

Strongly intensive fluctuations and correlations in ultrarelativistic nuclear collisions in the model with string fusion

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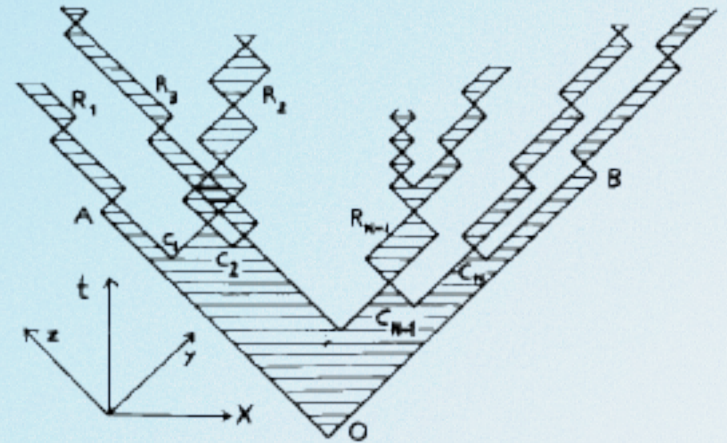
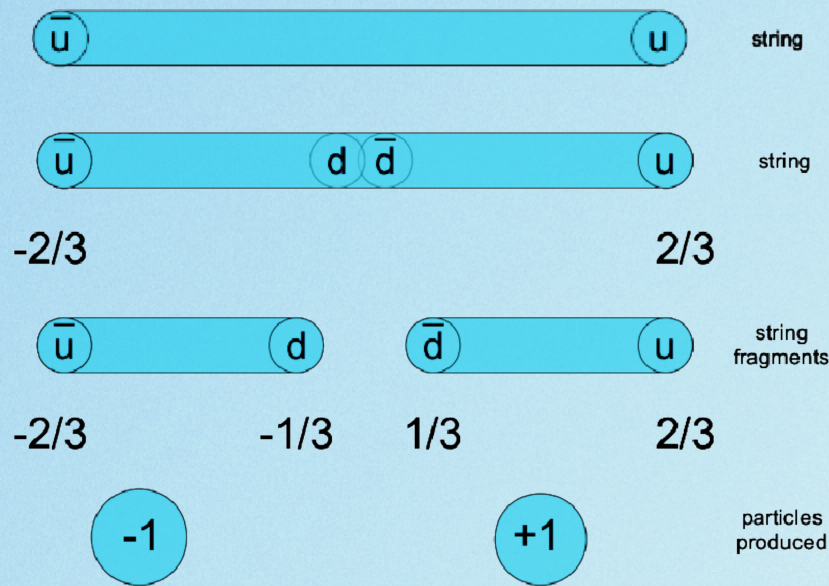
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Initial stages

- The soft QCD processes is not described by usual perturbation theory
- The model of quark-gluon strings, stretched between projectile and target partons – semiphenomenological approach to the multiparticle production



X. Artru and G. Mennessier, Nucl Phys B 70 (1974) 93
“String Model and Multiproduction”,

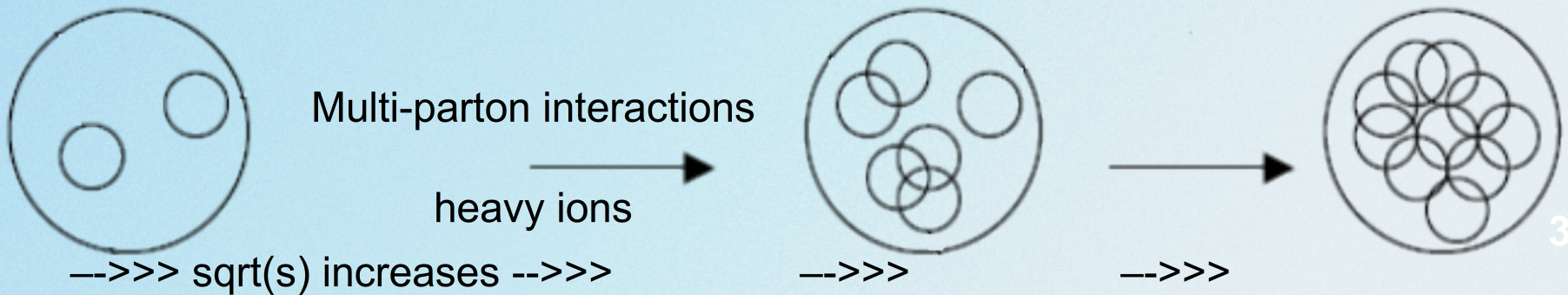
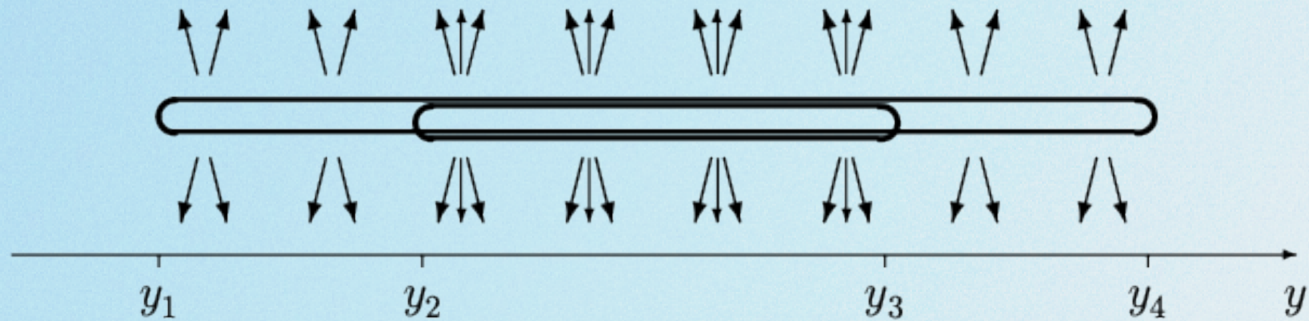
- Correlations play crucial role:
 - causality requires appearance of long-range correlations – if they exist – at the very early stages between particles detected in separated rapidity intervals

