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Confinement effects in heterogeneous nucleation of droplets on solid particles

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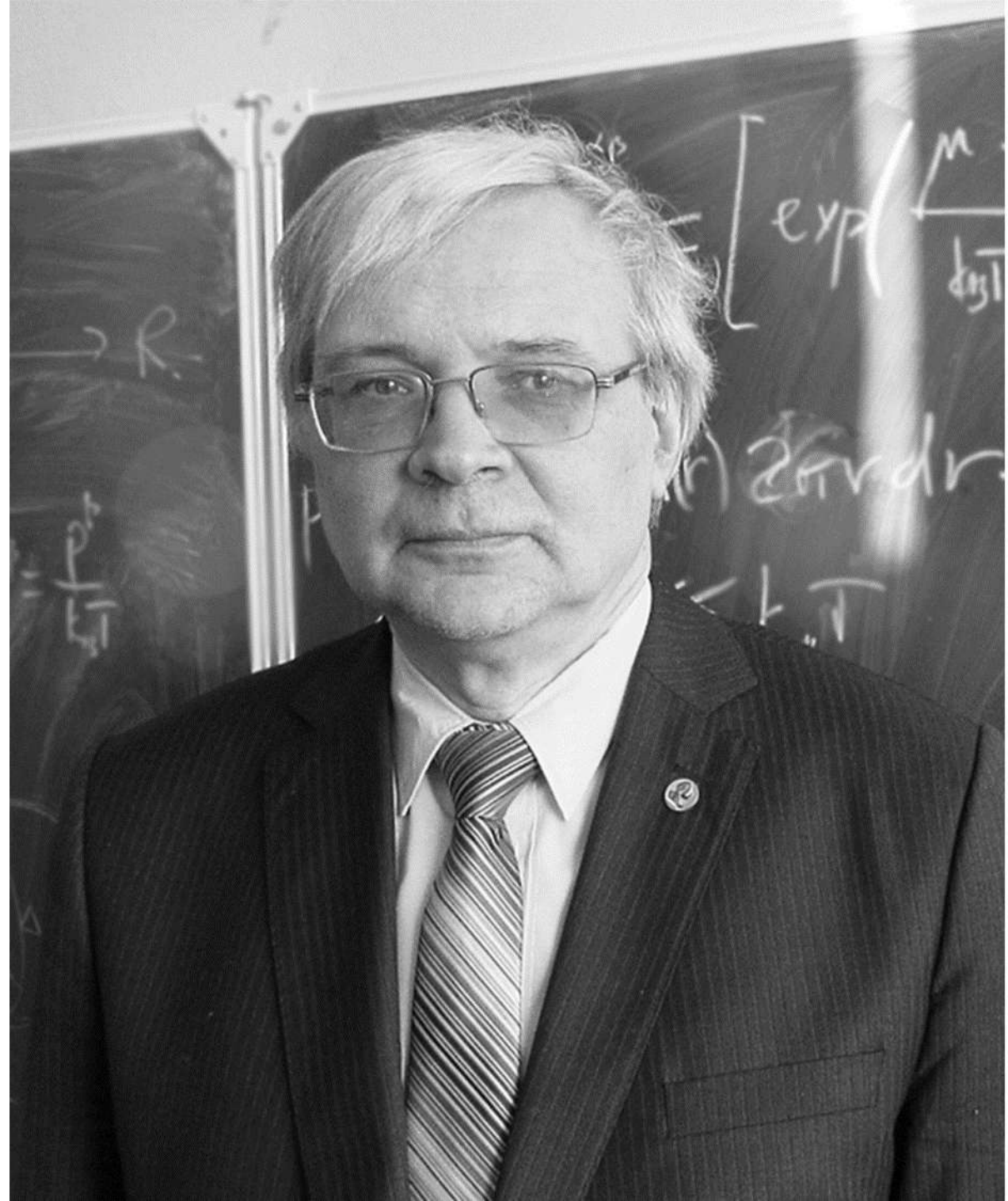


To the memory of

Dr.Sc., Corr. Member of RAS

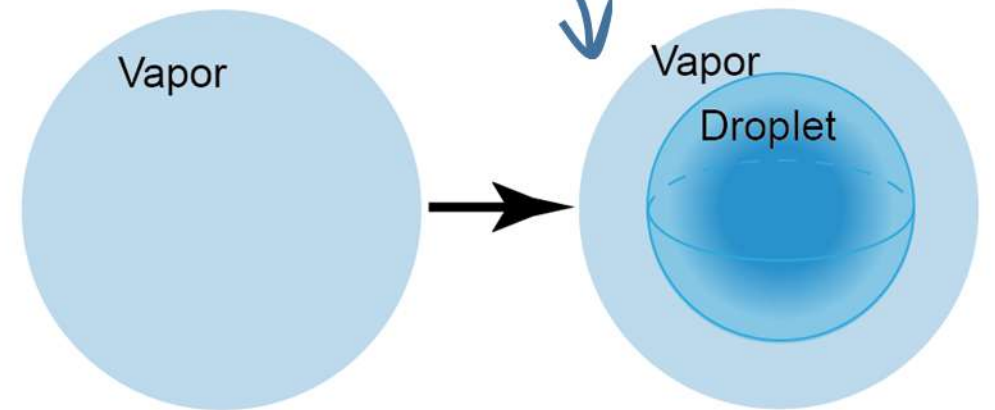
Alexander Kimovich Shchekin

28.07.1957 – 12.07.2023

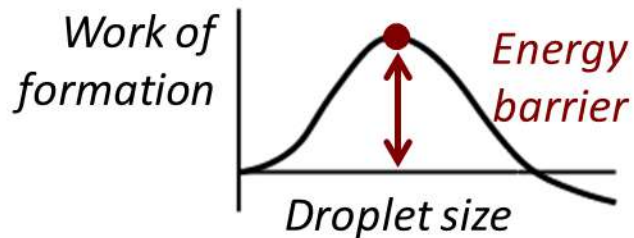


Confinement effects
in heterogeneous nucleation of droplets
on solid particles

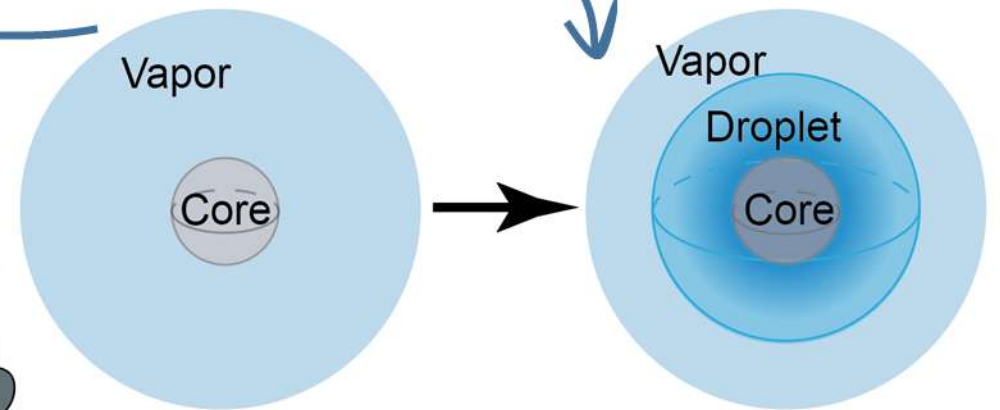
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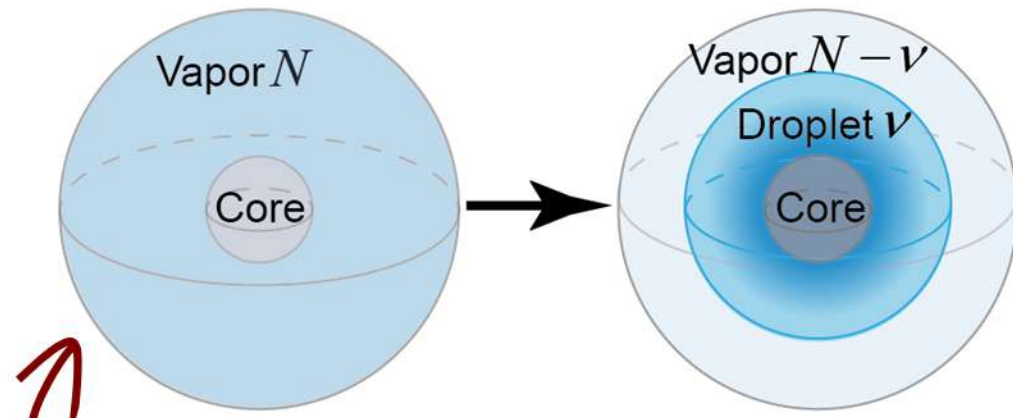


Confinement effects in heterogeneous nucleation of droplets on solid particles



That's what
happens
in nature!



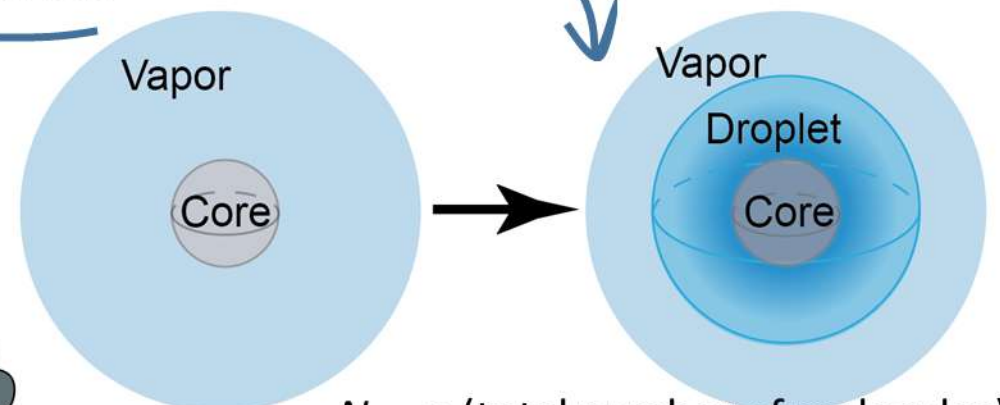


$N = \text{finite}$
 A closed system
 (Canonical Ensemble)

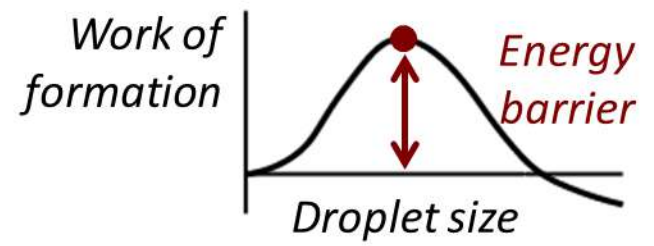
- in computer simulation
- in reality (in a small pore)

Confinement effects

in heterogeneous nucleation of droplets
on solid particles



$N = \infty$ (total number of molecules)
 An open system
 (Grand Canonical Ensemble)



That's what happens in nature!

