flow  
 Using 2-subevents doesn't work



Symbol “F2{Tp}(F1,F3)” means  $R_1$  calculated via  
 (4S resolution):

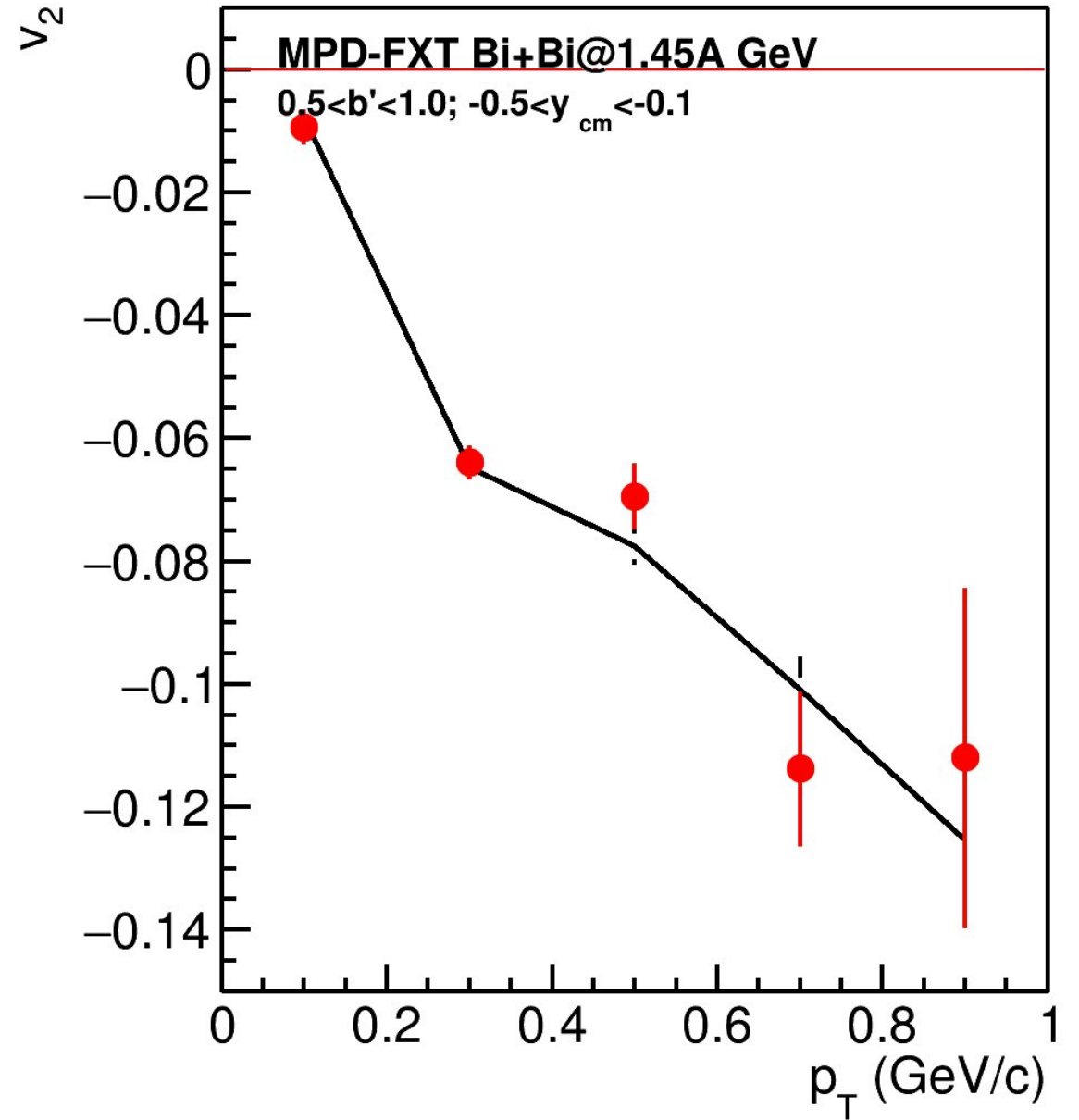
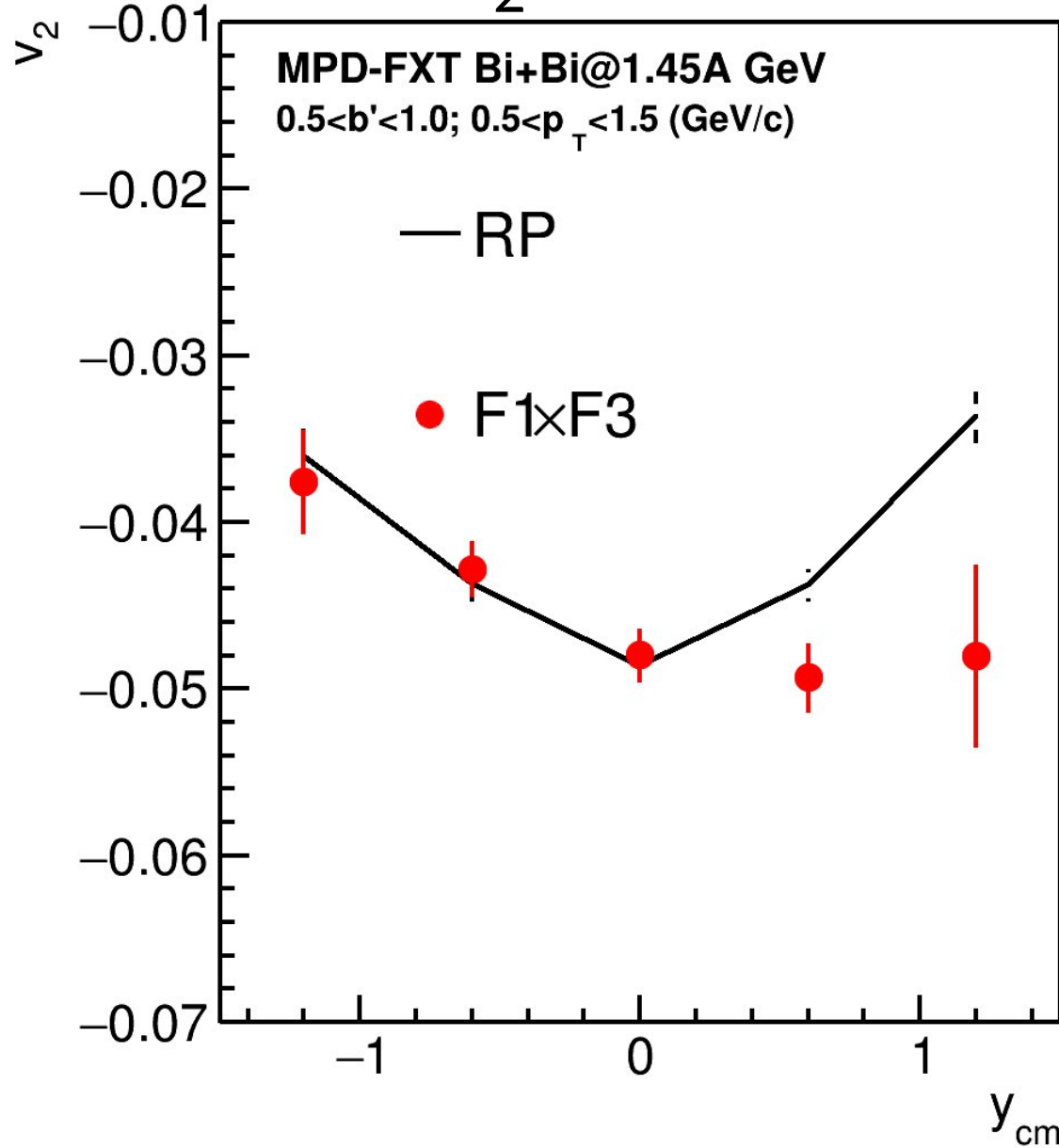
$$R_1^{F2\{Tp\}(F1,F3)} = \langle Q_1^{F2} Q_1^{Tp} \rangle \frac{\sqrt{\langle Q_1^{F1} Q_1^{F3} \rangle}}{\sqrt{\langle Q_1^{Tp} Q_1^{F1} \rangle \langle Q_1^{Tp} Q_1^{F3} \rangle}}$$