Strings in rapidity space

- Uniform and independent distribution of particles on rapidity from $y_{\min}$ to $y_{\max}$



- Can study string overlaps:



String fusion

$$Q^2(n) = \left(\sum_{i=1}^n \vec{Q}_i(1) \right)^2 = \sum_{i=1}^n Q_i^2(1) + \sum_{i \neq j} \vec{Q}_i(1) \cdot \vec{Q}_j(1)$$

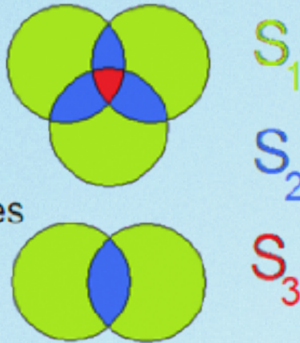
$$\langle Q^2(n) \rangle = n Q^2(1)$$

overlaps

SFM

$$C = \{S_1, S_2, \dots\}$$

S_k – area
covered k-times



String fusion mechanism predicts:

- decrease of multiplicity
- increase of p_T
- growth of p_T with multiplicity in pp, pA and AA collisions
- growth of strange particle yields

– results are in a good agreement with the experiment

$$\langle \mu \rangle_k = \mu_0 \sqrt{k} \frac{S_k}{\sigma_0}$$

$$\langle p_t^2 \rangle_k = p_0^2 \sqrt{k}$$

$$\langle p_t \rangle_k = p_0 \sqrt[4]{k}$$

S_k – area, where k strings are overlapping, σ_0 single string transverse area, μ_0 and p_0 – mean multiplicity and transverse momentum from one string

M. A. Braun, C. Pajares, Nucl. Phys. B 390 (1993) 542.

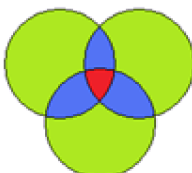
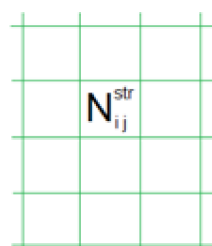
M. A. Braun, R. S. Kolevatov, C. Pajares, V. V. Vechernin, Eur. Phys. J. C 32 (2004) 535.

N.S. Amelin, N. Armesto, C. Pajares, D. Sousa, Eur.Phys.J.C22:149-163 (2001), arXiv:hep-ph/0103060

G. Ferreiro and C Pajares J. Phys. G: Nucl. Part. Phys. 23 1961 (1997)

String fusion

- Lattice realization of string fusion model

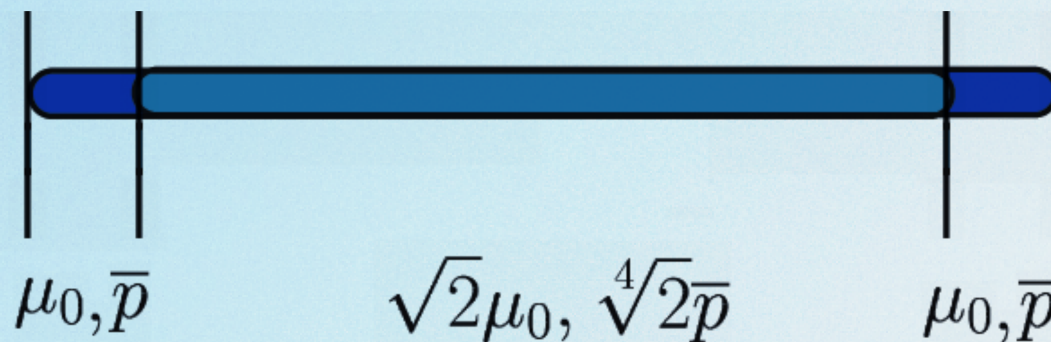
	"overlaps" (local fusion)		"overlaps" (local fusion)
SFM	○ $C = \{S_1, S_2, \dots\}$  S_1 S_2 S_3 S_k – area covered k-times	cellular analog of SFM	□ $C = \{N_{ij}^{str}\}$  N_{ij}^{str} $k_{ij} = N_{ij}^{str}$ – "occupation" numbers

V. V. Vechernin and R. S. Kolevaton, Vestn. SPb. Univ., Ser. Fiz. Khim., No. 2, 12 (2004); hep-ph/0304295.

V. V. Vechernin and R. S. Kolevaton, Vestn. SPb. Univ., Ser. Fiz. Khim., No. 4, 11 (2004); hep-ph/0305136.

I. A. Lakomov, V. V. Vechernin, PoS (Baldin ISHEPP XXI) 072 (2012).

- Fusion of finite rapidity strings



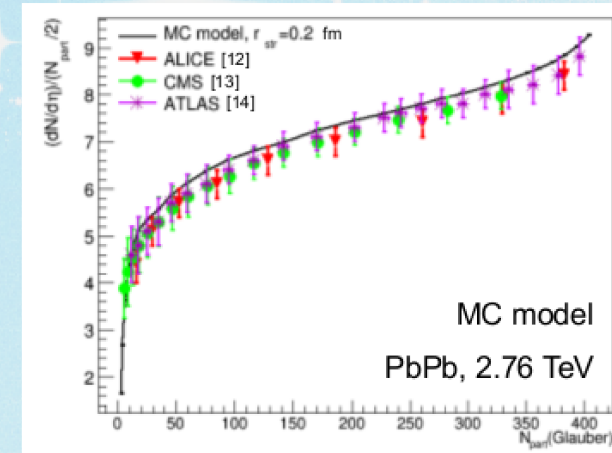
Monte Carlo model

- Partonic dipole-based picture of nucleons interaction.
- Energy and angular momentum conservation in the initial state of a nucleon.
- The probability of dipoles are defined by their transverse coordinates [7-8]:

$$f = \frac{\alpha_s^2}{2} \ln^2 \frac{|\vec{r}_1 - \vec{r}'_1| |\vec{r}_2 - \vec{r}'_2|}{|\vec{r}_1 - \vec{r}'_2| |\vec{r}_2 - \vec{r}'_1|}$$