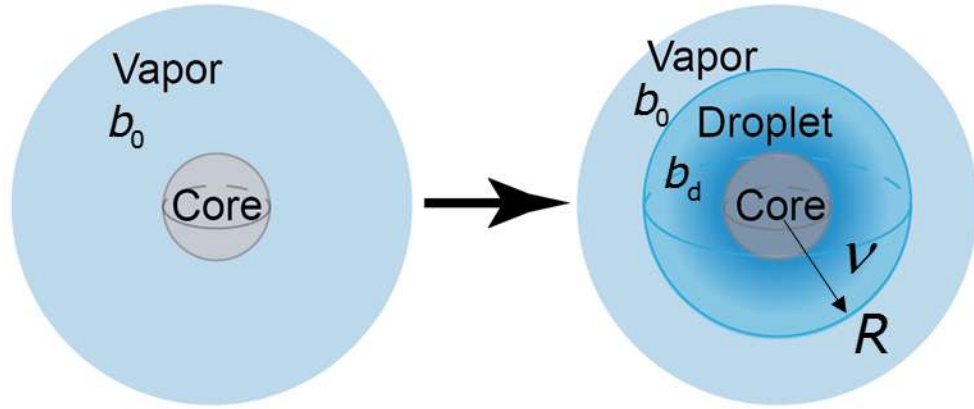


Plan

1. Thermodynamic “macro” description
2. Classical density functional theory “micro” calculations

Thermodynamic “macro” description

Open systems (Grand Canonical Ensemble)



$b_d(v) = b_0$ - at equilibrium

$$b_d \approx \frac{2\gamma}{k_B T n^\alpha R} - \frac{K \exp\left(-\frac{R - R_{\text{core}}}{\lambda}\right) R_{\text{core}}^2}{k_B T n^\alpha R^2}$$

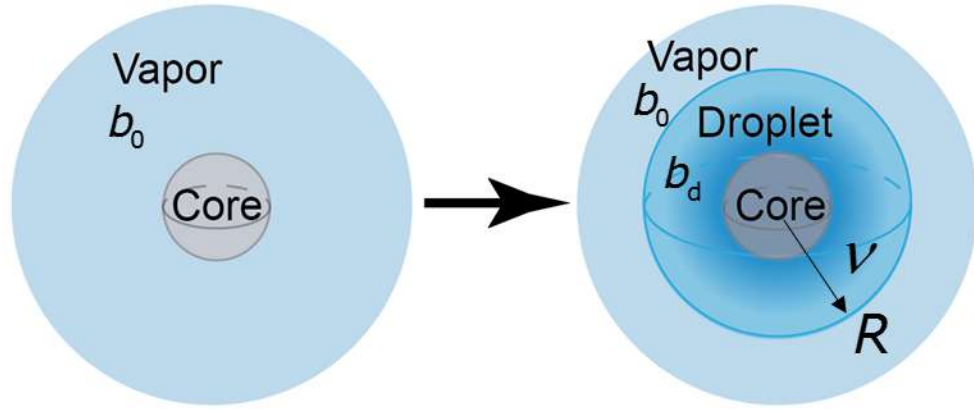
Surface tension Disjoining pressure

$$b \equiv \frac{\mu - \mu_\infty}{k_B T}$$

b_0 - initial vapor chemical potential

b_d - chemical potential in the droplet

Open systems (Grand Canonical Ensemble)



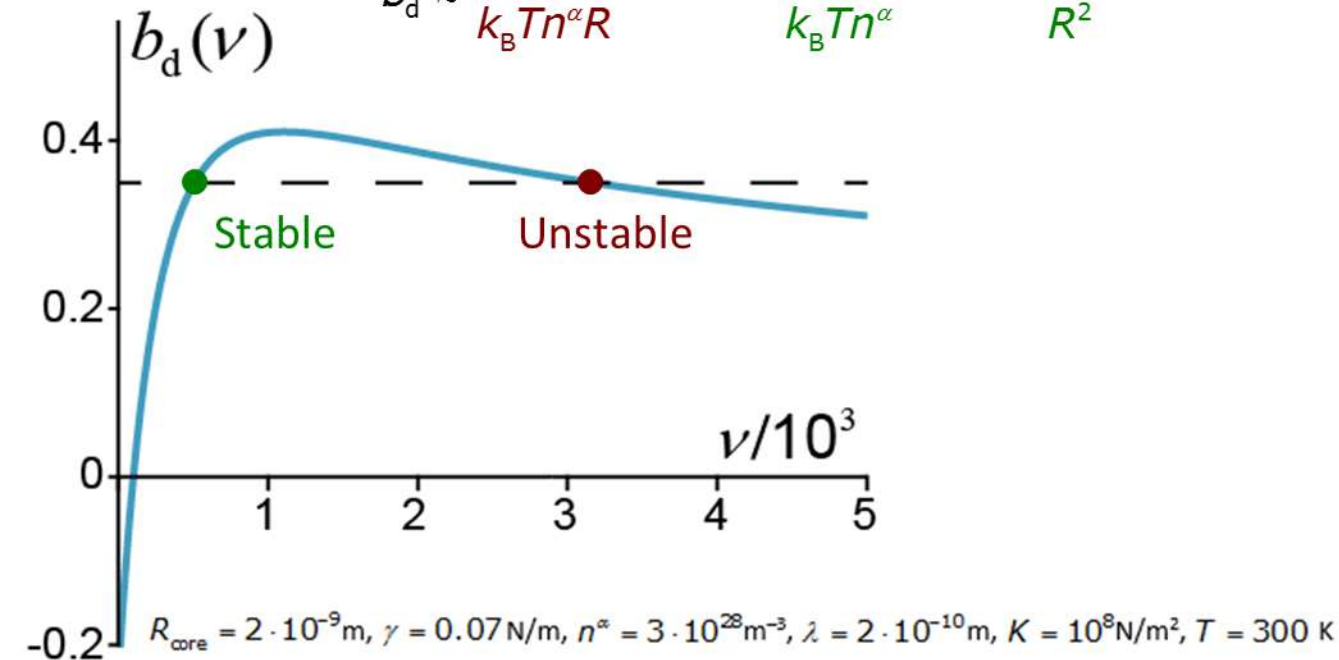
$$b \equiv \frac{\mu - \mu_\infty}{k_B T}$$

b_0 – initial vapor chemical potential

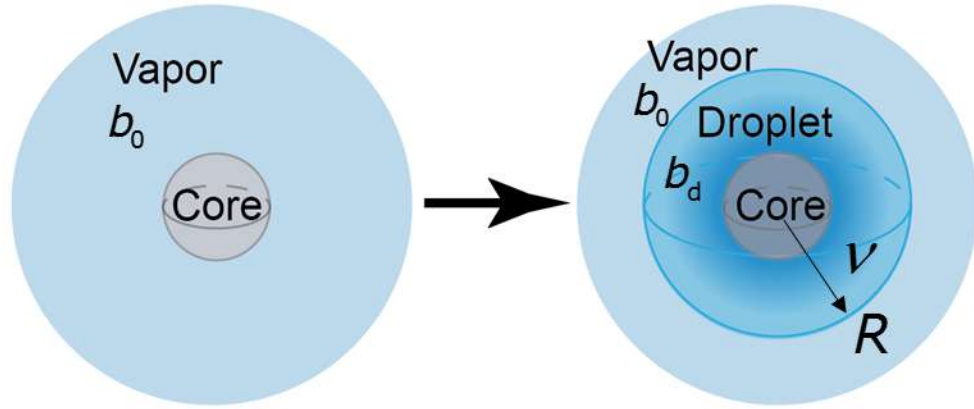
b_d – chemical potential in the droplet

$b_d(\nu) = b_0$ - at equilibrium

$$b_d \approx \frac{2\gamma}{k_B T n^\alpha R} - \frac{K \exp\left(-\frac{R - R_{\text{core}}}{\lambda}\right) R_{\text{core}}^2}{k_B T n^\alpha R^2}$$



Open systems (Grand Canonical Ensemble)

