

The simulation of interactions in the straw-based SPD
track detector and primary vertex reconstruction

Моделирование взаимодействий в трековом
детекторе SPD на основе строу-трубок и
реконструкция первичных вершин

S.D. Morozova^{a,1}, A.V. Shipilova^{a,b,2}

С.Д. Морозова^а, А.В. Шипилова^{а,б}

^a Samara National Research University, Samara, Russia

^b Joint Institute for Nuclear Research, Dubna, Russia

^a Самарский национальный исследовательский университет, Московское шоссе, 34, г. Самара, 443086, Россия

^b Объединенный институт ядерных исследований, ул. Жолио-Кюри, д.6, г. Дубна, Московская область, 141980, Россия

In this work we simulate the SPD NICA straw tracker detector response in the trigger-less regime using GEANT4 tools. We study the temporal structure of signals and investigate the vertex reconstruction efficiency using the simulation data. We develop a part of prototype software for event reconstruction at the stage of online data filtering.

В данной работе мы моделируем отклик строу-трекера SPD детектора коллайдера NICA в бестриггерном режиме с помощью инструментов GEANT4. Исследуется временная структура сигналов и эффективность реконструкции вершин по данным моделирования. Разработана часть прототипа программного обеспечения для реконструкции событий на онлайн этапе фильтрации данных.

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Introduction

The Spin Physics Detector (SPD) at JINR NICA complex is currently under construction [1]. The SPD is a universal facility for study of spin-related phenomena with deuteron and proton beams. The detector has a multi-layered structure, and one of its layers is a straw tracker, that should provide the measurement of the secondary and primary particles momenta with high precision. For the stage of the SPD active operation it is necessary to develop fast data processing algorithms for online data collection, event selection and primary vertices reconstruction.

¹E-mail: svtmorozova09@gmail.com

²E-mail: shipilova.av@ssau.ru

14 The straw tracker is the inner part of the SPD detector. Using the
 15 GEANT4 software package [2] we model a geometry of the SPD straw tracker
 16 (ST), its sensitive volumes and their response. We adopt a number of simpli-
 17 fications against the real ST geometry which would be insignificant for this
 18 stage of our study. We model the ST by a system of nested cylinders each
 19 constructed by one layer of parallel cylindrical tight-fitted straw tubes, as
 20 illustrated in the Fig. 1, right. The ST tubes have the outer polyethylene
 21 shell of thickness $R = 0.036$ mm and the inner cylindrical volume of radius
 22 $R = 4.934$ mm filled by $Ar(70\%) + CO_2(30\%)$ gas mixture, which includes
 23 tungsten wire (anode) of radius $R = 0.03$ mm, see Fig. 1, left. We assign a
 24 sensitive detector object of GEANT4 to the inner volume of each tube and
 25 adopt the tube numbering scheme where the unique number corresponds to
 26 the each tube.

27 The length of the time slice in the experiment is $10\mu s$, while the proton
 28 bunch crossings occur every 76 ns. The probability of proton-proton interac-
 29 tion in a one pp bunch crossing is simulated by Poisson distribution $f(k) =$
 30 $\frac{\lambda^k}{k!} e^{-\lambda}$ with expected value of $\lambda = 0.3$. The interaction point is placed into
 31 $(0, 0, z)$, where z is defined by Gaussian distribution $f(z) = \frac{1}{\sigma\sqrt{2\pi}} e^{-\frac{1}{2}(\frac{z-z_0}{\sigma})^2}$
 32 with $\sigma = 30$ cm and central value of $z_0 = 0$. The charged reaction products
 33 are modelled by muons carrying the energy of $E = 1$ GeV and the momentum
 34 p which direction is uniformly distributed in the 4π space. The number of
 35 muons produced in the pp collision is defined by a Poisson distribution with
 36 expected value of $\lambda = 7$.

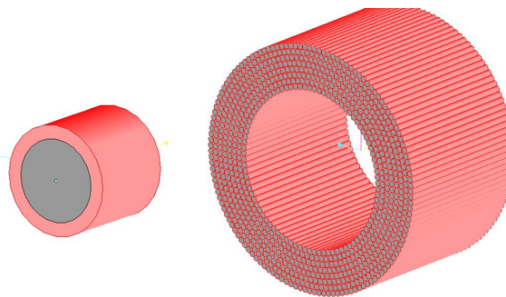


Fig. 1. Single straw (left) and 8 layers of straw detector model (right). Red layer – polyethylene, grey – gas, green – tungsten wire.

