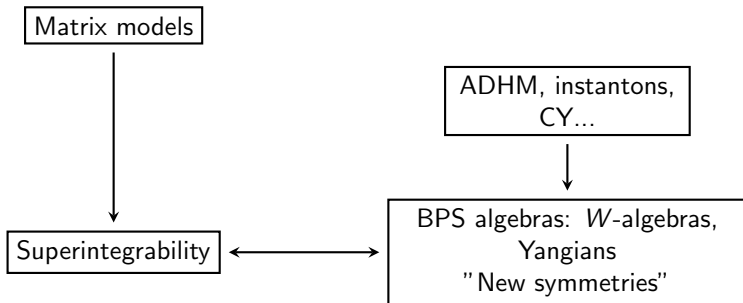


# Superintegrability of $\beta$ -deformed monomial matrix models and Uglov polynomials

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# Generalities



# Superintegrability of matrix models

Matrix models:

$$\langle f(X) \rangle = \int DX e^{V(X)} f(X)$$

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As an eigenvalue integral (for the Hermitian case)

$$\langle f(x) \rangle = \int \prod_{i=1}^N dx_i \Delta^2(x) \exp\left(\sum_{i=1}^N V(x_i)\right) f(x)$$

$$\Delta(x) = \prod_{i < j=1}^N (x_i - x_j)$$

# Superintegrability of matrix models

Sometimes expectation values can be computed exactly [[A.Mironov, A.Morozov \(2022\)](#)]:

$$\langle \text{Character} \rangle \sim \text{Character}$$

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$$\langle f(x) \rangle \sim f(*)$$

## Superintegrability monomial potentials

Among many - monomial matrix models are exactly solvable:

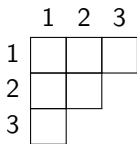
$$\int DX \exp\left(-\text{Tr} \frac{X^s}{s}\right) \dots$$



# Superintegrability monomial potentials

Simplest case - Gaussian matrix model,  $\text{Tr } V(X) = -\frac{1}{2} \text{Tr } X^2$ :

$$\begin{aligned} \left\langle \underbrace{S_R(\text{Tr } X^k)}_{\text{Schur functions}} \right\rangle &= \left( \prod_{(i,j) \in R} (N + j - i) \right) S_R \{p_k = \delta_{k,2}\} = \\ &= \frac{S_R \{p_k = N\}}{\underbrace{S_R \{p_k = \delta_{k,1}\}}_{\text{Schur functions}}} S_R \{p_k = \delta_{k,2}\} \end{aligned}$$



$$\rightarrow N^2(N+1)^2(N+2)(N-1)(N-2)$$

# Superintegrability of monomial potentials

Non-Gaussian models:

$$\int DX \exp\left(-\text{Tr} \frac{X^s}{s}\right) \dots$$

**?**

# Superintegrability of monomial potentials

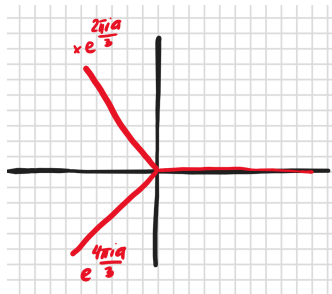
Non-Gaussian models:

$$\int_{C_{s,a}^{\otimes N}} \prod_{i=1}^N dx_i \Delta^2(x) \exp\left(-\sum_{i=1}^N \frac{x_i^s}{s}\right) \dots$$

$$C_{s,a} = \sum_{b=0}^{s-1} e^{\frac{2\pi iab}{s}} \times \left(0, e^{\frac{2\pi ib}{s}} \infty\right)$$

# Superintegrability of monomial potentials

Non-Gaussian models:



Contours in the cubic model

# Superintegrability of monomial potentials

Non-Gaussian models:

$$\int_{\mathbb{C}_{s,a}^{\otimes N}} \prod_{i=1}^N dx_i \Delta^2(x) \exp\left(-\sum_{i=1}^N \frac{x_i^s}{s}\right) \dots$$

$$\langle S_R(x) \rangle_{s,a} = \left( \prod_{(i,j) \in R} [[N+j-i]]_{s,a} [[N+j-i]]_{s,0} \right) S_R \{p_k = \delta_{k,s}\}$$

where  $[[n]]_{s,a} = n$  if  $n \bmod s = a$  and 1 otherwise. Also,  $N = 0, a \bmod s$  (for concreteness choose 0)