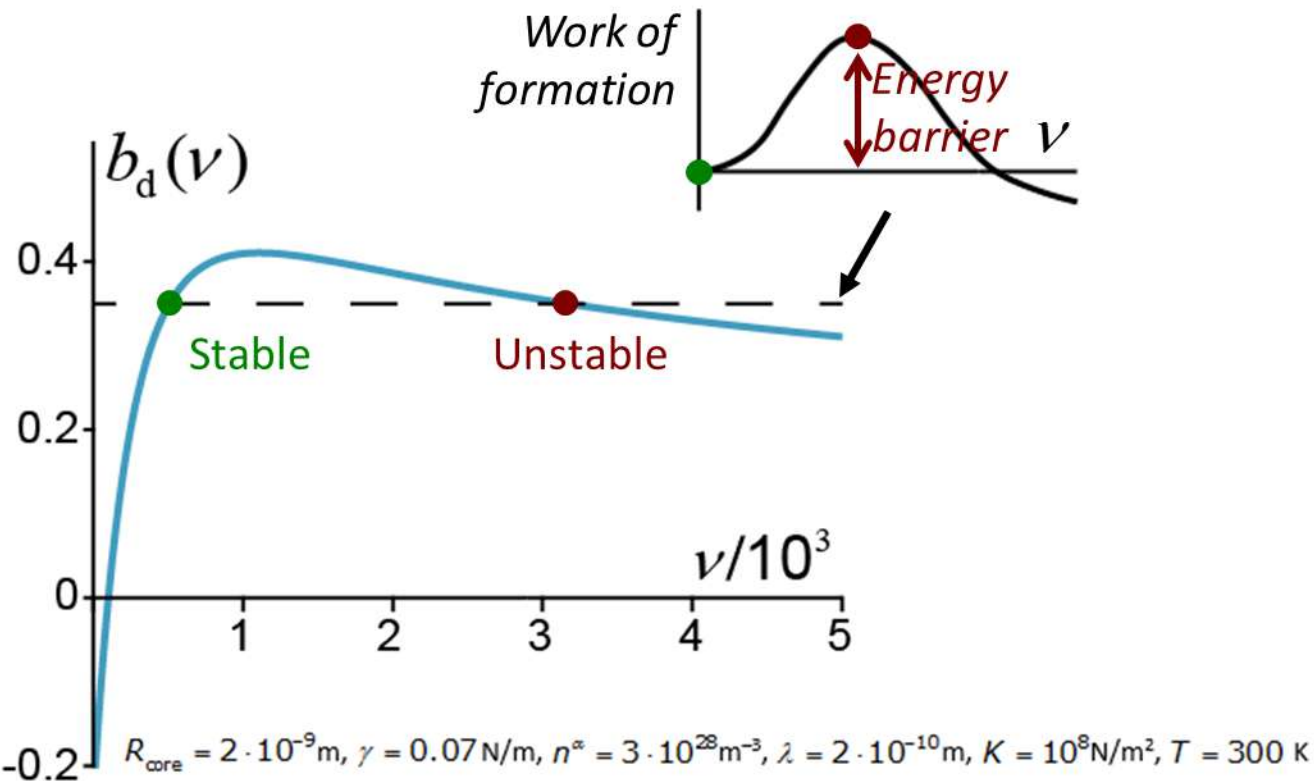


$$b \equiv \frac{\mu - \mu_\infty}{k_B T}$$

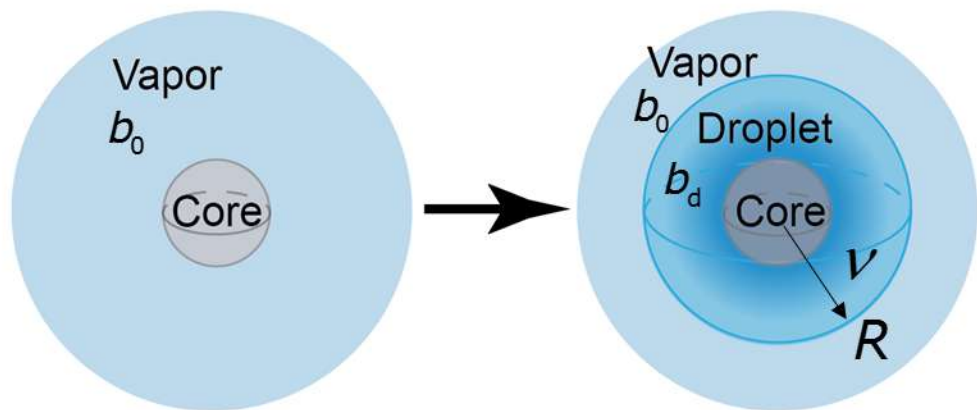
b_0 – initial vapor chemical potential

b_d – chemical potential in the droplet

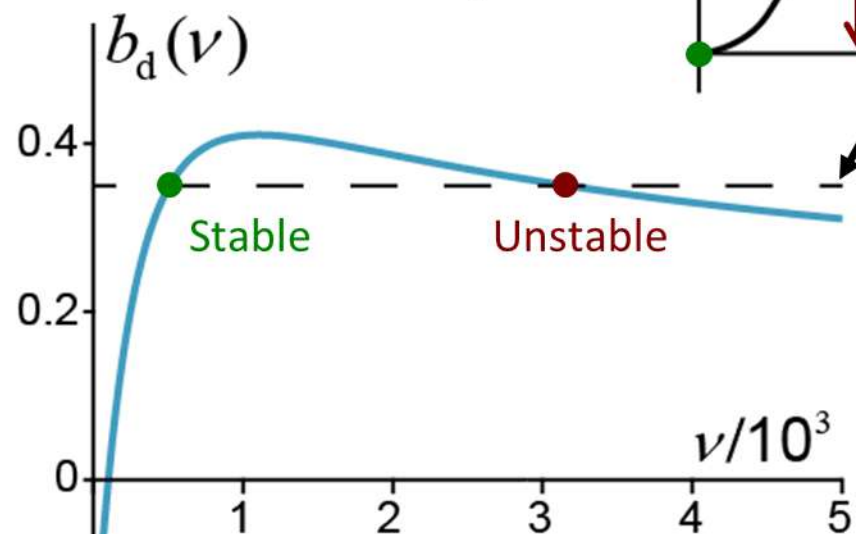
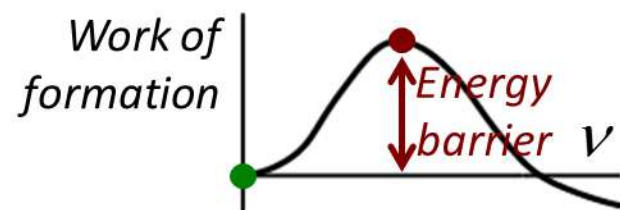
$b_d(v) = b_0$ - at equilibrium



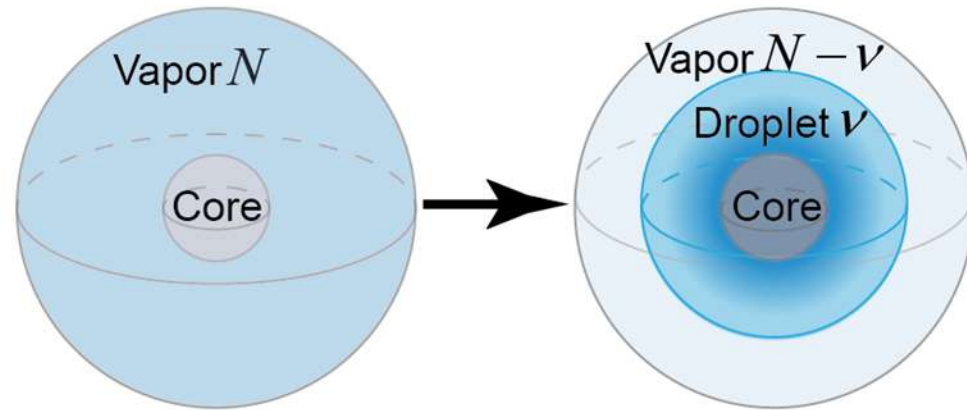
Open systems (Grand Canonical Ensemble)



$$b_d(v) = b_0 \text{ - at equilibrium}$$



Closed systems (Canonical Ensemble)

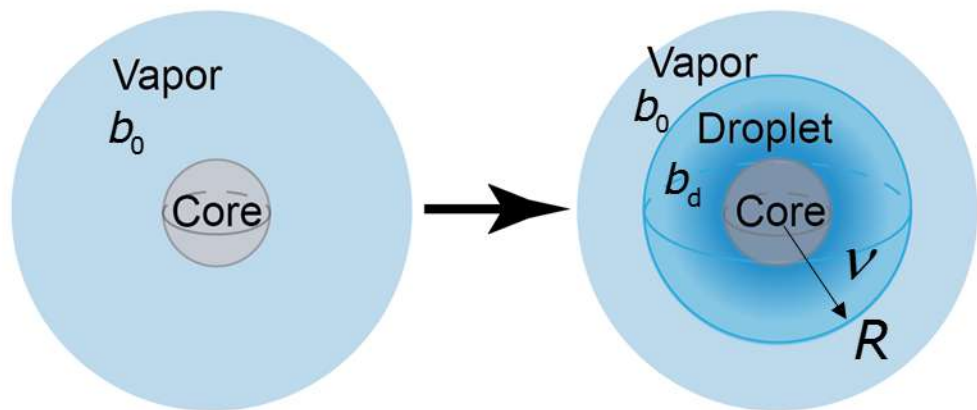


$$b_d = b_0 + \ln(1 - v / N)$$

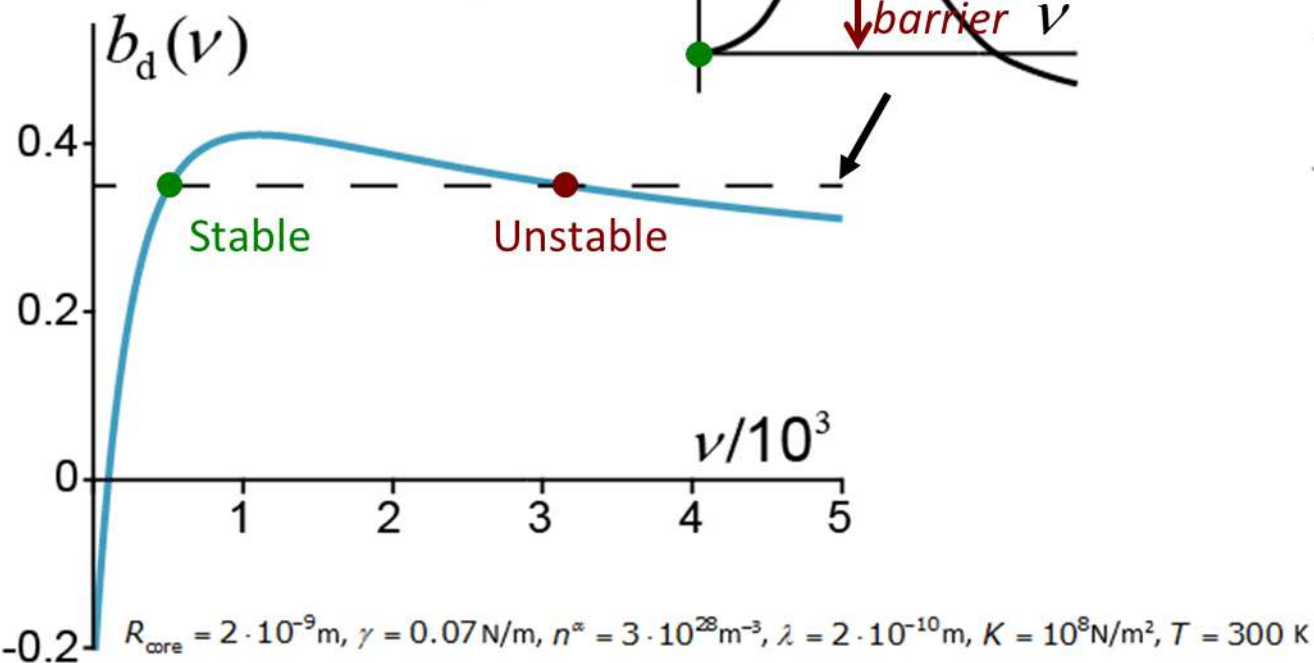
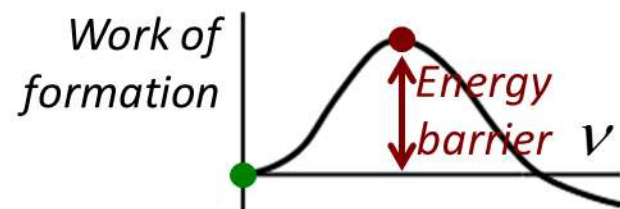
Ideal gas approximation for the vapor phase

