

On the azimuthal flow of protons in the heavy ion  
collisions at  $\sqrt{s_{NN}}=2-4$  GeV

Азимутальный поток протонов в столкновениях  
тяжелых ионов при энергиях  $\sqrt{s_{NN}}=2-4$  ГэВ

*M. Mamaev<sup>a,b,1</sup>*

*M.B. Mamaev<sup>a,b,1</sup>*

<sup>a</sup> National Research Nuclear University MEPhI (Moscow Engineering Physics  
Institute)

<sup>a</sup> Национальный исследовательский ядерный университет «МИФИ»

<sup>b</sup> Institute for Nuclear Research of the Russian Academy of Science

<sup>b</sup> Институт Ядерных Исследований Российской Академии Наук

Одной из основных целей экспериментов по столкновениям тяжелых ионов является изучение свойств сильно взаимодействующей материи, возникающей в области перекрытия сталкивающихся ионов в различных состояниях. При относительно низкой энергии в несколько ГэВ на пару нуклонов созданная материя может характеризоваться высокими барионными плотностями и относительно низкими температурами. Азимутальная анизотропия, возникающая при столкновении адронов, несёт ценную информацию о свойствах вещества в области перекрытия. В этой работе мы представляем результаты свойств масштабирования направленного потока протонов в зависимости от энергии столкновения, а также размера системы. Кроме того, 1  
приведено сравнение полученных результатов с теоретическими моделями.

One of the main goals of the heavy ion collisions experiments is studying the properties of strongly interacting matter created at different states in the overlap region of two intercepting ions. At a relatively low energy of several-GeV per nucleon pair created matter can be characterized by high net baryon densities and relatively low temperatures. The azimuthal anisotropy or produced in the collision hadrons is a valuable probe of the properties of the matter within the overlap region. In this work we present the observation of the scaling properties of the directed flow of protons on energy of the collision as well as the system size. We also elaborate on the directed flow of protons dependence on the geometry of the 2  
collision.

3 PACS: 44.25.+f; 44.90.+c

4 Introduction

5 The strongly interacting matter produced in heavy ion collisions at collision energies of  $\sqrt{s_{NN}}=2-4$  GeV is characterized by the net baryon densities 6  
of 2-5 times greater than the nuclear saturation density [1, 2]. This level of 7  
compression is achieved due to the presence of the colliding nuclei remnants 8

---

<sup>1</sup>E-mail: mam.mih.val@gmail.com

<sup>1</sup>E-mail: mam.mih.val@gmail.com

known as spectators [3]. We can estimate the time it took the accelerated nuclei to penetrate each other ( $t_{pass}$ ), as

$$t_{pass} = \frac{2R}{\sinh y_{beam}}, \quad (1)$$

where  $R$  is the radius of the colliding nuclei and  $y_{beam}$  is the beam rapidity. In non-central collisions the overlap region of two colliding nuclei is asymmetric. Via the interaction of the constituents of the matter within the overlap region as well as cold spectator matter this asymmetry transforms into azimuthal anisotropy of the produced particles [4]. This anisotropy can be quantified by decomposing the azimuthal angle distribution in a Fourier series with respect to plane spanned on vectors of beam direction and impact parameter (reaction plane) [5]:

$$\rho(\phi - \Psi_{RP}) = \frac{1}{2\pi} \left( 1 + 2 \sum_{n=1}^{\infty} v_n \cos n(\phi - \Psi_{RP}) \right), \quad (2)$$

where  $\phi$  is the azimuthal angle of particle momentum and  $\Psi_{RP}$  is the reaction plane. The coefficients of the decomposition  $v_n$  are defined as follows:

$$v_n = \langle \cos n(\phi - \Psi_{RP}) \rangle, \quad (3)$$

where the angle brackets denote the averaging over all particles in the collision as well as all collisions. The coefficients of the first and second orders —  $v_1$  and  $v_2$  — are usually referred to as directed flow and elliptic flow, respectively. The  $v_1$  and  $v_2$  signals are dominant in low energy heavy ion collisions and are sensitive probes of the properties of the matter created within the overlap region. The experimental measurements of the flow harmonics can lay a constraints on the macroscopic properties of the matter existing for the brief moments in the collision. In this work, we discuss the results of the measurements of  $v_1$  of protons in Ag+Ag collisions at  $\sqrt{s_{NN}}=2.4$  GeV ( $E_{kin}=1.23$  AGeV) and 2.55 GeV ( $E_{kin}=1.58$  AGeV), presented by the author at FAIR-NICA workshop in 2021. The preliminary results are based on the analysis of HADES data collected by the experiment in 2019. The detailed comparison with  $v_1$  results for Au+Au collisions at  $\sqrt{s_{NN}}=2.4$  GeV ( $E_{kin}=1.23$  AGeV) will be presented.

## Results

The directed ( $v_1$ ) flow at mid-rapidity can be quantified by its slope  $dv_1/dy|_{y=0}$ . It is determined as the linear term,  $dv_1/dy|_{y=0} = a_1$ , of a cubic ansatz  $v_1(y) = a_1 y + a_3 y^3$  which has been fitted to the measured  $v_1(y)$  data points. The left part of Fig. 1 shows the centrality dependence of the slope of  $v_1$  of protons at midrapidity  $dv_1/dy_{cm}|_{y_{cm}=0}$  for the case if one use the center-of-mass rapidity  $y = y_{cm}$ . The systematical errors correspond to effects not related to initial asymmetry of the collision (see [6, 7]). The slope

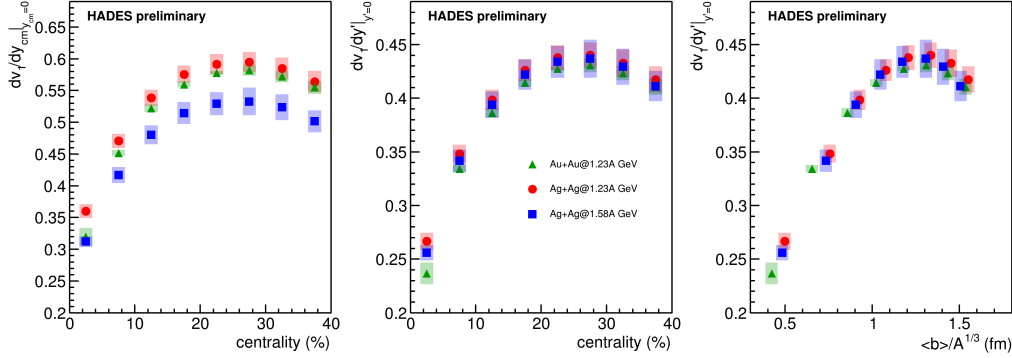


Fig. 1. Left: Directed flow slope of protons  $dv_1/dy|_{y_{cm}=0}$  at center-of-mass rapidity  $y_{cm} = 0$  as a function of centrality of the collision; Center: Directed flow slope of protons  $dv_1/dy'|_{y'=0}$  scaled on the beam rapidity as a function of centrality of the collision; Right: Directed flow slope of protons  $dv_1/dy'|_{y'=0}$  scaled on the beam rapidity as a function of the mean impact parameter divided over the cubic root of mass number of colliding nuclei  $\langle b \rangle / A^{1/3}$  (relative impact parameter of the collision).

