

XXV International Baldin Seminar on High Energy Physics Problems *Relativistic Nuclear Physics & Quantum Chromodynamics* September 18 - 23, 2023, Dubna, Russia

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Non-linear processes in photonand electron- laser interactions

## Definition: non-linear processes are those that cannot be described in terms of perturbative QED

## 1. First hopes:

**QED vacuum breakdown - spontaneous electron-positron pair creation** (Schwinger effect (previously predicted by Sautzer, Heizenberg&Euler and justified by Schwinger))

### Probability of $e^+e^-$ pair production

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## 2. Current status:

Non-linear multiphoton processes with (ultra) high power laser field:

## Non-linear multiphoton processes with (ultra) high power laser field:

Non-linear Breit-Wheeler  $e^+e^-$  pair production

$$\gamma + n\gamma_L \to e^+e^-$$
(L)

Non-linear Compton high energy photon production  $e^- + n\gamma_L \rightarrow e^- + \gamma_L (L)$ 

Non-linear trident processes

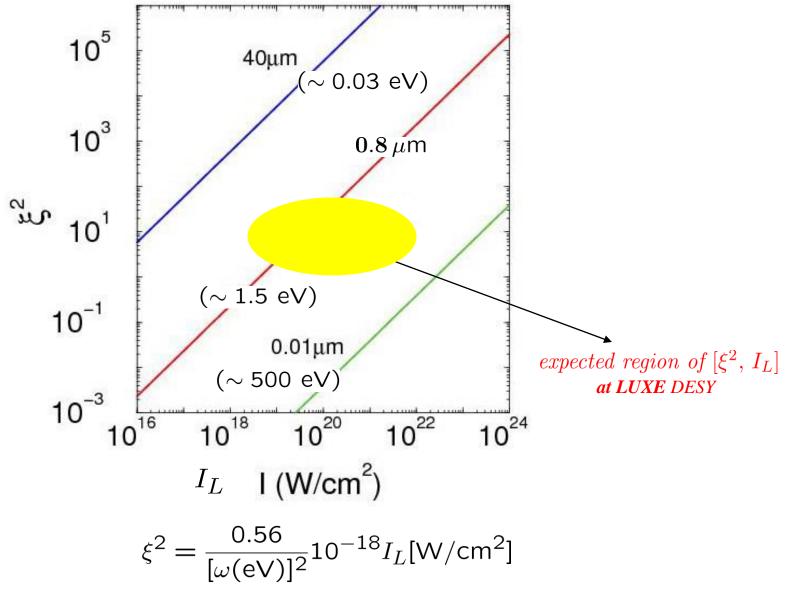
$$e^- + n\gamma_L \to e^- + e^+ + e^-$$
(L)

Reduced e.m. field intensity

$$\xi = \frac{e\mathcal{E}_{\mathrm{L}}}{m\omega_{\mathrm{L}}}$$

$$\xi^{2} = \frac{0.56 \, I [\text{W/cm}^{2}] \times 10^{-18}}{(\omega_{\text{L}} [\text{eV}])^{2}}$$

