

Differential equations method for expansion of hypergeometric functions

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Hypergeometric functions and Feynman integrals

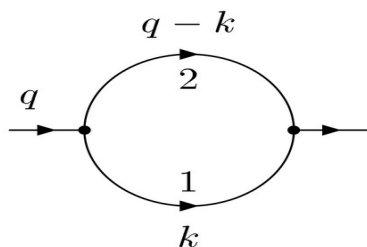
$${}_pF_q \left(\begin{matrix} a_1, \dots, a_p \\ b_1, \dots, b_q \end{matrix} \middle| z \right) = \sum_{n=0}^{\infty} c_n z^n, \quad \frac{c_{n+1}}{c_n} = \frac{(n+a_1)(n+a_2)\dots(n+a_p)}{(n+b_1)\dots(n+b_q)(n+1)}.$$

Methods for calculating Feynman integrals in terms of hypergeometry

Mellin-Barnes

dimensional reduction

Exact Frobenius



The diagram shows a bubble integral with two external lines. The left external line has momentum q entering. The right external line has momentum q leaving. The top internal line has momentum $q-k$ and is labeled with a '2'. The bottom internal line has momentum k and is labeled with a '1'.

$$= \int \frac{d^{4-2\varepsilon} k}{i\pi^{2-\varepsilon}} \frac{1}{(k^2 - m^2)((k-q)^2 - m^2)} = m^\varepsilon \Gamma(\varepsilon) {}_2F_1 \left(\begin{matrix} 1, \varepsilon \\ 3/2 \end{matrix} \middle| \frac{q^2}{4m^2} \right)$$

Polylogarithms

$$G(a_1, \dots, a_n; x) = \int_0^x \frac{G(a_2, \dots, a_n; x')}{x' - a_1} dx', \quad n > 0, \quad G(; x) = 1, \quad G(\vec{0}_n; x) = \frac{\log^n x}{n!}$$

A. B. Goncharov, Mathematical Research Letters 5, 497 (1998).

A. B. Goncharov, arXiv preprint math/0103059 (2001).

$$G(a; b) = \log \left(1 - \frac{b}{a} \right), \quad a \neq 0$$

$$\text{Li}_n(x) = -G \left(\vec{0}_{n-1}, \frac{1}{x}; 1 \right) = \int_0^x \frac{dx'}{x'} \text{Li}_{n-1}(x')$$

$$dG(a_1, \dots, a_n; x) = \sum_{i=1}^n G(a_1, \dots, a_{i-1}, a_{i+1}, \dots, a_n; x) d \log \left(\frac{a_{i-1} - a_i}{a_{i+1} - a_i} \right)$$

Polylogarithms form Hopf algebra

Differential equation

$${}_pF_q \left(\begin{matrix} a_1, \dots, a_p \\ b_1, \dots, b_q \end{matrix} \middle| z \right) = \sum_{n=0}^{\infty} \frac{(a_1)_n, \dots, (a_p)_n}{(b_1)_n, \dots, (b_q)_n} \frac{z^n}{n!}$$

$$\left[z \left(z \frac{d}{dz} + a_1 \right) \left(z \frac{d}{dz} + a_2 \right) \dots \left(z \frac{d}{dz} + a_p \right) - z \frac{d}{dz} \left(z \frac{d}{dz} + b_1 - 1 \right) \left(z \frac{d}{dz} + b_2 - 1 \right) \dots \left(z \frac{d}{dz} + b_q - 1 \right) \right] {}_pF_q = 0$$

$$J = \{f_0, f_1, \dots, f_q\}, \quad f_0 = {}_pF_q, \quad f_n = z^n \frac{d^n f_0}{dz^n}.$$

$$\frac{d}{dz} J = \left(\frac{\mathbf{A}}{z} + \frac{\mathbf{B}}{z-1} \right) J$$

$$\mathbf{C} = -\mathbf{A} - \mathbf{B}$$

$$\mathbf{A} : \{0, 1 - b_1, \dots, 1 - b_q\},$$

$$\mathbf{B} : \left\{ \underbrace{0, \dots, 0}_q, -q - \sum_{i=1}^p a_i + \sum_{i=1}^q b_i \right\},$$

$$\mathbf{C} : \{a_1, a_2, \dots, a_p\}.$$

DE solution and Lee algorithm

$$\frac{dJ}{dx} = \mathcal{M}(x, \varepsilon)J, \quad J = T(x, \varepsilon)\tilde{J}, \quad \frac{d\tilde{J}}{dx} = \tilde{\mathcal{M}}\tilde{J} = \left[T^{-1}\mathcal{M}T - T^{-1}\frac{d}{dx}T \right] \tilde{J}$$

$$\tilde{\mathcal{M}}(x, \varepsilon) = \varepsilon \sum_r \frac{\mathcal{M}_r}{x - x_r} = \varepsilon \bar{\mathcal{M}}(x)$$

$$U(x) = \text{Pexp} \left[\varepsilon \int_{x_0}^x \bar{\mathcal{M}}(x') dx' \right] = \sum_n \varepsilon^n \int_{x > x_n \dots > x_1 > x_0} \dots \int dx_n \dots dx_1 [\bar{\mathcal{M}}(x_n) \dots \bar{\mathcal{M}}(x_1)].$$

