

Bjorken Sum Rule With New Analytic Coupling

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Abstract: We analyze recently obtained experimental data for the polarized Bjorken sum rule in the region of small values of Q^2 . Our investigation is based on a new form of coupling constant which doesn't contain the Landau pole. We found an excellent agreement between the experimental data and the predictions of analytic QCD, as well as a strong difference between these data and the results obtained in the framework of perturbative QCD.

Introduction

We introduce the derivatives (in the k -order of perturbation theory (PT)) [1]

$$\tilde{a}_{n+1}^{(k)}(Q^2) = \frac{(-1)^n d^n a_s^{(k)}(Q^2)}{n! (dL)^n}, \quad a_s^{(k)}(Q^2) = \frac{\beta_0 \alpha_s^{(k)}(Q^2)}{4\pi},$$

where $L = \ln \frac{Q^2}{\Lambda^2}$, $Q^2 = -q^2$, q^2 – transferred momentum in the Euclidean domain for spacelike processes. The series of derivatives $\tilde{a}_n(Q^2)$ can successfully replace the corresponding series of $a_s(Q^2)$ -powers.

Analytic Coupling

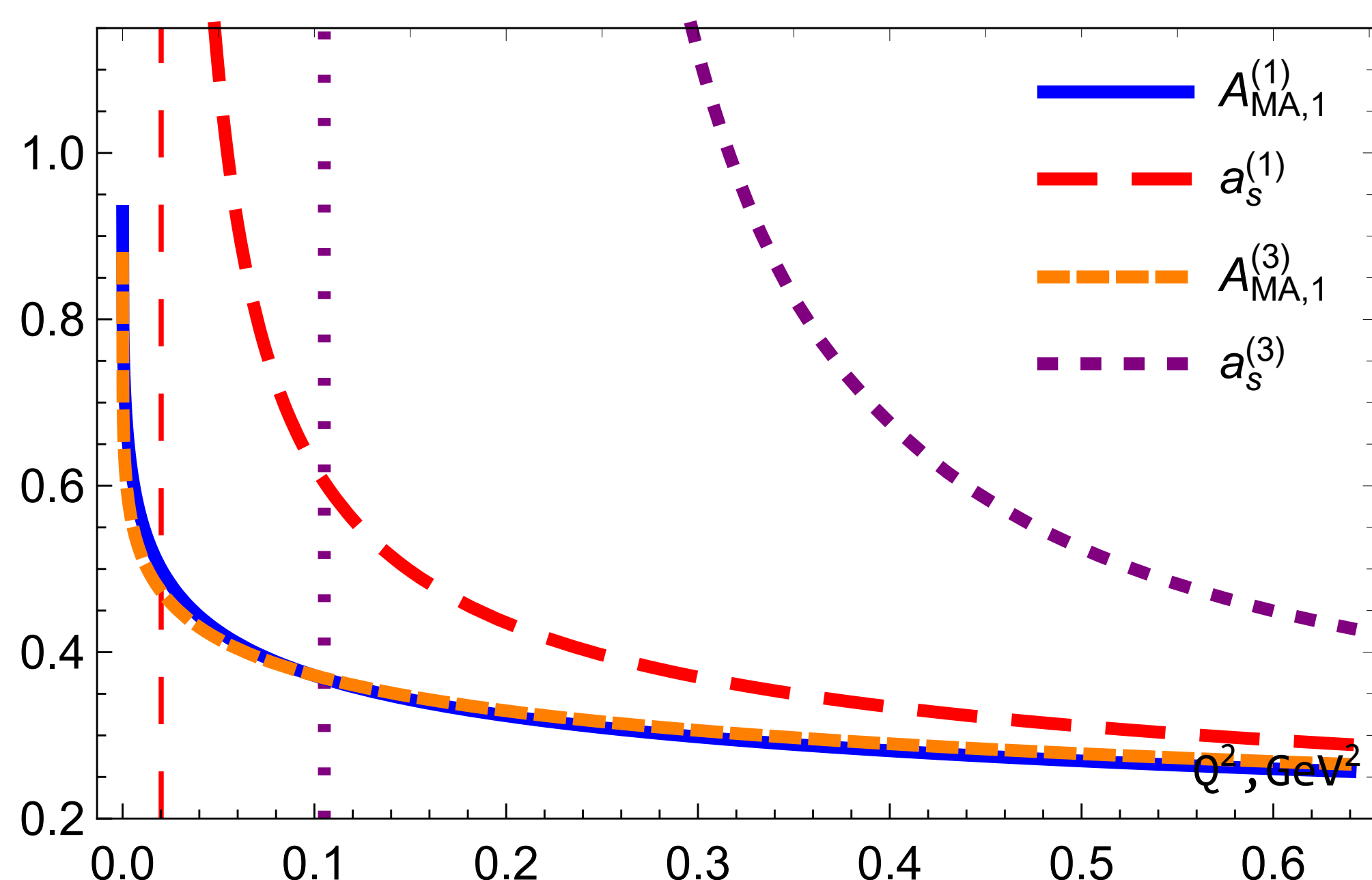


Fig.1 Comparison $a_s^{(i)}(Q^2)$ and $A_{MA,i}^{(i)}(Q^2)$. Vertical lines indicate the appropriate values of Λ_i .