# Reduced field intensity $\xi^2$ vs. laser intensity



## Original Breit-Wheeler process $\gamma' + \gamma \rightarrow e^+e^-$

**DECEMBER 15, 1934** 

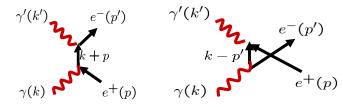
PHYSICAL REVIEW

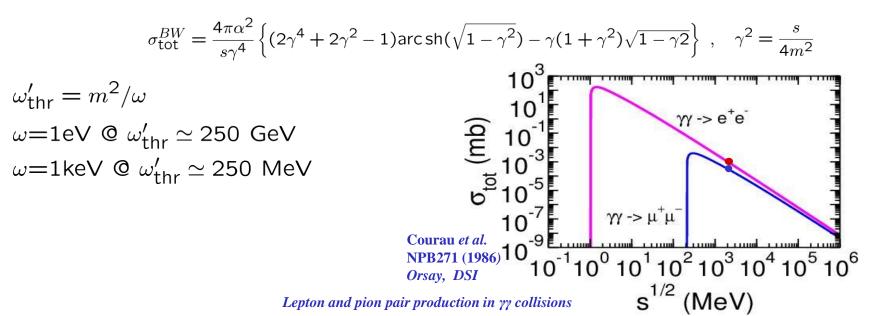
VOLUME 46

#### Collision of Two Light Quanta

G. BREIT\* AND JOHN A. WHEELER,\*\* Department of Physics, New York University (Received October 23, 1934)

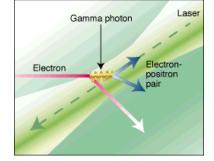
The recombination of free electrons and free positrons and its connection with the Compton effect have been treated by Dirac before the experimental discovery of the positron. In the present note are given analogous calculations for the production of positron electron pairs as a result of the collision of two light quanta. The angular distribution of the ejected pairs is calculated for different polarizations, and formulas are given for the angular distribution of photons due to recombination. The results are applied to the collision of high energy photons of cosmic radiation with the temperature radiation of interstellar space. The effect on the absorption of such quanta is found to be negligibly small.





SLAC (E-144) experiment D. Burke et al., PRL 79 (1997)  $\gamma' + L \rightarrow e^+ e^-$  generalized multi-photon process

Kinematics of BW - process



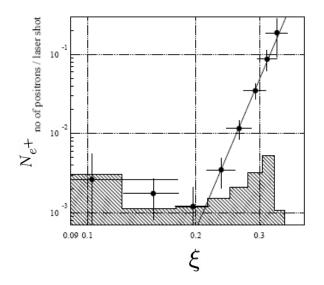
$$s_{thr}(\gamma'\gamma) = (k+k')^2 = 4\omega\omega' = 4m^2$$
  $\omega'_{thr} = \frac{m_e^2}{\omega} \simeq \frac{0.26 \cdot 10^{12} (eV^2)}{2.35 \, eV \, (SLAC)} \simeq 111 \, \text{GeV}$ 

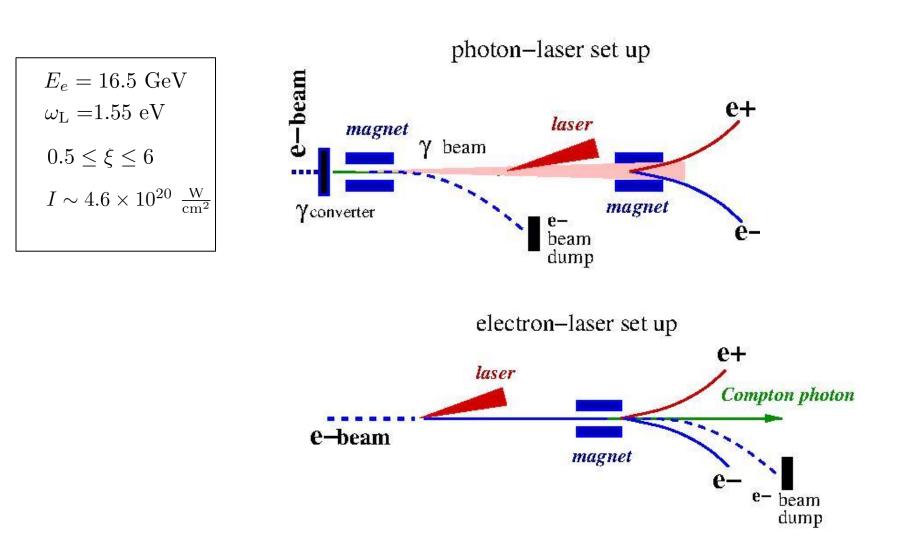
$$\omega'_{\text{SLAC}} \simeq 29 \text{ GeV} \implies \frac{\omega'_{thr}}{\omega'_{Bremmst}} \simeq 3.83$$

$$\gamma' + n\gamma \to e^+e^- \to n_{\min} \ge 4$$

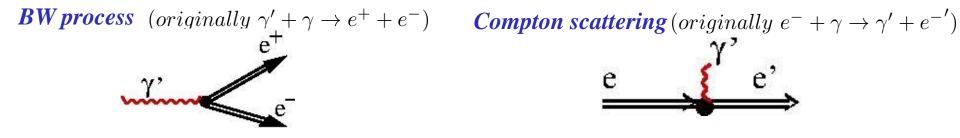
$$I \sim 2 \times 10^{18} \text{ W/cm}^2 \rightarrow 0.1 < \xi < 0.36$$