# MPD-FXT: $v_1$ for $\pi^+$



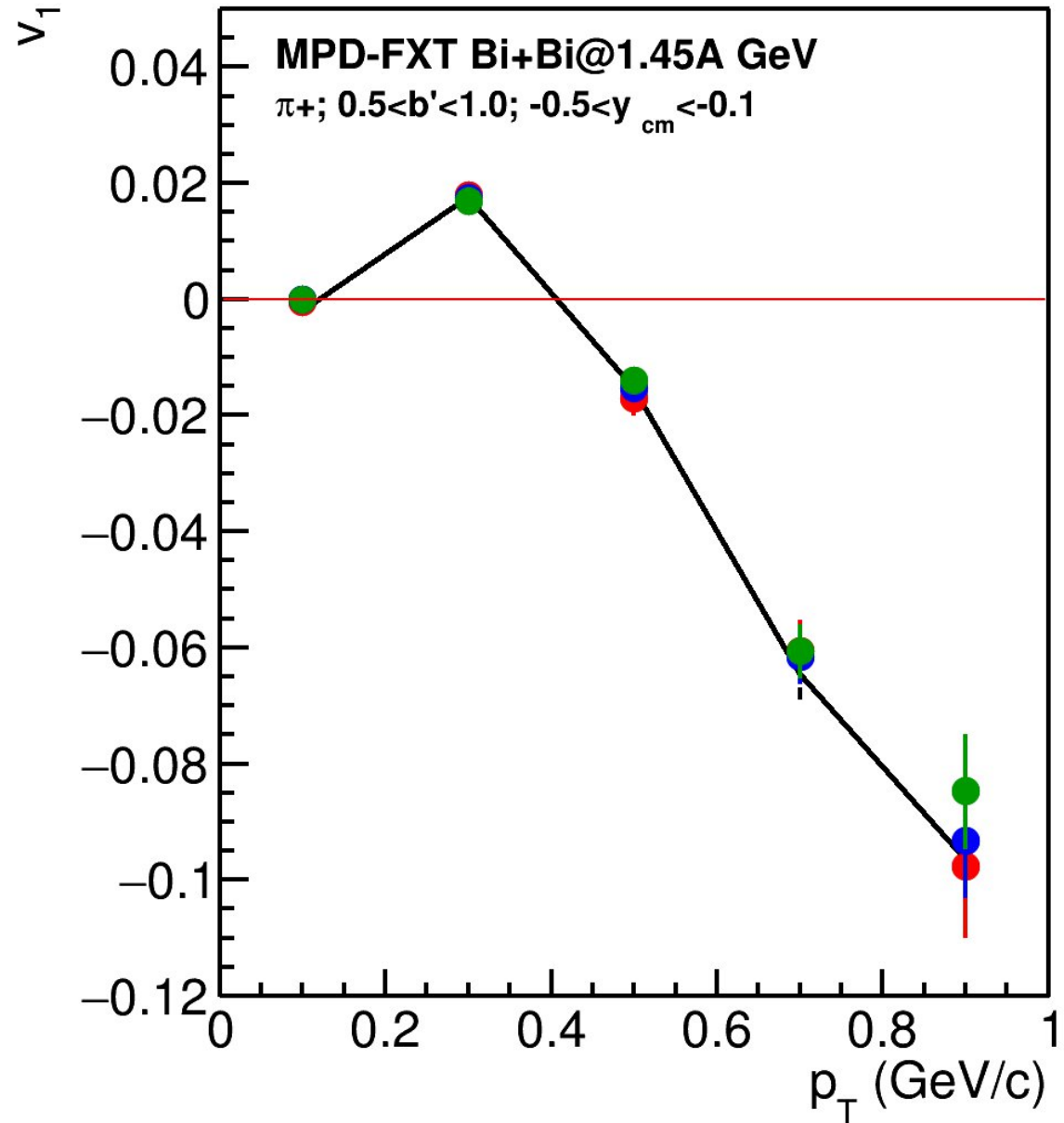
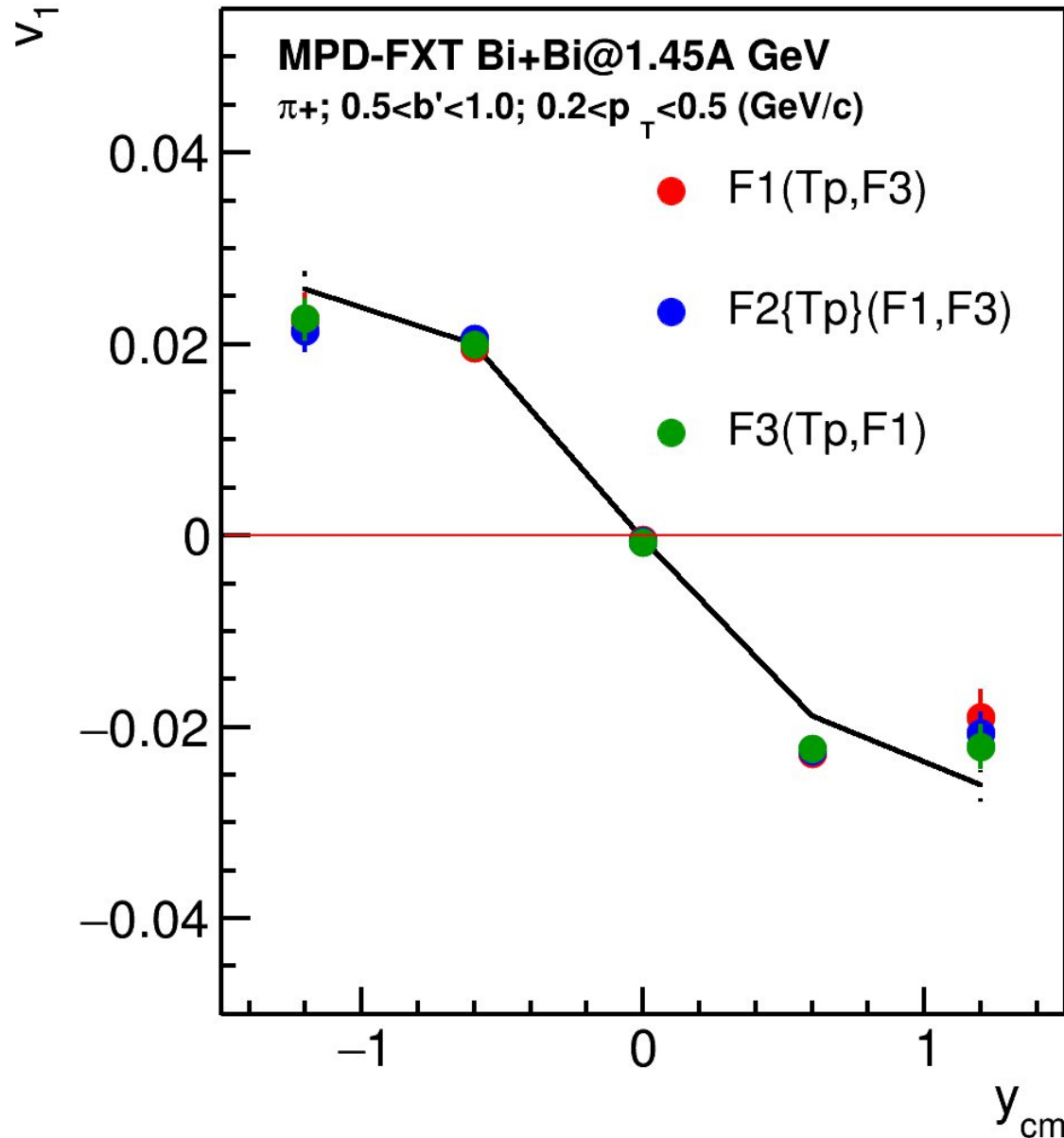
$v_1$  is consistent with model signal for  $y < 1$ ; We need more statistics

# MPD-FXT: $v_2$ for $\pi^+$



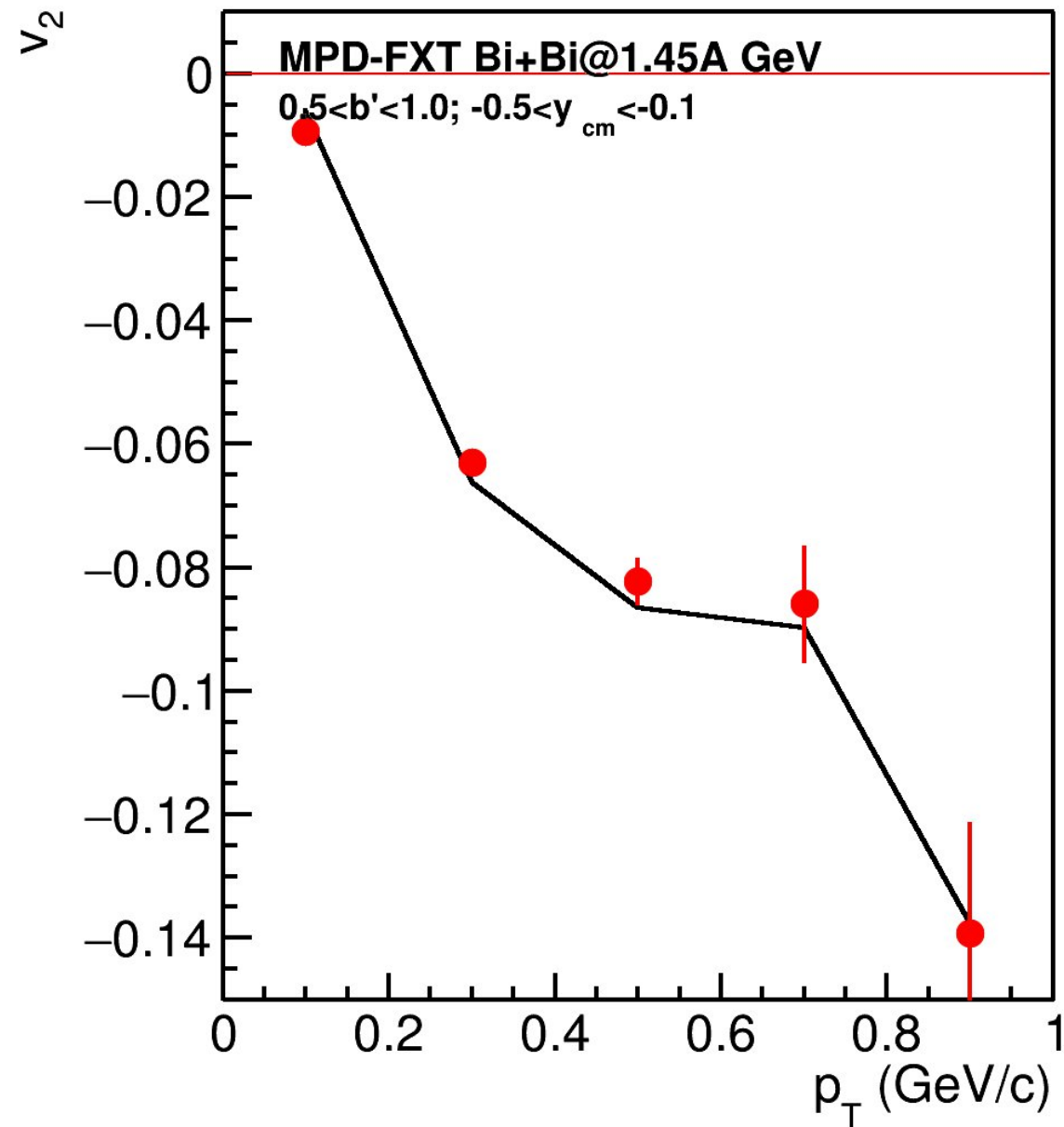
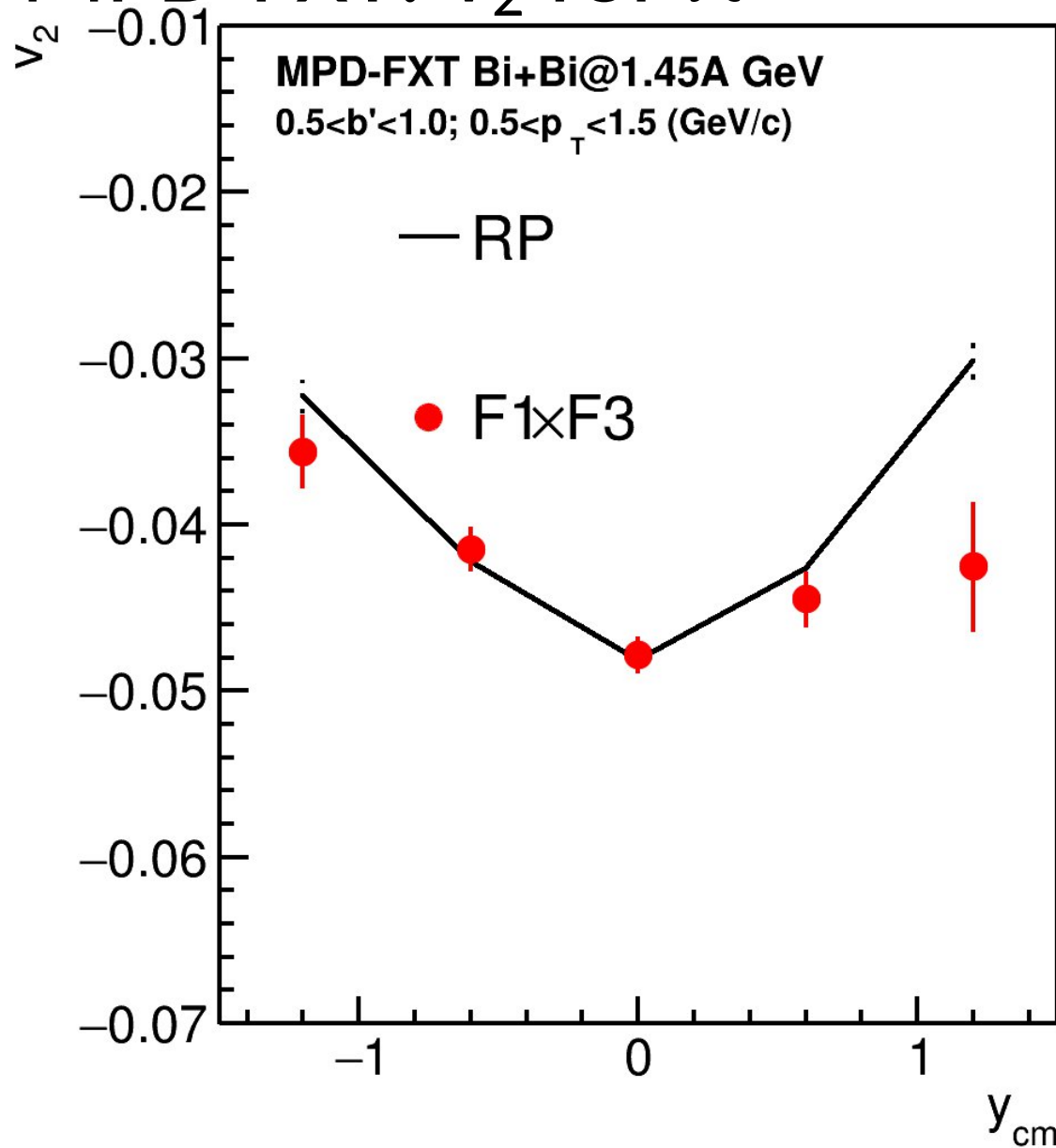
$v_2$  is consistent with model signal for  $y < 0$ ; We need more statistics