Multiplicity and transverse momentum are obtained in the approach of colour strings, stretched between projectile and target partons

- The interaction of strings is realized in the accordance with the **string fusion** model
- Multiplicity from one string is distributed according to Poisson



V. N. Kovalenko.. Phys. Atom. Nucl. 76, 1189 (2013), arXiv:1211.6209 [hep-ph]

V. Kovalenko, V. Vechernin., PoS (Baldin ISHEPP XXI) 077, arXiv:1212.2590 [nucl-th], 2012

Forward-backward correlations

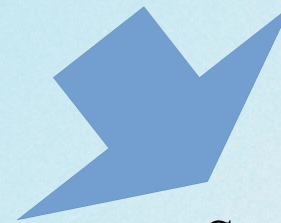


Select 2 variables in windows:

F, B:
$$\left\{ \begin{array}{ll} n & \text{-- number of charged particles in the window (Q=n}_+ \text{ - n}_-) \\ p_t = \frac{1}{n} \sum_{i=1}^n p_{ti} & \text{-- mean (in the event!) transverse momentum of charged particles in the given window} \end{array} \right.$$

correlation coefficient

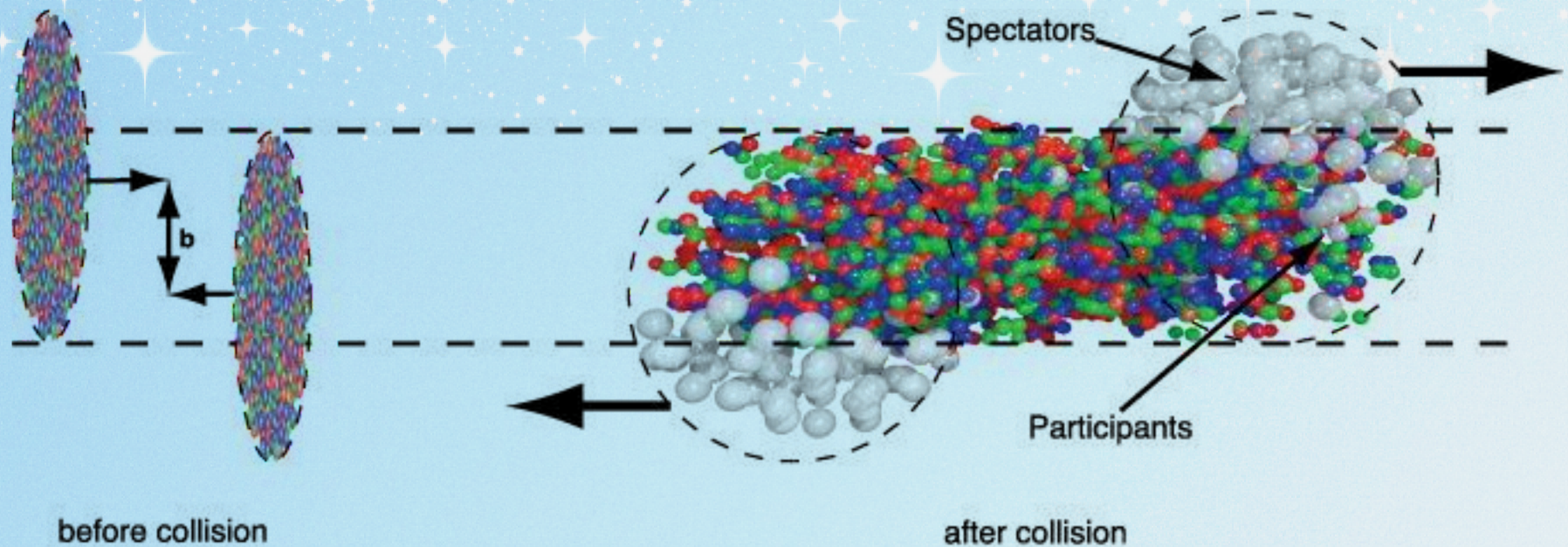
$$b = \frac{\langle FB \rangle - \langle F \rangle \langle B \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$



Several types of correlation coefficient:

n-n , pt-n, **pt-pt**, Q-Q

Centrality of AA collisions



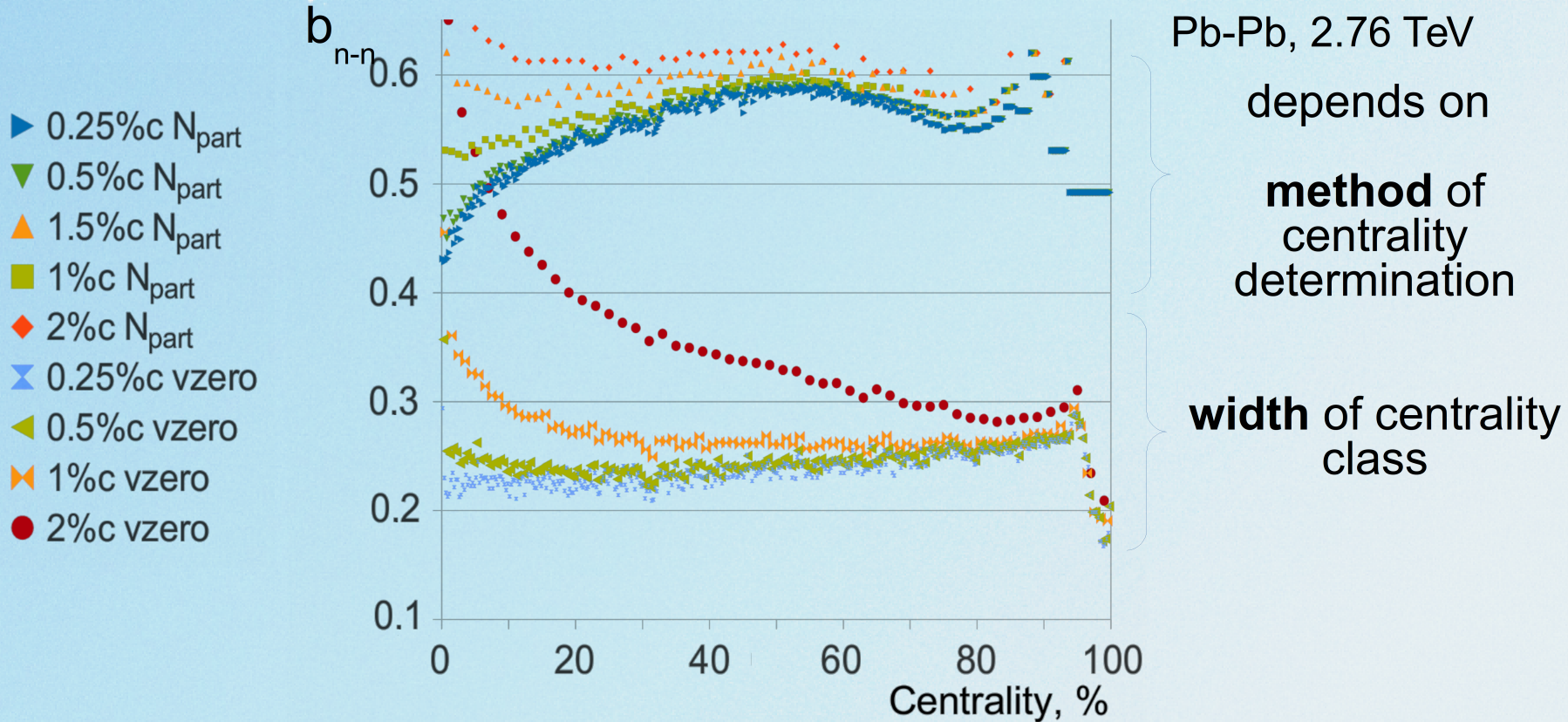
- Nucleon-participants N_{part} - nucleons collided at least once
- Nucleon-spectators N_{spect} - nucleons, which didn't interact
- Number of nucleon-nucleon collisions N_{coll}
- Multiplicity of charge particles N_{ch}

Forward-backward correlations in AA

multiplicity – extensive variable

n-n correlation coefficient

$$b = \frac{\langle FB \rangle - \langle F \rangle \langle B \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$



Mean pt-pt correlations in SF model

mean event transverse momentum – **intensive** variable

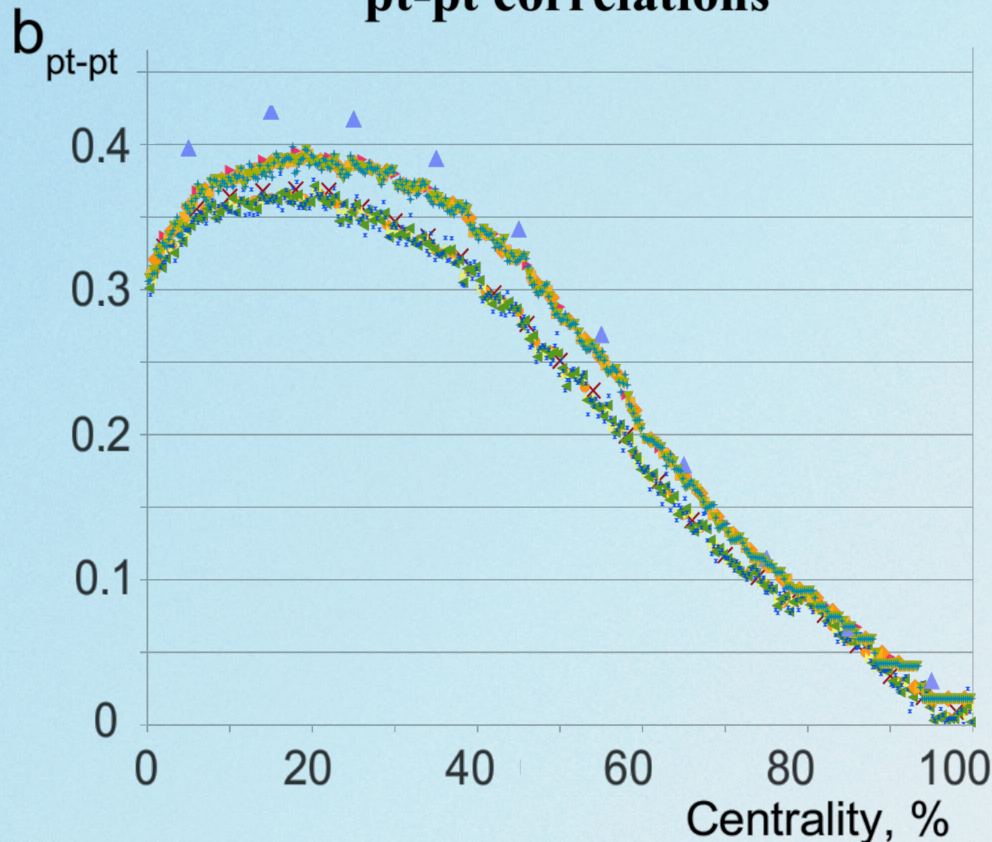
pt-pt correlation
coefficient

$$b = \frac{\langle FB \rangle - \langle F \rangle \langle B \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$

