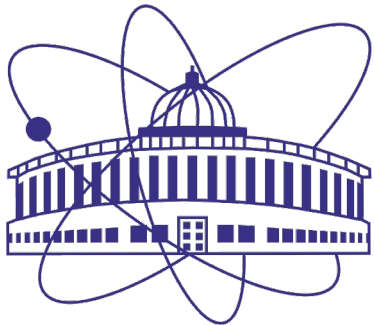


# Performance study of the anisotropic flow measurements with fixed-target mode of the MPD experiment at NICA

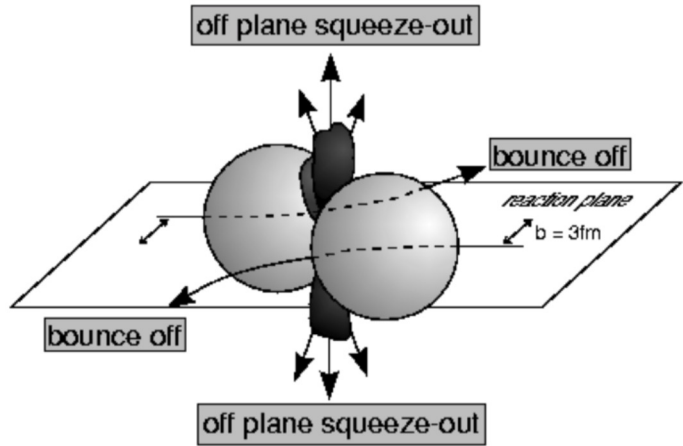
P. Parfenov, M. Mamaev and A. Taranenko  
(JINR, NRNU MEPhI)

LXXIV International Conference Nucleus-2024:  
Fundamental Problems and Applications  
1-5 July 2024

The work has been supported by the Ministry of Science and Higher Education of the Russian Federation, Project "Fundamental and applied research at the NICA megascience experimental complex" № FSWU-2024-0024



# Anisotropic flow & spectators



The azimuthal angle distribution is decomposed in a Fourier series relative to reaction plane angle:

$$\rho(\varphi - \Psi_{RP}) = \frac{1}{2\pi} \left( 1 + 2 \sum_{n=1}^{\infty} v_n \cos n(\varphi - \Psi_{RP}) \right)$$

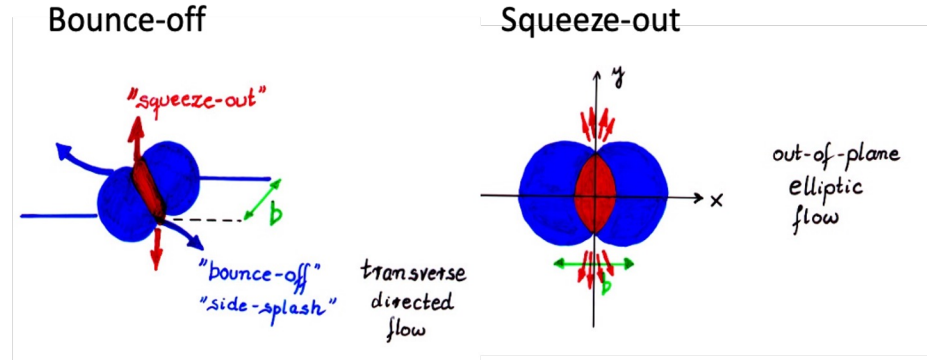
Anisotropic flow:

$$v_n = \langle \cos [n(\varphi - \Psi_{RP})] \rangle$$

$v_1$  - directed flow,  $v_2$  - elliptic flow

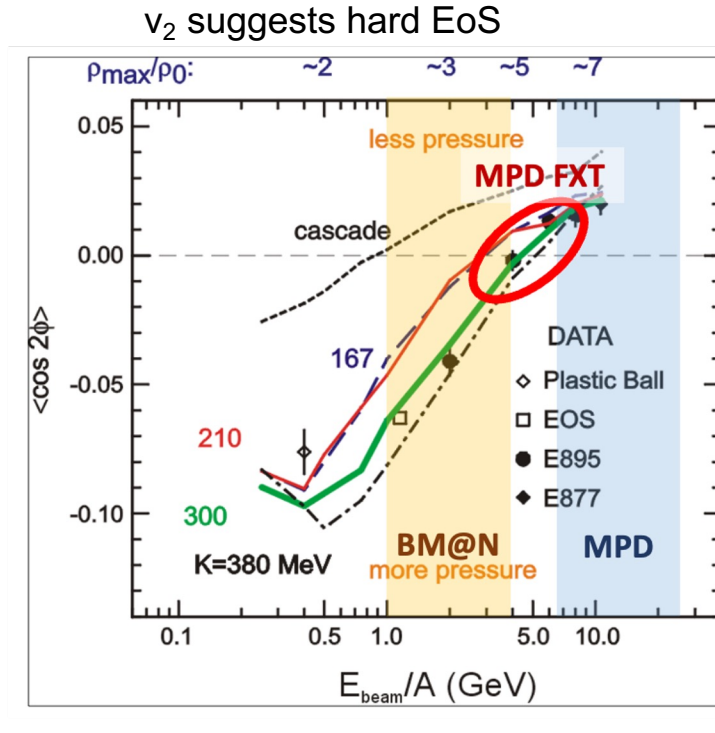
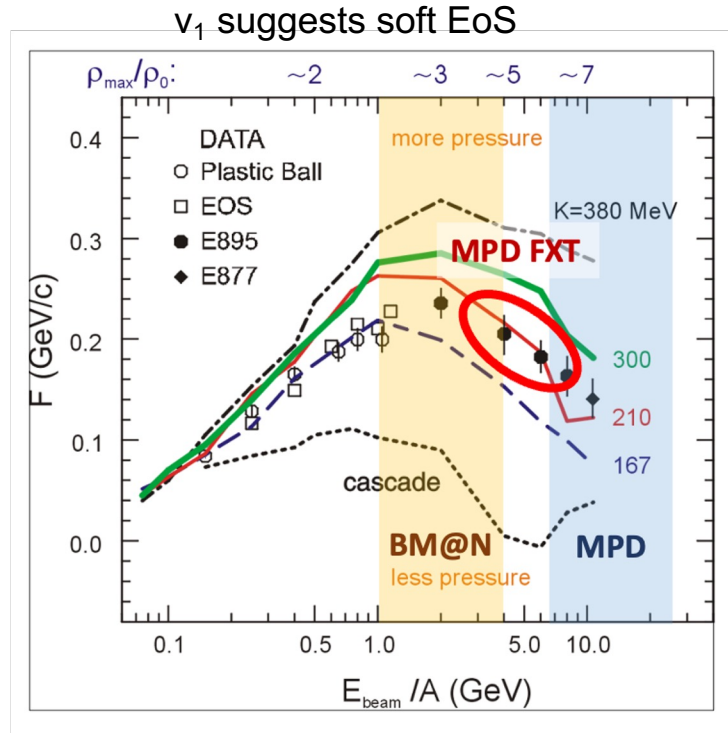
## Anisotropic flow is sensitive to:

- Compressibility of the created matter  
 $(t_{exp} = R/c_s, c_s = c\sqrt{dp/d\varepsilon})$
- Time of the interaction between overlap region and spectators  
 $(t_{pass} = 2R/\gamma_{CM}\beta_{CM})$