$R_{\text{core}} = 2 \cdot 10^{-9} \text{m}$, $\gamma = 0.07 \text{N/m}$, $n^\infty = 3 \cdot 10^{28} \text{m}^{-3}$, $\lambda = 2 \cdot 10^{-10} \text{m}$, $K = 10^8 \text{N/m}^2$, $T = 300 \text{K}$

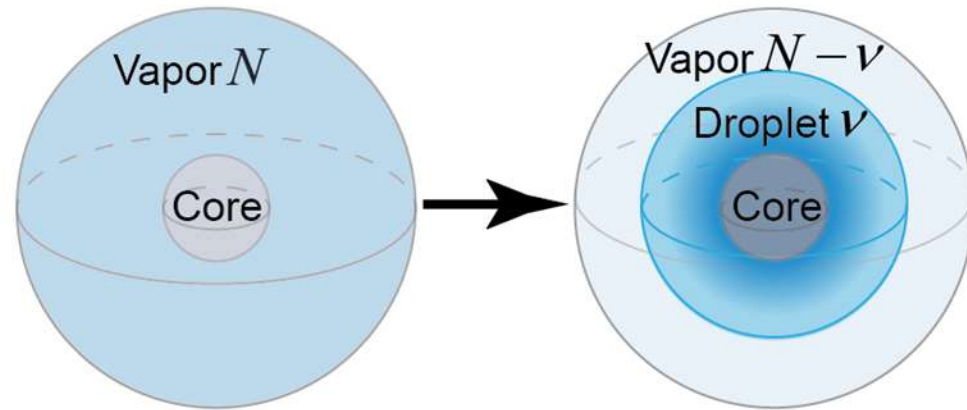
Open systems (Grand Canonical Ensemble)



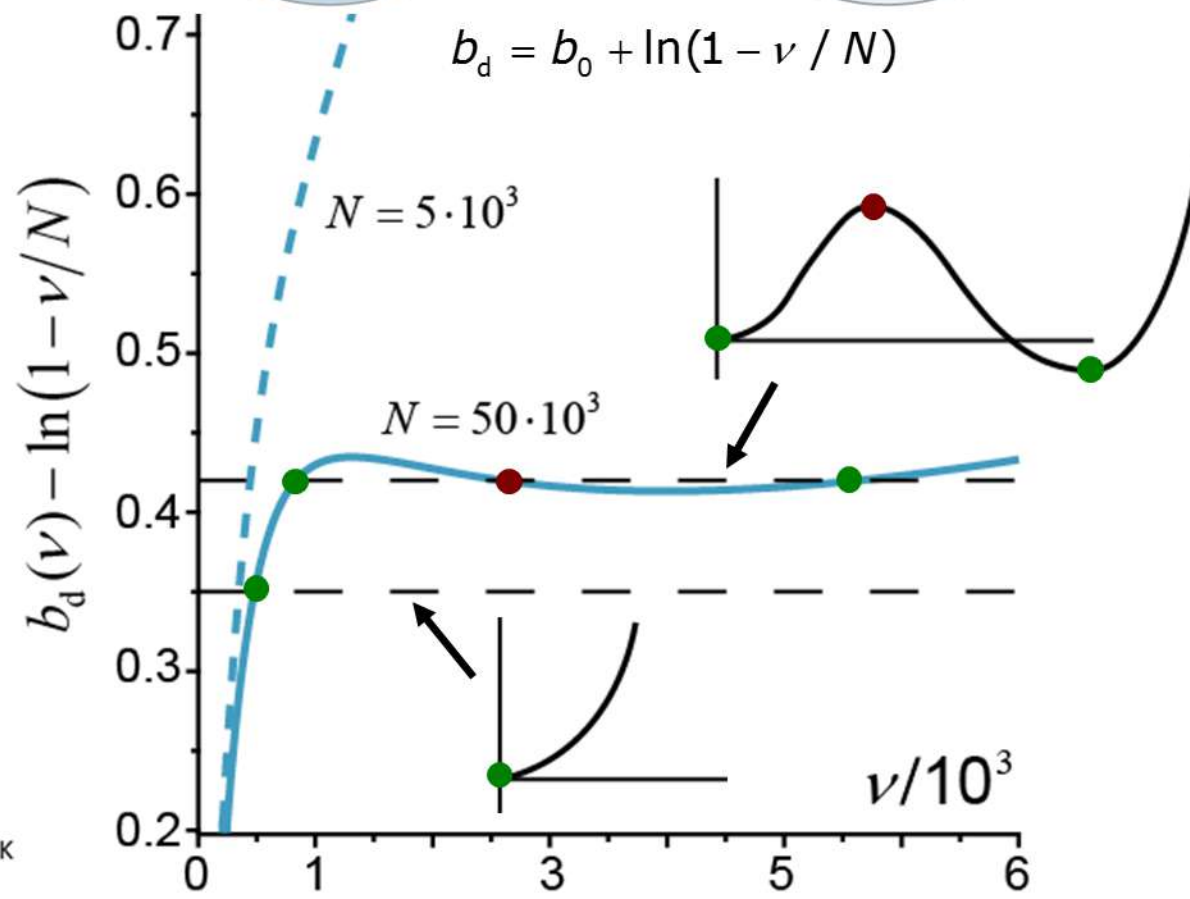
$$b_d(v) = b_0 \text{ - at equilibrium}$$



Closed systems (Canonical Ensemble)



$$b_d = b_0 + \ln(1 - v / N)$$



Classical density functional theory (DFT)

Density functional theory (DFT)

$\frac{\delta F[\rho]}{\delta \rho} = 0 \Rightarrow$ an equation for the density profile $\rho(\vec{r})$ at equilibrium $\left\{ \begin{array}{l} \text{stable – minimum} \\ \text{unstable – saddle point} \end{array} \right.$

$F = ?$ various models...

Density functional theory (DFT)

$$\frac{\delta F[\rho]}{\delta \rho} = 0 \Rightarrow \text{an equation for the density profile } \rho(\vec{r}) \text{ at equilibrium} \begin{cases} \text{stable – minimum} \\ \text{unstable – saddle point} \end{cases}$$

$F = ?$ various models...

1. Square-gradient DFT *R. Evans (1979)*

$$F[\rho(\vec{r})] = F_{\text{ideal}} + \int d\vec{r} \left(-a\rho(\vec{r})^2 + \frac{C}{2}(\nabla\rho(\vec{r}))^2 \right) + k_B T \int d\vec{r} \rho(\vec{r}) \frac{4\eta(\vec{r}) - 3\eta(\vec{r})^2}{(1 - \eta(\vec{r}))^2} + \int d\vec{r} v_{\text{ext}}(\vec{r})\rho(\vec{r}), \quad \eta \equiv \frac{\pi d^3}{6} \rho$$

Long-range attraction

Hard-sphere repulsion
(Carnahan–Starling)

+ assumption of spherical symmetry

2. DFT + Fundamental Measure Theory *R. Roth et al. (2002)*

$$F[\rho(\vec{r})] = F_{\text{ideal}} + \frac{1}{2} \iint d\vec{r} d\vec{r}' \rho(\vec{r}) u(|\vec{r} - \vec{r}'|) \rho(\vec{r}') + k_B T \int d\vec{r} \Phi(\{n_\alpha(\vec{r})\}) + \int d\vec{r} v_{\text{ext}}(\vec{r})\rho(\vec{r})$$

Long-range attraction

Hard-sphere repulsion
(Fundamental Measure Theory)

$$n_\alpha(\vec{r}) \equiv \int d\vec{r}' \rho(\vec{r}') \omega_\alpha(\vec{r} - \vec{r}')$$