$$B, F, \rightarrow p_t = \frac{1}{n} \sum_{i=1}^n p_{ti}$$

Pb-Pb, 2.76 TeV

pt-pt correlations



robust

- ✦ 0.25% c N_{part}
- ▼ 0.5% c N_{part}
- ▲ 1.5% c N_{part}
- 1% c N_{part}
- ◆ 2% c N_{part}
- ▶ 4% c N_{part}
- ▲ 10% c N_{part}
- ✕ 0.25% c v_{zero}
- ▼ 0.5% c v_{zero}
- ✦ 1% c v_{zero}
- 2% c v_{zero}
- ✕ 4% c v_{zero}

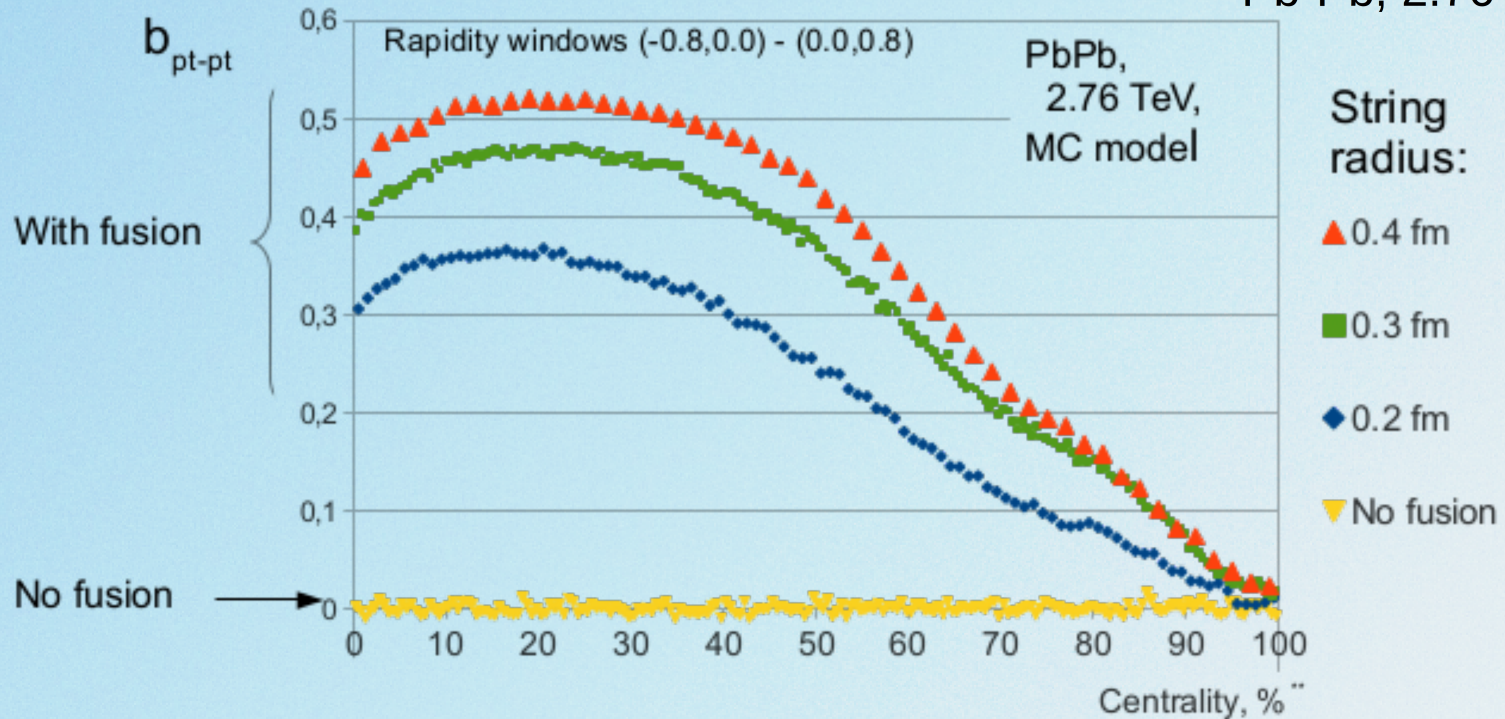
Mean pt-pt correlations in SF model

correlation
coefficient

$$b = \frac{\langle FB \rangle - \langle F \rangle \langle B \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$