[C.Cordova, B. Heidenreich, A. Popolitov, S. Shakirov (2018)]

# Superintegrability of monomial potentials

Non-Gaussian models:

$$[[N + j - i]]_{3,0} : \begin{array}{|c|c|c|c|c|} \hline \text{tan} & & & \text{tan} & \\ \hline & \text{tan} & & & \\ \hline & & \text{tan} & & \\ \hline \text{tan} & & & & \\ \hline & \text{tan} & & & \\ \hline \end{array} \rightarrow N^3(N+3)(N-3)^2$$

$$[[N + j - i]]_{3,1} : \begin{array}{|c|c|c|c|} \hline & \text{dark purple} & & \text{dark purple} \\ \hline & & \text{dark purple} & \\ \hline \text{dark purple} & & & \\ \hline & \text{dark purple} & & \\ \hline & & & \\ \hline \end{array} \rightarrow (N+1)^2(N+4)(N-2)^2$$

## $\beta$ -deformation

The Gaussian model is known to survive a deformation of the measure:

$$\Delta^2(x) \rightarrow \Delta^{2\beta}(x)$$

This corresponds to:

$$S_R \rightarrow J_R - \text{Jack polynomials}$$

and

$$\langle J_R(x) \rangle = \frac{J_R\{N\}}{\underbrace{J_R\{\delta_{k,1}\}}_{\prod_{(i,j) \in R} (\beta N + j - \beta i)}} J_R\{\delta_{k,2}\}$$

## $\beta$ -deformation

$\beta$ -deformation is consistent among many subjects :

- "Temperature" in statistical physics, interpolation between  $\beta = 1/2, 1, 2$
- Central charge in CFT  $c = 1 - 6 \left( \sqrt{\beta} - \frac{1}{\sqrt{\beta}} \right)^2$
- AGT-related equivariant parameters in Nekrasov functions:  
 $\beta = -\frac{\epsilon_1}{\epsilon_2}$
- Non trivial coupling  $g = \beta(\beta - 1)$  in Calogero model. Jack polynomials - eigenfunctions.
- more ...



# $\beta$ -deformation

Is there a  $\beta$ -deformation of Non-Gaussian superintegrability?

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$$\langle S_R(x) \rangle_{s,a} = \left( \prod_{(i,j) \in R} [[N+j-i]_{s,a} [[N+j-i]_{s,0}] \right) S_R \{p_k = \delta_{k,s}\}$$

# $\beta$ -deformation

Naive Guess does not work!

$$\Delta^2 \rightarrow \Delta^{2\beta}$$

$$\langle J_R(x) \rangle_{s,a} = \text{complete mess}$$

## $\beta$ -deformation

The problem is we are doing an *uneducated guess* but need to do an *educated guess*

We need some symmetry argument. The symmetry appears to be not the *internal* for the matrix model, but the symmetry of it's *moduli space*

## $\beta$ -deformation

We need to cook up a " deformed" partition function

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In QFT terms:

$$Z(T_{O_i}) = \int D\Phi \exp \left( -S[\Phi] + \sum_i \int T_{O_i} O_i[\Phi(x)] \right)$$

## $\beta$ -deformation

We need to cook up a " deformed" partition function

In matrix models:

$$\begin{aligned} Z(p_k) &= \int DX \exp \left( \text{Tr} V(X) + \sum_{k=1}^{\infty} \frac{p_k}{k} \text{Tr} X^k \right) = \\ &= \sum_R \langle S_R(\text{Tr} X^k) \rangle_{S_R(p)} \end{aligned}$$

## $\beta$ -deformation

Symmetry acts on the  $p_k$  parameters. It's main realization is the  $W$ -representation:

$$Z(p_k) = e^{W\left[p_k, \frac{\partial}{\partial p_k}\right]} \cdot 1$$



## $\beta$ -deformation

It is intimately connected to superintegrability

$$W_n[g] \cdot S_R(p) = \sum_{\substack{Q: \\ |Q|=|R|+n}} \prod_{(i,j) \in Q/R} g(j-i) \langle p_n S_R | S_Q \rangle S_Q(p)$$

$$\langle S_R(x) \rangle = \left( \prod_{(i,j) \in R} g(j-i) \right) S_R \{p_k = \delta_{k,n}\}$$

