

ARieL: Physics at Future e^+e^- colliders

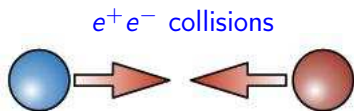
Renat Sadykov
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on behalf of the SANC team

Helmholtz International Summer School
"Modern Colliders - Theory and Experiments 2018"
&
International Workshop
"Calculations for Modern and Future Colliders (CALC2018)"
Dubna, 23.07.2018

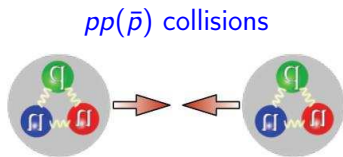
Outline

- Motivation and goals
- Description of the project
- Polarized Bhabha scattering $e^+e^- \rightarrow e^+e^-$ at NLO EW
- Preliminary results at NLO EW for other processes with polarized e^+e^- beams
- Conclusions and plans

Motivation: choice of collider types



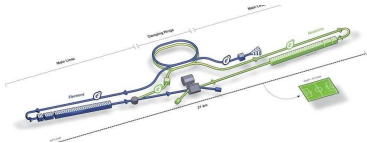
- Point-like particles
- Total annihilation: initial state known
- Decent background
- Limited in energy, but — **precision!**



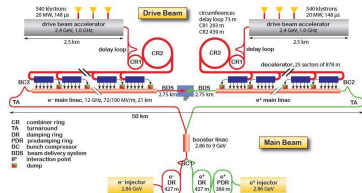
- Composite particles
- Random energy of the hard interaction
- High background
- High energy — **discovery!**

Projects of e^+e^- collider

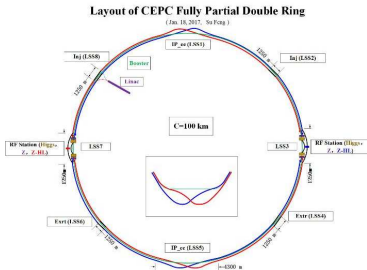
ILC (2030) 250 GeV



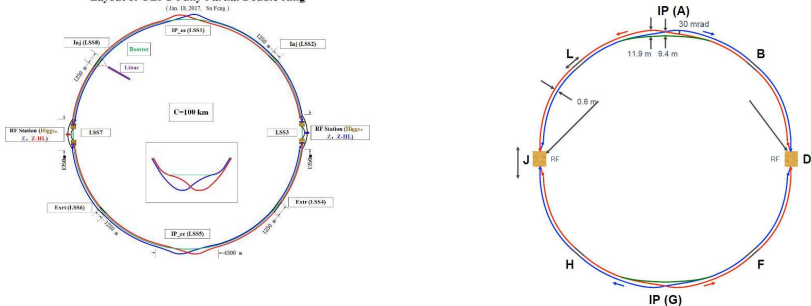
CLIC (2035) 380-3000 GeV



CEPC (2030) 240 GeV



FCC-ee (2039) 350 GeV



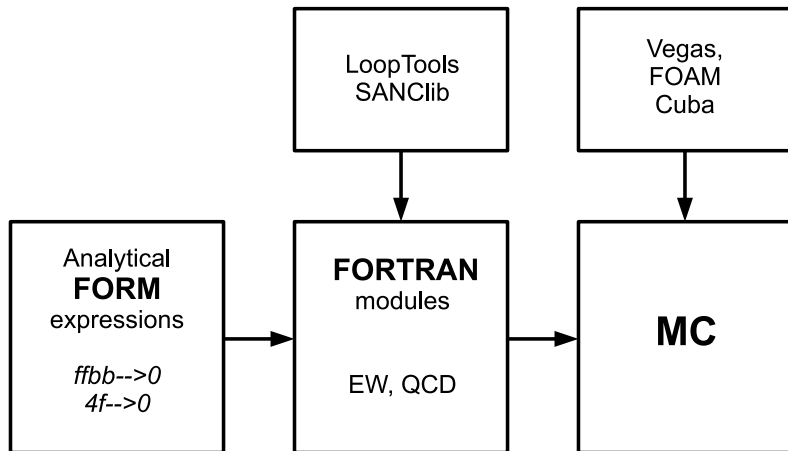
Why ARIeL?