# MPD-FXT: $v_1$ for $\pi^-$



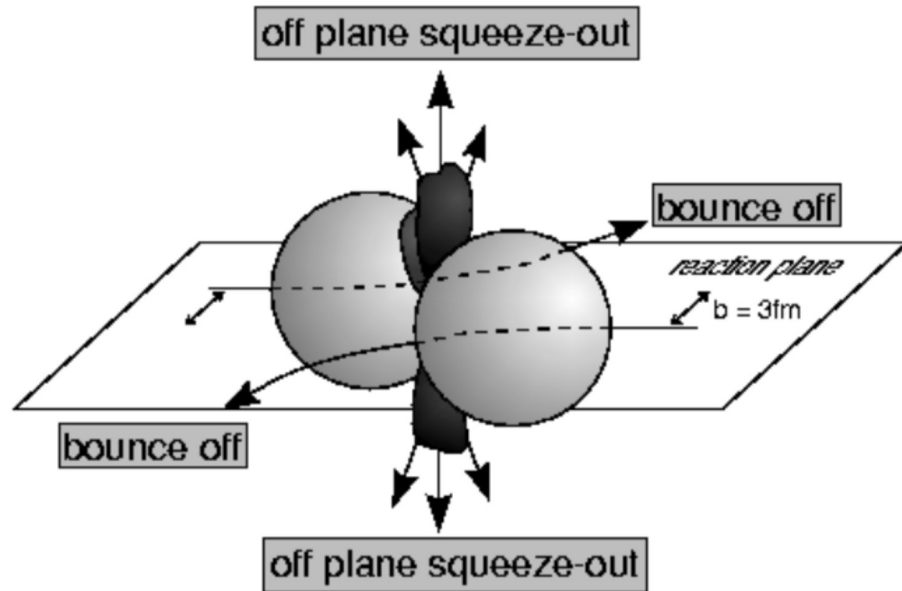
$v_1$  is consistent with model signal for  $y < 1$ ; We need more

# MPD-FXT: $v_2$ for $\pi^-$



$v_2$  is consistent with model signal for  $y < 0$ ; We need more statistics

# Anisotropic flow & spectators



The azimuthal angle distribution is decomposed in a Fourier series relative to reaction plane angle:

$$\rho(\varphi - \Psi_{RP}) = \frac{1}{2\pi} \left( 1 + 2 \sum_{n=1}^{\infty} v_n \cos n(\varphi - \Psi_{RP}) \right)$$

Anisotropic flow:

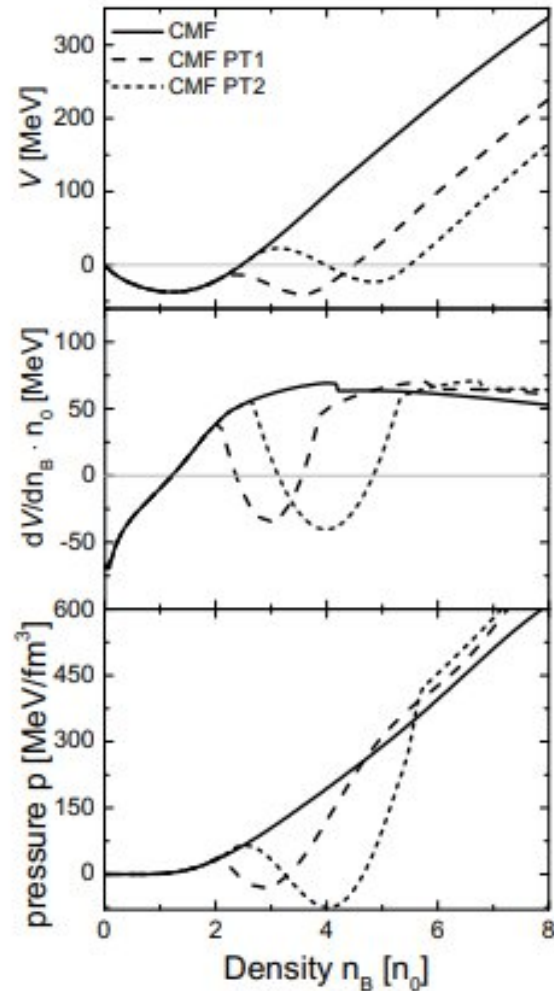
$$v_n = \langle \cos [n(\varphi - \Psi_{RP})] \rangle$$

Anisotropic flow is sensitive to:

- Time of the interaction between overlap region and spectators
- Compressibility of the created matter

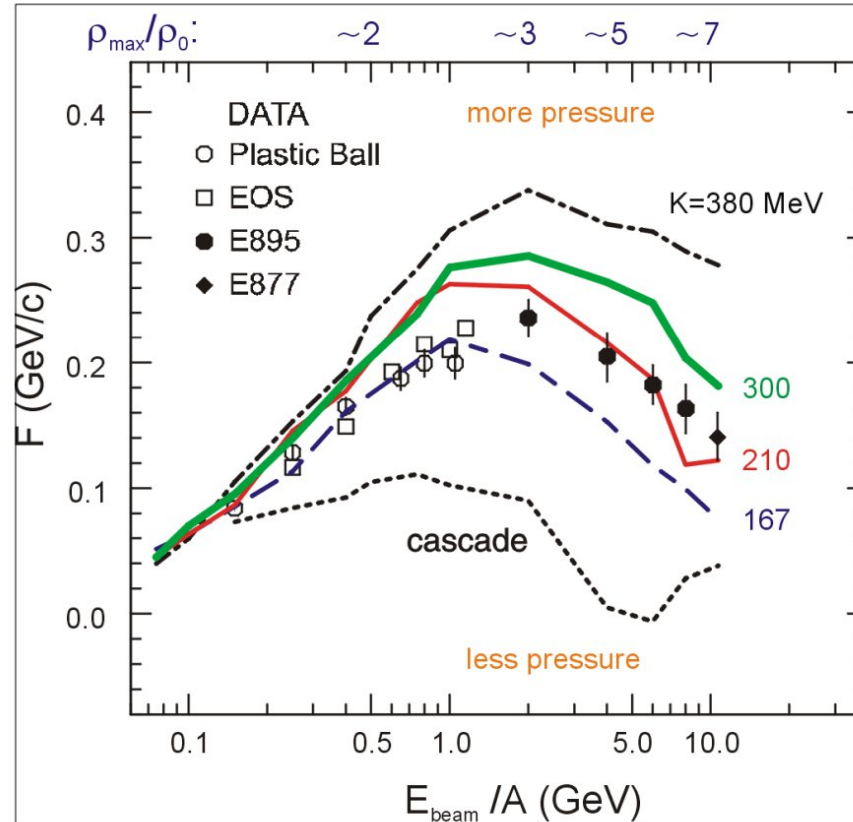
# $v_n$ as a function of collision energy

P. DANIELEWICZ, R. LACEY, W. LYNCH  
[10.1126/science.1078070](https://doi.org/10.1126/science.1078070)

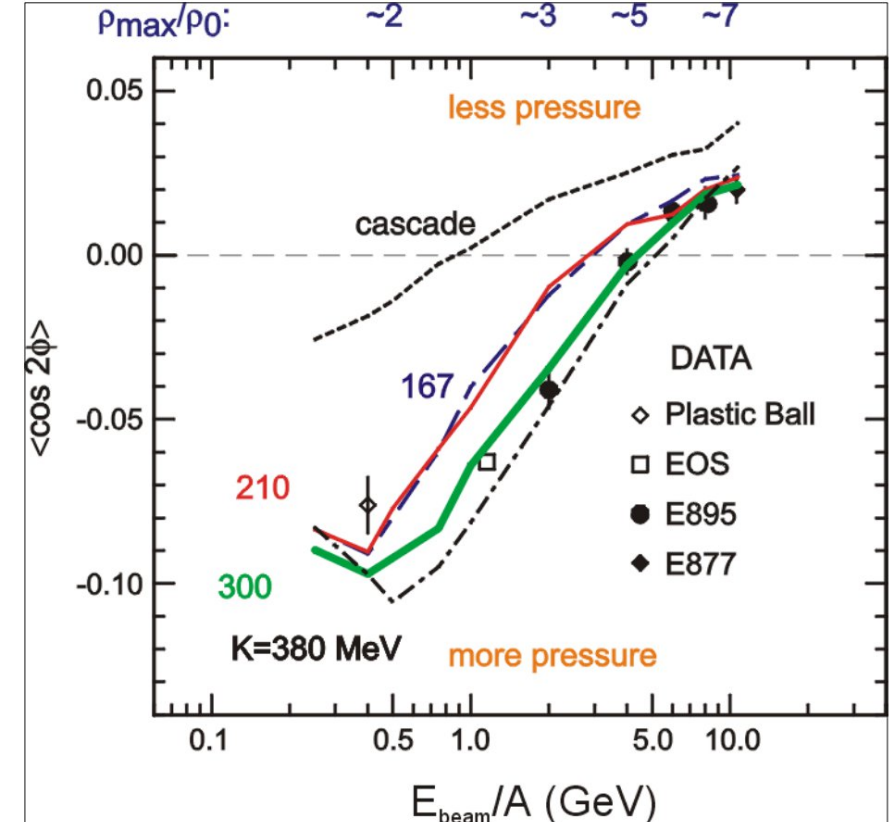


EPJ Web of Conferences 276, 01021 (2023)