Balance transformation

$$\mathcal{B}(\mathcal{P}, x_1, x_2; x) = \mathcal{I} - \mathcal{P} + \frac{x - x_2}{x - x_1} \mathcal{P}$$

J. M. Henn, Physical review letters, vol. 110, no. 25, p. 251601, 2013.

R. N. Lee, JHEP, vol. 04, p. 108, 2015.

Case	A eigenvalues	B eigenvalues	C eigenvalues	variable change
<i>A</i>	$m_a + q_a \varepsilon$	$m_b + q_b \varepsilon$	$m_c + q_c \varepsilon$	nun
<i>B</i>	$\frac{k_a}{n} + m_a + q_a \varepsilon$	$\frac{k_b}{n} + m_b + q_b \varepsilon$	$m_c + q_c \varepsilon$	$z \rightarrow \frac{z_1^n}{1+z_1^n}$
<i>C</i>	$\frac{k_a}{n} + m_a + q_a \varepsilon$	$m_b + q_b \varepsilon$	$\frac{k_c}{n} + m_c + q_c \varepsilon$	$z \rightarrow z_2^n$
<i>D</i>	$m_a + q_a \varepsilon$	$\frac{k_b}{n} + m_b + q_b \varepsilon$	$\frac{k_c}{n} + m_c + q_c \varepsilon$	$z \rightarrow 1 - z_3^n$
<i>E</i>	$\frac{1}{2} + m_a + q_a \varepsilon$	$\frac{1}{2} + m_b + q_b \varepsilon$	$\frac{1}{2} + m_c + q_c \varepsilon$	$z \rightarrow -\frac{4z_4^2}{(z_4^2-1)^2}$
<i>F*</i>	$\frac{k_a}{n} + m_a + q_a \varepsilon$	$\frac{k_a}{n} + m_b + q_b \varepsilon$	$\frac{k_a}{n} + m_c + q_c \varepsilon$	$z \rightarrow \frac{z_1^n}{1+z_1^n}$

$$k = \#\{a_i | i = 1, \dots, q; a_i \notin \mathbb{Z}\}, \quad l = \#\{b_j | j = 1, \dots, p; b_j \notin \mathbb{Z}\}$$

1. $n = 2$ and $|k - l| \geq 2$.
2. $n > 2$ and $(k \geq 2$ or $l \geq 2)$ and $\{a_i - a_1, b_j - a_1 | i = 1, q, j = 1, p\} \not\subset \mathbb{Z}$.
3. $n > 2$, $k = l = 1$ and $a_i - b_j \notin \mathbb{Z}$ where a_i and b_j are two non-integer indices.

A basis $J = \{f_0, f_1, \dots, f_q\}$ with a DE $dJ/dx = MJ$ forms a vector bundle on the Riemannian sphere $\mathbb{C} \cup \infty$, M is the connection of the bundle.

According to the Birkhoff-Grothendieck theorem, any vector bundle on a Riemannian sphere can be represented as a sum of one-dimensional bundles, each of which characterized by an integer.

The entire bundle is characterized by a set of numbers
 $D = \text{diag}(d_1, \dots, d_p)$ (vector bundle splitting)

It can be shown that a system of differential equations can be reduced to epsilon form only if the splitting of the corresponding bundle is trivial

R. N. Lee and A. A. Pomeransky, arXiv:1707.07856 [hep-th]

$${}_2F_1 \left(\begin{array}{c} \frac{1}{2} + \alpha\varepsilon, -\frac{1}{2} + \beta\varepsilon \\ 1 + \gamma\varepsilon \end{array} \middle| x \right)$$

$$D = \text{diag}(1, -1)$$

$${}_3F_2 \left(\begin{array}{c} \alpha\varepsilon, \beta\varepsilon, \gamma\varepsilon \\ \frac{3}{2} + \psi\varepsilon, \frac{1}{2} + \omega\varepsilon \end{array} \middle| x \right)$$

$$D = \text{diag}(1, 0, -1)$$

Main steps of the algorithm

Determine a basis of p elements for the chosen hypergeometric functions (or with indices shifted by an integer) ${}_pF_q$ and write down the corresponding DE system.



Determine whether a suitable variable replacement exists and whether a given function can be expanded in terms of MPLs



Use the Lee algorithm to reduce the DE system to epsilon form



Integrate the resulting system in terms of polylogarithms



If necessary, apply a sequence of differential operators to restore the original function