- **A**dvanced
- **R**esearch of
- **I**nteractions in
- e^+e^-
- **coL**lisions

Goals

- Preparation of CLIC research program:
 - Precision study of $e^+e^- \rightarrow \gamma\gamma$ and setting limits on the NP models
 - Precision measurement of the Higgs boson mass M_H
 - Determination of top quark polarization
 - Measurement of $\gamma\gamma \rightarrow W^+W^-$ and $\gamma\gamma \rightarrow ZZ$ and search for anomalous quaring coupling
- Theoretical support:
 - Create e^+e^- Monte Carlo generator with polarization at **complete one-loop EW** and **leading multiloop** for processes $e^+e^- \rightarrow e^+e^- (\mu\mu, \tau\tau, tt, HZ, H\gamma, Z\gamma, ZZ, H\nu\nu, H\mu\mu, ff\gamma, \gamma\gamma)$
 - Interface NLO EW RC to PYTHIA8
 - Implement a single-resonance approach to complex processes
 - Elaborate the standard procedure for $2 \rightarrow 3, 4$ helicity amplitudes
 - Create building blocks for complete EW 2-loops and QCD 3-loops, plus leading EW 3-loops and QCD 4-loops

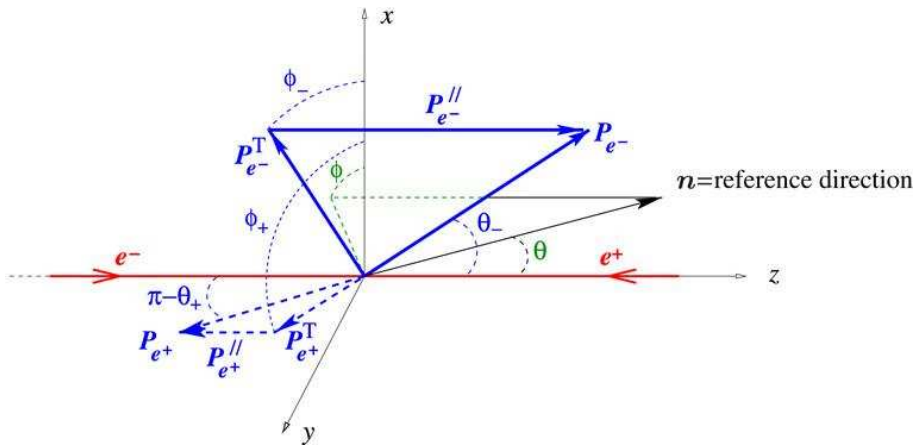
The calculation framework scheme



Processes with polarized beams in MCSANC

- NLO EW corrections for polarized e^+e^- scattering:
 - Bhabha scattering (D.Bardin et al. Phys.Rev. D98 (2018) no.1, 013001)
 - $e^+e^- \rightarrow \mu^+\mu^-$, $e^+e^- \rightarrow \tau^+\tau^-$ (preliminary results)
 - $e^+e^- \rightarrow Z\gamma$ (preliminary results)
 - $e^+e^- \rightarrow t\bar{t}$ (in progress)
 - $e^+e^- \rightarrow ZH$ (in progress)
 - $e^+e^- \rightarrow \gamma\gamma$ (in progress)
 - $e^+e^- \rightarrow ZZ$ (in progress)
 - $e^+e^- \rightarrow f\bar{f}\gamma$ (future plans)
 - $e^+e^- \rightarrow f\bar{f}H$ (future plans)
- NLO EW corrections for polarized $\gamma\gamma$ scattering:
 - $\gamma\gamma \rightarrow \gamma\gamma$ (future plans)
 - $\gamma\gamma \rightarrow Z\gamma$ (future plans)
 - $\gamma\gamma \rightarrow ZZ$ (future plans)