$v_1$  suggests softer EOS



$v_2$  suggests harder EOS

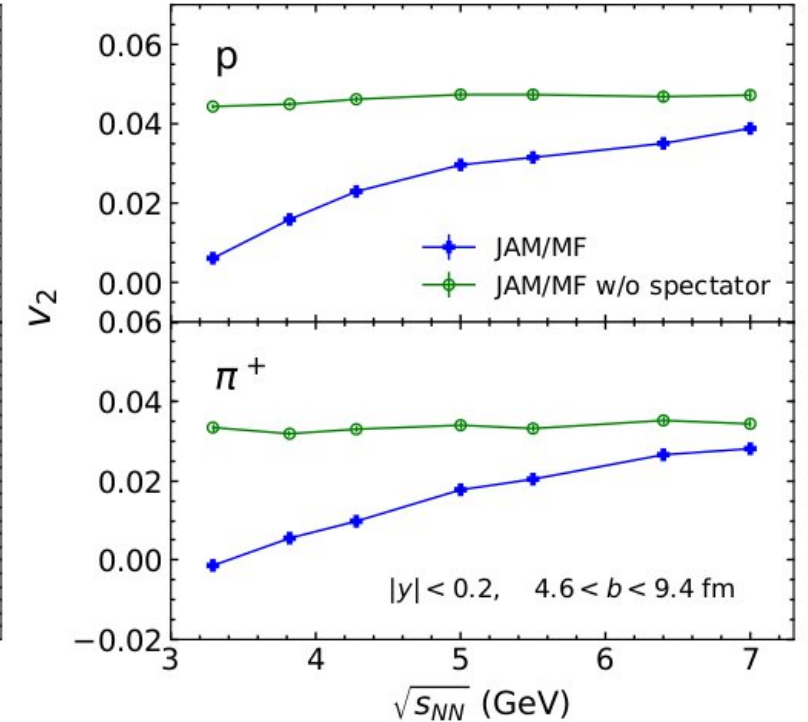
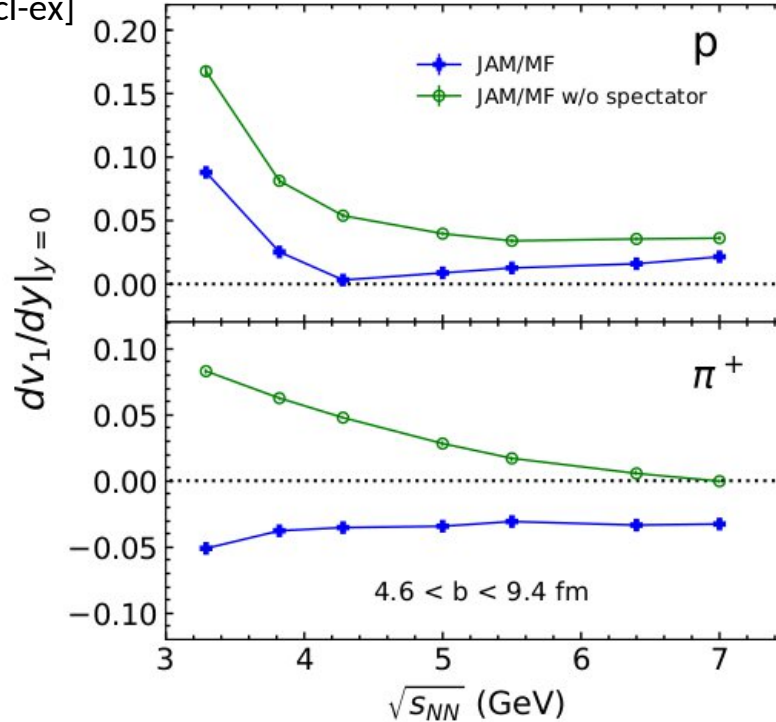
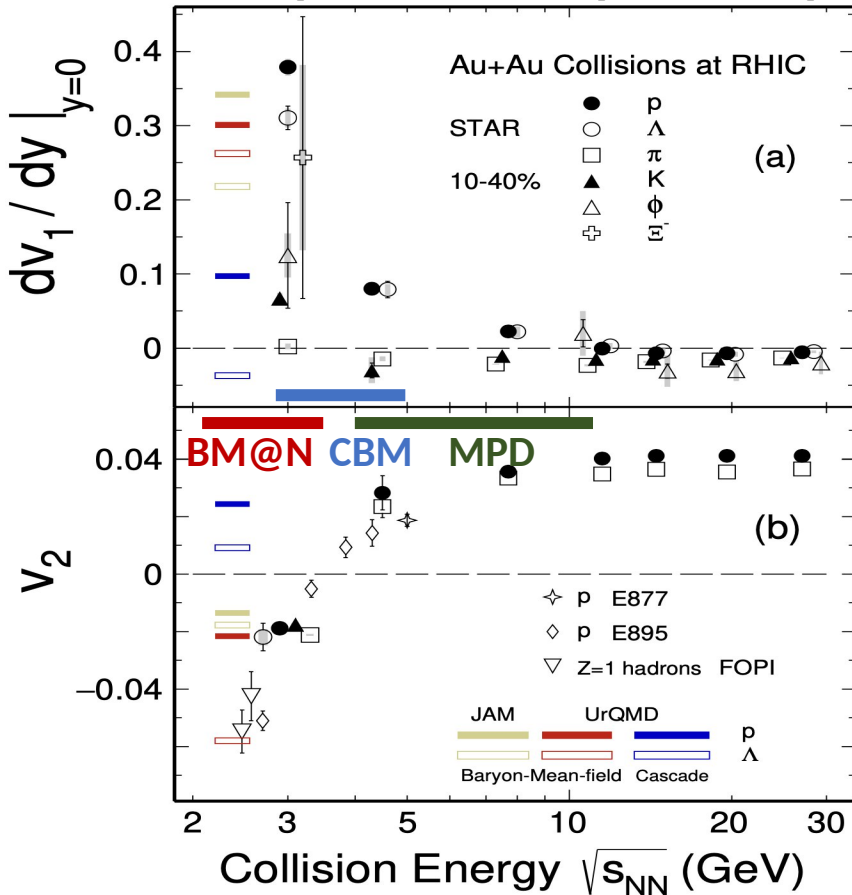


Describing the high-density matter using the mean field  
 Flow measurements constrain the mean field

Discrepancy is probably due to non-flow correlations

# Anisotropic flow in Au+Au collisions at Nuclotron-NICA energies

M. Abdallah et al. [STAR Collaboration] 2108.00908 [nucl-ex]



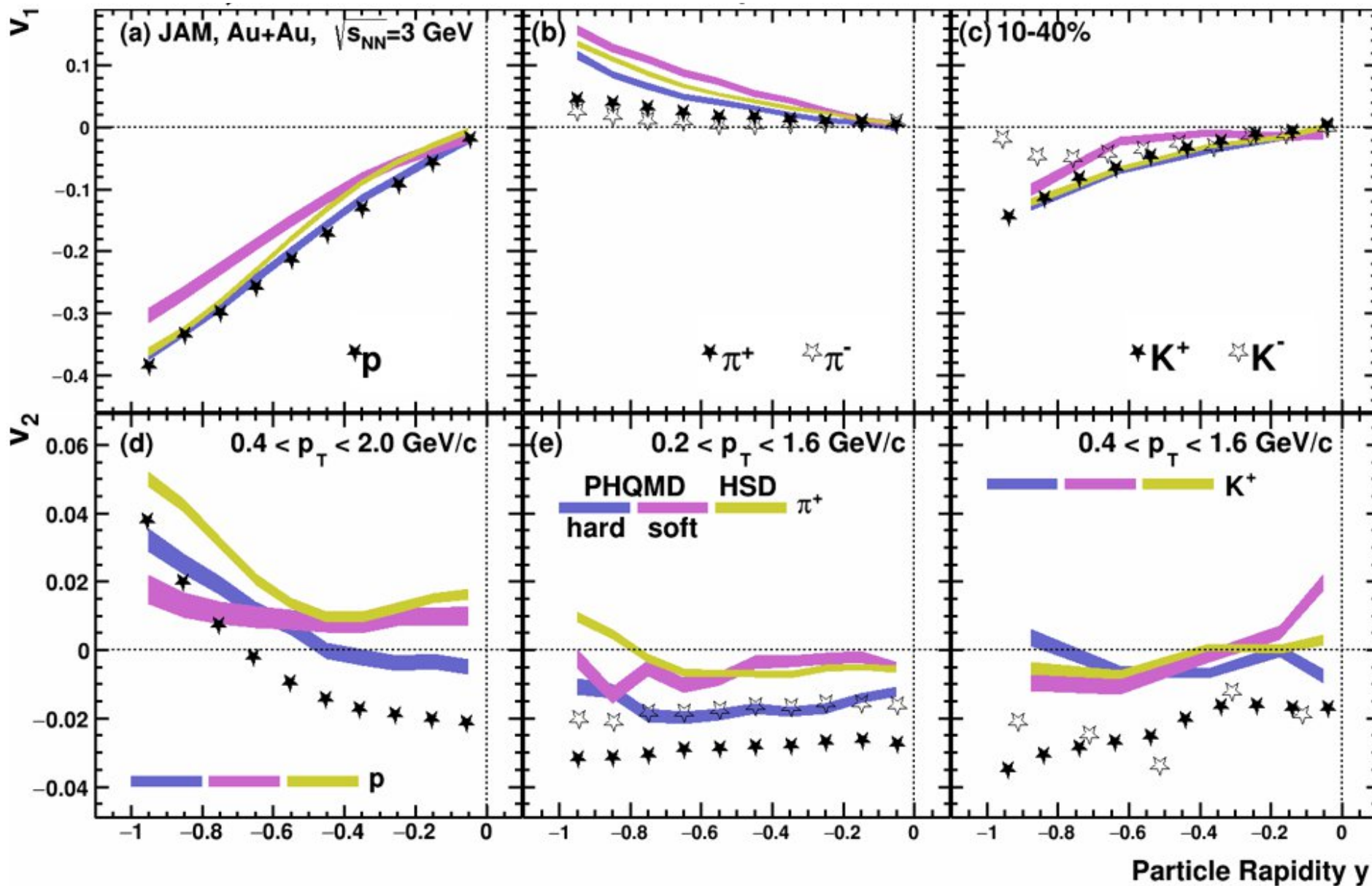
$$\frac{dN}{d\phi} \propto 1 + 2 \sum_{n=1} v_n \cos [n(\phi - \Psi_{RP})], \quad v_n = \langle \cos [n(\phi - \Psi_{RP})] \rangle$$