In the frame of analytical perturbation theory (APT) [2] one can construct new holomorphic couplant $A_{MA}^{(i)}(Q^2)$

$$A_{MA}^{(i)}(Q^2) = \frac{i}{\pi} \int_0^{+\infty} \frac{d\sigma}{(\sigma + Q^2)} r_{pt}^{(i)}(\sigma), \quad r_{pt}^{(i)}(\sigma) = \text{Im } a_s^{(i)}(-\sigma - i\epsilon).$$

The final expressions for $\tilde{a}_\nu^{(i+1)}(Q^2)$ are represented as the sum of LO expression and the high order corrections [3]

$$\tilde{a}_\nu^{(i+1)}(Q^2) = \tilde{a}_\nu^{(1)}(Q^2) + \sum_{m=1}^i \frac{\Gamma(\nu + m)}{m! \Gamma(\nu)} \left(\hat{R}_m \tilde{a}_{\nu+m}^{(i+1)}(Q^2) \right),$$

where \hat{R}_m – differential operators ($\sim d^m/d\nu^m$). The structure of spectral integral allows to perform the same operation for $\tilde{A}_{MA}^{(i)}(Q^2)$:

$$\tilde{A}_{MA,\nu}^{(i+1)}(Q^2) = \tilde{A}_{MA,\nu}^{(1)}(Q^2) + \sum_{m=1}^i \frac{\Gamma(\nu + m)}{m! \Gamma(\nu)} \left(\hat{R}_m \tilde{A}_{MA,\nu+m}^{(1)}(Q^2) \right).$$

Bjorken Sum Rule

The definition of polarized Bjorken sum rule (BSR)

$$\Gamma_1^{p-n}(Q^2) = \int_0^1 dx [g_1^p(x, Q^2) - g_1^n(x, Q^2)].$$

BSR in the OPE form (twist-2+massive twist-4) reads [4]

$$\Gamma_1^{p-n}(Q^2) = \frac{g_A}{6} (1 - D_{BS}(Q^2)) + \frac{\hat{\mu}_4 M^2}{Q^2 + M^2}$$

and another form of twist-4 term for small Q^2 values [5]

$$\Gamma_1^{p-n}(Q^2) = \frac{g_A}{6} (1 - D_{BS}(Q^2)) + \frac{\hat{\mu}_4 M^2 (Q^2 + M^2)}{(Q^2 + M^2)^2 + M^2 \sigma^2}.$$

The twist-2 term $D_{BS}(Q^2)$ in PT and APT takes the form

$$D_{BS}^{(k)}(Q^2) = \frac{4}{\beta_0} \left(\tilde{a}_1^{(k)} + \sum_{m=2}^k \tilde{d}_{m-1} \tilde{a}_m^{(k)} \right),$$

$$D_{MA,BS}^{(k)}(Q^2) = \frac{4}{\beta_0} \left(A_{MA}^{(k)} + \sum_{m=2}^k \tilde{d}_{m-1} \tilde{A}_{MA,\nu=m}^{(k)} \right).$$