# $v_n$ at Nuclotron-NICA energies

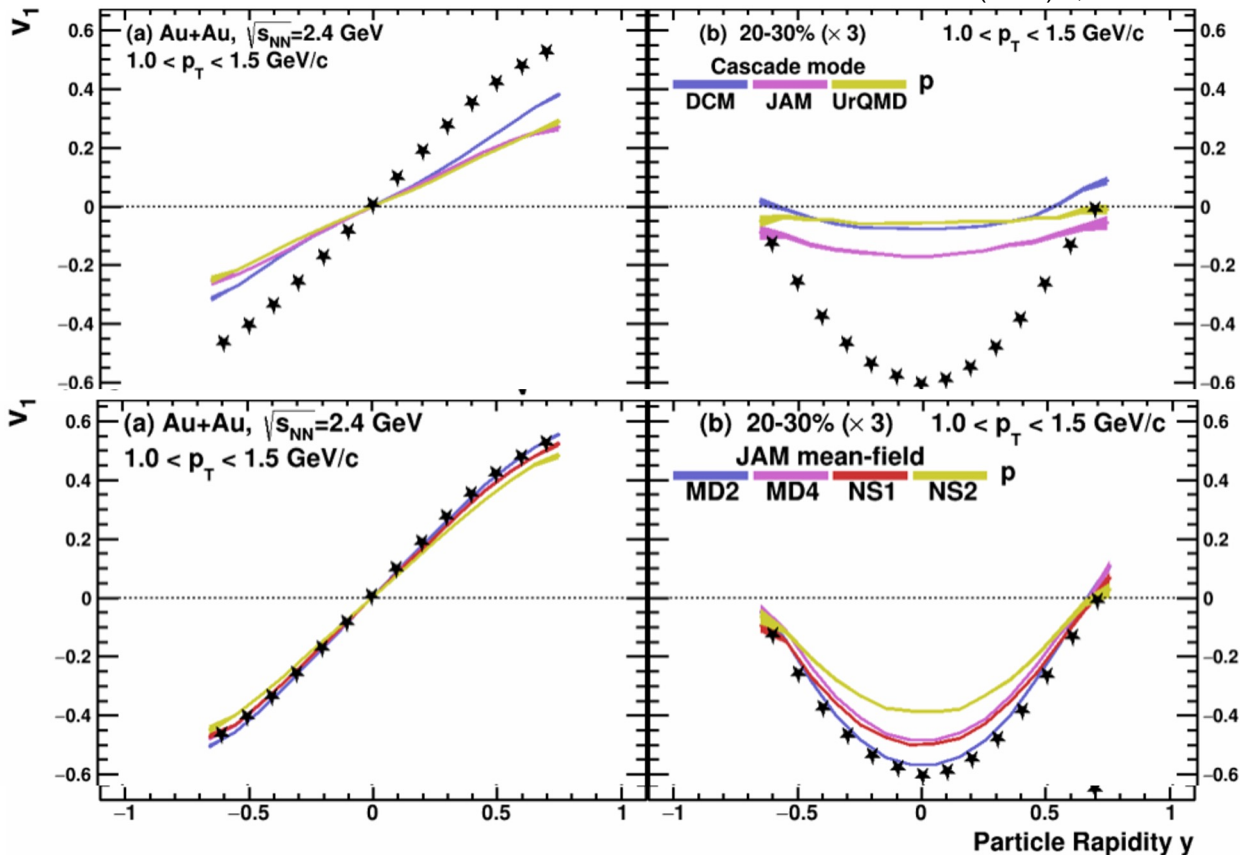
P. DANIELEWICZ, R. LACEY, W. LYNCH  
[10.1126/science.1078070](https://doi.org/10.1126/science.1078070)



- $v_n$  results from the E895 experiment are ambiguous:
  - $v_1$  suggests soft EoS and  $v_2$  suggests hard EoS
- Additional experimental data are required to address this discrepancy

# Selecting the model

P.Parfenov Particles 5 (2022) 4, 561-579

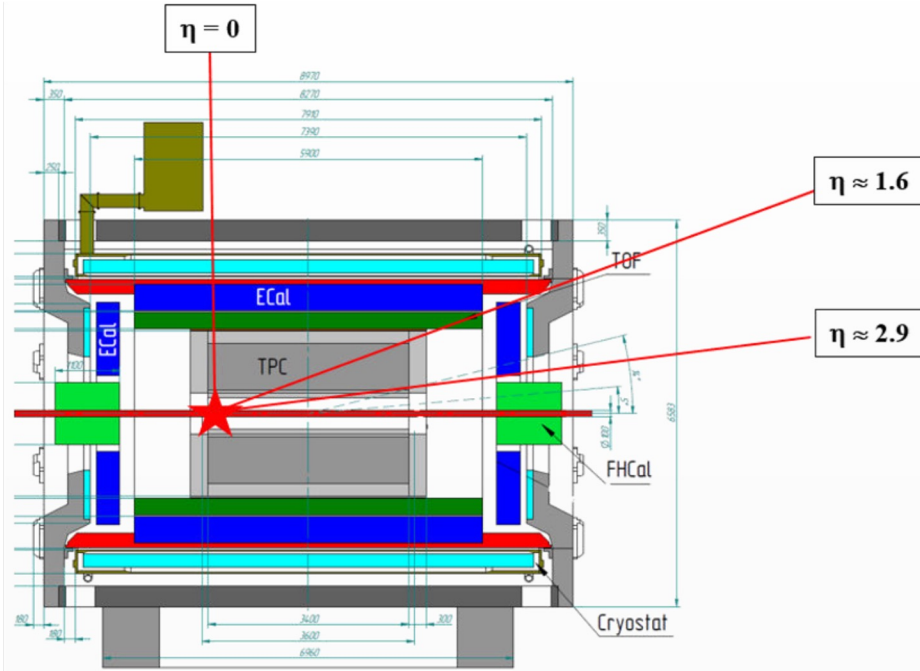


Cascade models fail to reproduce  $v_n$  at low-energy heavy-ion collision

Mean field models reproduce the  $v_n$  rather well

# MPD in Fixed-Target Mode (FXT)

## MPD-FXT



- Model used: UrQMD mean-field
  - Bi+Bi,  $E_{\text{kin}}=1.45$  AGeV ( $\sqrt{s_{\text{NN}}}=2.5$  GeV)
  - Bi+Bi,  $E_{\text{kin}}=2.92$  AGeV ( $\sqrt{s_{\text{NN}}}=3.0$  GeV)
  - Bi+Bi,  $E_{\text{kin}}=4.65$  AGeV ( $\sqrt{s_{\text{NN}}}=3.5$  GeV)
- Point-like target at  $z = -115$  cm
- GEANT4 transport
- Multiplicity-based centrality determination
- PID using information from TPC and TOF
- Primary track selection:  $\text{DCA} < 1$  cm
- Track selection:
  - $N_{\text{hits}} > 27$  (protons),  $N_{\text{hits}} > 22$  (pions)

# The Bayesian inversion method ( $\Gamma$ -fit)

Relation between multiplicity  $N_{ch}$  and impact parameter  $b$  is defined by the fluctuation kernel:

$$P(N_{ch}|c_b) = \frac{1}{\Gamma(k(c_b))\theta^k} N_{ch}^{k(c_b)-1} e^{-N_{ch}/\theta} \quad \frac{\sigma^2}{\langle N_{ch} \rangle} = \theta \approx const, k = \frac{\langle N_{ch} \rangle}{\theta}$$

$$c_b = \int_0^b P(b') db' - \text{centrality based on impact parameter}$$

Mean multiplicity as a function of  $c_b$  can be defined as follows:

$$\langle N_{ch} \rangle = N_{knee} \exp\left(\sum_{j=1}^3 a_j c_b^j\right) \quad N_{knee}, \theta, a_j - 5 \text{ parameters}$$

Fit function for  $N_{ch}$  distribution:

$b$ -distribution for a given  $N_{ch}$  range:

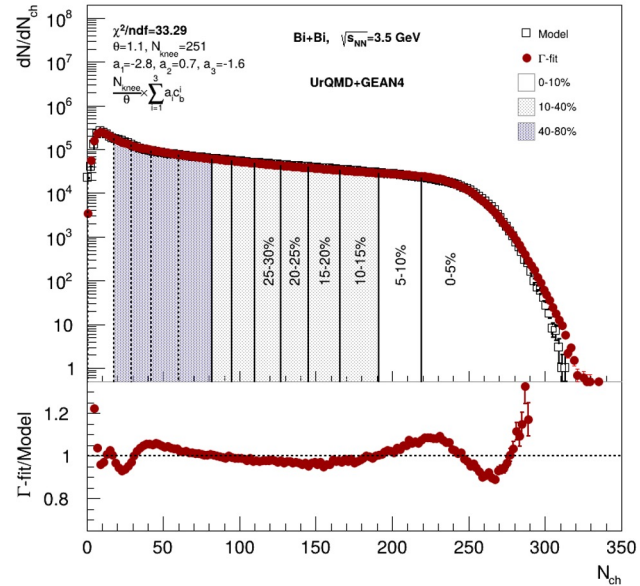
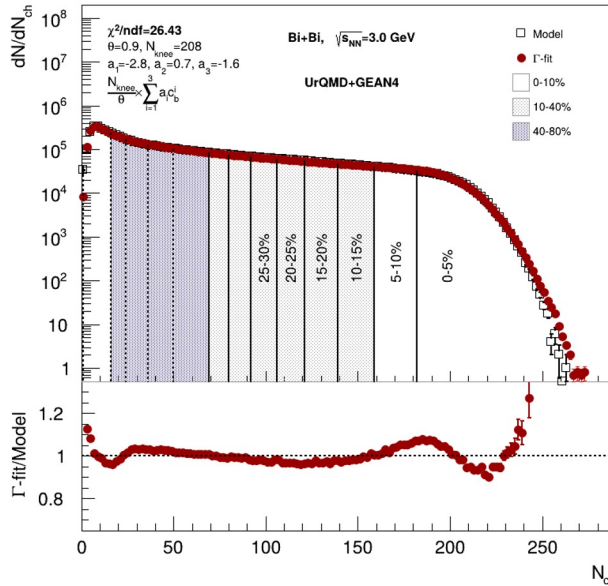
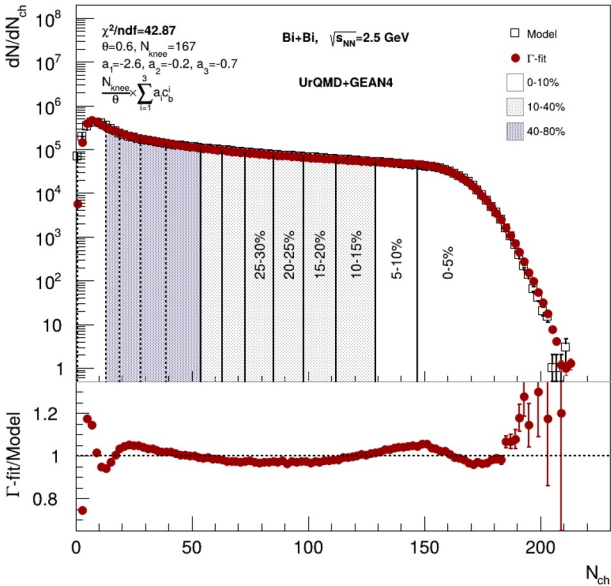
$$P(N_{ch}) = \int_0^1 P(N_{ch}|c_b) dc_b \quad P(b|n_1 < N_{ch} < n_2) = P(b) \frac{\int_{n_1}^{n_2} P(N_{ch}|b) dN_{ch}}{\int_{n_1}^{n_2} P(N_{ch}) dN_{ch}}$$

**2 main steps of the method:**

Fit experimental (model) distribution with  $P(N)$

Construct  $P(b|E)$  using Bayes' theorem:  
 $P(b|N) = P(b)P(N|b)/P(N)$

# Centrality determination: multiplicity fit



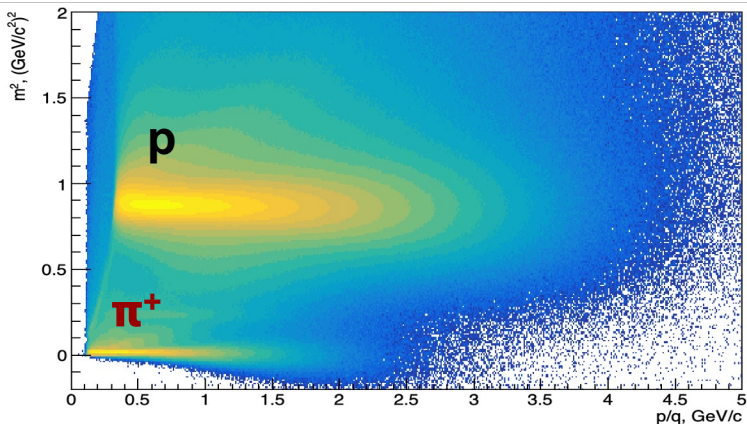
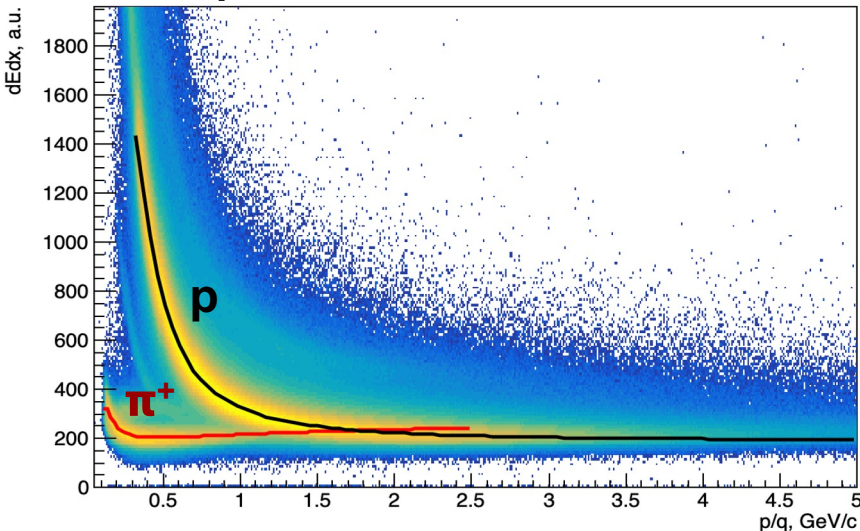
Cuts on tracks:

- $N_{hits} > 16$
- $0 < \eta < 2$

Good agreement between fit and data

Multiplicity-based centrality determination ( $\Gamma$ -fit) was used

# PID procedure



Fit  $dE/dx$  distributions with Bethe-Bloch parametrization:

$$f(\beta\gamma) = \frac{p_1}{\beta^{p_4}} \left( p_2 - \beta^{p_4} - \ln \left( p_3 + \frac{1}{(\beta\gamma)^{p_5}} \right) \right)$$

$$\beta^2 = \frac{p^2}{m^2 + p^2}, \beta\gamma = \frac{p}{m}$$

$p_i$  - fit parameters

Fit  $(dE/dx - f(\beta\gamma))/f(\beta\gamma)$  with gaus in the slices of  $p/q$  and get  $\sigma_p(dE/dx)$

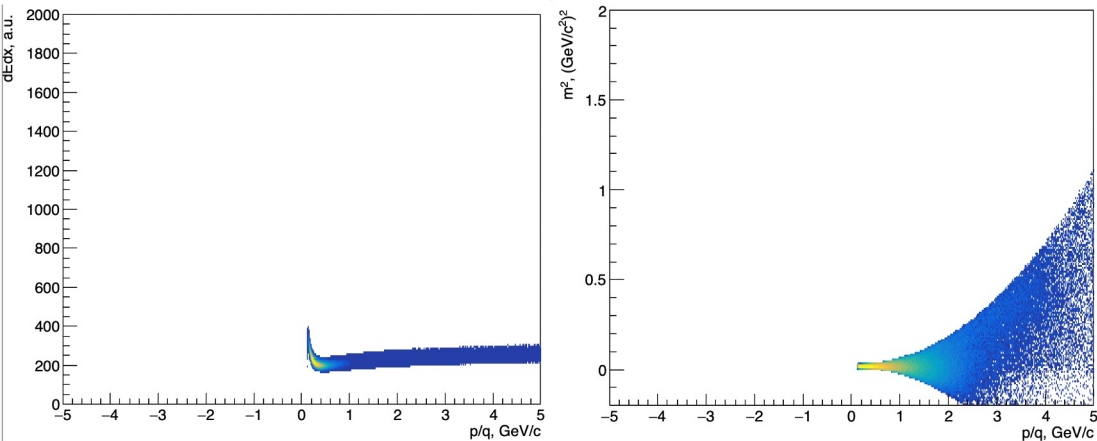
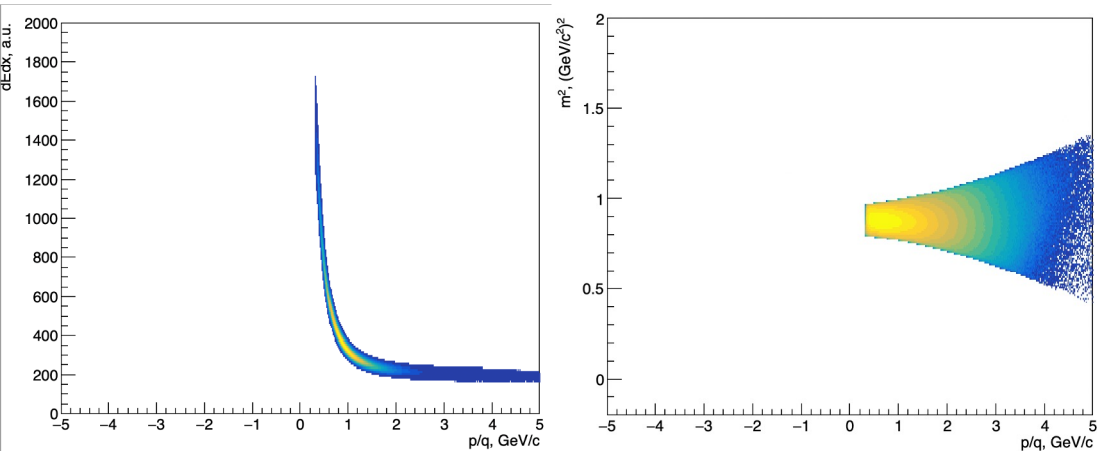
Fit  $m^2$  with gaus in the slices of  $p/q$  and get  $\sigma_p(m^2)$