Decomposition of the e^\pm polarization vector



G. Moortgat-Pick et al. Phys. Rept. 460 (2008) 131–243

Matrix element squared

$$\begin{aligned}
 |\mathcal{M}|^2 = & L_{e^-}^{\parallel} R_{e^+}^{\parallel} |\mathcal{H}_{-+}|^2 + R_{e^-}^{\parallel} L_{e^+}^{\parallel} |\mathcal{H}_{+-}|^2 + L_{e^-}^{\parallel} L_{e^+}^{\parallel} |\mathcal{H}_{--}|^2 + R_{e^-}^{\parallel} R_{e^+}^{\parallel} |\mathcal{H}_{++}|^2 \\
 & - \frac{1}{2} P_{e^-}^T P_{e^+}^T \operatorname{Re} \left[e^{i(\Phi_+ - \Phi_-)} \mathcal{H}_{++} \mathcal{H}_{--}^* + e^{i(\Phi_+ + \Phi_-)} \mathcal{H}_{+-} \mathcal{H}_{-+}^* \right] \\
 & + P_{e^-}^T \operatorname{Re} \left[e^{i\Phi_-} \left(L_{e^+}^{\parallel} \mathcal{H}_{+-} \mathcal{H}_{--}^* + R_{e^+}^{\parallel} \mathcal{H}_{++} \mathcal{H}_{-+}^* \right) \right] \\
 & - P_{e^+}^T \operatorname{Re} \left[e^{i\Phi_+} \left(L_{e^-}^{\parallel} \mathcal{H}_{-+} \mathcal{H}_{--}^* + R_{e^-}^{\parallel} \mathcal{H}_{++} \mathcal{H}_{+-}^* \right) \right],
 \end{aligned}$$

where

$$L_{e^\pm}^{\parallel} = \frac{1}{2}(1 - P_{e^\pm}^{\parallel}), \quad R_{e^\pm}^{\parallel} = \frac{1}{2}(1 + P_{e^\pm}^{\parallel}), \quad \Phi_{\pm} = \phi_{\pm} - \phi,$$

$\mathcal{H}_{--}, \mathcal{H}_{++}, \mathcal{H}_{-+}, \mathcal{H}_{+-}$ — helicity amplitudes.

Longitudinally-polarized beams

With longitudinally-polarized beams, cross-sections at an e^+e^- collider can be subdivided into four parts:

$$\sigma_{P_{e^-}P_{e^+}} = R_{e^-}R_{e^+}\sigma_{RR} + L_{e^-}L_{e^+}\sigma_{LL} + R_{e^-}L_{e^+}\sigma_{RL} + L_{e^-}R_{e^+}\sigma_{LR},$$

where σ_{RL} stands for the cross-section if the e^- beam is completely right-handed polarized ($P_{e^-} = +1$) and the e^+ beam is completely left-handed polarized ($P_{e^+} = -1$). The cross-sections σ_{LR} , σ_{RR} and σ_{LL} are defined analogously.

Polarized Bhabha scattering at one-loop

Here we present complete one-loop EW corrections to Bhabha scattering $e^+e^- \rightarrow e^+e^-$ with longitudinally polarized initial particles.

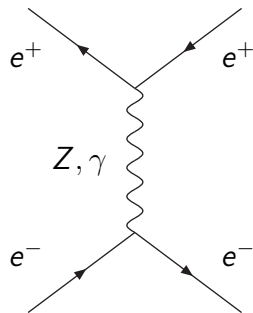
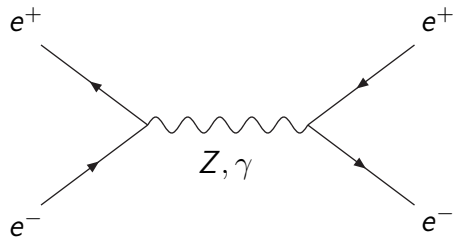
The cross-section of this process at one-loop can be divided into four parts:

$$\sigma^{1\text{-loop}} = \sigma^{\text{Born}} + \sigma^{\text{virt}}(\lambda) + \sigma^{\text{soft}}(\lambda, \omega) + \sigma^{\text{hard}}(\omega),$$

where σ^{Born} — Born level cross-section, σ^{virt} — contribution of virtual(loop) corrections, σ^{soft} — contribution due to soft photon emission, σ^{hard} — contribution due to hard photon emission (with energy $E_\gamma > \omega \frac{\sqrt{s}}{2}$).

Auxiliary parameters λ ("photon mass") and ω cancel out after summation.