38 The propagation of the charged particle through the gas leads to its en-
 39 ergy loss. We declare the hits collection object to store the characteristics
 40 of particle energy loss points. To simulate the ST response time we should
 41 define the shortest distance from the particle track to the anode (e.g. tube
 42 axis). In case there are several energy loss points in the same logical volume

43 we adopt the approximation where the only first and last points are consid-
 44 ered, then the shortest distance is calculated by crossing lines formula. The
 45 dependence of the electron avalanche drift time on the distance was simulated
 46 using Garfield simulation software [3], [4], see TDR [5]. We approximated
 47 this dependence by the analytic formula and used it to obtain the time distri-
 48 butions of ST response. The results of our simulation are presented in the
 49 fig. 2, where all the histograms were created using CERN ROOT tools [6].
 50 We found a significant overlap of the ST response times for particles produced
 51 in different bunch crossings from the same time slice. This fact points out to
 52 the problem of signal decoding for event reconstruction when collecting data
 53 in a real experiment.

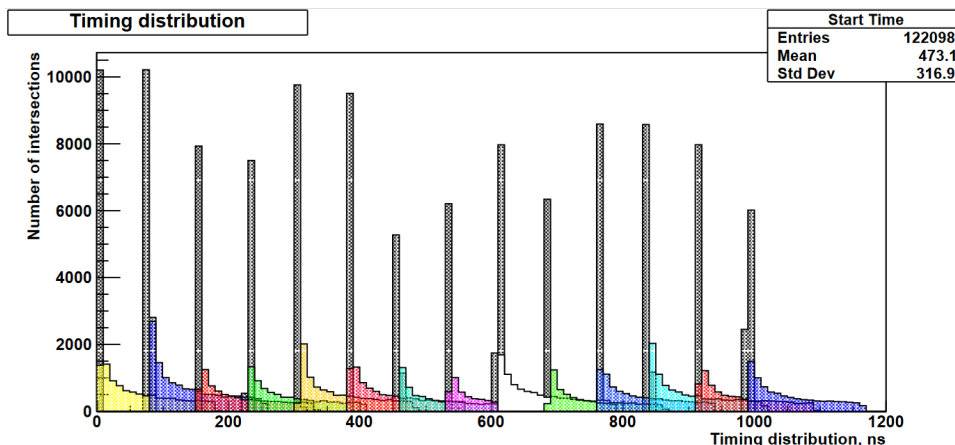


Fig. 2. Timing distribution averaged by 1000 time slices. Grey area – the sample particle intersection time of the sensitive volume. Coloured area – ST response time distribution.

54 Primary vertex reconstruction

55 Using the hits collection data one can perform a reconstruction of parti-
 56 cle tracks. We skipped the hits recovery step assuming the coordinates
 57 of the particles energy loss points are already known. We suppose a uni-
 58 form magnetic field along z axis of $B = 1$ T without endcup effects. Thus,
 59 we approximate the sample charged particle trajectories in the XOY plane
 60 transverse to the field by parabolic function $y = a_1x^2 + a_2x + a_3$, where
 61 the coefficients a_i , $i = 1, 2, 3$ are determined from the hits data using the
 62 least-squares method. The simulated tracks and hits in the XOY plane
 63 from the primary particles of one time slice are illustrated in the fig. 3,
 64 left, while the corresponding example of track-approximating curve is shown
 65 in the right. The coefficients a_i determine the $z(l)$ dependence for each
 66 primary particle, where l is the arc length of the parabolic segment. It
 67 can be calculated by the simple formula $l = \int_0^{x_0} \sqrt{1 + (2a_1x + a_2)^2} dx =$
 68 $\frac{1}{4a_1} \ln \left(\left| \sqrt{(2a_1x + a_2)^2 + 1} + 2a_1x + a_2 \right| \right) + (2a_1x + a_2) \sqrt{(2a_1x + a_2)^2 + 1}.$

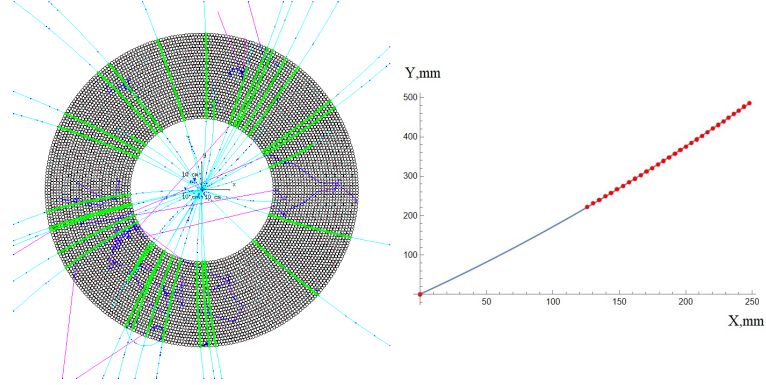


Fig. 3. The hits (green) and tracks (cyan) of one time slice simulation, XOY plane (left). Example of one track approximation in XOY plane (right).