Anisotropic flow at FAIR/NICA energies is a delicate balance between:

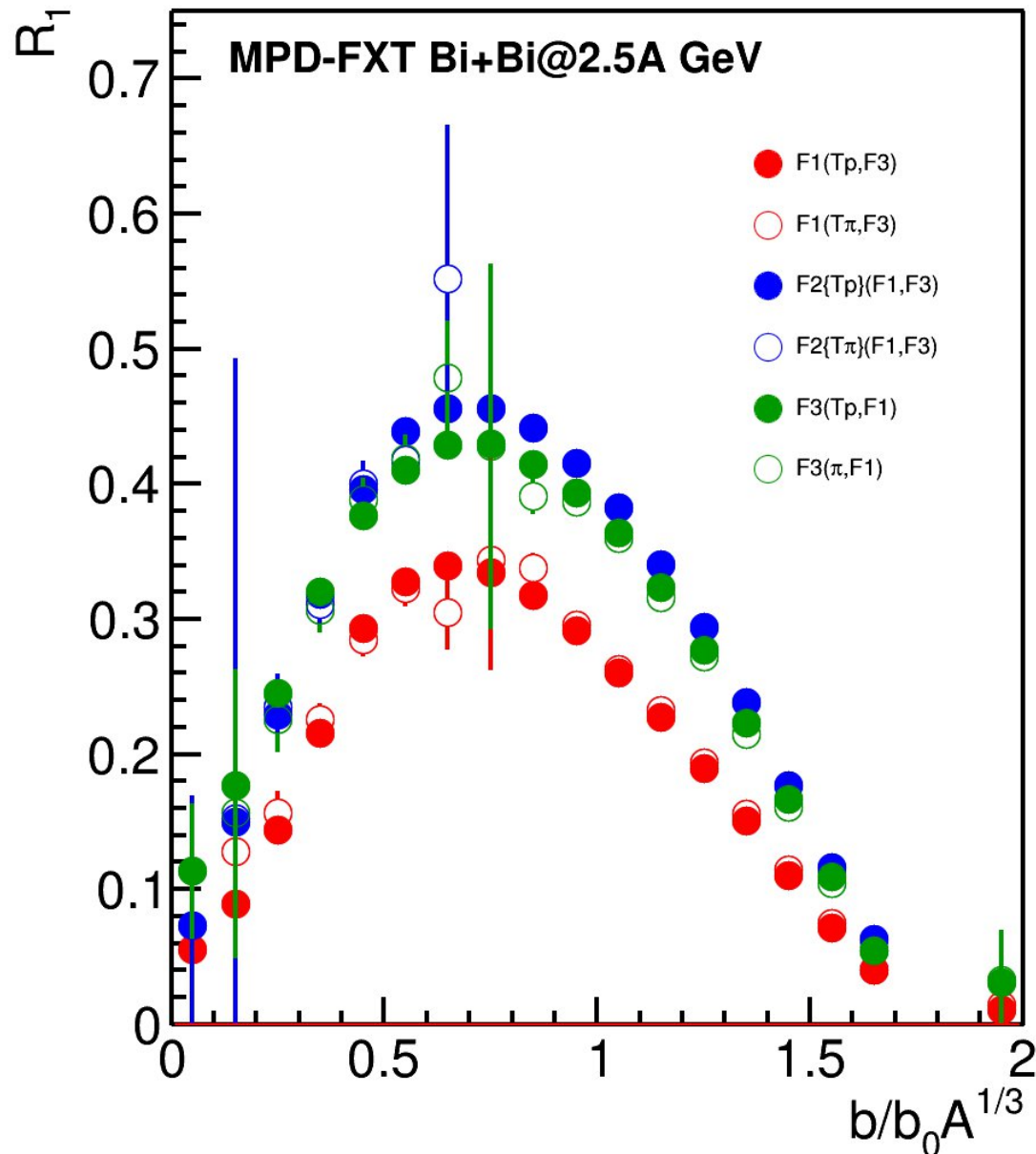
- I. The ability of pressure developed early in the reaction zone ( $t_{exp} = R/c_s$ ,  $c_s = c \sqrt{dp/d\varepsilon}$ ) and
- II. The passage time for removal of the shadowing by spectators ( $t_{pass} = 2R/\gamma_{CM}\beta_{CM}$ )



# STAR-FXT vs JAM

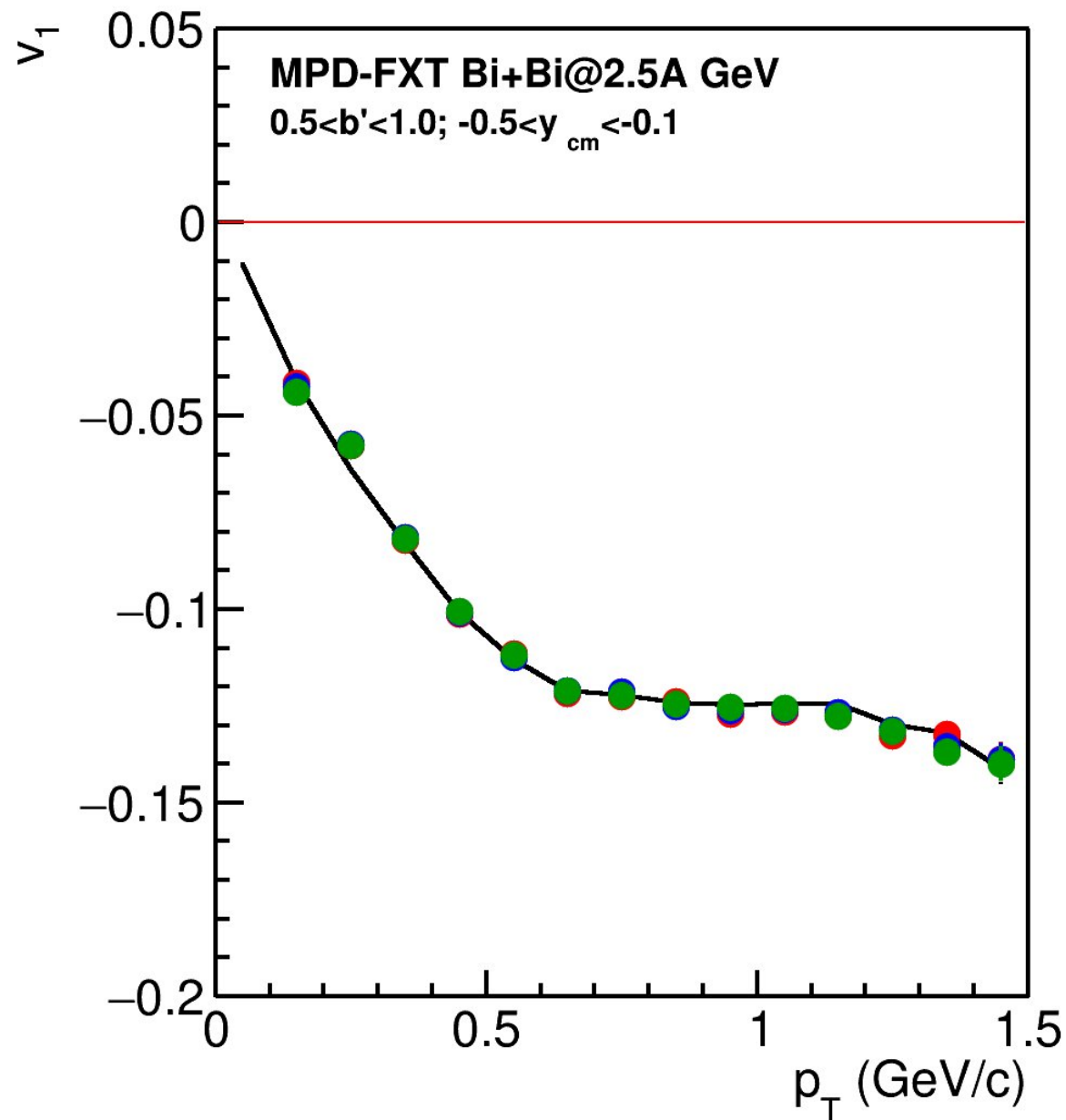
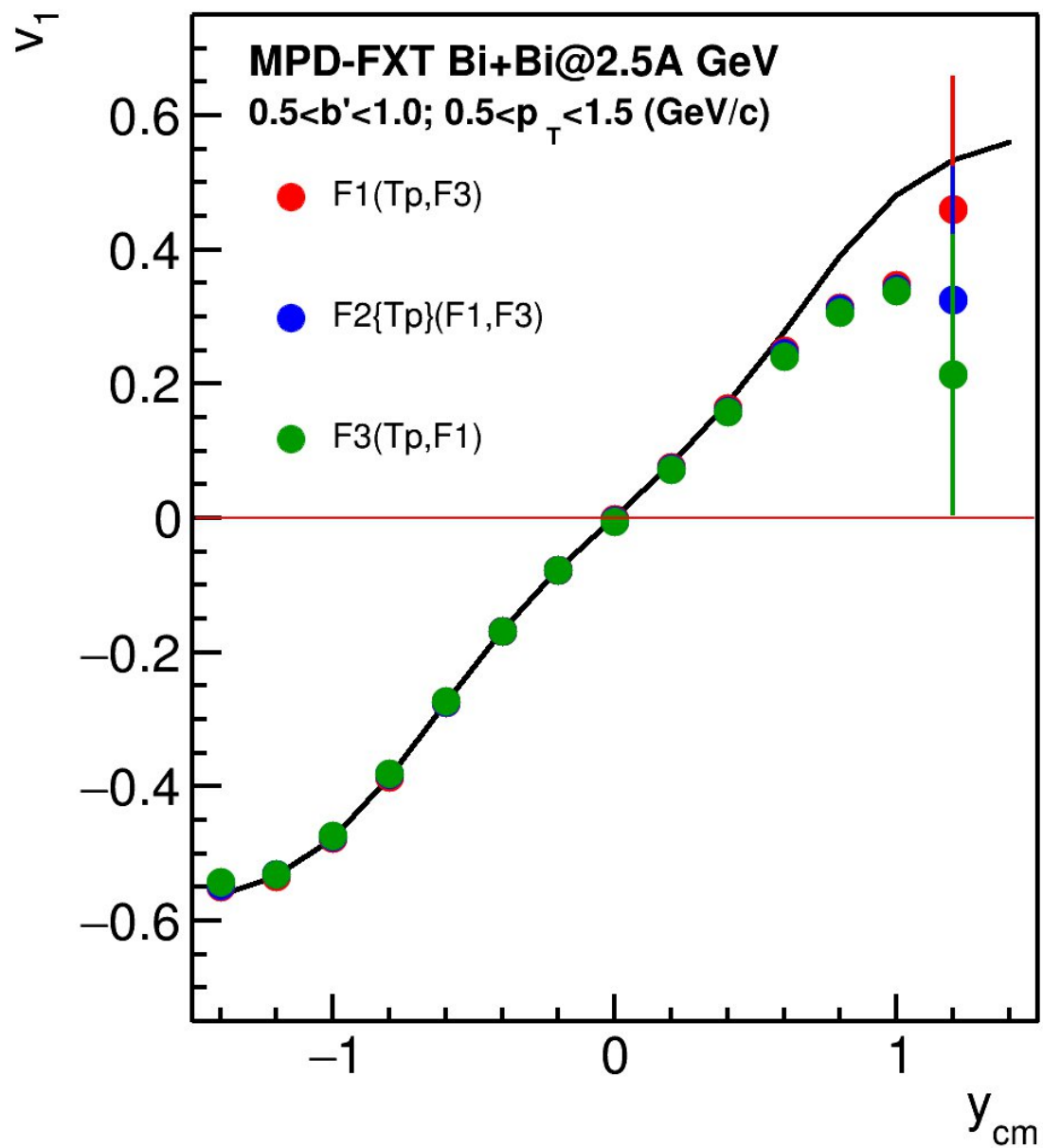