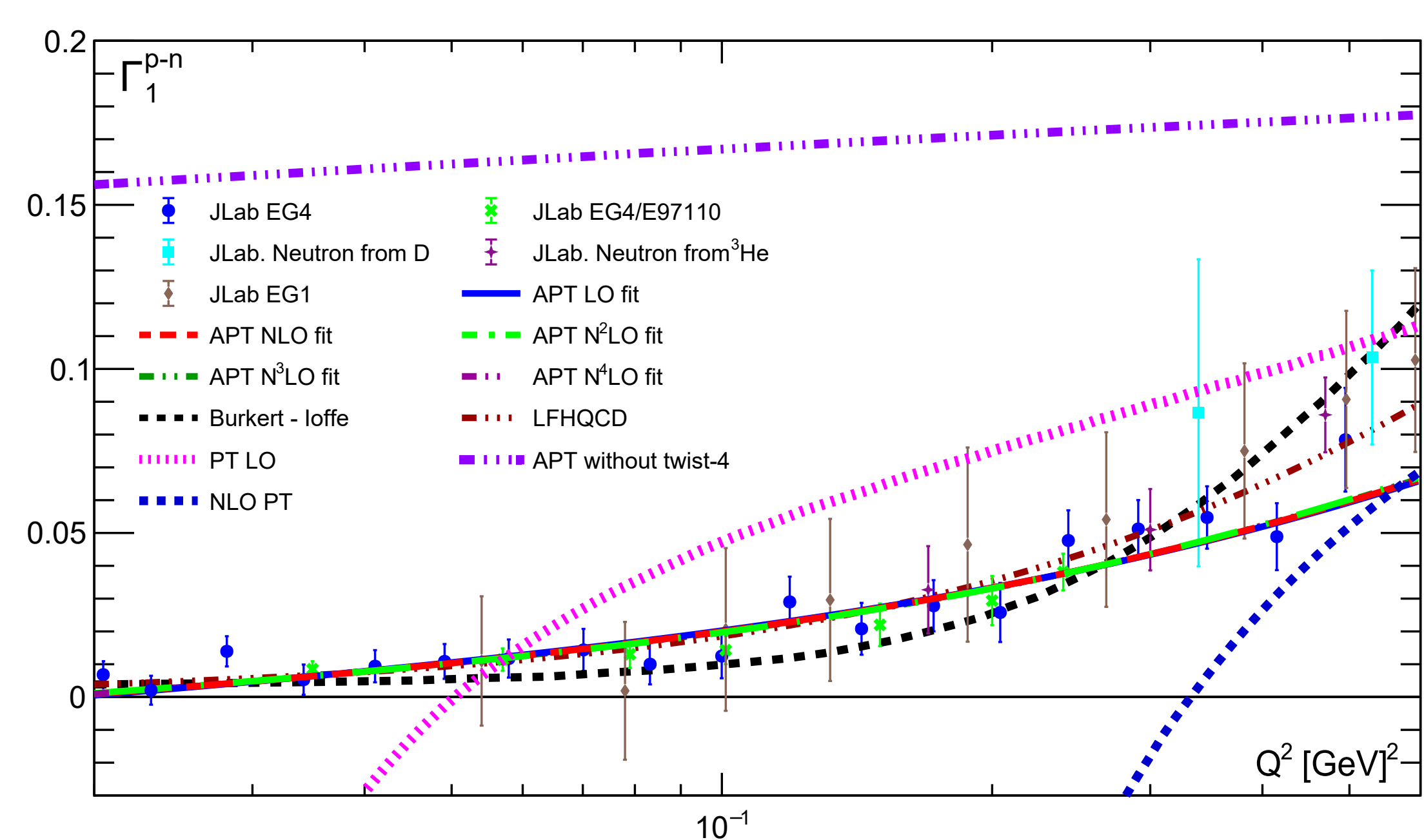


Fig.2 The results for $\Gamma_1^{p-n}(Q^2)$ in the first four orders of APT with $\sigma = \sigma_\rho$.

The values of the fit parameters with $\sigma = \sigma_\rho = 145$ MeV (the ρ -meson decay width) and $\sigma = 0$:

	M^2 for $\sigma = \sigma_\rho$ (for $\sigma = 0$)	$\hat{\mu}_{MA,4}$ for $\sigma = \sigma_\rho$ (for $\sigma = 0$)	$\chi^2/(\text{d.o.f.})$ for $\sigma = \sigma_\rho$ (for $\sigma = 0$)
LO	1.592 ± 0.300 (1.631 ± 0.301)	-0.168 ± 0.002 (-0.166 ± 0.001)	0.788 (0.789)
NLO	1.505 ± 0.286 (1.545 ± 0.287)	-0.157 ± 0.002 (-0.155 ± 0.001)	0.755 (0.757)
N ² LO	1.378 ± 0.242 (1.417 ± 0.241)	-0.159 ± 0.002 (-0.156 ± 0.002)	0.728 (0.728)
N ³ LO	1.389 ± 0.247 (1.429 ± 0.248)	-0.159 ± 0.002 (-0.157 ± 0.002)	0.747 (0.747)
N ⁴ LO	1.422 ± 0.259 (1.462 ± 0.259)	-0.159 ± 0.002 (-0.157 ± 0.001)	0.754 (0.754)

Summary

- The calculation results taking into account only statistical uncertainties.
- The cases $\sigma = 0$ and $\sigma = \sigma_\rho$ lead to very similar values for the fitting parameters and χ^2 -factor.
- The quality of the APT fits is very good (as evidenced quantitatively by the values of $\chi^2/(\text{d.o.f.})$) and much better than PT fits.

Acknowledgments

This work was supported in part by the Foundation for the Advancement of Theoretical Physics and Mathematics ‘‘BASIS’’.

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