Bhabha scattering: Born-level diagrams



Polarized Bhabha scattering: HA for Born and Virtual parts

At one-loop level we have six non-zero HAs (four independent):

$$\begin{aligned}
 \mathcal{H}_{++++} &= \mathcal{H}_{----} = -2e^2 \frac{s}{t} \left(\mathcal{F}_{QQ}^{(\gamma, Z)}(t, s, u) - \chi_Z^t \delta_e \mathcal{F}_{QL}^Z(t, s, u) \right), \\
 \mathcal{H}_{+--+} &= \mathcal{H}_{-+-+} = -2e^2 \frac{t}{s} \left(\mathcal{F}_{QQ}^{(\gamma, Z)}(s, t, u) - \chi_Z^s \delta_e \mathcal{F}_{QL}^Z(s, t, u) \right), \\
 \mathcal{H}_{+---} &= 2e^2 \left(\frac{u}{s} \left[\mathcal{F}_{QQ}^{(\gamma, Z)}(s, t, u) + \chi_Z^s (\mathcal{F}_{LL}^Z(s, t, u) - 2\delta_e \mathcal{F}_{QL}^Z(s, t, u)) \right] \right. \\
 &\quad \left. + \frac{u}{t} \left[\mathcal{F}_{QQ}^{(\gamma, Z)}(t, s, u) + \chi_Z^t (\mathcal{F}_{LL}^Z(t, s, u) - 2\delta_e \mathcal{F}_{QL}^Z(t, s, u)) \right] \right), \\
 \mathcal{H}_{-++-} &= 2e^2 \left(\frac{u}{s} \mathcal{F}_{QQ}^{(\gamma, Z)}(s, t, u) + \frac{u}{t} \mathcal{F}_{QQ}^{(\gamma, Z)}(t, s, u) \right),
 \end{aligned}$$

where

$$\chi_Z^s = \frac{1}{4s_W^2 c_W^2} \frac{s}{s - M_Z^2 + iM_Z \Gamma_Z}, \quad \chi_Z^t = \frac{1}{4s_W^2 c_W^2} \frac{t}{t - M_Z^2}, \quad \delta_e = v_e - a_e = 2s_W^2,$$

$$\mathcal{F}_{QQ}^{(\gamma, Z)}(a, b, c) = \mathcal{F}_{QQ}^\gamma(a, b, c) + \chi_Z^a \delta_e^2 \mathcal{F}_{QQ}^Z(a, b, c).$$

We get the Born level HAs by replacing $\mathcal{F}_{LL}^Z \rightarrow 1$, $\mathcal{F}_{QL}^Z \rightarrow 1$, $\mathcal{F}_{QQ}^Z \rightarrow 1$ and $\mathcal{F}_{QQ}^\gamma \rightarrow 1$.

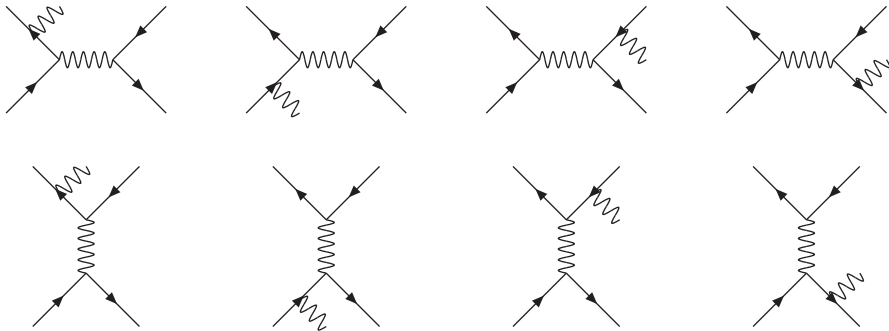
Polarized Bhabha scattering: soft photon contribution

The soft photon contribution contains the infrared divergences which compensate the infrared divergences of the one-loop QED corrections.