## $\beta$ -deformation

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$$W_n[g] \cdot S_R(p) = \sum_{\substack{Q: \\ |Q|=|R|+n}} \prod_{(i,j) \in Q/R} g(j-i) \langle p_n S_R | S_Q \rangle S_Q(p)$$

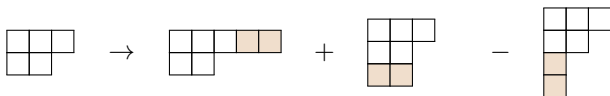
These operators are nicely  $\beta$  deformed :

$$W_n[g] \cdot J_R(p) = \sum_{\substack{Q: \\ |Q|=|R|+n}} \prod_{(i,j) \in Q/R} g(j-\beta i) \langle p_n J_R | J_Q \rangle J_Q(p)$$

# $\beta$ -deformation

In the Gaussian model:

$$W_2 S_R = \sum_{R+[2]} (N + j_{\square} - i_{\square})(N + j_{\square} + 1 - i_{\square}) S_{R+[2]} - \sum_{R+[1,1]} (N + j_{\square} - i_{\square})(N + j_{\square} - i_{\square} - 1) S_{R+[1,1]}$$



## $\beta$ -deformation

$\beta$ -deformation of  $W$ -operators can be deduced from algebra.

For the Gaussian model it corresponds to the deformation:

$$\mathcal{W}_{1+\infty} \longrightarrow Y(\hat{\mathfrak{gl}}_1)$$

## $\beta$ -deformation

Generically it is known, at least for Schur functions, how to reverse engineer operator from function:

$$\{g(j-i), n\} \longrightarrow W_n[g] \in \mathcal{W}_{1+\infty}$$

We can try to do this for non-Gaussian model. What we find is quite complicated, but the important part is that it corresponds to quantum toroidal DIM algebra, with  $(q, t)$  parameters at:

$$(q, t) = (\omega_s, \omega_s), \quad \omega_s = e^{\frac{2\pi i}{s}}$$

## $\beta$ -deformation

$$V(q) = \oint dz : \frac{e^{\varphi(qz) - \varphi(z)}}{q - 1} : \quad - \text{ at } q = \omega_s$$

$$\varphi(z) = \sum_{k=1}^{\infty} z^{-k} p_k + \sum_{k=1}^{\infty} z^k \frac{1}{k} \frac{\partial}{\partial p_k}$$

$$V(q)S_R = \left( \sum_{(i,j) \in R} q^{j-i} \right) S_R$$

## $\beta$ -deformation

For the DIM algebra the relevant polynomials are Macdonald polynomials.

At  $q = t$  Macdonald polynomials always reduce to Schur! They are insensitive to the value of  $q = t$ .

But the operators and the deformation are!

Hence we should deform from this special point

## $\beta$ -deformation

Finally, the clue! One should take a deformation near the root of unity point in the  $(q, t)$  plane.

$$q = \omega_r u, \quad t = \omega_r u^\beta, \quad u \rightarrow 1$$



# $\beta$ -deformation

This hints us both at the measure and the symmetric functions.

Measure

$$\Delta_{(q,t)}(x) \rightarrow \Delta_{\text{Uglouv}}^{(r)}(x) = \prod_{i < j}^N (x_i^r - x_j^r)^{\frac{2(\beta-1)}{r}} (x_i - x_j)^2$$

Reduce to ordinary  $\beta$ -deformation and Jack polynomials at  $r = 1$

## $\beta$ -deformation

This hints us both at the measure and the symmetric functions.