$$B, F, \rightarrow p_t = \frac{1}{n} \sum_{i=1}^n p_{ti}$$

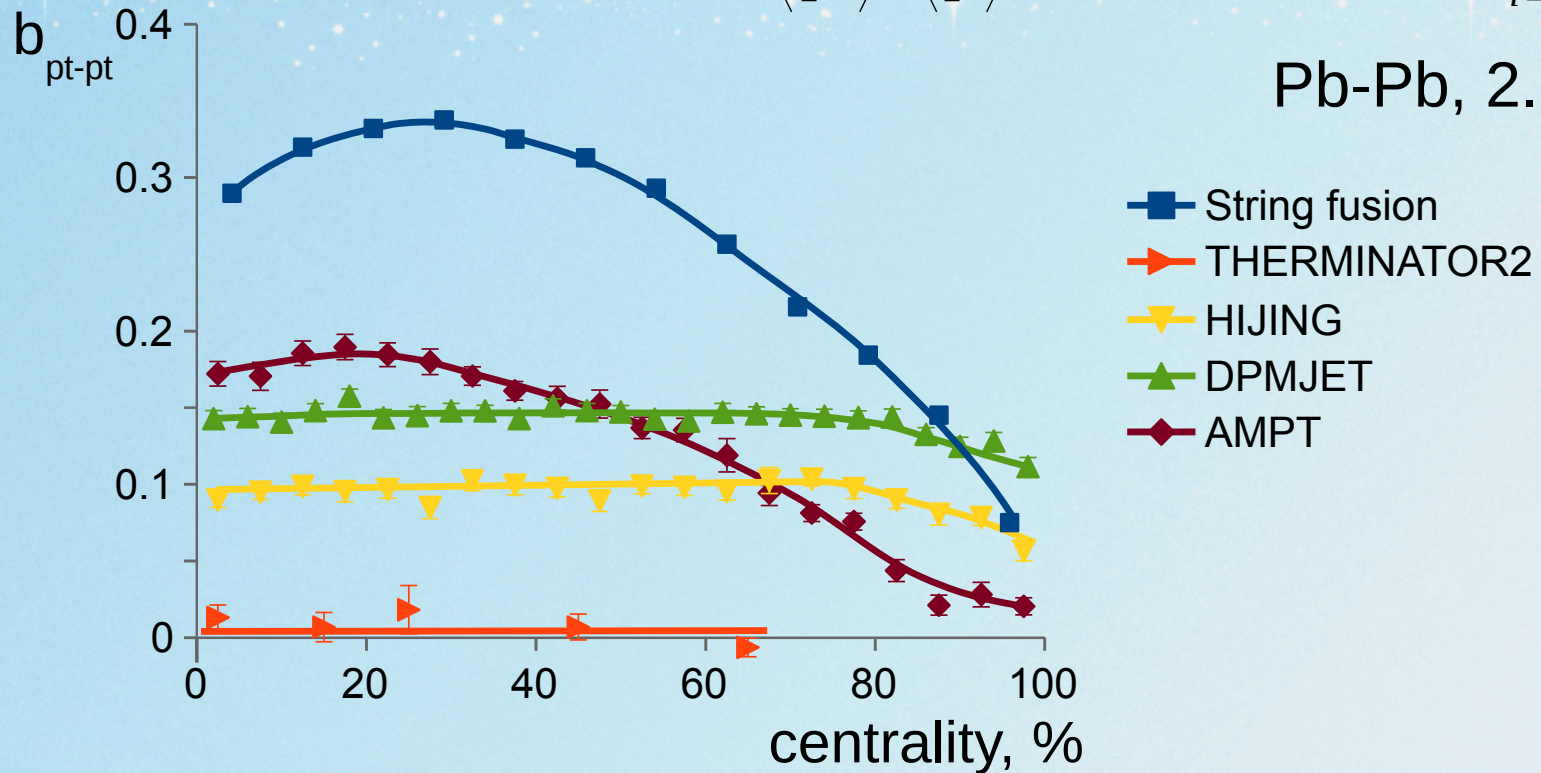
Pb-Pb, 2.76 TeV



Mean pt-pt correlations in different models

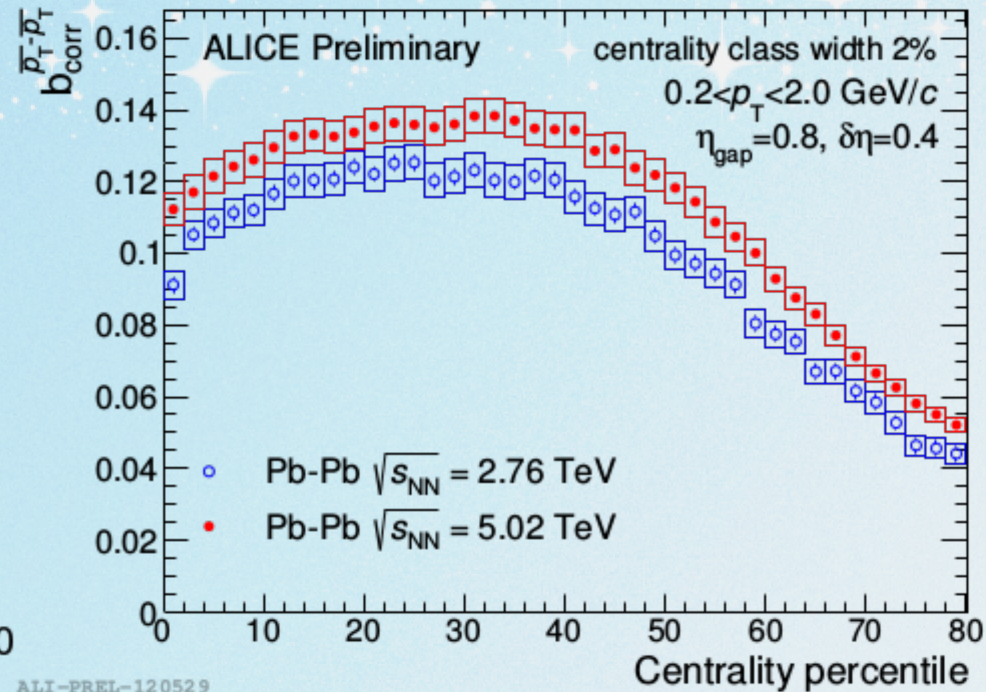
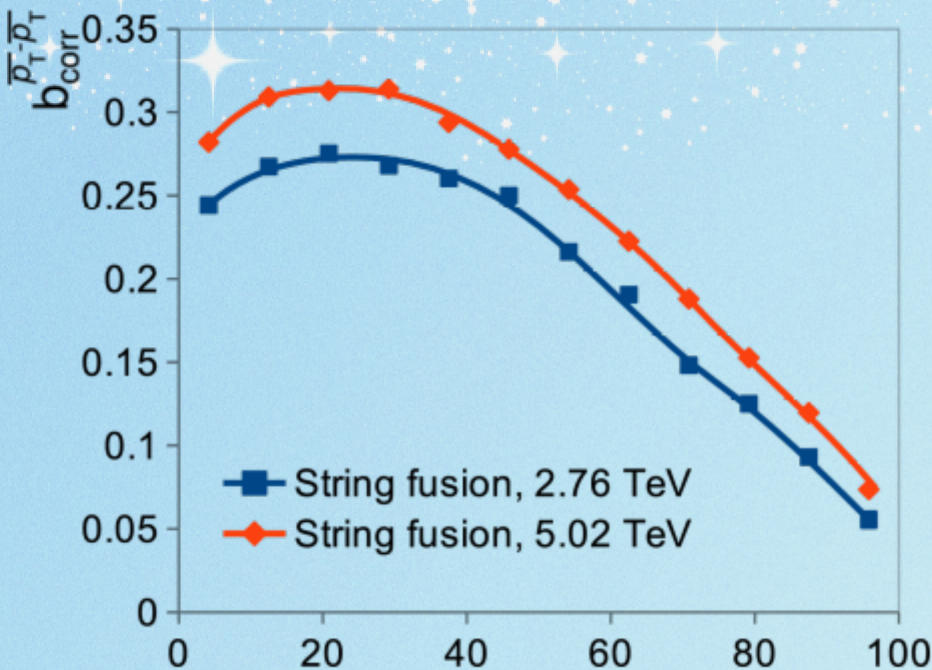
$$b = \frac{\langle FB \rangle - \langle F \rangle \langle B \rangle}{\langle F^2 \rangle - \langle F \rangle^2} \quad B, F, \rightarrow p_t = \frac{1}{n} \sum_{i=1}^n p_{ti}$$