This soft photon correction can be calculated analytically and to be factorized to Born cross section. The polarization dependence is contained in σ^{Born} .

$$\begin{aligned}\sigma^{\text{soft}}(\lambda, \omega) &= -\sigma^{\text{Born}} \frac{\alpha}{\pi} \left\{ \left(1 + \ln \left(\frac{m_e^2}{s} \right) \right)^2 + \ln \left(-\frac{u}{s} \right)^2 - \ln \left(-\frac{t}{s} \right)^2 \right. \\ &\quad - 2\text{Li}_2 \left(-\frac{u}{s} \right) + 2\text{Li}_2 \left(-\frac{t}{s} \right) + 4\text{Li}_2(1) \\ &\quad \left. - 1 + 2 \ln \left(\frac{4\omega^2}{\lambda} \right) \left[1 + \ln \left(\frac{m_e^2}{s} \right) - \ln \left(\frac{t}{u} \right) \right] \right\}.\end{aligned}$$

Polarized Bhabha scattering: hard photon contribution



For HA for hard bremsstrahlung see presentation by [Yahor Dydyshka](#)

Polarized Bhabha scattering: Monte Carlo generator

We created Monte Carlo generator of unweighted events for the polarized Bhabha scattering $e^+e^- \rightarrow e^+e^-$ with complete one-loop EW corrections.

This generator uses adaptive algorithm **mFOAM** [CPC 177:441-458,2007] which is a part of **ROOT** [<https://root.cern.ch>] program.

It will be interfaced to **PYTHIA8** [CPC 178 (2008) 852–867] program.

Setup for tuned comparison

We performed a tuned comparison of our results for polarized Born and hard Bremsstrahlung with the results **WHIZARD** [Eur.Phys.J.C71 (2011) 1742] program. The contributions of soft and virtual parts were compared with the results of **a1TALC** [CPC 174 (2006) 71-82] program

Input parameters:

$$\alpha^{-1}(0) = 137.03599976,$$

$$M_W = 80.4514958 \text{ GeV}, \quad M_Z = 91.1876 \text{ GeV}, \quad \Gamma_Z = 2.49977 \text{ GeV},$$

$$m_e = 0.51099907 \text{ MeV}, \quad m_\mu = 0.105658389 \text{ GeV}, \quad m_\tau = 1.77705 \text{ GeV},$$

$$m_d = 0.083 \text{ GeV}, \quad m_s = 0.215 \text{ GeV}, \quad m_b = 4.7 \text{ GeV},$$

$$m_u = 0.062 \text{ GeV}, \quad m_c = 1.5 \text{ GeV}, \quad m_t = 173.8 \text{ GeV}.$$

Cuts:

$$|\cos\theta| < 0.9,$$

$$E_\gamma > 1 \text{ GeV} \quad (\text{for comparison of hard Bremsstrahlung}).$$

$e^+e^- \rightarrow e^+e^-$: **WHIZARD** vs **MCSANC** (Born)

P_{e^-}, P_{e^+}	0, 0	-0.8, 0	-0.8, -0.6	-0.8, 0.6
$\sqrt{s} = 250 \text{ GeV}$				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	56.677(1)	57.774(1)	56.272(1)	59.276(1)
$\sigma_{e^+e^-}^{\text{Born}}$, pb	56.677(1)	57.775(1)	56.272(1)	59.275(1)
$\sqrt{s} = 500 \text{ GeV}$				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	14.379(1)	15.030(1)	12.706(1)	17.355(1)
$\sigma_{e^+e^-}^{\text{Born}}$, pb	14.379(1)	15.030(1)	12.706(1)	17.354(1)
$\sqrt{s} = 1000 \text{ GeV}$				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	3.6792(1)	3.9057(1)	3.0358(1)	4.7756(1)
$\sigma_{e^+e^-}^{\text{Born}}$, pb	3.6792(1)	3.9057(1)	3.0358(1)	4.7755(1)

$e^+e^- \rightarrow e^+e^-$: **WHIZARD** vs **MCSANC** (hard)

P_{e^-}, P_{e^+}	0, 0	-0.8, 0	-0.8, -0.6	-0.8, 0.6
$\sqrt{s} = 250$ GeV				
$\sigma_{e^+e^-}^{\text{hard}}$, pb	48.62(1)	49.58(1)	48.74(1)	50.40(1)
$\sigma_{e^+e^-}^{\text{hard}}$, pb	48.65(1)	49.56(1)	48.78(1)	50.44(1)
$\sqrt{s} = 500$ GeV				
$\sigma_{e^+e^-}^{\text{hard}}$, pb	15.14(1)	15.81(1)	13.54(1)	18.07(1)
$\sigma_{e^+e^-}^{\text{hard}}$, pb	15.12(1)	15.79(1)	13.55(1)	18.11(2)
$\sqrt{s} = 1000$ GeV				
$\sigma_{e^+e^-}^{\text{hard}}$, pb	4.693(1)	4.976(1)	3.912(1)	6.041(1)
$\sigma_{e^+e^-}^{\text{hard}}$, pb	4.694(1)	4.975(1)	3.913(1)	6.043(1)