Example

$$\begin{aligned}
 {}_3F_2 \left(\begin{array}{c} 1, \frac{\varepsilon+1}{2}, \frac{\varepsilon}{2} \\ \frac{1-\varepsilon}{2}, \frac{\varepsilon+3}{2} \end{array} \middle| z \right) &= 1 + \varepsilon \left(-\frac{(z_4^2+1)G(-1, z_4)}{2z_4} + \frac{(z_4^2+1)G(1, z_4)}{2z_4} + 1 \right) \\
 &+ \varepsilon^2 \left(\left(2z_4 + \frac{2}{z_4}\right) G(f_4^1, -1, z_4) - \frac{2(z_4^2+1)G(f_4^1, 1, z_4)}{z_4} + \dots \right) + \mathcal{O}(\varepsilon^3).
 \end{aligned}$$

$$\begin{aligned}
 {}_3F_2 \left(\begin{array}{c} 1, \frac{\varepsilon+1}{2}, \frac{\varepsilon+2}{2} \\ \frac{1-\varepsilon}{2}, \frac{\varepsilon+3}{2} \end{array} \middle| z \right) &= \left(\frac{2z}{\varepsilon} \frac{d}{dz} + 1 \right) {}_3F_2 \left(\begin{array}{c} 1, \frac{\varepsilon+1}{2}, \frac{\varepsilon}{2} \\ \frac{1-\varepsilon}{2}, \frac{\varepsilon+3}{2} \end{array} \middle| z \right) = \\
 &\frac{(z_4^2 - 1)^2 (G(-1, z_4) - G(1, z_4))}{2(z_4^3 + z_4)} + \frac{\varepsilon(z_4^2 - 1)}{2(z_4^3 + z_4)} \left(-4G(f_4^1, -1, z_4) + \right. \\
 &\left. + 4G(f_4^1, 1, z_4) + G(-1, z_4) - G(1, z_4) + G(-1, -1, z_4) + \dots \right) + \mathcal{O}(\varepsilon^2).
 \end{aligned}$$

Appell functions

$$F_1(\alpha, \beta_1, \beta_2, \gamma; x, y) = \sum_{m,n=0}^{\infty} \frac{(\alpha)_{m+n}(\beta_1)_m(\beta_2)_n}{(\gamma)_{m+n}m!n!} x^m y^n, \quad |x| < 1, |y| < 1$$

$$\left[x(1-x) \frac{\partial^2}{\partial x^2} + y(1-x) \frac{\partial^2}{\partial x \partial y} + [\gamma - (\alpha + \beta_1 + 1)x] \frac{\partial}{\partial x} - \beta_1 y \frac{\partial}{\partial y} - \alpha \beta_1 \right] F_1 = 0,$$

$$\left[y(1-y) \frac{\partial^2}{\partial y^2} + x(1-y) \frac{\partial^2}{\partial x \partial y} + [\gamma - (\alpha + \beta_2 + 1)y] \frac{\partial}{\partial y} - \beta_2 x \frac{\partial}{\partial x} - \alpha \beta_2 \right] F_1 = 0.$$

$$J_1 = \left\{ F_1, x \frac{\partial}{\partial x} F_1, y \frac{\partial}{\partial y} F_1 \right\}, \quad \frac{\partial}{\partial x} J_1 = \left(\frac{\mathbf{A}_0}{x} + \frac{\mathbf{A}_1}{x-1} + \frac{\mathbf{A}_y}{x-y} \right) J_1$$

Lauricella functions

$$F_A^{(n)}(\alpha; \beta_1, \dots, \beta_n; \gamma_1, \dots, \gamma_n; x_1, \dots, x_n) = \sum_{m_1, \dots, m_n=0}^{\infty} \frac{(\alpha)_{m_1+\dots+m_n} (\beta_1)_{m_1} \dots (\beta_n)_{m_n}}{(\gamma_1)_{m_1} \dots (\gamma_n)_{m_n} m_1! \dots m_n!} x_1^{m_1} \dots x_n^{m_n},$$

$$F_B^{(n)}(\alpha_1, \dots, \alpha_n; \beta_1, \dots, \beta_n; \gamma; x_1, \dots, x_n) = \sum_{m_1, \dots, m_n=0}^{\infty} \frac{(\alpha_1)_{m_1} \dots (\alpha_n)_{m_n} (\beta_1)_{m_1} \dots (\beta_n)_{m_n}}{(\gamma)_{m_1+\dots+m_n} m_1! \dots m_n!} x_1^{m_1} \dots x_n^{m_n},$$

$$F_D^{(n)}(\alpha; \beta_1, \dots, \beta_n; \gamma; x_1, \dots, x_n) = \sum_{m_1, \dots, m_n=0}^{\infty} \frac{(\alpha)_{m_1+\dots+m_n} (\beta_1)_{m_1} \dots (\beta_n)_{m_n}}{(\gamma)_{m_1+\dots+m_n} m_1! \dots m_n!} x_1^{m_1} \dots x_n^{m_n}.$$

$$\left\{ \theta_{x_{j_1}} \dots \theta_{x_{j_k}} F_i \mid 0 \leq k \leq n, j_1 < j_2 < \dots < j_k \right\}, \quad i = A, B, \quad \theta_x = x \frac{d}{dx}$$

$$\left\{ F_D, \theta_{x_j} F_D \mid j = 1, \dots, n \right\},$$

Conclusions

In this work we study hypergeometric functions and their expansions in terms of polylogarithms.

Of particular importance here are non-trivial variable changes

All algorithms described here for calculating the expansion of single and multiple hypergeometric series were implemented in the **Diogenes** package (Coming soon) written in Wolfram Mathematica language.

Thank you for your attention!