# $R_1$ for FHCaI spectator plane

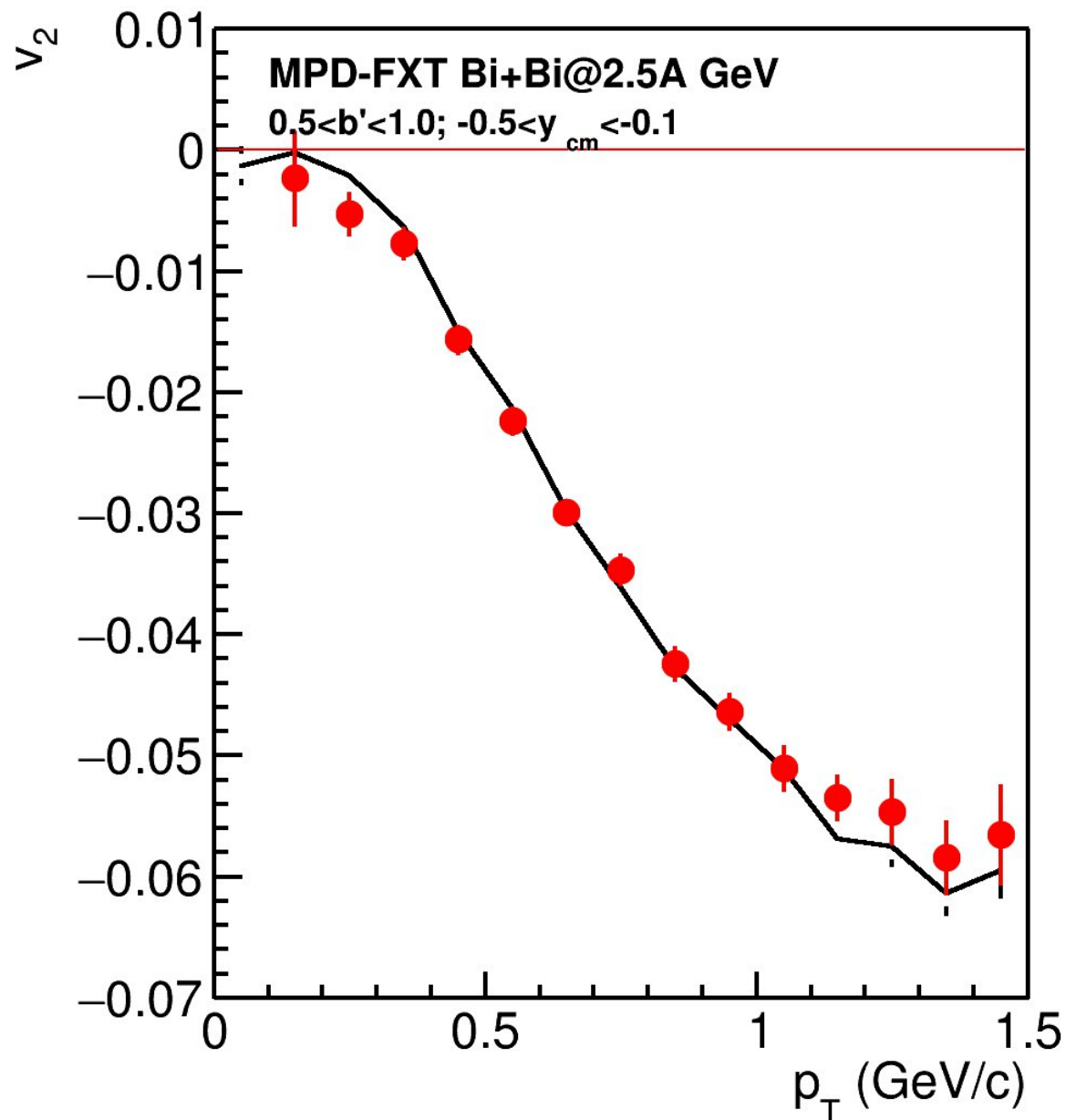
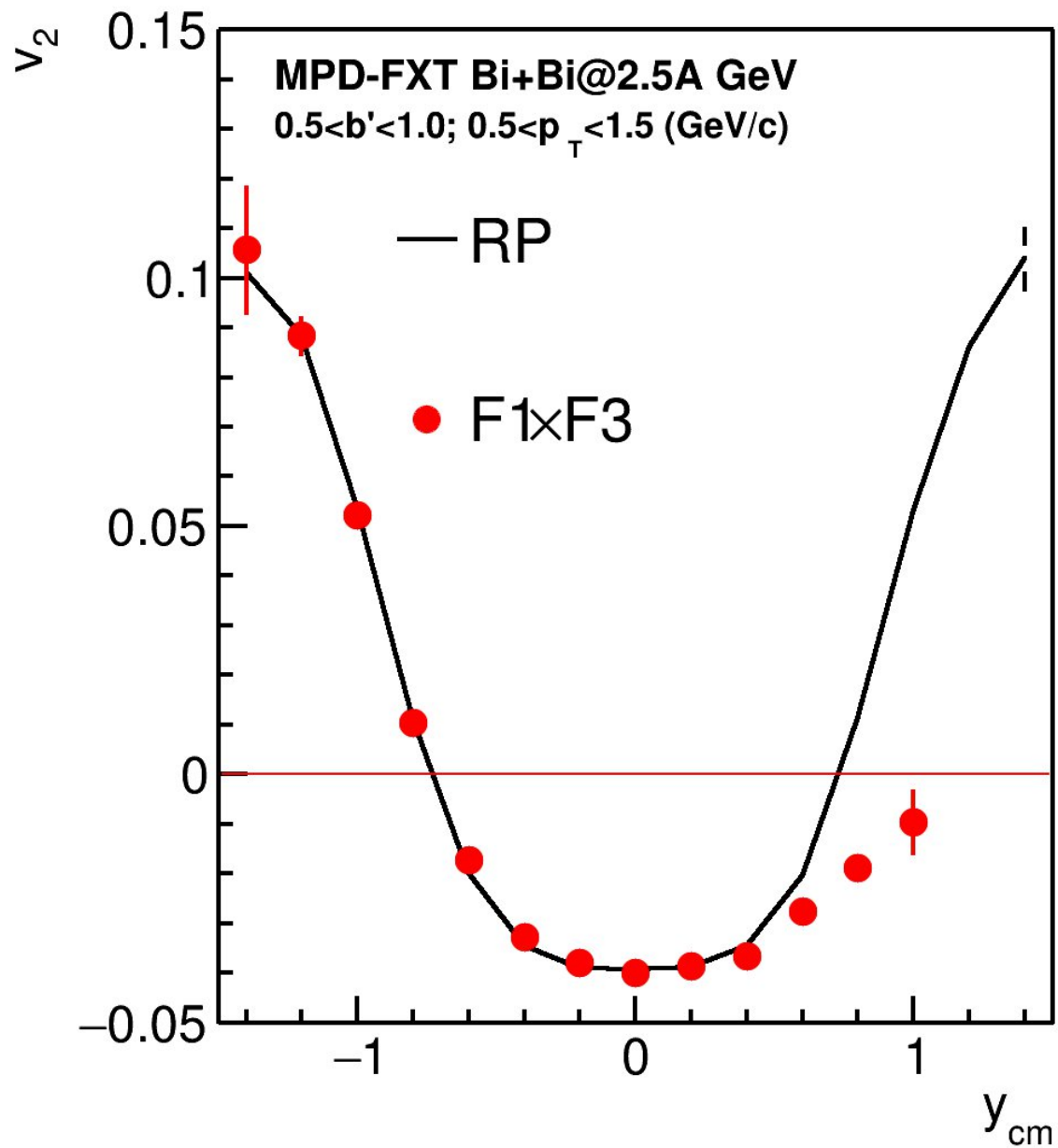


Good agreement between  $R_1$  calculated using different combinations of Q-vectors with significant rapidity separation