$e^+e^- \rightarrow e^+e^-$: **a**i**TALC** vs **MCSANC** (virtual+soft)

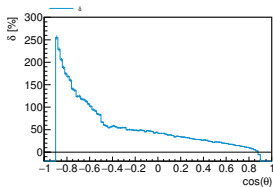
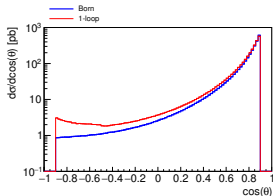
$\cos \theta$	$\sigma_{e^+e^-}^{\text{Born}}$, pb	$\sigma_{e^+e^-}^{\text{Born+virt+soft}}$, pb
-0.9	$2.16999 \cdot 10^{-1}$ $2.16999 \cdot 10^{-1}$	$1.93445 \cdot 10^{-1}$ $1.93445 \cdot 10^{-1}$
-0.5	$2.61360 \cdot 10^{-1}$ $2.61360 \cdot 10^{-1}$	$2.38707 \cdot 10^{-1}$ $2.38707 \cdot 10^{-1}$
0	$5.98142 \cdot 10^{-1}$ $5.98142 \cdot 10^{-1}$	$5.46677 \cdot 10^{-1}$ $5.46677 \cdot 10^{-1}$
+0.5	$4.21273 \cdot 10^0$ $4.21273 \cdot 10^0$	$3.81301 \cdot 10^0$ $3.81301 \cdot 10^0$
+0.9	$1.89160 \cdot 10^2$ $1.89160 \cdot 10^2$	$1.72928 \cdot 10^2$ $1.72928 \cdot 10^2$
+0.99	$2.06556 \cdot 10^4$ $2.06555 \cdot 10^4$	$1.90607 \cdot 10^4$ $1.90607 \cdot 10^4$
+0.999	$2.08236 \cdot 10^6$ $2.08236 \cdot 10^6$	$1.91624 \cdot 10^6$ $1.91624 \cdot 10^6$
+0.9999	$2.08429 \cdot 10^8$ $2.08429 \cdot 10^8$	$1.91402 \cdot 10^8$ $1.91402 \cdot 10^8$

$e^+e^- \rightarrow e^+e^-$: Born vs 1-loop

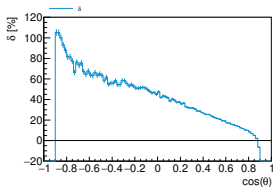
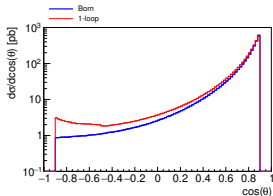
P_{e^-}, P_{e^+}	0, 0	-0.8, 0	-0.8, -0.6	-0.8, 0.6
$\sqrt{s} = 250 \text{ GeV}$				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	56.6763(1)	57.7738(1)	56.2725(4)	59.2753(5)
$\sigma_{e^+e^-}^{\text{1-loop}}$, pb	61.731(6)	62.587(6)	61.878(6)	63.287(7)
δ , %	8.92(1)	8.33(1)	9.96(1)	6.77(1)
$\sqrt{s} = 500 \text{ GeV}$				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	14.3789(1)	15.0305(1)	12.7061(1)	17.3550(2)
$\sigma_{e^+e^-}^{\text{1-loop}}$, pb	15.465(2)	15.870(2)	13.861(1)	17.884(2)
δ , %	7.56(1)	5.59(1)	9.09(1)	3.05(1)
$\sqrt{s} = 1000 \text{ GeV}$				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	3.67921(1)	3.90568(1)	3.03577(3)	4.77562(5)
$\sigma_{e^+e^-}^{\text{1-loop}}$, pb	3.8637(4)	3.9445(4)	3.2332(3)	4.6542(7)
δ , %	5.02(1)	0.99(1)	6.50(1)	-2.54(1)