**$(dE/dx, m) \rightarrow (x, y)$  coordinates for PID:**

$$x_p = \frac{(dE/dx)^{meas} - (dE/dx)_p^{fit}}{(dE/dx)_p^{fit} \sigma_p^{dE/dx}}, \quad y_p = \frac{m^2 - m_p^2}{\sigma_p^{m^2}}$$



# PID procedure: Results



$$x_p = \frac{(dE/dx)^{meas} - (dE/dx)_p^{fit}}{(dE/dx)_p^{fit} \sigma_p^{dE/dx}}$$

$$y_p = \frac{m^2 - m_p^2}{\sigma_p^{m^2}}$$

Protons:

$$\sqrt{x_p^2 + y_p^2} < 2, \sqrt{x_\pi^2 + y_\pi^2} > 3$$

Pions ( $\pi^+$ ):

$$\sqrt{x_\pi^2 + y_\pi^2} < 2, \sqrt{x_p^2 + y_p^2} > 3$$

Pions ( $\pi^-$ ):

charge < 0

# (y-pt) distribution, efficiency and $\delta p_T$ (protons)

$$\text{eff} = \frac{\frac{dN}{dydp_T}(\text{reco})}{\frac{dN}{dydp_T}(\text{sim})}$$

$$\Delta p_T = \frac{|p_T^{\text{reco}} - p_T^{\text{mc}}|}{p_T^{\text{mc}}}$$

Bi+Bi  $\sqrt{s_{\text{NN}}=2.5 \text{ GeV}}$

Cuts for reco tracks:

- Nhits>27
- DCA< 1 cm
- PID (TPC+TOF)
- Primary (DCA<1 cm)

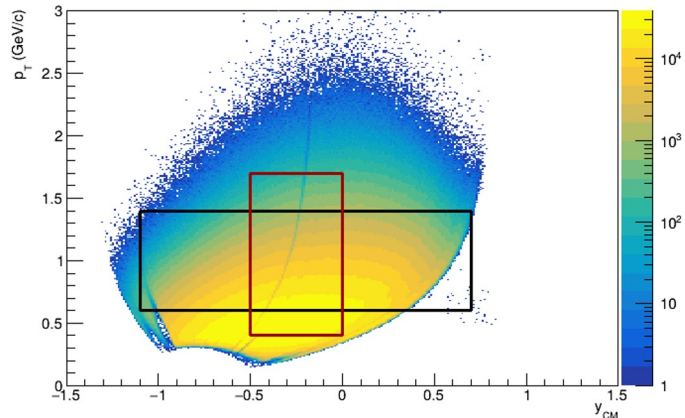
Cuts for sim particles:

- PID (pdg code)
- Primary (motherId)

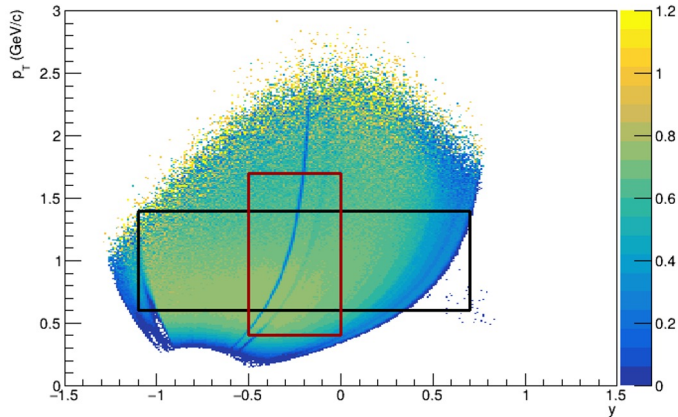
**Black box:** acceptance window for  $v_n(y)$

**Red box:** acceptance window for  $v_n(p_T)$

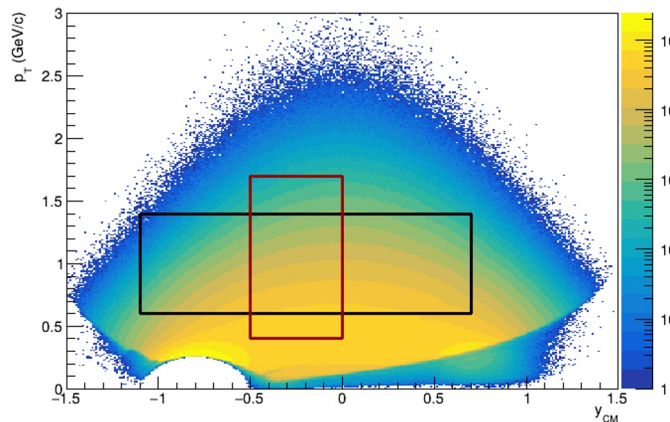
Reconstructed protons Ycm-pT



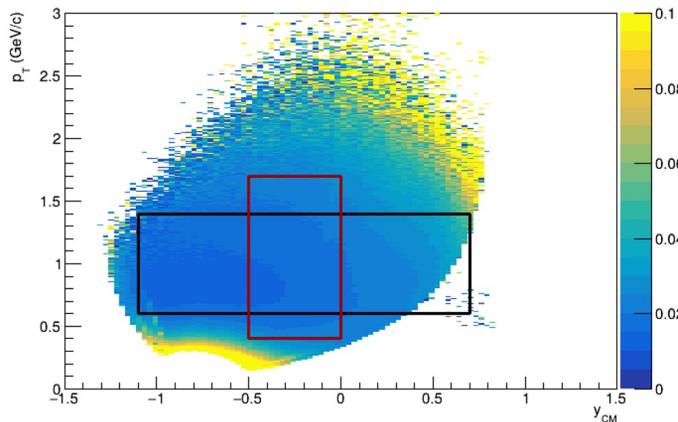
Efficiency (Y-pT) of primary protons



Simulated protons Ycm-pT



Pt-resolution for reconstructed protons in Ycm-pT plane



# Flow vectors

From momentum of each measured particle define a  $u_n$ -vector in transverse plane:

$$u_n = e^{in\phi}$$

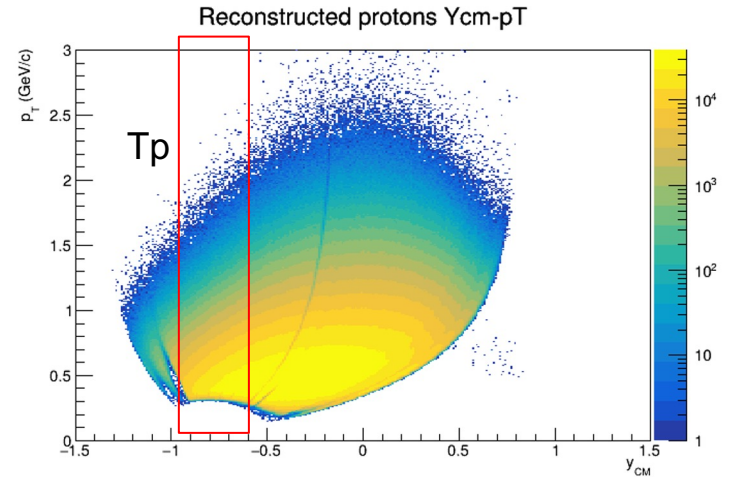
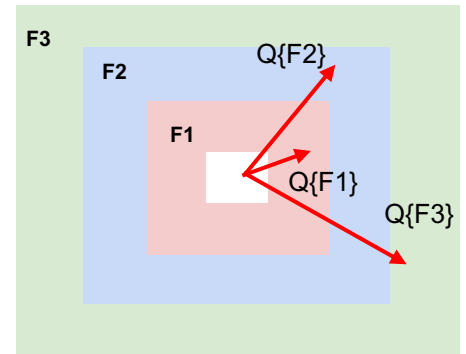
where  $\phi$  is the azimuthal angle