69 Then, $z(l)$ can be approximated by the linear function, which should be ex-
 70 trapolated to the intersection with z -axis to determine the primary vertex
 71 position. The simulated tracks and hits in the ZOY plane from the primary
 particles of one time slice are illustrated in the fig. 4.

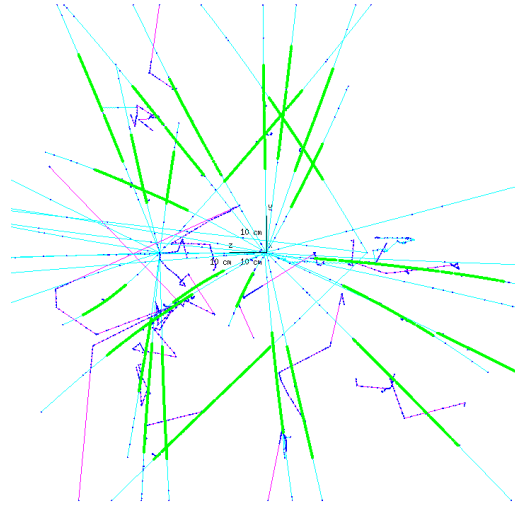


Fig. 4. The hits (green) and tracks (cyan) of one time slice simulation, ZOY plane.

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73 After recovering the starting z -position of each track, name z_0 , it is nec-
 74 essary to separate the tracks into clusters with common vertices and evaluate
 75 the vertex recovery efficiency. The set of initial z_0 values was sorted in as-
 76 cending order, then divided into the groups starting from some boundary z_0
 77 with step H . In the beginning, the boundary z_0 was the smallest z_0 in the
 78 whole set, and then the smallest one from the remaining set. Within a track
 79 group we determine the mean value of z_0 , name z_c , this will be the expected
 80 vertex of the group of tracks which z_0 lie close enough to the mean value.
 81 Within a step we can distinguish only one vertex. Then we calculated the
 82 distance from the z_c to the true vertices z_T of the tracks belonging to this
 83 group. The efficiency was calculated as the ratio of the number of correctly
 84 recovered true vertices to the total number of true vertices. A correctly re-

85 constructed vertex is considered to be the vertex that can be distinguished
 86 on the interval H , and all the assigned tracks actually belong to this vertex.
 87 The distribution of the distance between the reconstructed vertex and the
 88 true vertex is presented in fig. 5. We analysed a set of step values H to
 89 achieve the maximum efficiency. At the current moment the achieved recov-
 90 ery efficiency value is 61%. However, we imposed the most strong conditions
 91 on the track recovery purity.

92 In practice, the requirements on the purity of vertex recovery can be
 93 reduced if we impose restrictions on the energies of the particles which tra-
 94 jectories to be fitted, that will improve the efficiency. Thus, we can separate
 95 the tracks by their z_0 and then combine them into the clusters with common
 vertices, which will correspond to the separate bunch crossings points.

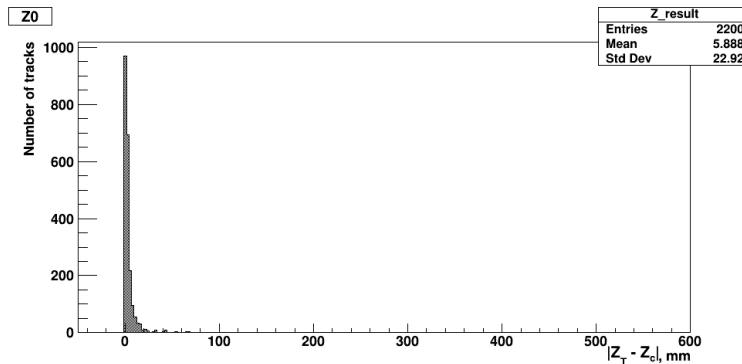


Fig. 5. The distribution of the distance between the reconstructed vertex (z_c) and the true vertex (z_T), the step $H = 6$ cm.

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97 Conclusions

98 In this study we created a simplified model of the straw tracker of SPD
 99 NICA detector using GEANT4 software tools. Introducing the hit collec-
 100 tions, we studied the temporal structure of the events, and found a significant
 101 overlap of straw tubes response times from different bunch crossings. Using
 102 the hits collection data we developed an algorithm for primary vertex recov-
 103 ery with current efficiency of 61%, obtained with the most strict conditions
 104 on the purity of reconstruction. These results are to be a part of prototype
 105 for the online data processing software.

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