$e^+e^- \rightarrow e^+e^-$: distributions in $\cos\theta$

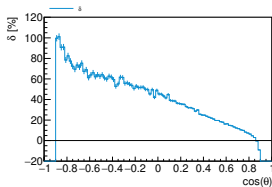
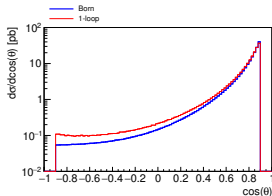
$\sqrt{s} = 250$ GeV



$\sqrt{s} = 500$ GeV



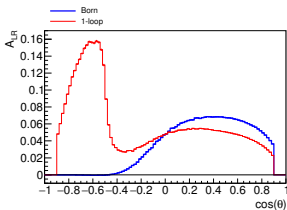
$\sqrt{s} = 1000$ GeV



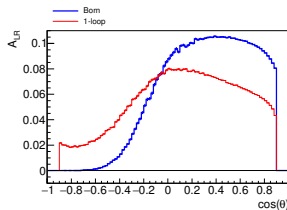
$e^+e^- \rightarrow e^+e^-$: A_{LR} distributions in $\cos\theta$

$$A_{LR} = \frac{\sigma_{LR} - \sigma_{RL}}{\sigma_{LR} + \sigma_{RL}}$$

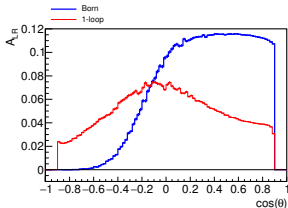
$\sqrt{s} = 250$ GeV



$\sqrt{s} = 500$ GeV



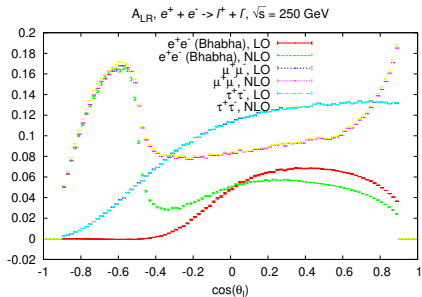
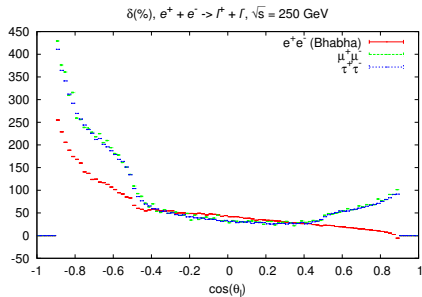
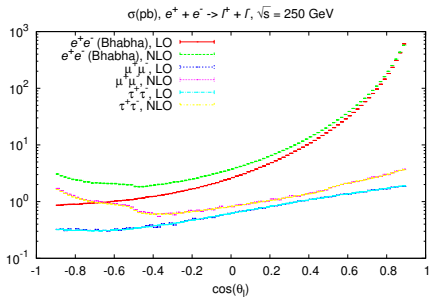
$\sqrt{s} = 1000$ GeV



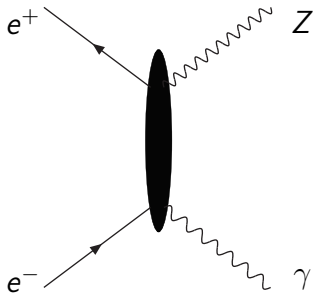
$e^+e^- \rightarrow \ell^+\ell^-$: Born vs 1-loop (preliminary)