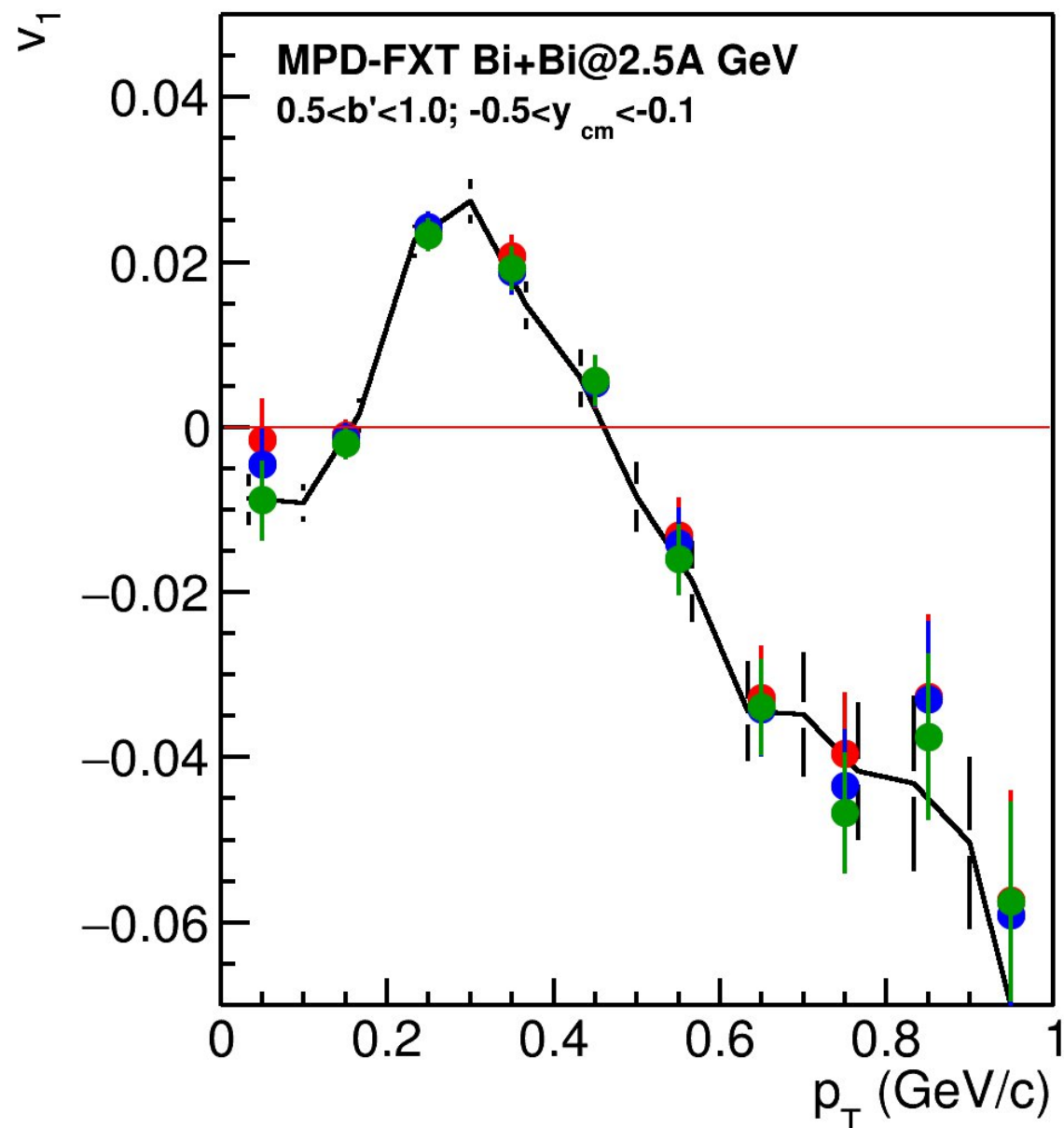
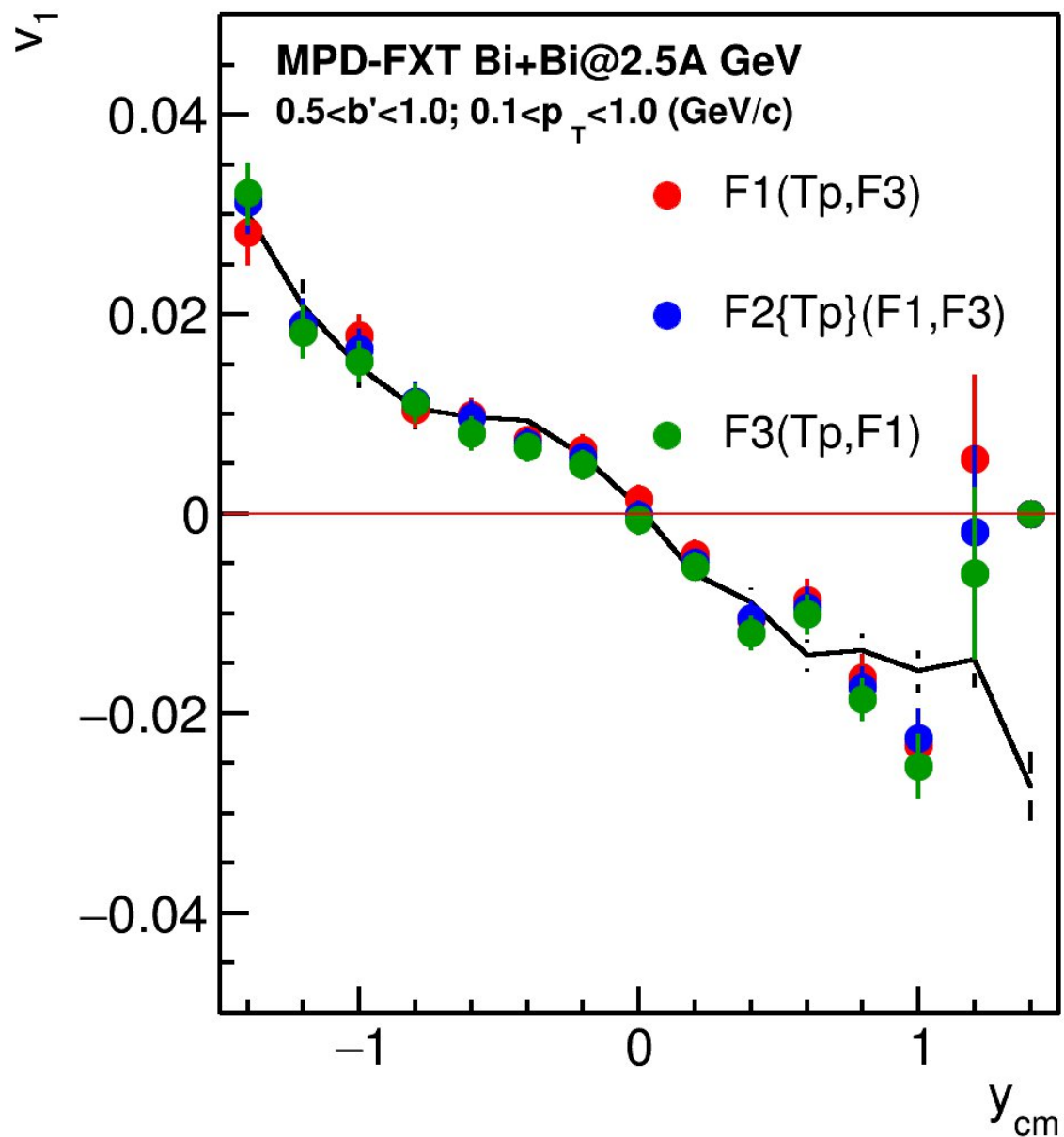
# $v_1$ for protons



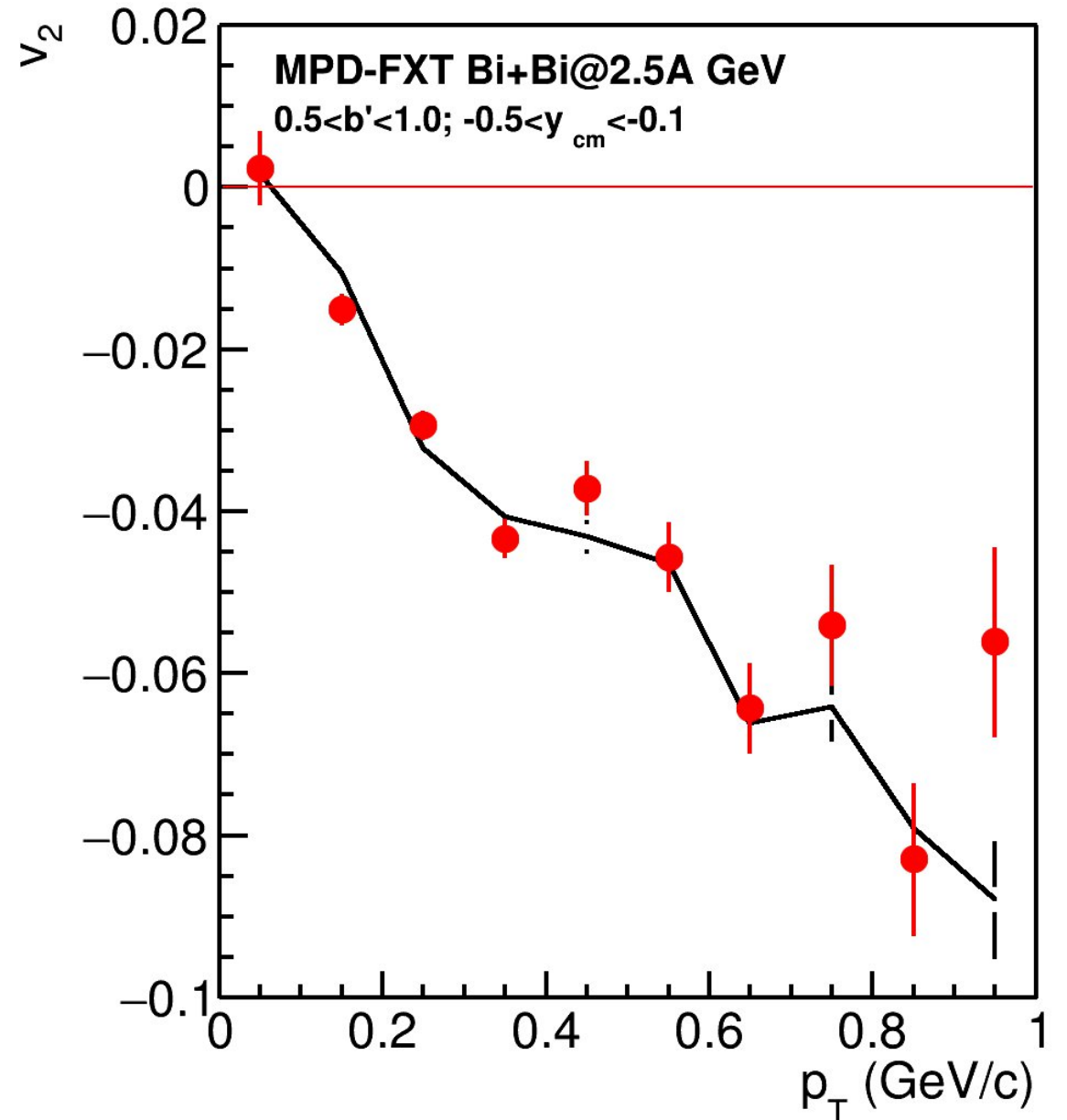
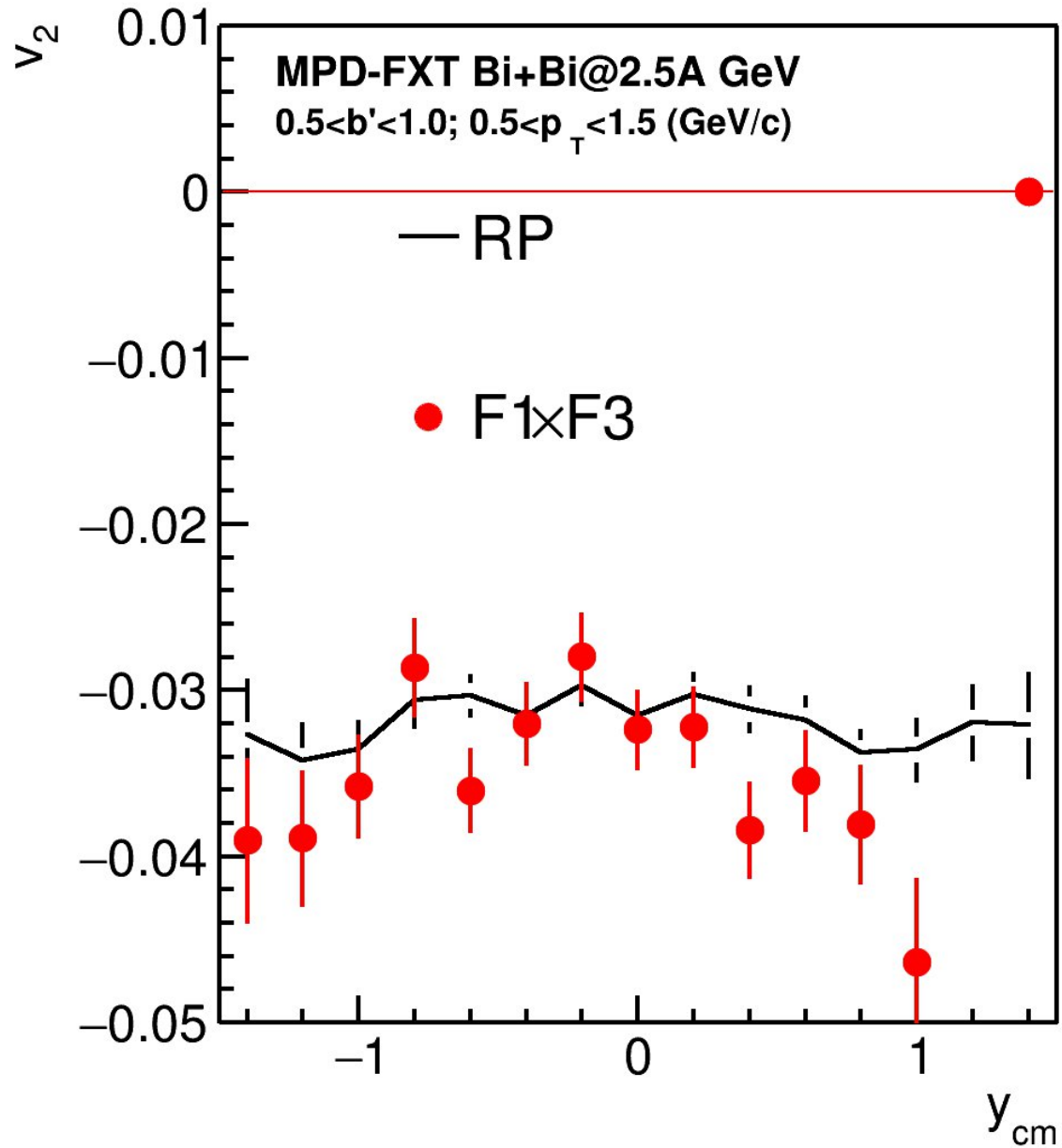
# $v_2$ for protons