 $(\omega = 2.35 \text{ eV})$ 





## Interaction of charge particles with background field is considered in Furry picture



Volkov solution

$$\psi_p(\phi) = \left[1 + \frac{e(\gamma \cdot k)(\gamma \cdot A)}{2(k \cdot p)}\right] \frac{u_{p'}}{\sqrt{2p_0}} e^{-ip \cdot x} \exp\left[-i \int\limits_{-\infty}^{\phi} \left(\frac{e(p \cdot A)}{(k \cdot p)} - \frac{e^2 A^2}{2(k \cdot p)}\right) d\phi'\right]$$

 $A(\phi)-e.m.$  background field with  $\phi=kx$ 

$$S_{fi} = -ie \int d^4x \langle f|\gamma \cdot \varepsilon(k')|i\rangle \frac{\mathrm{e}^{-ik \cdot x}}{\sqrt{2\omega'}}$$
$$M_{fi}(kx) = \int_{-\infty}^{\infty} d\ell \mathrm{e}^{-i\ell kx} M_{fi}(\ell) \Rightarrow S_{fi} = \int_{\ell_{\min}}^{\infty} d\ell M_{fi}(\ell) \,\delta^4(k' + \ell k - p_{e^-} - p_{e^+})$$

 $l\omega$  is the energy of the pulse involved into process

Non-linear BW process  $\gamma' + L \rightarrow e^+ + e^-$ 

A.T., Kämpfer, PRL **108**(2013), EPJD **68** (2014), EPJD **74** (2020)

 $A = (0, \mathbf{A}), \ \mathbf{A}(\phi) = f(\phi) [\mathbf{a}\cos(\phi)] \qquad \mathcal{E} = -\frac{\partial \mathbf{A}}{\partial t} \qquad \phi = kx \text{ - invariant phase,}$ *e.m. potential:* with  $\mathbf{a} = \mathbf{e}_0 \frac{\xi m}{2}$  $f(\phi)$  - envelope function  $\sum_{i} |S_{fi}|^2 \propto \int dl |M(\mathbf{e}_0, \mathbf{e}', k, k', \xi, f(\phi))|^2 (2\pi)^4 \delta(\ell k + k' - q - q')$  $W \sim \int dl w_l = \int dl \left[ \xi^2 u(\widetilde{A}_1(l)^2 - \widetilde{A}_0(l)\widetilde{A}_2(l)) - [|p \cdot e_i''\widetilde{A}_0(l) + ea \cdot e_i''\widetilde{A}_1(l)|^2]/m^2 \right] ,$  $e'' = e' - k'(k \cdot e')/k \cdot k'$  $\widetilde{A}_m(\ell) = \frac{1}{2\pi} \int d\phi f^m(\phi) \cos^m(\phi) e^{i\ell\phi - i\mathcal{P}^{(lin)}(\phi)}$  $\mathcal{P}^{(lin)}(\phi) = \tilde{\alpha}(\phi) - \tilde{\beta}(\phi) ,$  $\tilde{\alpha}(\phi) = z \cos \phi_e \int^{\phi} d\phi' f(\phi') \cos(\phi') ,$  $\tilde{\beta}(\phi) = \frac{u\ell\xi^2}{u_\ell} \int^{\phi} d\phi' f^2(\phi') \cos^2(\phi') .$  $u = (kk')/4(kp)(kp'), u_{\ell} = \ell/\zeta, z = 2\ell\xi((u/u_{\ell})(1-(u/u_{\ell}))^{1/2})$ 