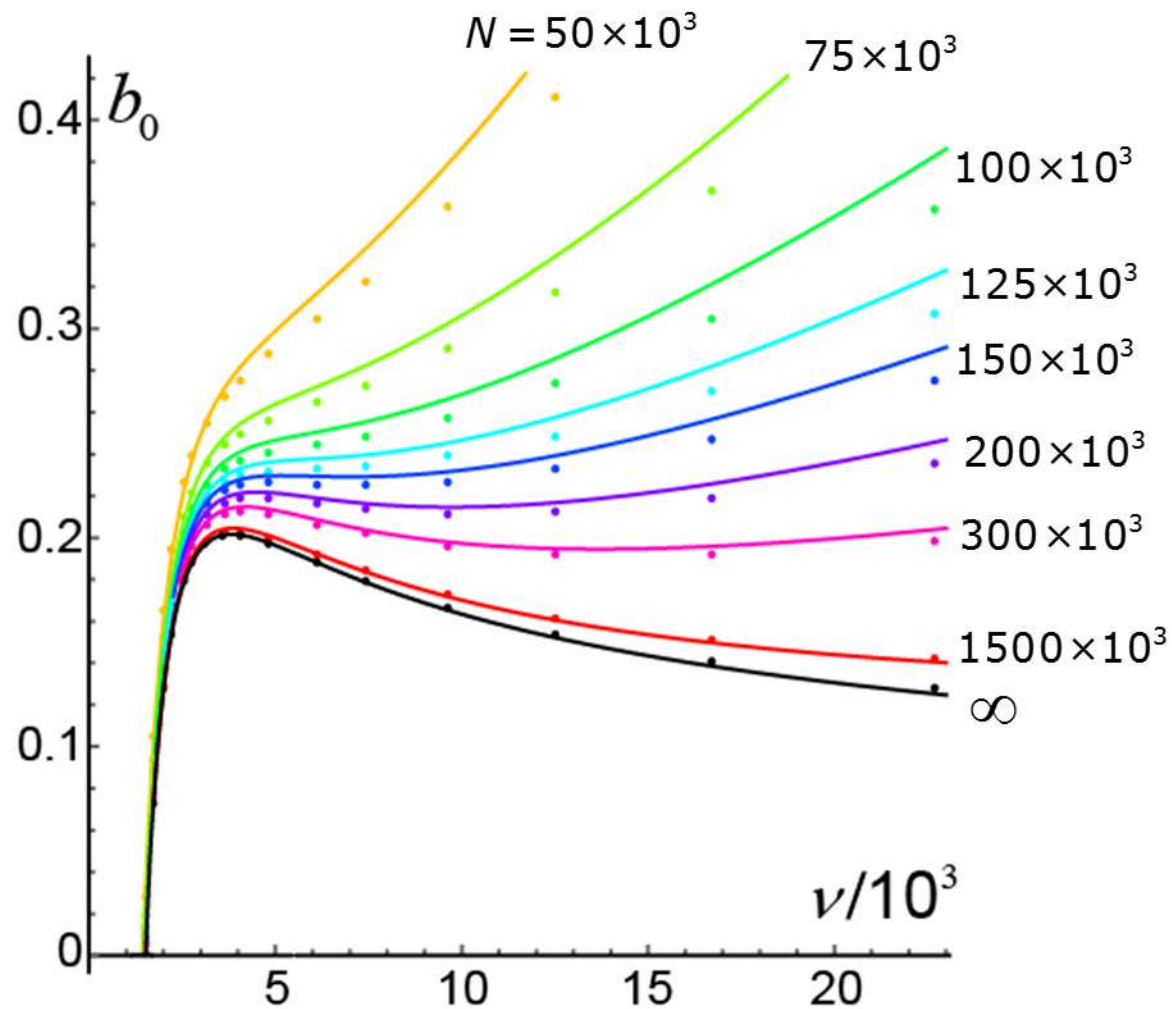
*A library for DFT calculations:
J. Lutsko, classicalDFT.*

github.com/jimlutsko/classicalDFT.

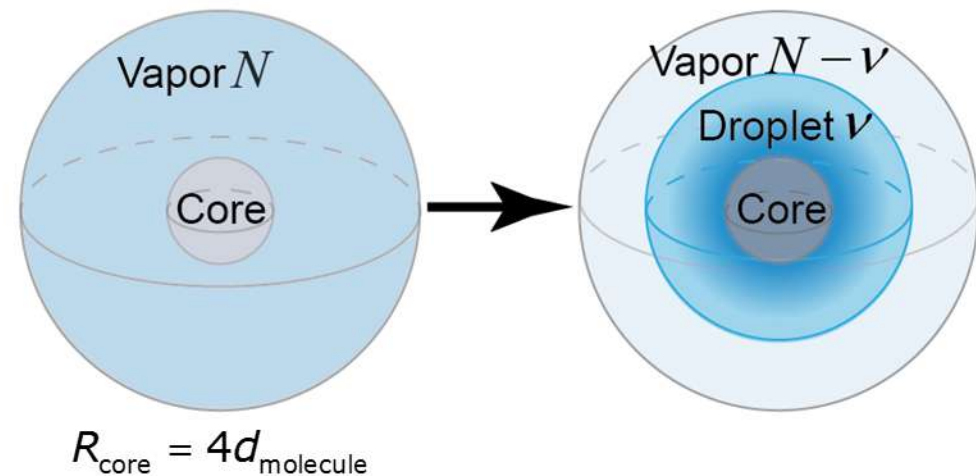
DFT results for closed systems

Results

1. Square-gradient DFT, large N



v is not close to $N \Rightarrow$ The ideal-gas approximation works



Dots – DFT

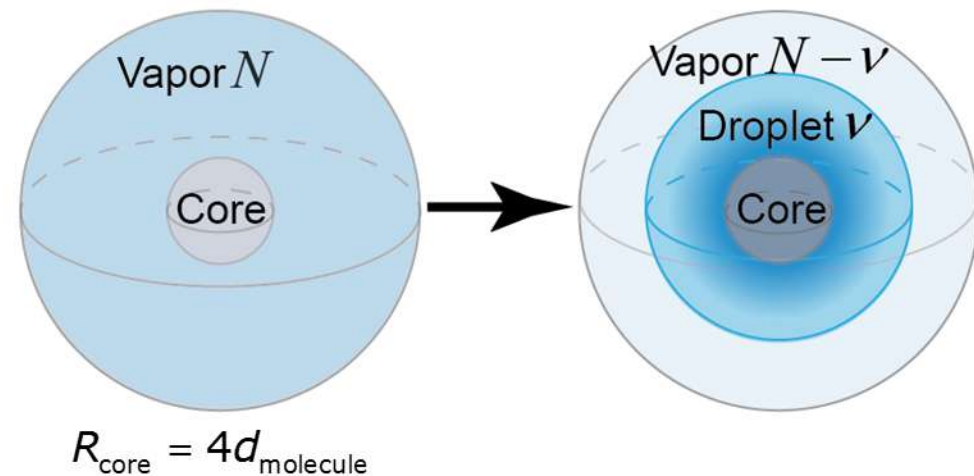
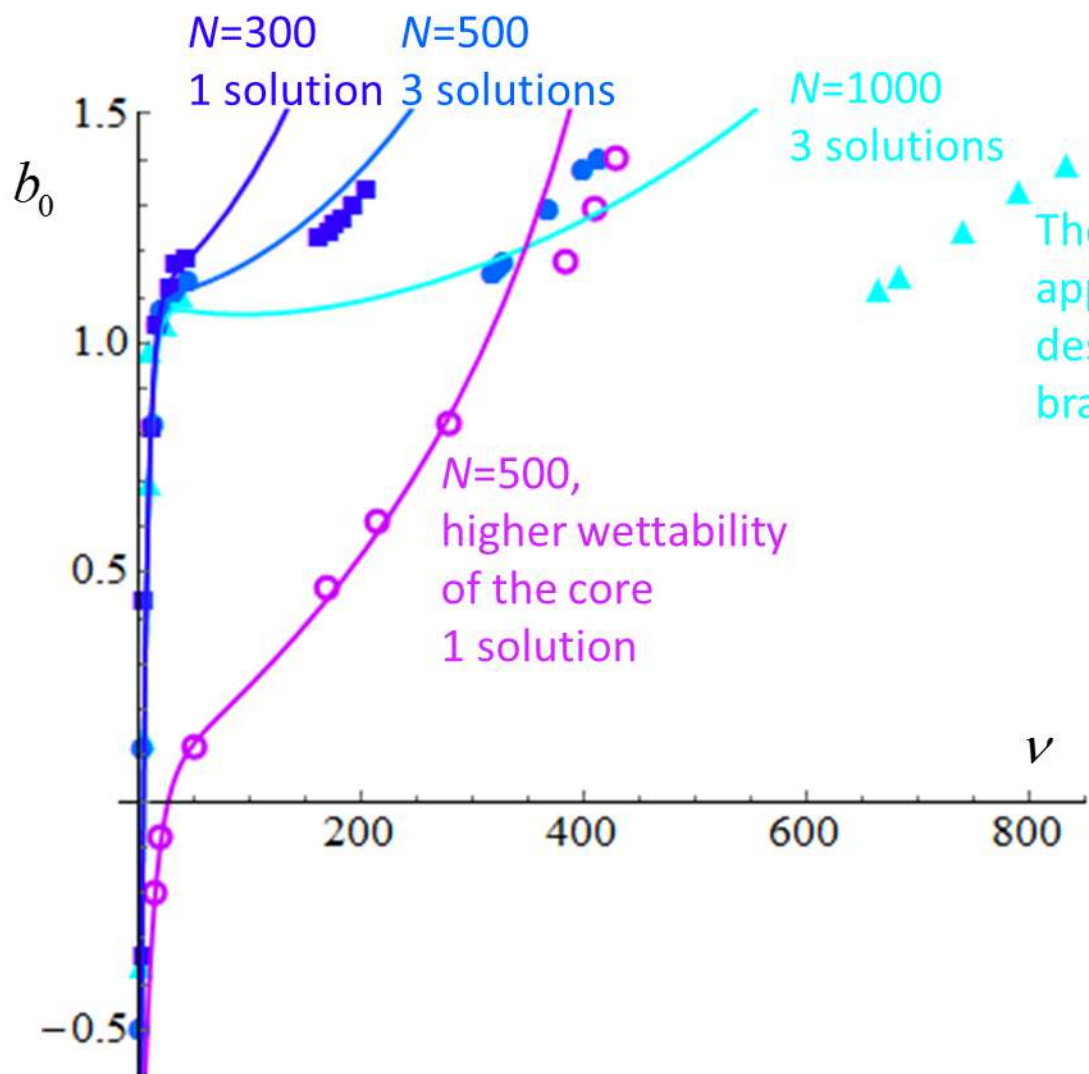
Curves – thermodynamics (fitting)

The dependence of the initial vapor chemical potential b_0 on the number v of droplet's molecules with different total numbers N of argon-like fluid molecules in a closed system.

$$T = 90\text{K}, a = 3.58 \cdot 10^{-49} \text{ J}\cdot\text{m}^3, C = 0.551d^5k_{\text{B}}T, d = 3.14 \cdot 10^{-10} \text{ m}$$