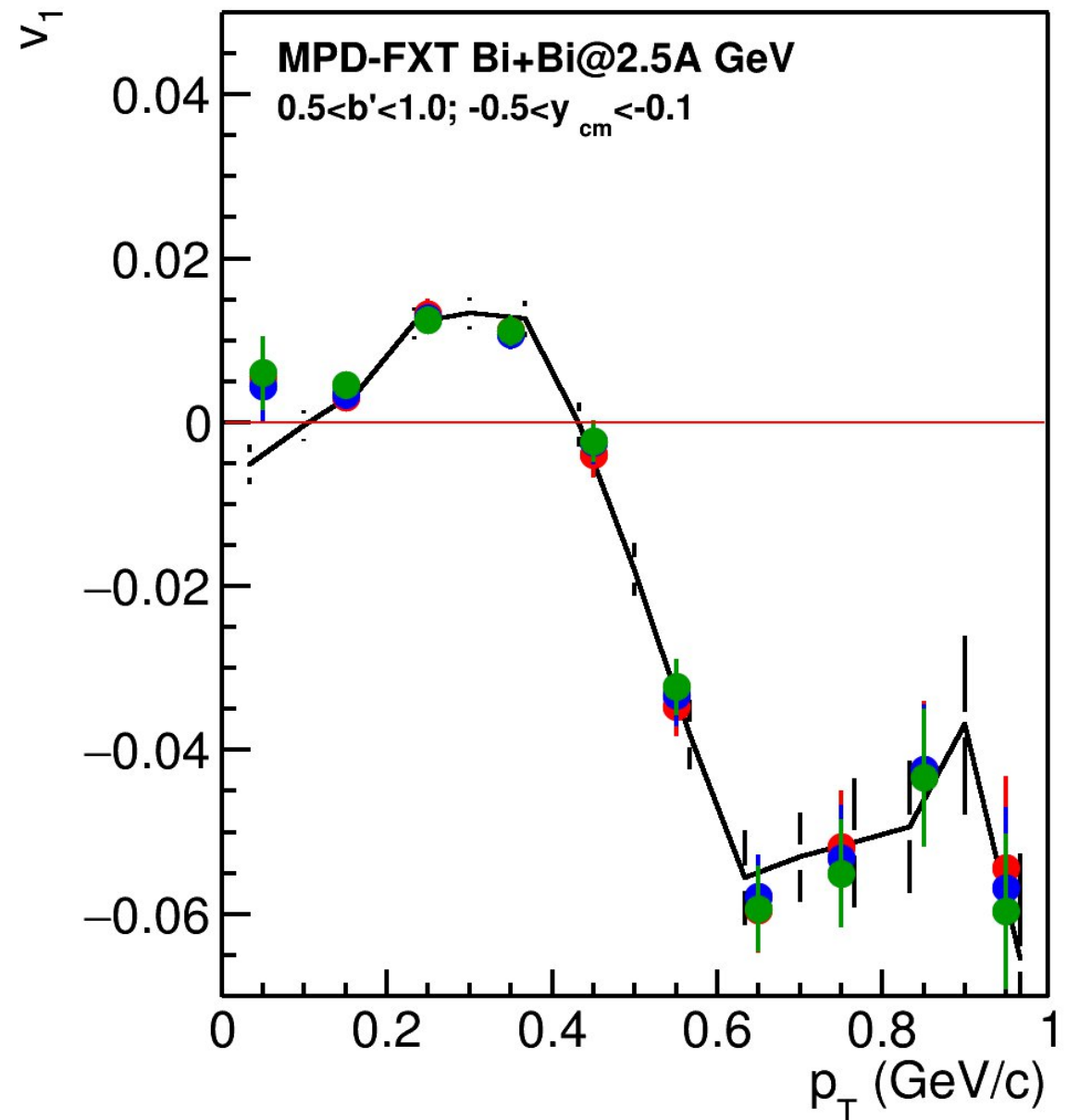
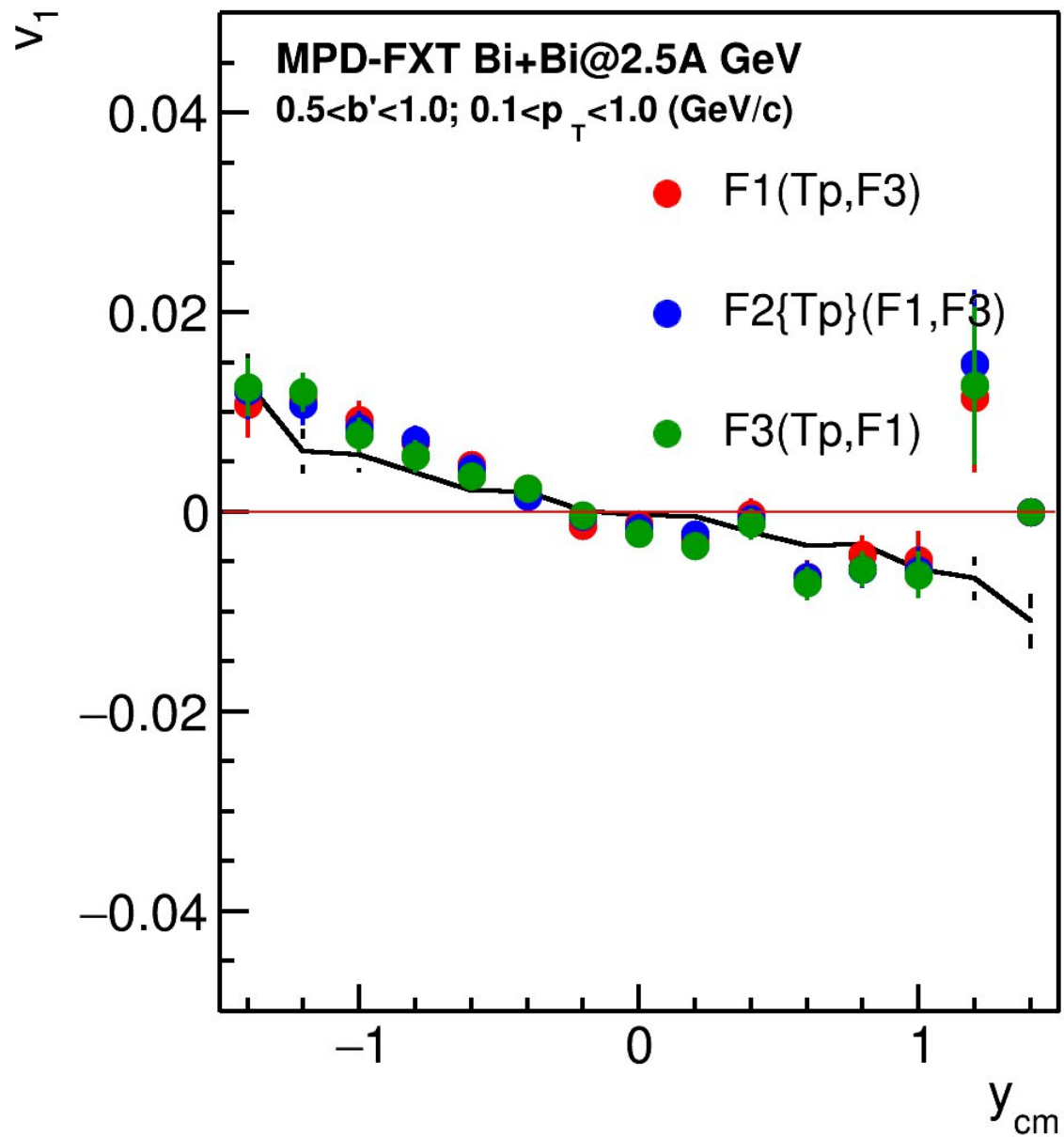
# $v_1$ for $\pi^+$



# $v_2$ for $\pi^+$



# $v_1$ for $\pi^-$



# $v_2$ for $\pi^-$