$$w(\ell) = w_{\perp}(\ell) + w_{\parallel}(\ell)$$

$$w_{\parallel}(\ell) = \xi^{2}(u-1) \left( |\tilde{A}_{1}(\ell)|^{2} - [\tilde{A}_{0}(\ell)\tilde{A}_{2}^{*}(\ell) \right) + (1+\tau^{2}) |\tilde{A}_{0}(\ell)|^{2}, \qquad perpendicular polarization \\ \psi_{\perp}(\ell) = \xi^{2} u \left( \tilde{A}_{1}(\ell)|^{2} - [\tilde{A}_{0}(\ell)\tilde{A}_{2}^{*}(\ell) \right) - \tau^{2} |\tilde{A}_{0}(\ell)|^{2}, \qquad \tau^{2} = (u_{\ell}/u-1) \sin^{2}\varphi_{c}.$$

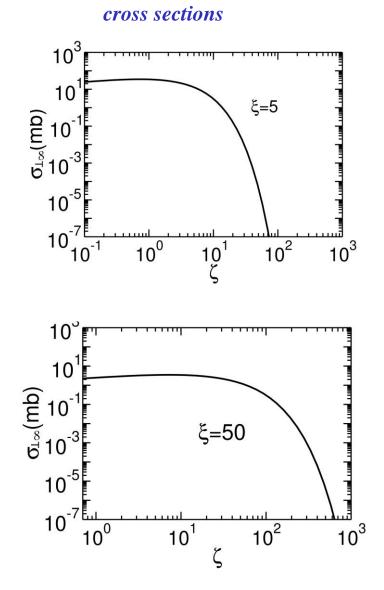
$$f_{\perp} = \sigma_{\perp} + \sigma_{\parallel}$$

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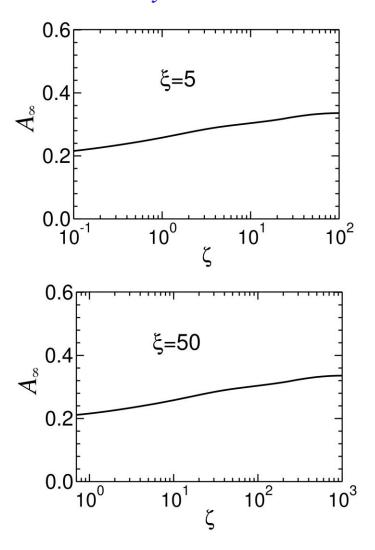
$$f_{\perp} = \sigma_{\perp} + \sigma_{\perp} + \sigma_{\perp}$$

$$f_{\perp} = \sigma_{\perp} + \sigma_{\perp}$$

an analog of "cumulative number"



asymmetries



*Non-linear Compton scattering*  $e^- + L \rightarrow \gamma' + e^{-\prime}$ 

A.T., Kämpfer, Phys.Rev.A103(2021), Phys.Part.Nucl.47(2016) A.T., arXiv 2307.00621

 $M = \sum_{a=1,2} e'^*_a M_a$ ,  $e'_{1,2}$  are the polarization vectors of  $\gamma'$ (linear polarization)

$$\rho_{ab}^{f} = \frac{M_{a}M_{b}^{*}}{\sum_{a} |M_{a}|^{2}} \text{ spin-density matrix}$$
  
$$\xi_{3}^{f} = \operatorname{Sp}(\rho\sigma_{3})) = \frac{|M_{1}|^{2} - |M_{2}|^{2}}{|M_{1}|^{2} + |M_{2}|^{2}} = \mathcal{A}$$