Results

2. Fundamental Measure DFT, small N



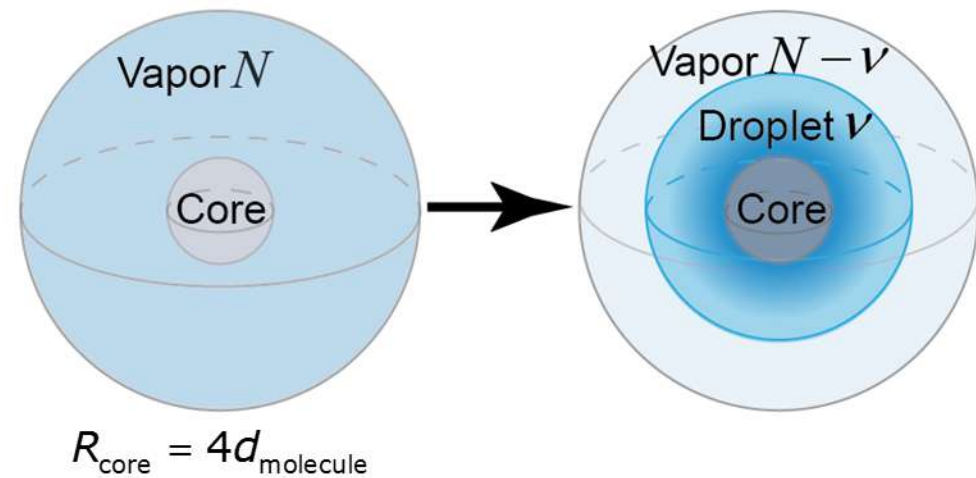
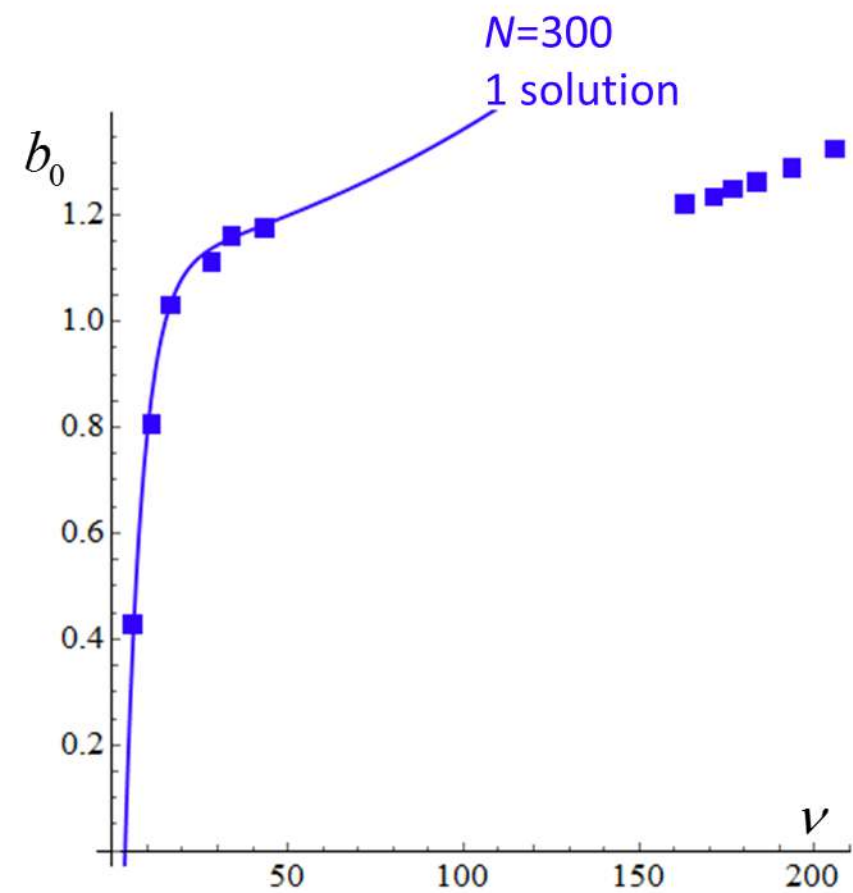
The ideal-gas approximation does not describe the "large drop" branch properly

Dots – DFT
Curves – thermodynamics (fitting)

The dependence of the initial vapor chemical potential b_0 on the number ν of droplet's molecules with different total numbers N of argon-like fluid molecules in a closed system.
Only **stable** solutions are depicted.

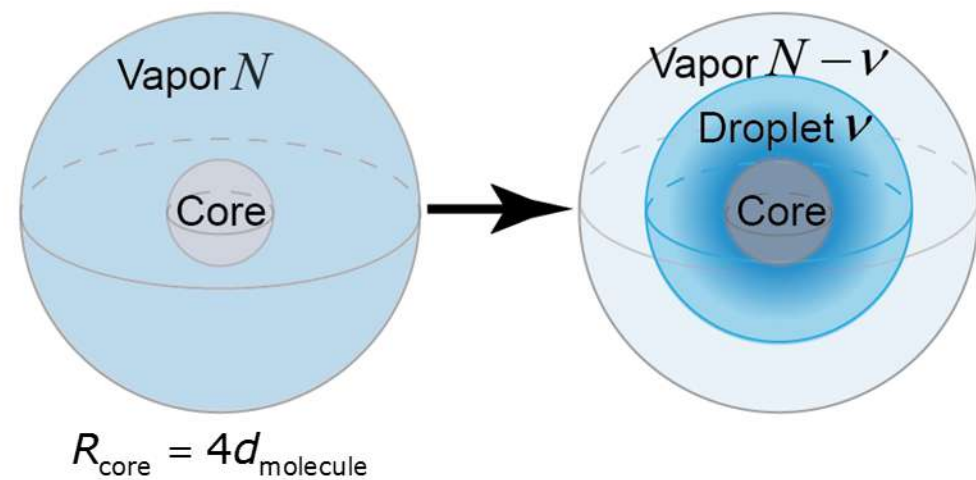
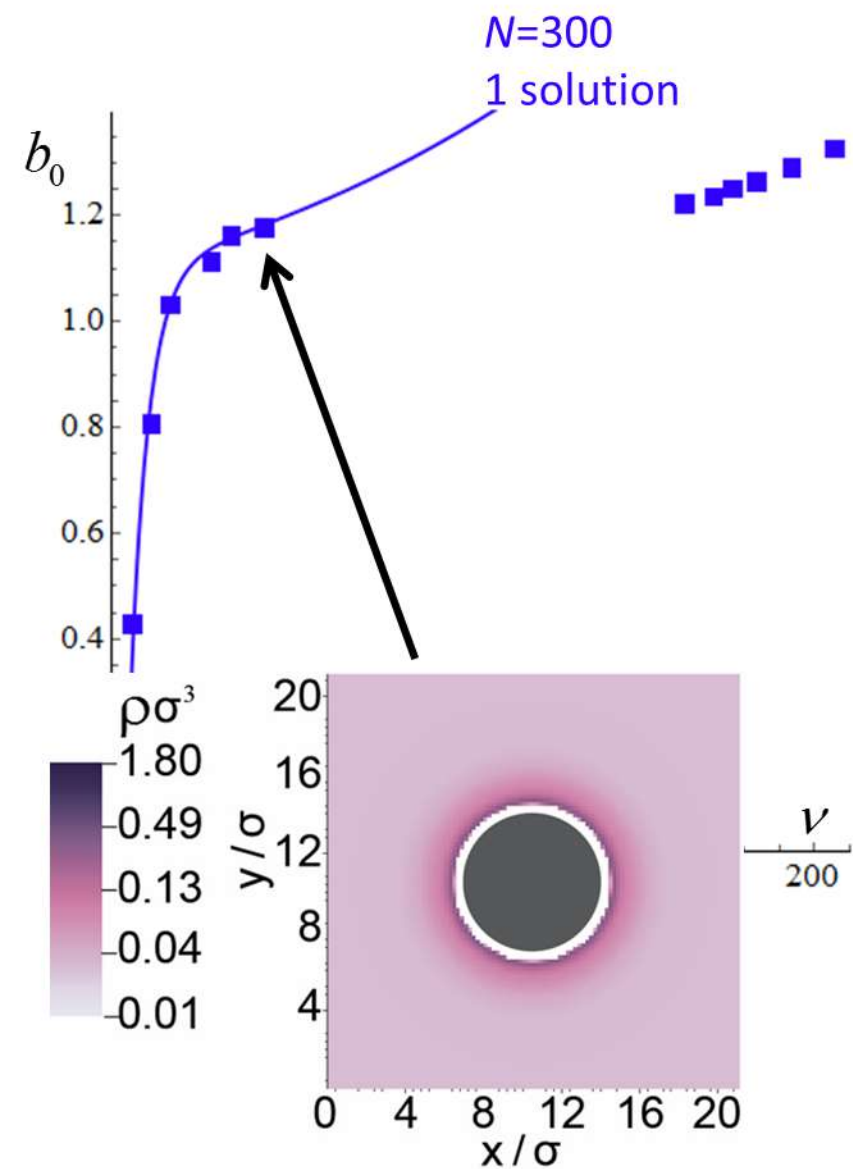
Results