Uglov polynomials:

$$\lim_{\omega_r} M_R(p_k) = U_R^r(p_k)$$

For example:

$$U_{[2]}^{(2)}(p_k) = \frac{\beta p_1^2 + p_2}{\beta + 1} \quad U_{[2]}^{(1)}(p_k) = \frac{p_1^2 + \beta p_2}{\beta + 1}$$

Reduce to ordinary  $\beta$ -deformation and Jack polynomials at  $r = 1$

# $\beta$ -deformation

Just as Jack polynomials, Uglov polynomials are nice objects:

- Eigenfunctions for  $\mathfrak{gl}_r$  spin-Calogero system [D.Uglov (1998)]
- Label highest weight vectors in superVirasoro algebra (at  $r = 2$ ) [M.Bershtein, A.Vargulevich (2022)]
- SCFT [V.Belavin, A.Zhakenov (2020)]
- Vectors in representation of  $Y(\hat{\mathfrak{gl}}_r)$  [N.Tselousov, D.Galakhov, A.Morozov (2024)]
- Uglov limit itself is related to ADHM on  $\mathbb{C}^2/\mathbb{Z}_r$  [T. Kimura (2012)].

## $\beta$ -deformation of the monomial matrix model

Finally. Combine everything and put  $r = s$ :

$$\begin{aligned} \left\langle U_R^{(s)}(x) \right\rangle_{s,a}^\beta &= \int_{\mathcal{C}_{s,a}^{\otimes N}} \prod_{i=1}^N dx_i \Delta_{\text{Uglöv}}^{(s)}(x) \exp \left( - \sum_{i=1}^N \frac{x_i^s}{s} \right) U_R^{(s)}(x) = \\ &= \left( \prod_{(i,j) \in R} [[\beta N + j - \beta i]]_{s,a} [[\beta N + (j-1) - \beta(i-1)]]_{s,0} \right) U_R^{(s)} \{ \delta_{k,s} \} \end{aligned}$$

## Discussion and questions?

- Looking at peculiar symmetries of matrix models we are able to deduce the correct basis for a non-trivial extension of superintegrability
- The type of special functions are tied not only to the measure, but also to the potential
- Can this result be deduced from superintegrability in some  $(q, t)$ -model?
- For the Gaussian model we have  $r = 2$  and  $r = 1$ , why is that?
- Somehow we still land at some very special functions
- Limit is related to instanton counting on  $\mathbb{C}^2/\mathbb{Z}_r$  orbifold. Is there a well-defined limit and special functions related to the  $\mathbb{C}^2/\Gamma_{r,\nu}$ -orbifold ?

Extra slides

## The Gaussian model

For  $V(X) = \frac{X^2}{2}$  we actually have two options  $r = s = 2$  and  $r = 1$ .  
Both are superintegrable.

## Example of superintegrability

$$\langle \text{Schur}_R (\text{Tr } X^k) \rangle = \frac{\text{Schur}_R \{N\}}{\text{Schur}_R \{\delta_{k,1}\}} \text{Schur}_R \{\delta_{k,2}\}$$

Where  $\text{Schur}_R$  are Schur functions, for example

$$\text{Schur}_{[2]}(p) = \frac{p_2}{2} + \frac{p_1^2}{2}, \quad p_k = \text{Tr } X^k$$

The r.h.s are Schur functions evaluated at special points

$$\frac{\text{Schur}_R \{p_k = N\}}{\text{Schur}_R \{p_k = \delta_{k,1}\}} = \prod_{(i,j) \in R} (N + j - i)$$



# Schur polynomials

- Generating function:

$$\exp\left(\sum_{k=1}^{\infty} \frac{z^k p_k}{k}\right) = \sum_{k=1}^{\infty} s_k(p) z^k$$

- Determinant form

$$S_R(p) = \det_{i,j} s_{\lambda_i - i + j}(p)$$

- As symmetric functions:

$$p_k = \sum_{i=1}^N x_i^k$$

- As characters for  $GL(N)$ :

$$S_R(p_k = \text{Tr } X^k) = \text{Tr } \rho_R(X)$$

## Jack polynomials

$$J_{[2]} = \frac{\beta p_1^2 + p_2}{\beta + 1}$$

$$J_{[1,1]} = \frac{1}{2} (p_1^2 - p_2)$$

$$J_{[3]} = \frac{\beta^2 p_1^3 + 3\beta p_2 p_1 + 2p_3}{(\beta + 1)(\beta + 2)}$$

$$J_{[2,1]} = \frac{\beta p_1^3 - \beta p_2 p_1 + p_2 p_1 - p_3}{2\beta + 1}$$