43  $dv_1/dy|_{y_{cm}=0}$  for the collision of Au+Au and Ag+Ag at the same energy is  
 44 consistent with the systematic error. The slope  $dv_1/dy|_{y_{cm}=0}$  decreases with  
 45 increasing the collision energy. In this energy range, the anisotropic flow is  
 46 strongly affected by the presence of cold spectators due to sizable passage  
 47 time  $t_{pass}$ . The observed change in the slope  $dv_1/dy|_{y_{cm}=0}$  can be attributed  
 48 to the reduction of the shadowing effects by the spectator matter due to de-  
 49 crease in the  $t_{pass}$ . The rapidity dependence of  $v_1$  of protons becomes less  
 50 complicated if one uses the scaled rapidity  $y' = y_{cm}/y_{beam}$ , since for the col-  
 51 liding beams one then always has  $y'_{beam} = \pm 1$  in the center-of-mass frame, see  
 52 middle panel of Fig 1. The scaled rapidity ( $y' = y_{cm}/y_{beam}$ ) dependence of  $v_1$   
 53 may reflect the partial scaling of  $v_1$  with  $t_{pass}$  in this energy range. The slope  
 54 results for all three data-sets are in agreement within the systematic error  
 55 except for the central-most collisions. To compare flow results for different  
 56 colliding systems, it was suggested to use the scaled impact parameter  $b_0$ ,  
 57 defined by  $b_0 = b/b_{max}$ , taking  $b_{max} = 1.15(A_P^{1/3} + A_T^{1/3})$  fm [10]. Since we  
 58 study the symmetric colliding systems ( $A_T = A_P = A$ ) we use mean impact  
 59 parameter  $\langle b \rangle$  at corresponding centrality class normalized to a cubic root  
 60 of colliding ion mass number  $\langle b \rangle / A^{1/3}$ . The right panel of Fig. 1 shows the  
 61 slope  $dv_1/dy'|_{y'=0}$  as a function of  $\langle b \rangle / A^{1/3}$ . The scaled directed flow slope  
 62 of protons  $dv_1/dy'|_{y'=0}$  shows the same dependency on the relative impact  
 63 parameter of the collision  $\langle b \rangle / A^{1/3}$  in all three cases. Figure 1 shows, that  
 64 the use of the scaled variables may simplify the comparison of  $v_n$  results for  
 65 different colliding systems and collision energies.

66  
 67 The directed flow slope of protons  $dv_1/dy|_{y=0}$  as a function of collision en-  
 68 ergy  $\sqrt{s_{NN}}$  is shown in Fig. 2. The  $v_1$  results for E895 experiment are taken  
 69 from [8], for the STAR-FXT program from [9] and FOPI experiment results

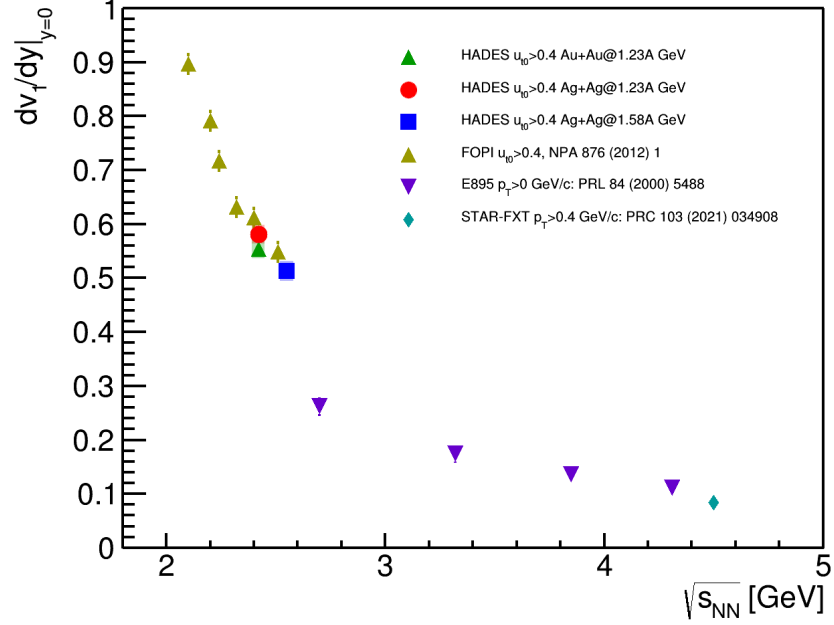


Fig. 2. The directed flow slope of protons  $dv_1/dy|_{y=0}$  as a function of collision energy  $\sqrt{s_{NN}}$ : comparison with the world data. The results were taken from: the E895 experiment [8], the STAR-FXT program [9] and the FOPI measurements [10].