2. Fundamental Measure DFT, small N



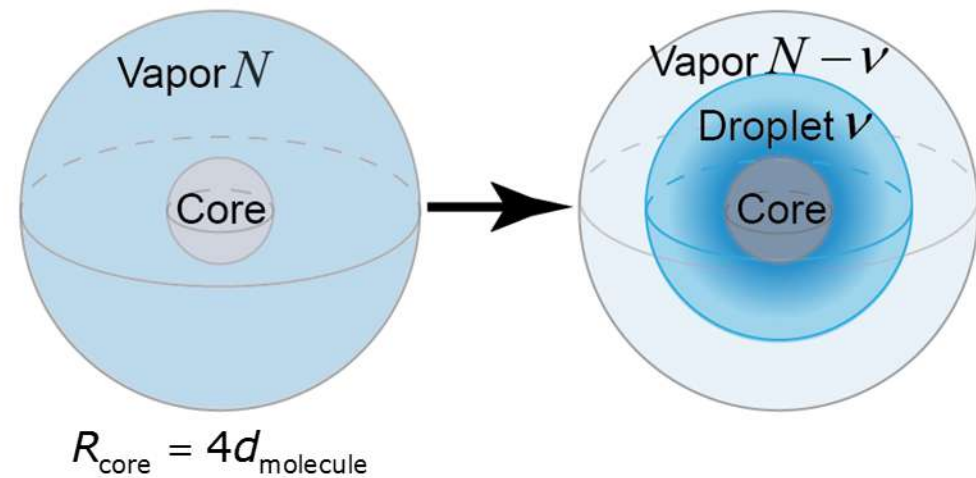
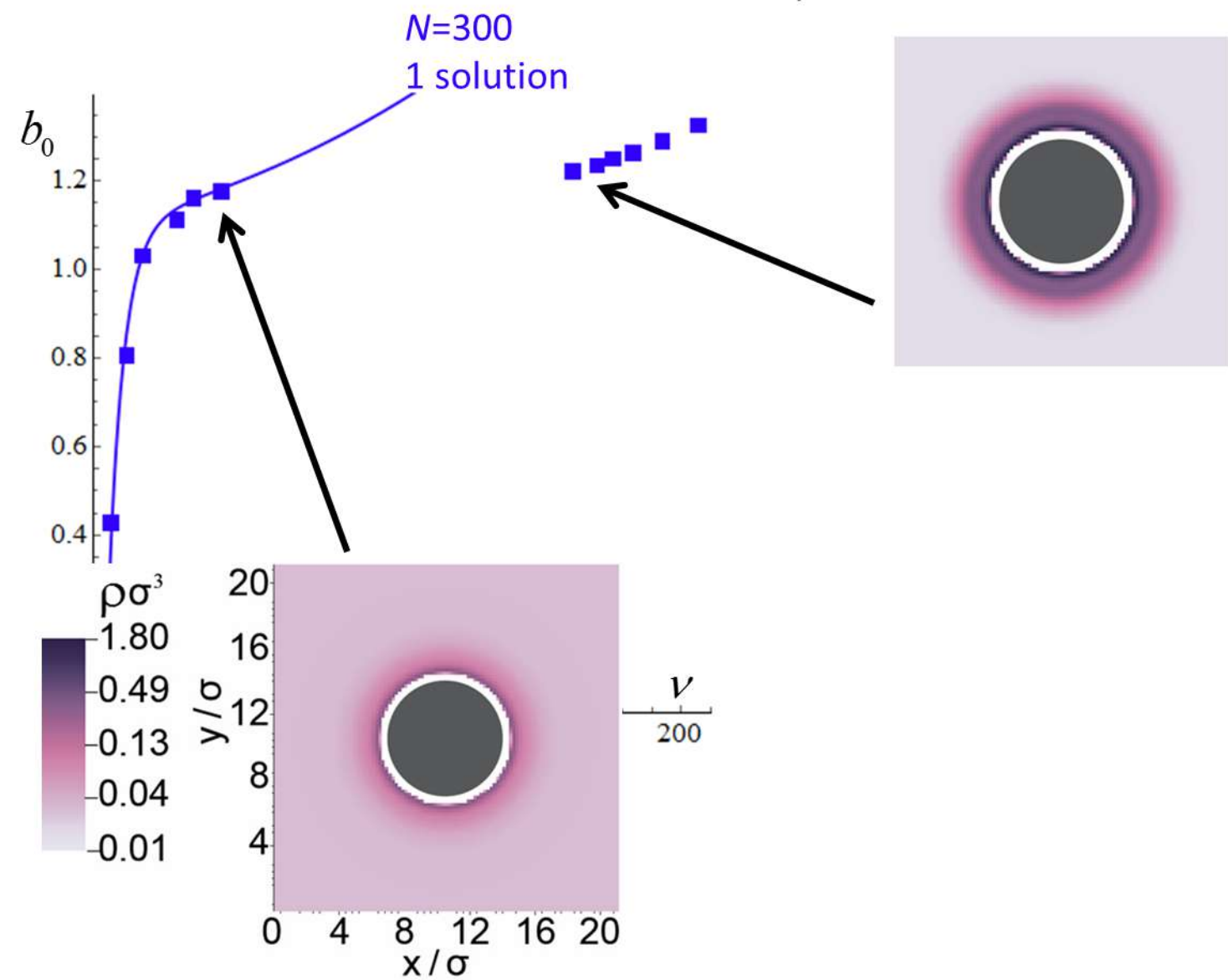
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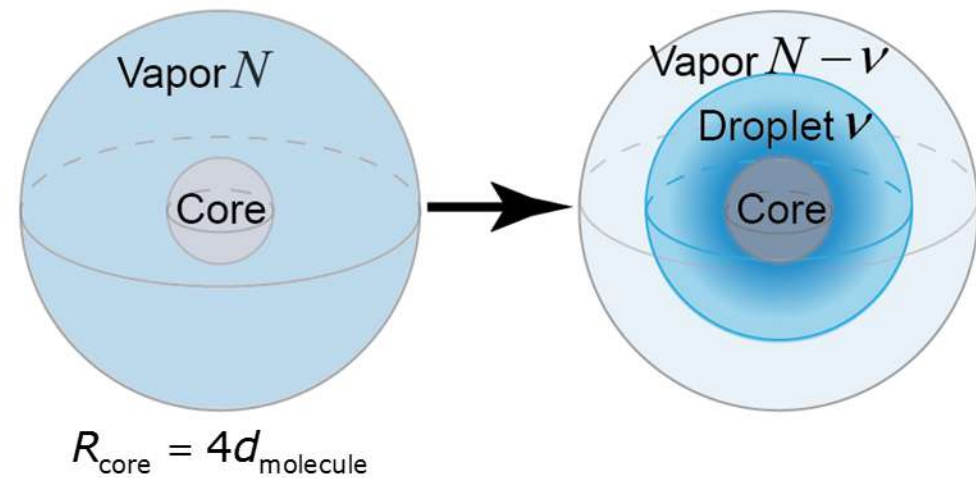
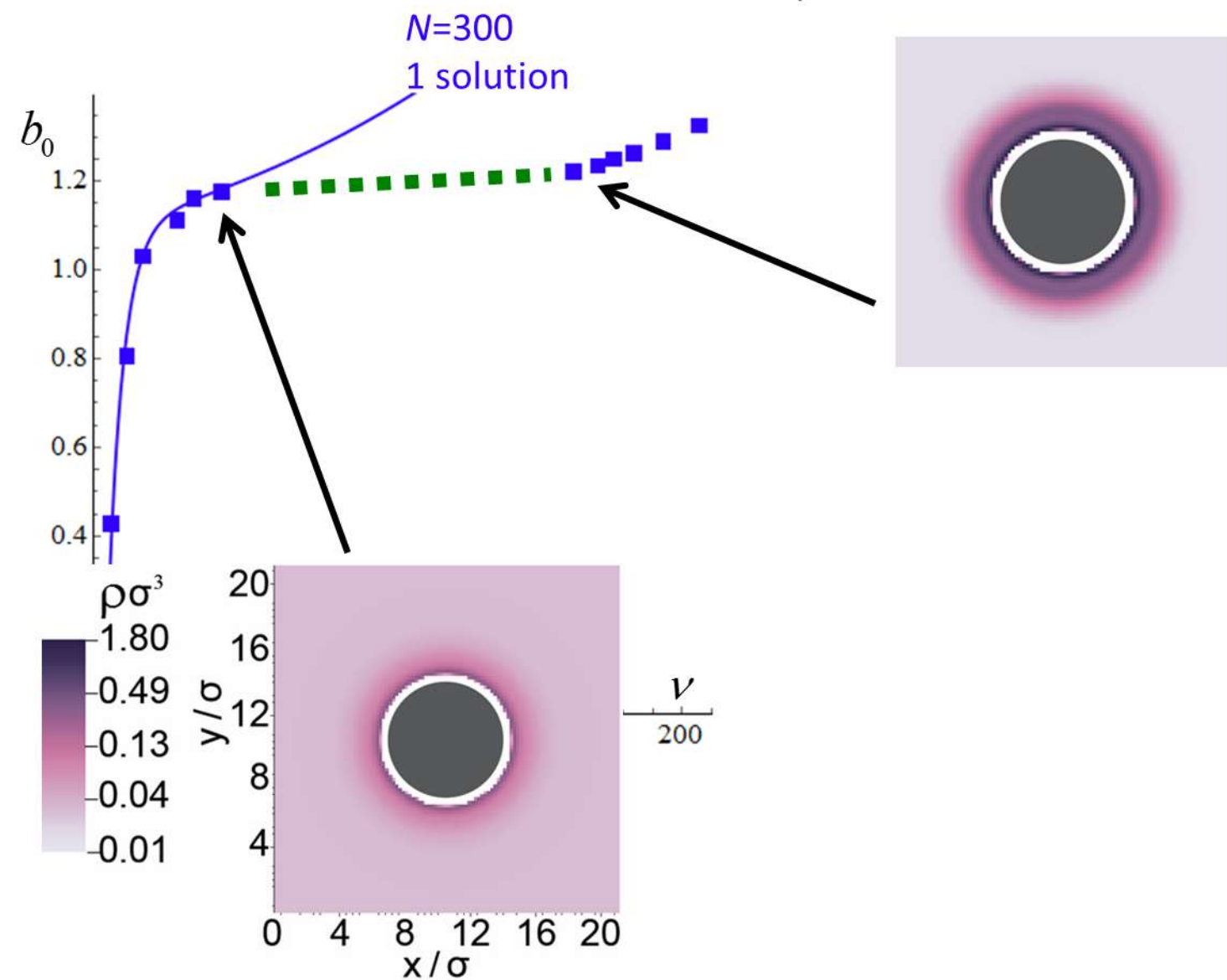
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2. Fundamental Measure DFT, small N