P_{e^-}, P_{e^+}	0, 0	-0.8, 0	-0.8, -0.6	-0.8, 0.6
$\sqrt{s} = 250 \text{ GeV}$				
$e^- + e^+ \rightarrow e^- + e^+$ (Bhabha)				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	56.6763(1)	57.7738(1)	56.2725(4)	59.2753(5)
$\sigma_{e^+e^-}^{\text{1-loop}}$, pb	61.731(6)	62.587(6)	61.878(6)	63.287(7)
δ , %	8.92(1)	8.33(1)	9.96(1)	6.77(1)
$e^- + e^+ \rightarrow \mu^- + \mu^+$				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	1.4174(1)	1.5462(1)	0.7690(2)	2.3231(2)
$\sigma_{e^+e^-}^{\text{1-loop}}$, pb	2.3987(2)	2.616(2)	1.3015(1)	3.929(1)
δ , %	69.22(2)	69.13(1)	69.18(2)	69.12(1)
$e^- + e^+ \rightarrow \tau^- + \tau^+$				
$\sigma_{e^+e^-}^{\text{Born}}$, pb	1.4174(1)	1.5461(1)	0.7692(1)	2.3230(1)
$\sigma_{e^+e^-}^{\text{1-loop}}$, pb	2.3609(1)	2.5773(1)	1.2817(1)	3.8728(2)
δ , %	66.56(1)	66.69(1)	66.62(1)	66.71(1)

$e^+e^- \rightarrow l^+l^-$: distributions in $\cos\theta_l$ (preliminary)



$$e^+e^- \rightarrow Z\gamma$$



The expressions for HA and the results for unpolarized case were published in [Eur.Phys.J.C54:187-197,2008](#)

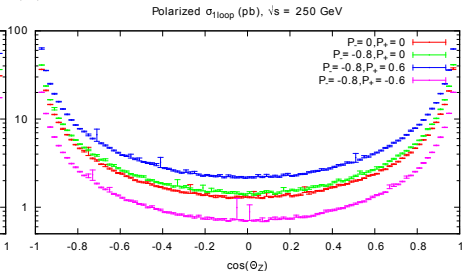
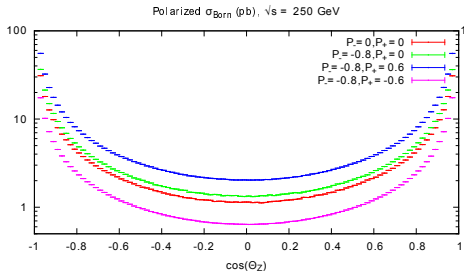
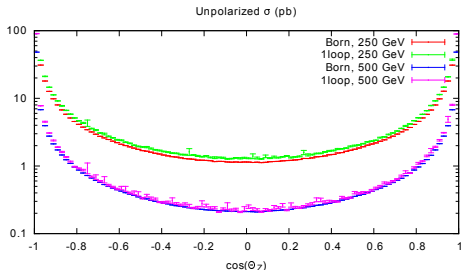
Here we present the preliminary results for 1-loop corrections taking into account the effect of polarization of e^+e^- beams.

$e^+e^- \rightarrow Z\gamma$: Born vs 1-loop (preliminary)

\sqrt{s}	250	500	1000
$\sigma_{Z\gamma}^{\text{Born}}$, pb	15.7038(6)	3.3858(3)	0.81958(3)
$\sigma_{Z\gamma}^{\text{1-loop}}$, pb	24.37(1)	5.23(6)	1.237(3)
δ , %	55.20(6)	54.4(2)	50.9(4)

P_{e^-}, P_{e^+}	0, 0	-0.8, 0	-0.8, 0.6	-0.8, -0.6
$\sqrt{s} = 250$ GeV				
$\sigma_{Z\gamma}^{\text{Born}}$, pb	15.7038(6)	18.520(4)	28.174(2)	8.870(3)
$\sigma_{Z\gamma}^{\text{1-loop}}$, pb	24.37(1)	28.00(1)	42.13(2)	13.53(1)
δ , %	55.20(6)	51.11(8)	49.55(7)	52.57(9)

$e^+e^- \rightarrow Z\gamma$: distributions in $\cos\theta_Z$ (preliminary)



Conclusions and plans

- We are building a powerful team of experimentalists and theoreticians that will prepare the physics research at the next generation e^+e^- collider
- We created the FORTRAN modules for polarized Bhabha scattering with complete one-loop EW corrections. Based on these modules Monte Carlo generator [[Phys.Rev. D98 \(2018\) no.1, 013001](#)] of unweighted events was created with possibility to produce events in Les Houches Event format
- Preliminary results for polarized $e^+e^- \rightarrow \mu^+\mu^-$, $e^+e^- \rightarrow \tau^+\tau^-$ and $e^+e^- \rightarrow Z\gamma$ were presented. More processes will be included in MC generator in the future