Stoks parameter for intrinsic spin of recoil photon  $~\gamma^{\prime}$ 

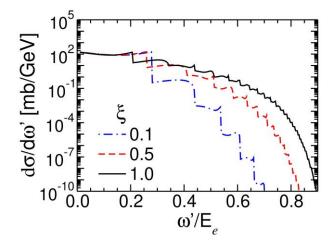
spin asymmetry  $\mathcal{A}$  (or  $\xi_3^f$ ) shows direction and degree of  $\gamma'$  spin polarization

$$\mathbf{e}_{1}^{\prime} = \frac{[\mathbf{k}, \mathbf{k}^{\prime}]}{|[\mathbf{k}, \mathbf{k}^{\prime}]|}, \qquad \mathbf{e}_{2}^{\prime} = \frac{|\mathbf{k}^{\prime}, \mathbf{e}_{1}^{\prime}|}{|\mathbf{k}^{\prime}|}, \qquad \longrightarrow \qquad \begin{array}{c} \mathbf{e}_{1}^{\prime} = -\mathbf{x}\sin\varphi + \mathbf{y}\cos\varphi & \text{A.I.Ahiezer, V.B. Berestetsky. QED} \\ \mathbf{e}_{2}^{\prime} = -\mathbf{x}\cos\theta\cos\varphi - \mathbf{y}\cos\theta\sin\varphi + \mathbf{z}\sin\theta & \text{.} \end{array}$$

$$d^{2}\sigma_{1} = d\varphi d\omega' \mathcal{K} \sum_{l=1}^{\infty} \left[\xi^{2} \widetilde{A}_{1}^{2} \sin^{2}\varphi + \xi^{2} \frac{u^{2}}{4(1+u)} \left(\widetilde{A}_{1}^{2} - \widetilde{A}_{0}\widetilde{A}_{2}\right)\right], \quad where \quad \mathcal{K} = \frac{4\alpha^{2}}{\xi\chi m^{2} E_{e}}$$
$$d^{2}\sigma_{2} = d\varphi d\omega' \mathcal{K} \sum_{\ell=1}^{\infty} \left[-\widetilde{A}_{0}^{2} - \xi^{2} \widetilde{A}_{1}^{2} \sin^{2}\varphi + \xi^{2} \left(1 + \frac{u^{2}}{4(1+u)} \left(\widetilde{A}_{1}^{2} - \widetilde{A}_{0}\widetilde{A}_{2}\right)\right)\right].$$

$$\begin{split} d^2\sigma &= d^2\sigma_1 + d^2\sigma_2 & \text{average asymmetry} \\ \mathcal{A}(\varphi, \omega') &= \frac{d^2\sigma_1 - d^2\sigma_2}{d^2\sigma} \ . & \langle \mathcal{A}(\omega') \rangle_{\varphi} = \frac{1}{2\pi} \int_{0}^{2\pi} d\varphi \, \mathcal{A}(\varphi, \omega') \ , \end{split}$$

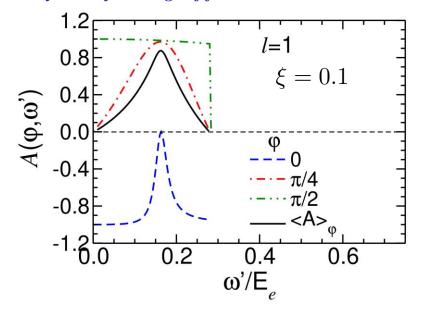
#### unpolarized cross sections



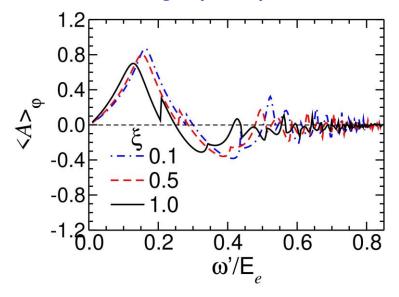
$$E_e = 16.5 \text{ GeV}$$

$$\omega_{\max}'(\ell) = \frac{2\omega E_e}{2E_e |\cos\theta + \ell\omega(1 - \cos\theta)|}$$