Sum over a group of  $u_n$ -vectors in one event forms  $Q_n$ -vector:

$$Q_n = \frac{\sum_{k=1}^N w_n^k u_n^k}{\sum_{k=1}^N w_n^k} = |Q_n| e^{in\Psi_n^{EP}}$$

$\Psi_n^{EP}$  is the event plane angle

Modules of FHCAL divided into 3 groups



**Additional subevents from tracks not pointing at FHCAL:**  
**Tp:** p;  $-1.0 < y < -0.6$ ;

# Flow methods for $v_n$ calculation

Tested in HADES: M Mamaev et al 2020 PPNuclei 53, 277–281  
M Mamaev et al 2020 J. Phys.: Conf. Ser. 1690 012122

Scalar product (SP) method:

$$v_1 = \frac{\langle u_1 Q_1^{F1} \rangle}{R_1^{F1}} \quad v_2 = \frac{\langle u_2 Q_1^{F1} Q_1^{F3} \rangle}{R_1^{F1} R_1^{F3}}$$

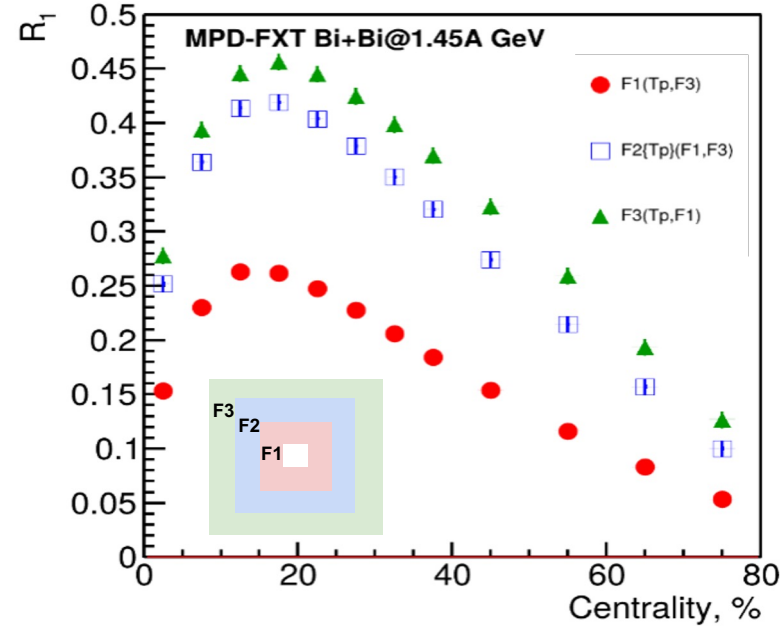
Where  $R_1$  is the resolution correction factor

$$R_1^{F1} = \langle \cos(\Psi_1^{F1} - \Psi_1^{RP}) \rangle$$

Symbol “F2(F1,F3)” means  $R_1$  calculated via  
(3S resolution):

$$R_1^{F2(F1,F3)} = \frac{\sqrt{\langle Q_1^{F2} Q_1^{F1} \rangle \langle Q_1^{F2} Q_1^{F3} \rangle}}{\sqrt{\langle Q_1^{F1} Q_1^{F3} \rangle}}$$

$$R_1^{F2\{Tp\}(F1,F3)} = \langle Q_1^{F2} Q_1^{Tp} \rangle \frac{\sqrt{\langle Q_1^{F1} Q_1^{F3} \rangle}}{\sqrt{\langle Q_1^{Tp} Q_1^{F1} \rangle \langle Q_1^{Tp} Q_1^{F3} \rangle}}$$

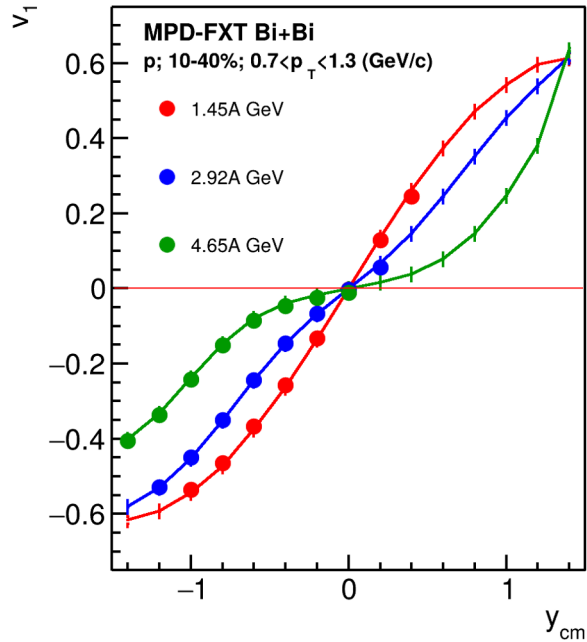


Symbol “F2{Tp}(F1,F3)” means  $R_1$   
calculated via (4S resolution):