Pb-Pb, 2.76 TeV



Vladimir Kovalenko, Vladimir Vechernin, J. Phys. Conf. Ser. 798, 012053 (2017),
arXiv:1611.07274 [nucl-th]

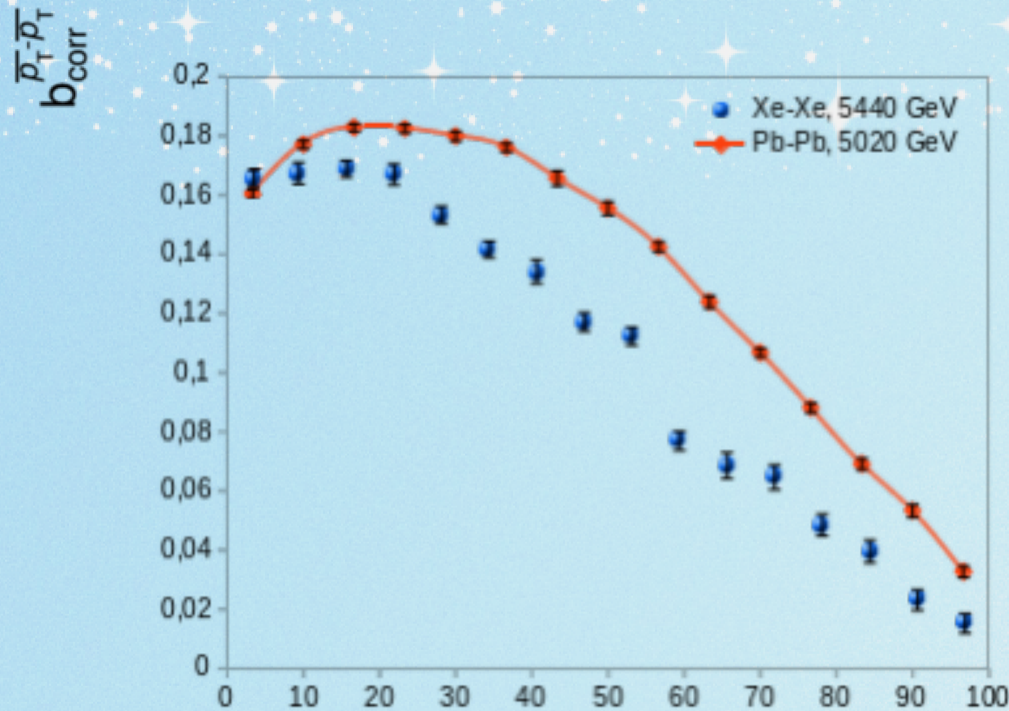
Energy dependence of pt-pt correlation coefficient



- String fusion model predicts the increase of $b_{\text{pt-pt}}$ from 2.76 TeV to 5.02 TeV at the level of 20%, that agrees with the experiment

Preliminary experimental data – I. Altsybeev, “KnE Energy and Physics, pages 304-312, 2018 DOI: 10.18502/ken.v3i1.1759 arXiv:1711.04844 [nucl-ex]