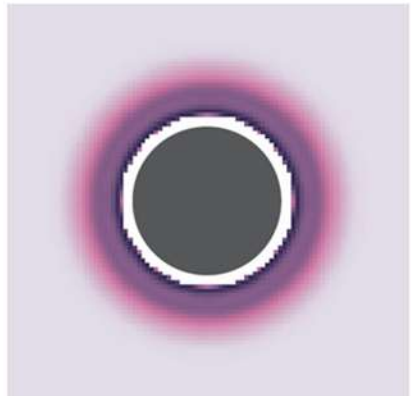
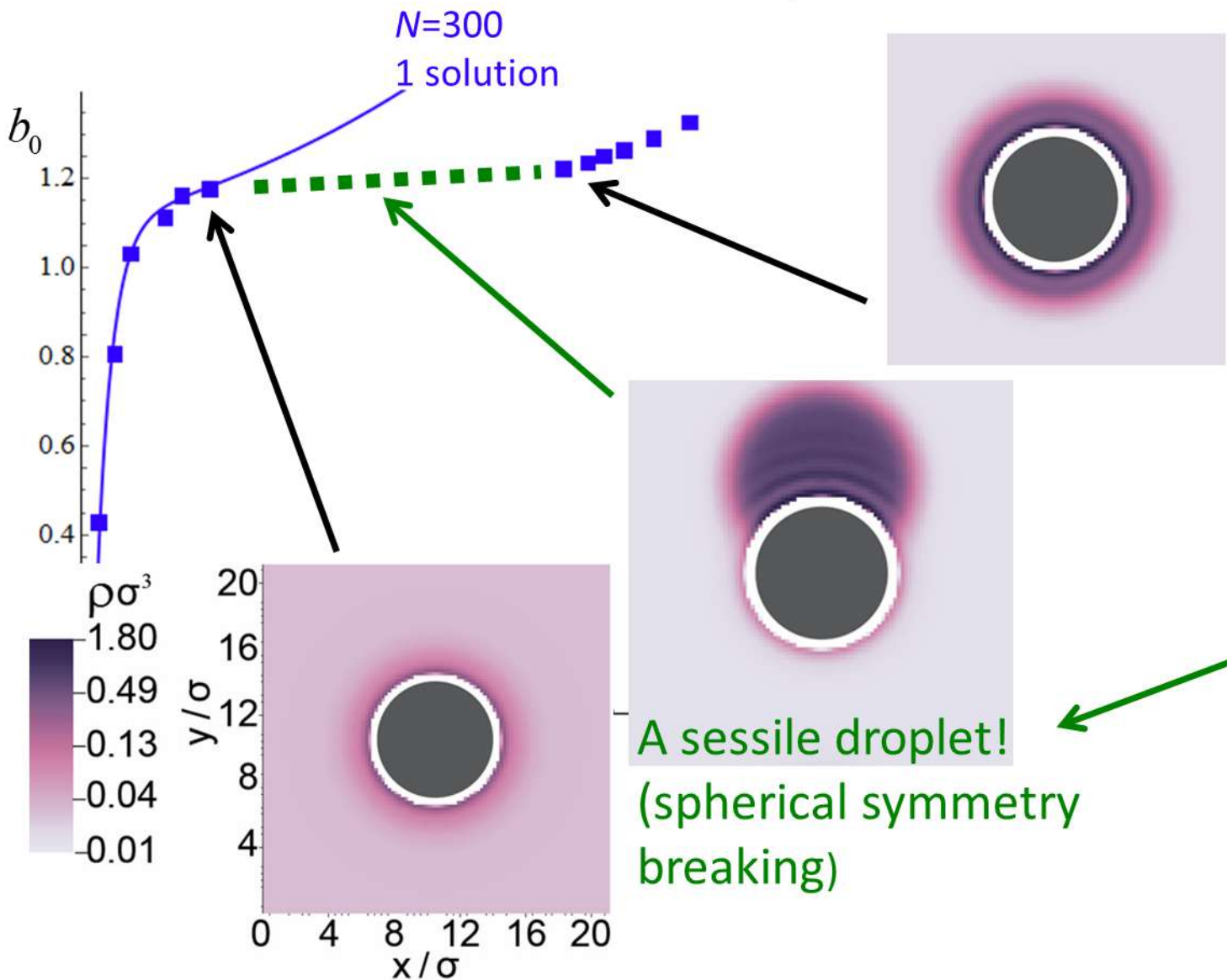
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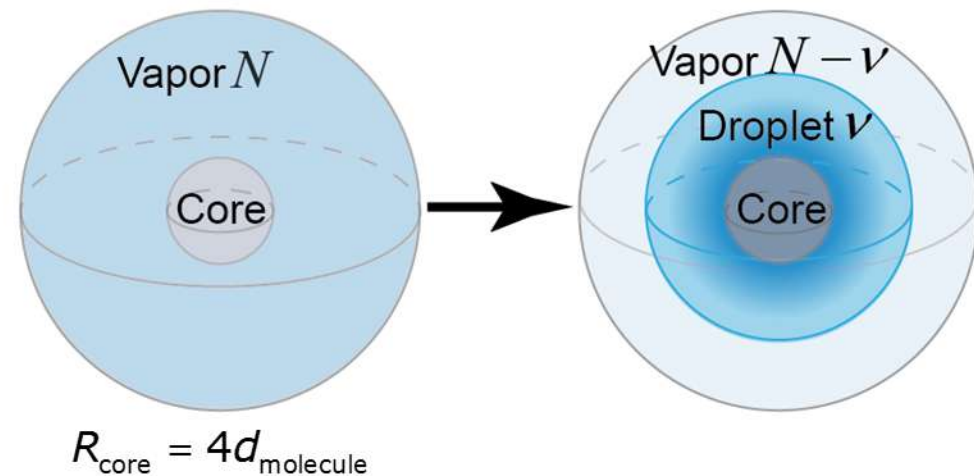


Results

2. Fundamental Measure DFT, small N

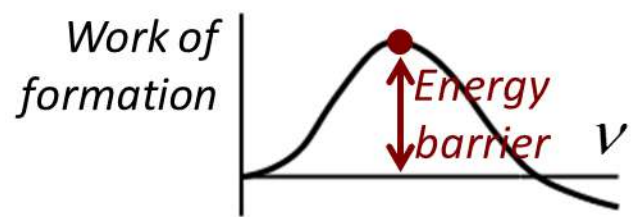


A sessile droplet!
(spherical symmetry breaking)



It's interesting not only for closed systems:

Stable in CE \leftrightarrow unstable in GCE



\Rightarrow in some region of parameters, the transition over the sessile droplet state is more energy-efficient than the transition over the spherical droplet state

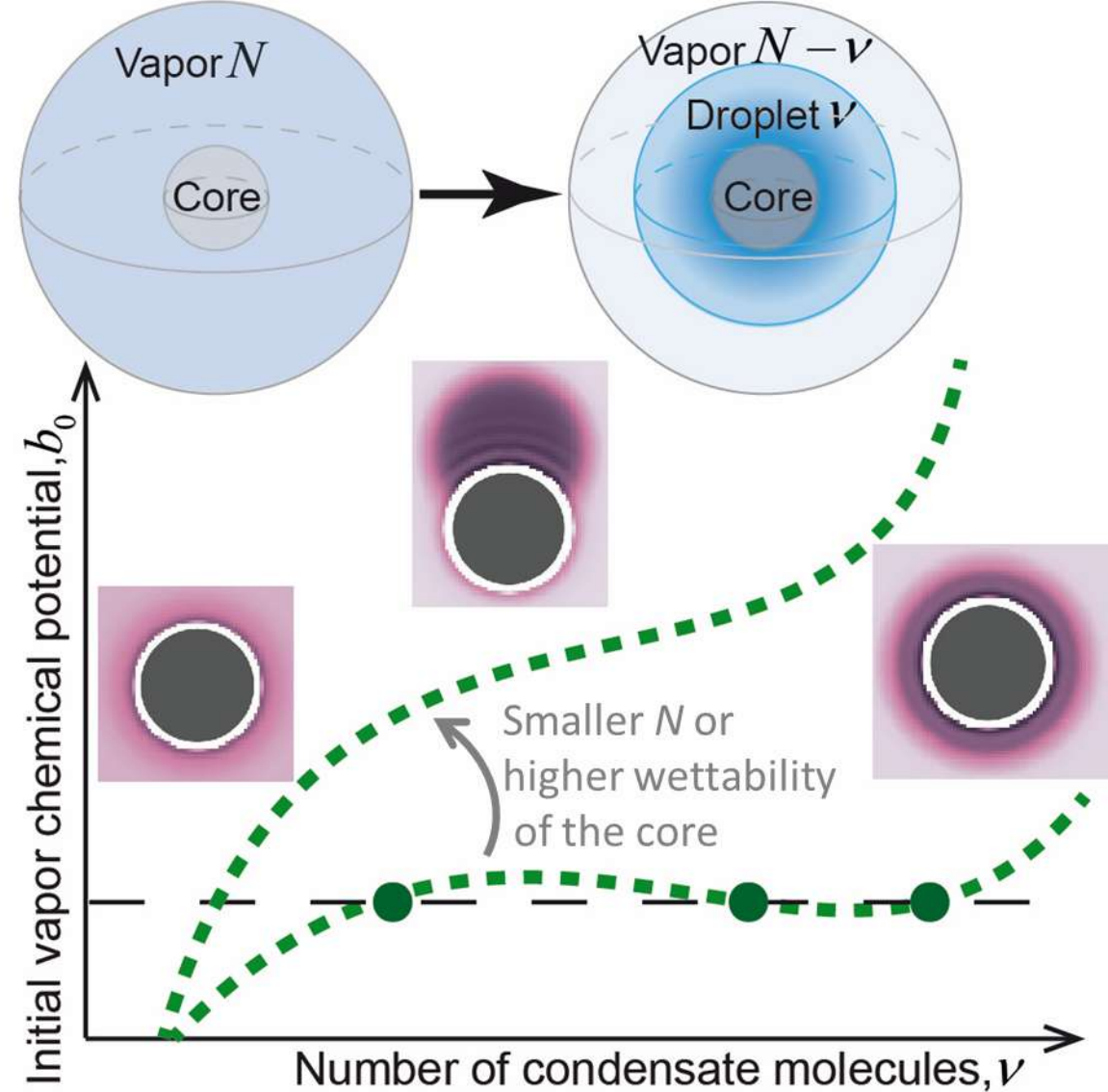
Conclusions

1. Confinement \rightarrow 1 solution OR 3 solutions

2. The ideal-gas approximation for the vapor phase describes only the left branch of the chemical potential (small droplets)

3. The sessile droplet formation (spherical symmetry breaking) in the case of low molecule number N and high initial supersaturation

4. It corresponds to the unstable solution in GCE (an easy way to find a saddle point – to find a minimum in the corresponding CE)



Thanks for your attention!

References

1. Ø. Wilhelmsen, D. Bedeaux, S. Kjelstrup, D. Reguera, J. Chem. Phys. 140, 024704 (2014).
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7. L. Gosteva, A. Shchekin, Physics of Particles and Nuclei Letters, 20 (5), 1084–1087 (2023).
8. A. Shchekin, L. Gosteva, T. Lebedeva, D. Tatyanyenko, Langmuir, *under review*.