from [10]. The values are presented without the scaling on the beam rapidity  
 and the FOPI-similar momentum criteria ( $u_{t0} > 0.4$ ) is applied, which cor-  
 responds to the  $p_T > 0.3$  GeV at  $E_{kin} = 1.23A$  GeV and  $p_T > 0.35$  GeV at  
 $E_{kin} = 1.58A$  GeV. The points for Au+Au at  $E_{kin} = 1.23A$  GeV and Ag+Ag  
 at  $E_{kin} = 1.23A$  and  $1.58A$  GeV follow the world data trend for  $v_1$ .

## Summary

The results for the directed flow slope of protons at midrapidity  $dv_1/dy|_y = 0$  are presented as a function of centrality and relative impact parameter of the collision for Au+Au at  $E_{kin} = 1.23A$  GeV and Ag+Ag at  $E_{kin} = 1.23A$  and  $1.58A$  GeV. Scaling the results on the beam rapidity reduces the dependency of  $dv_1/dy|_y = 0$  on the energy of the collision. Scaled directed flow slope of protons  $dv_1/dy'|_{y'} = 0$  shows the same dependency on the relative impact parameter of the collision in all three data-sets. The obtained results for  $dv_1/dy|_y = 0$  follow the existing world data trend as a function of collision energy.

## REFERENCES

1. Bzdak A., Esumi S., Koch V., Liao J., Stephanov M., Xu N. Mapping the Phases of Quantum Chromodynamics with Beam Energy Scan // Phys. Rept. — 2020. — V. 853. — P. 1–87. — arXiv:1906.00936.

- 89 2. *Xu N., others.* Nuclear Matter at High Density and Equation of State. —  
90 2022.
- 91 3. *Danielewicz P., Lacey R., Lynch W.G.* Determination of the equation of  
92 state of dense matter // Science. — 2002. — V. 298. — P. 1592–1596. —  
93 arXiv:nucl-th/0208016.
- 94 4. *Herrmann N., Wessels J.P., Wienold T.* Collective flow in heavy ion  
95 collisions // Ann. Rev. Nucl. Part. Sci. — 1999. — V. 49. — P. 581–632.
- 96 5. *Voloshin S.A., Poskanzer A.M., Snellings R.* Collective phenomena in  
97 non-central nuclear collisions // Landolt-Bornstein. — 2010. — V. 23. —  
98 P. 293–333. — arXiv:0809.2949 [nucl-ex].
- 99 6. *Mamaev M. et al.* [HADES Collaboration] Directed flow of protons with  
100 the event plane and scalar product methods in the HADES experiment  
101 at SIS18 // J. Phys. Conf. Ser. — 2020. — V. 1690, no. 1. — P. 012122.
- 102 7. *Abgaryan V. et al.* [MPD Collaboration] Status and initial physics per-  
103 formance studies of the MPD experiment at NICA // Eur. Phys. J. A. —  
104 2022. — V. 58, no. 7. — P. 140. — arXiv:2202.08970.
- 105 8. *Liu H. et al.* [E895 Collaboration] Sideward flow in Au + Au collisions  
106 between 2-A-GeV and 8-A-GeV // Phys. Rev. Lett. — 2000. — V. 84. —  
107 P. 5488–5492. — arXiv:nucl-ex/0005005.
- 108 9. *Abdallah M.S. et al.* [STAR Collaboration] Disappearance of partonic  
109 collectivity in sNN=3GeV Au+Au collisions at RHIC // Phys. Lett. B. —  
110 2022. — V. 827. — P. 137003. — arXiv:2108.00908.
- 111 10. *Reisdorf W. et al.* [FOPI Collaboration] Systematics of azimuthal asym-  
112 metries in heavy ion collisions in the 1 A GeV regime // Nucl. Phys. A. —  
113 2012. — V. 876. — P. 1–60. — arXiv:1112.3180 [nucl-ex].