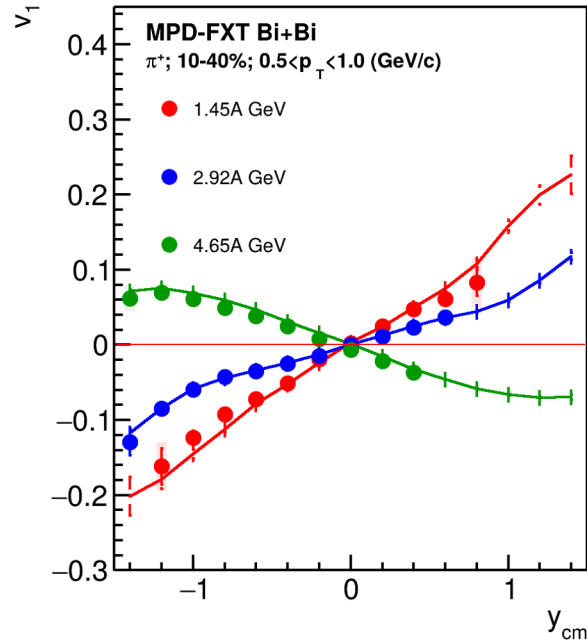
# Results: $v_1(y)$

Systematics: xx, yy, F1, F2, F3

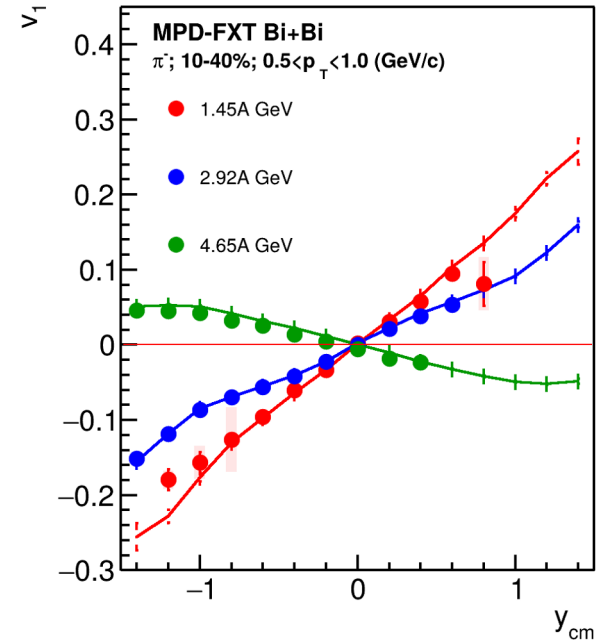
**p**



**$\pi^+$**



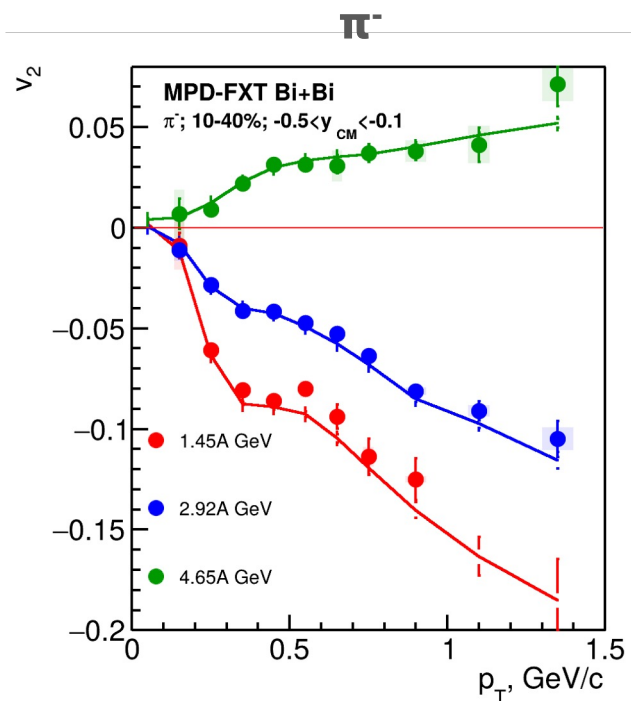
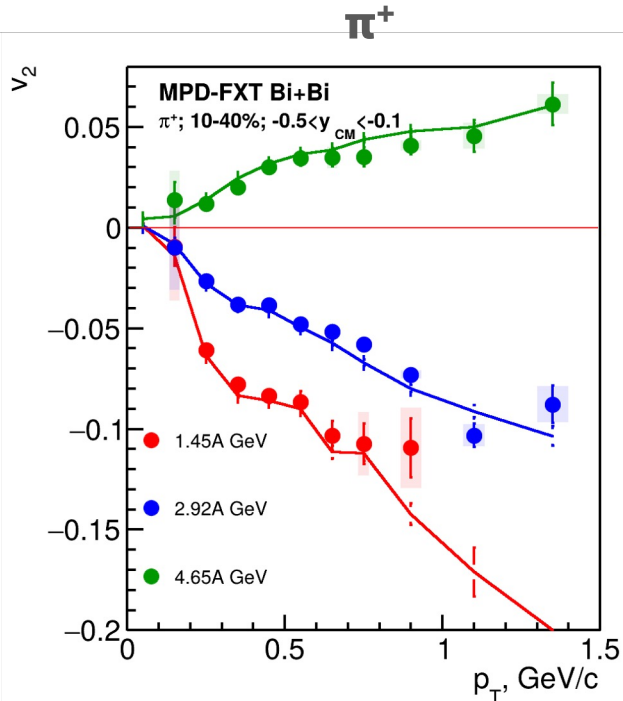
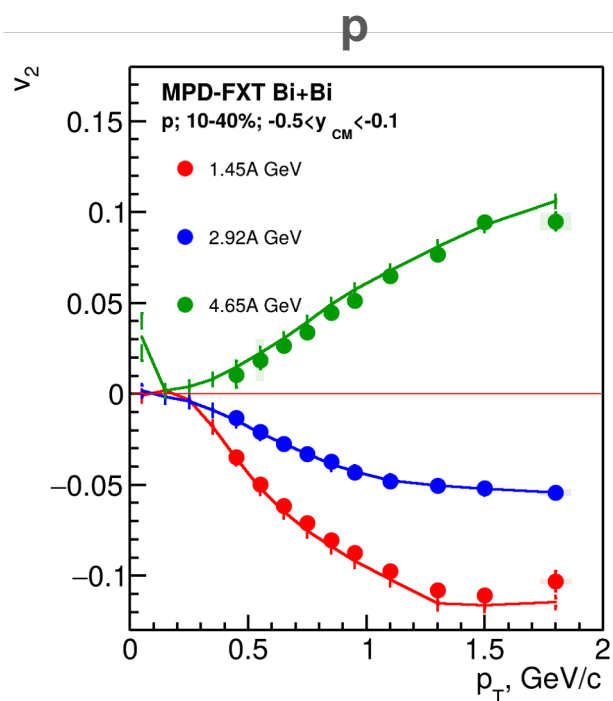
**$\pi^-$**