asymmetry in range of first harmonic with *l*=1

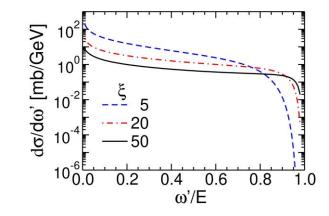


average asymmetry with all l



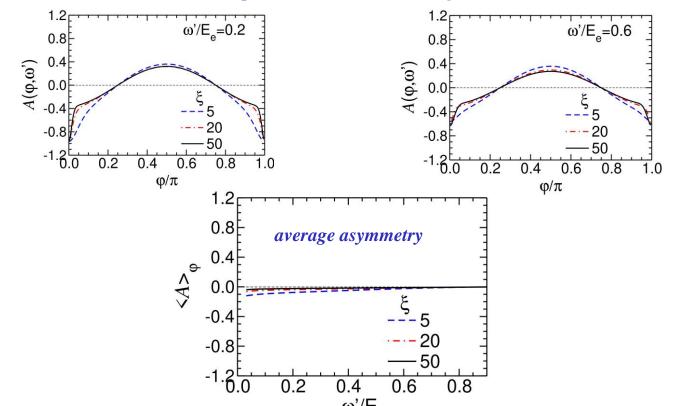
large and ultra-large field intensity  $\xi \gg 1$ 

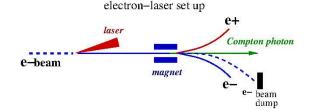
#### unpolarized cross sections



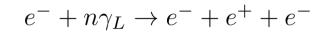
 $E_e = 16.5 \text{ GeV}$ 

dependence on azimuthal angle  $\varphi$ 



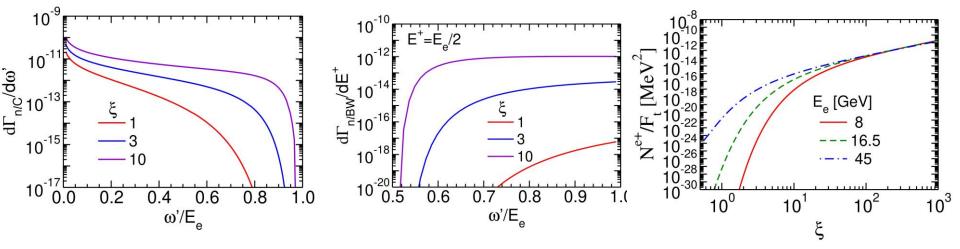


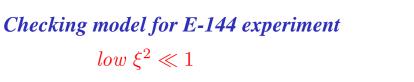
## Non-linear trident process

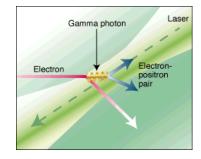


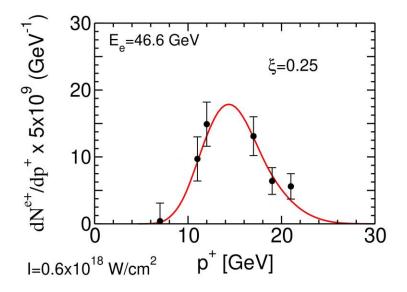
A.I. T., Acosta, Kampfer, PRA(104)2022 EPJST 230 (2021)

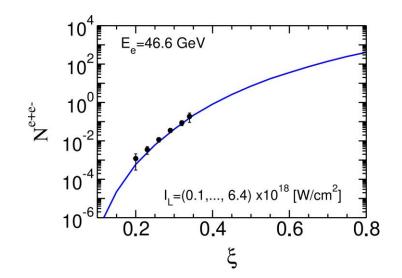
$$\frac{dN^{e^+}}{dE^+} = F_t \int_{E^+}^{E_e} d\omega' \frac{d\Gamma_{n\ell C}(\omega')}{d\omega'} \frac{d\Gamma_{n\ell BW}(\omega', E^+)}{dE^+}$$
$$F_t = \frac{1}{2} \left(\frac{2\pi}{\omega_{\rm L}}\right)^2 N_e \simeq 8.2 \times 10^{21} \text{ MeV}^2$$
$$N_e = 10^9$$











# **Summary**

1. We made predictions for unpolarized cross sections of non-linear BW and Compton processes in a wide region of e.m. field intensity

2. The main patterns of spin observables in BW and Compton scattering have been studied

3. The yield of electron-positron pairs at LUXE kinematic have been performed

4. The model was checked successfully for SLAC E-144 experiment



# THE END

Thank you for attention !

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