Good agreement with MC data

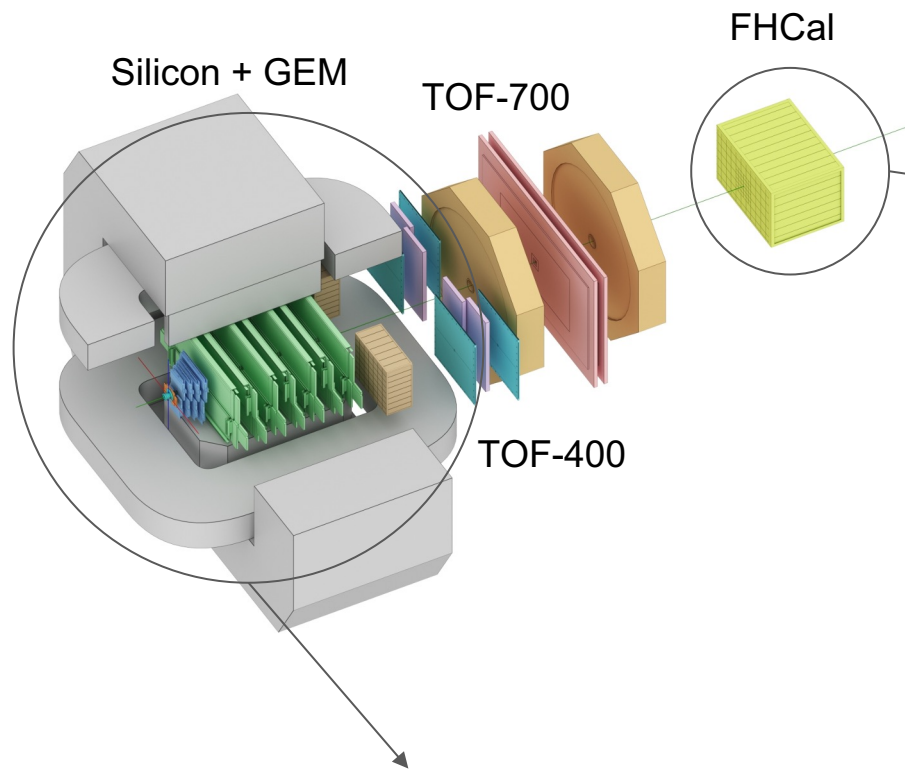
# Results: $v_2(p_T)$

Systematics: xxx, xyy

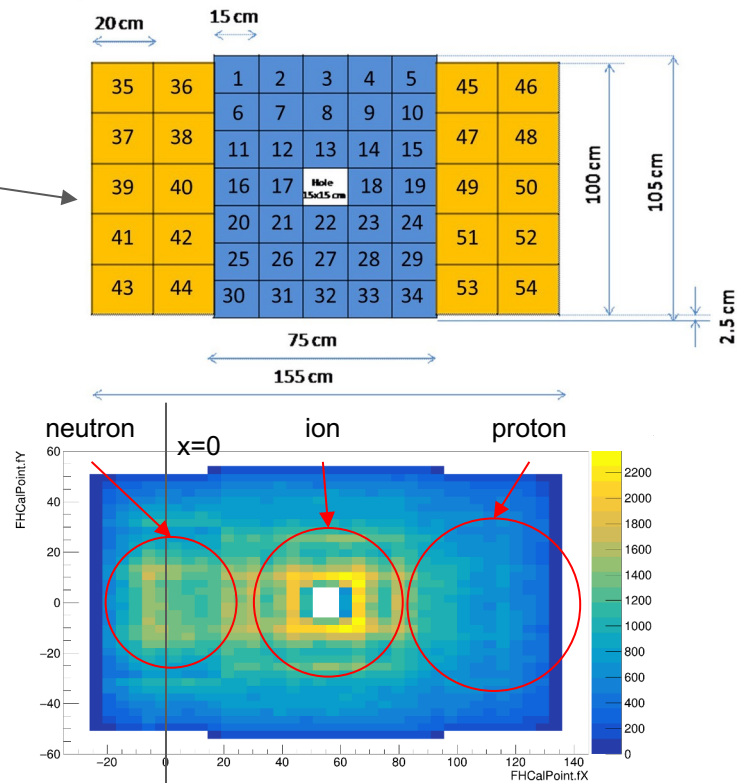


Good agreement with MC data

# The BM@N experiment (GEANT4 simulation for RUN8)

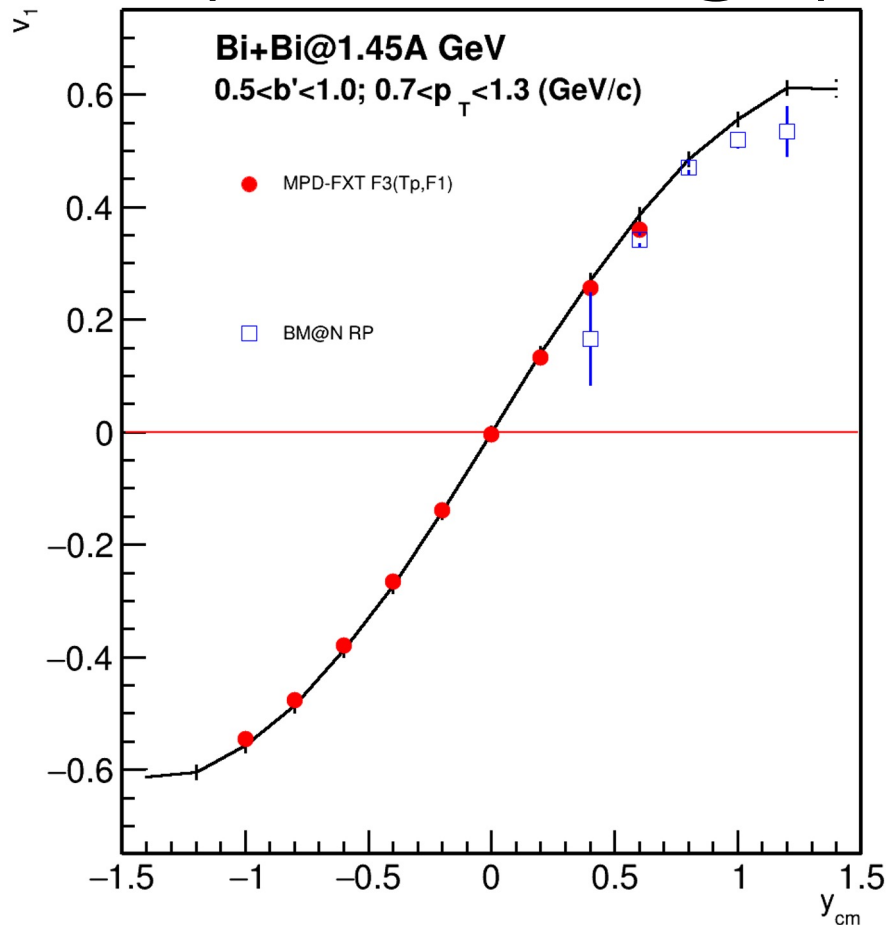


Square-like tracking system within the magnetic field deflecting particles along X-axis



Charge splitting on the surface of the FHCAL is observed due to magnetic field

# Comparison with BM@N performance



BM@N TOF system (TOF-400 and TOF-700) has poor midrapidity coverage at  $\sqrt{s_{NN}} = 2.5$  GeV

- One needs to check higher energies ( $\sqrt{s_{NN}} = 3, 3.5$  GeV)
- More statistics are required due to the effects of magnetic field in BM@N:
  - Only “yy” component of  $\langle uQ \rangle$  and  $\langle QQ \rangle$  correlation can be used

Despite the challenges, both MPD-FXT and BM@N can be used in  $v_n$  measurements:

- To widen rapidity coverage
- To perform a cross-check in the future



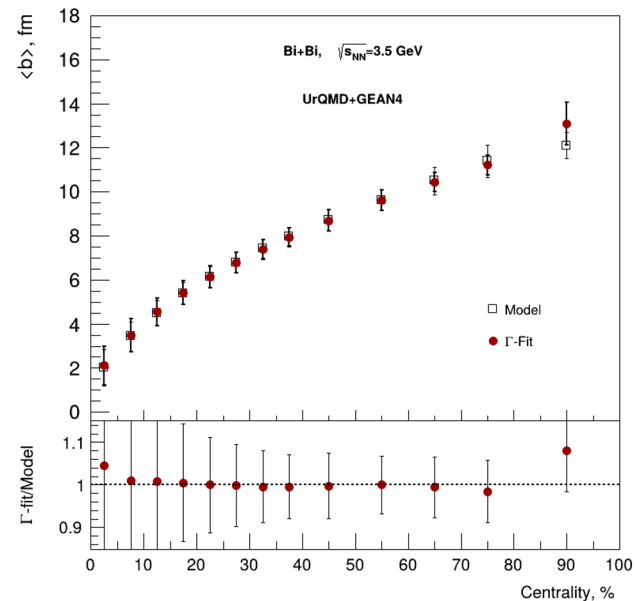
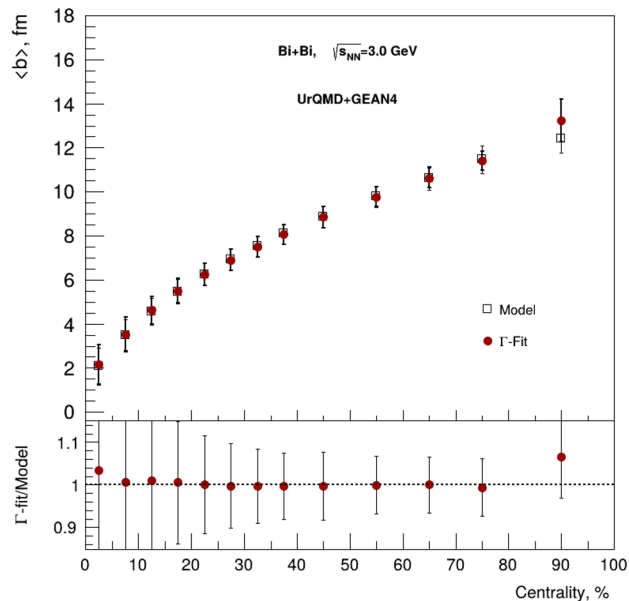
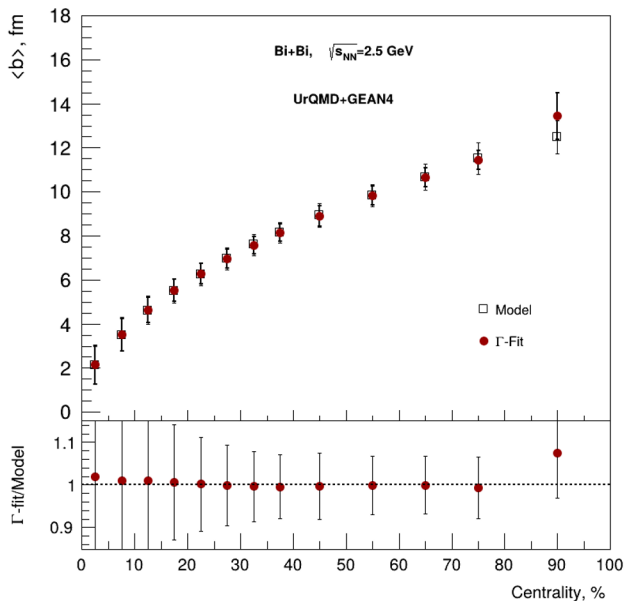
# Summary

- **Performance study for the anisotropic flow measurements was shown for the MPD-FXT using realistic procedures for centrality determination, primary track selection and PID:**
  - Multiplicity-based centrality determination using  $\Gamma$ -fit shows good agreement between fit and data
  - Overall good agreement between the estimated fit and impact parameter with the corresponding values taken directly from the model
- **Basic PID was performed using  $dE/dx$  from TPC and  $m^2$  from TOF**
- **Directed and elliptic flow of protons and pions were measured for  $\sqrt{s_{NN}} = 2.5, 3, 3.5$  GeV:**
  - Good agreement between reconstructed and model data within corresponding acceptance windows for all particle species
- **Both MPD-FXT and BM@N can complement each other in terms of  $v_n$ :**
  - Cross-checks can be performed to test the implemented flow measurement techniques
  - Using results from both experiments can widen the rapidity coverage - **no single fixed target experiment can achieve that!**

**New data from the BM@N and MPD (MPD-FXT) is required to address existing discrepancies in the experimental data and provide further constraints for the EoS in the models**

# Backup

# Centrality determination: $\langle b \rangle$ vs Centrality



Cuts on tracks:

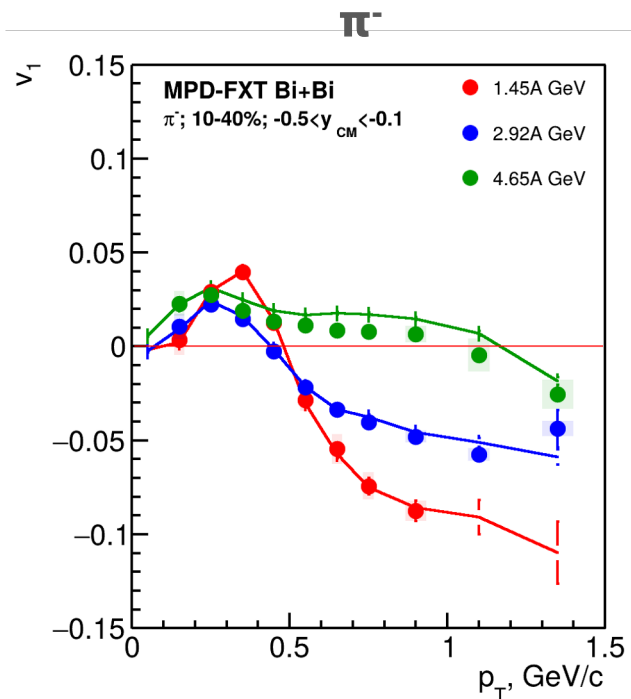
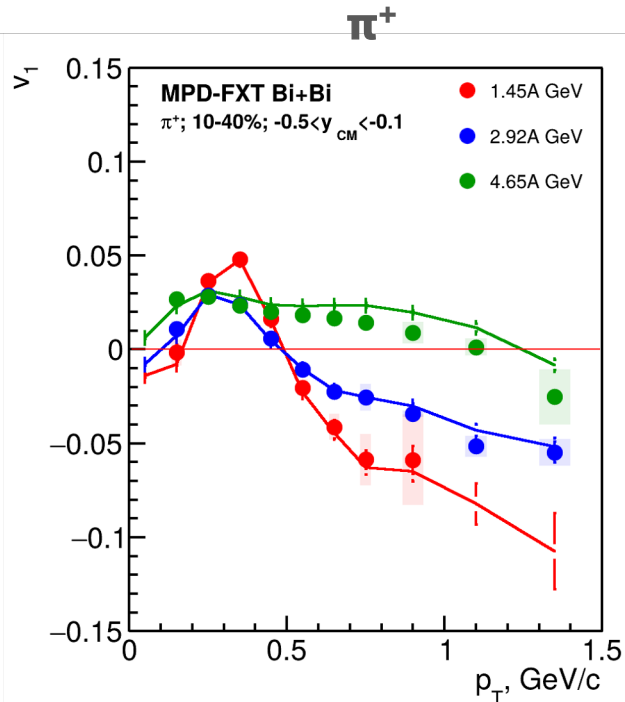
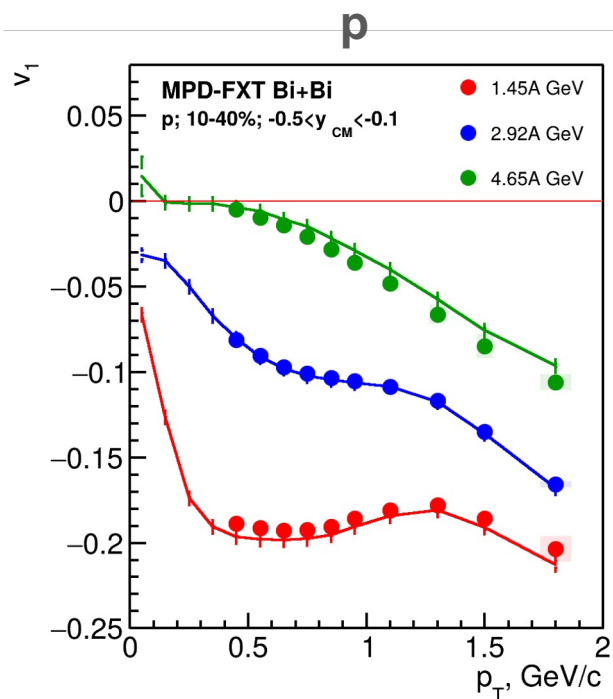
- $N_{\text{hits}} > 16$
- $0 < \eta < 2$

Good agreement between fit and data

Multiplicity-based centrality determination using inverse Bayes was used

# Results: $v_1(p_T)$

Systematics: xx, yy, F1, F2, F3

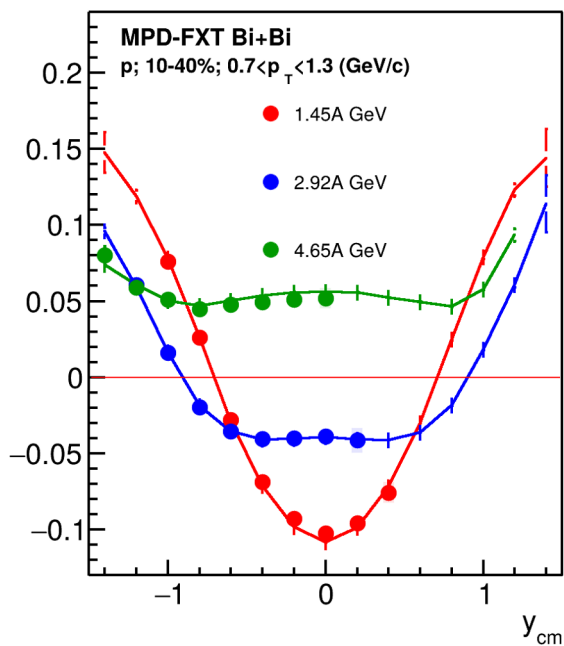


Good agreement with MC data

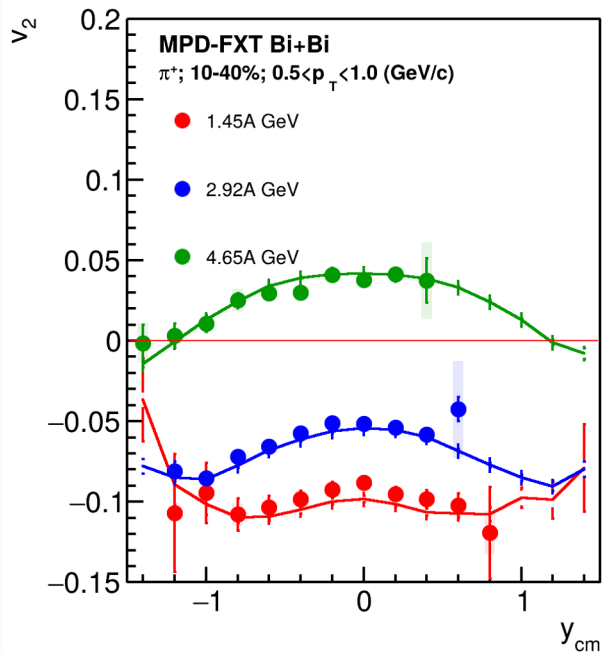
# Results: $v_2(y)$

Systematics: xxx, xyy

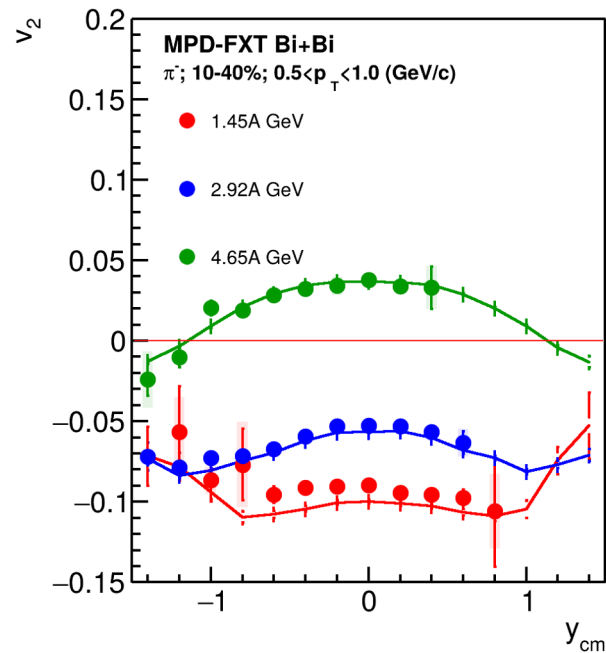
**p**



**$\pi^+$**



**$\pi^-$**



Good agreement with MC data