System size dependence of pt-pt correlation coefficient

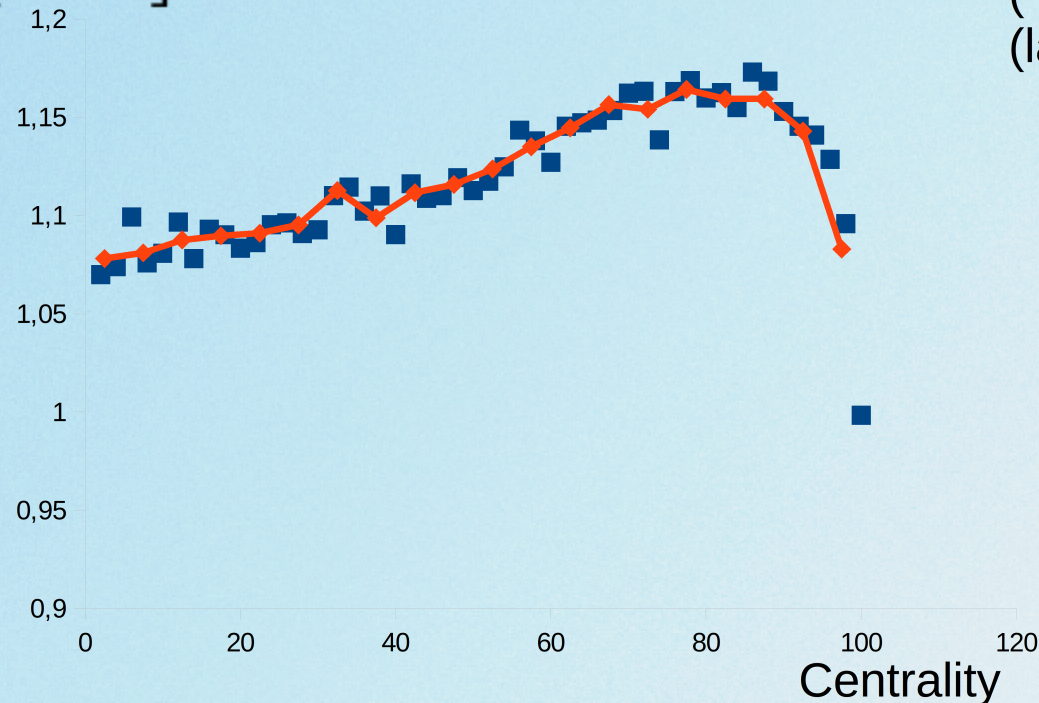


- Correlation coefficients $b_{\text{pt-pt}}$ shows decrease in very central collisions only in Pb-Pb, in smaller system Xe-Xe it reaches plateau.

Long-range Strongly Intensive multiplicity fluctuations

$$\Sigma[n_F, n_B] = \frac{\langle n_B \rangle \omega[n_F] + \langle n_F \rangle \omega[n_B] - 2(\langle n_B n_F \rangle - \langle n_B \rangle \langle n_F \rangle)}{\langle n_B \rangle + \langle n_F \rangle}$$

$\Sigma[n_F, n_B]$

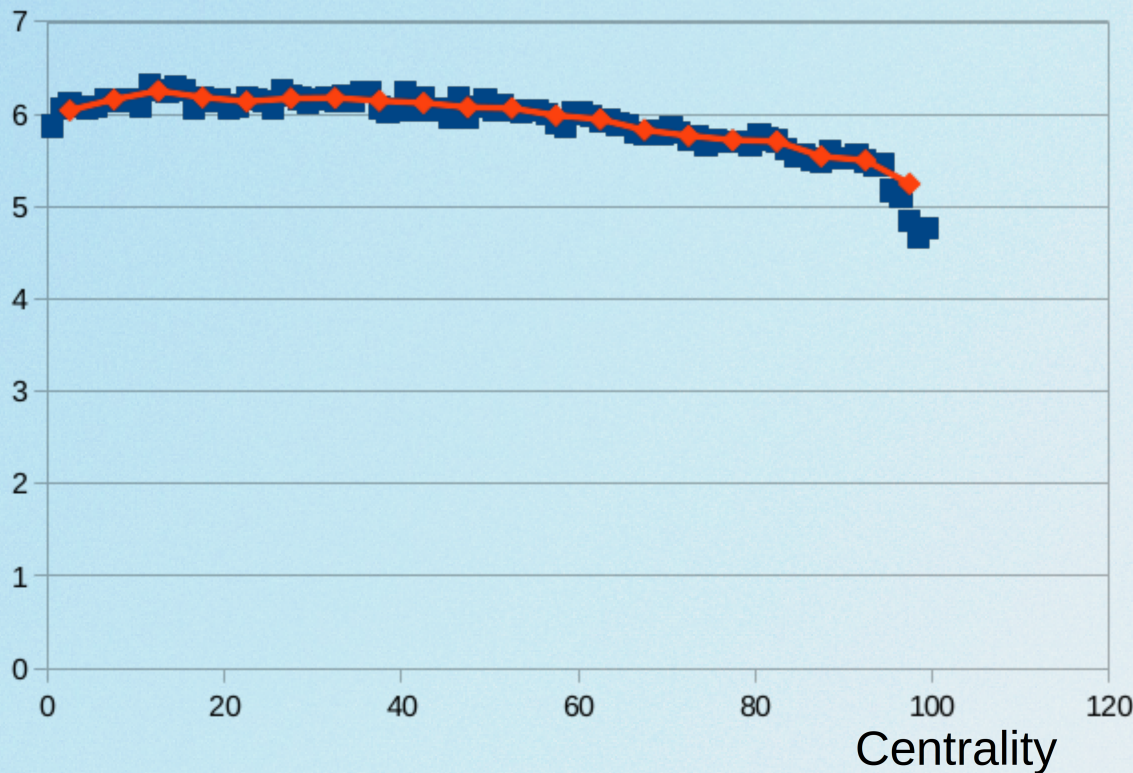


Long-range SI pt-multiplicity fluctuations

$$\Sigma[A, B] = \frac{1}{C_\Sigma} (\langle B \rangle_\omega [A] + \langle A \rangle_\omega [B] - 2(\langle AB \rangle - \langle A \rangle \langle B \rangle))$$

$$A = P_T = \sum_{i=1}^N p_{T_i}, \quad B = N \quad \longrightarrow \quad C_\Delta = C_\Sigma = \langle N \rangle_\omega [p_T]$$

$$\Sigma[P_T | N]$$

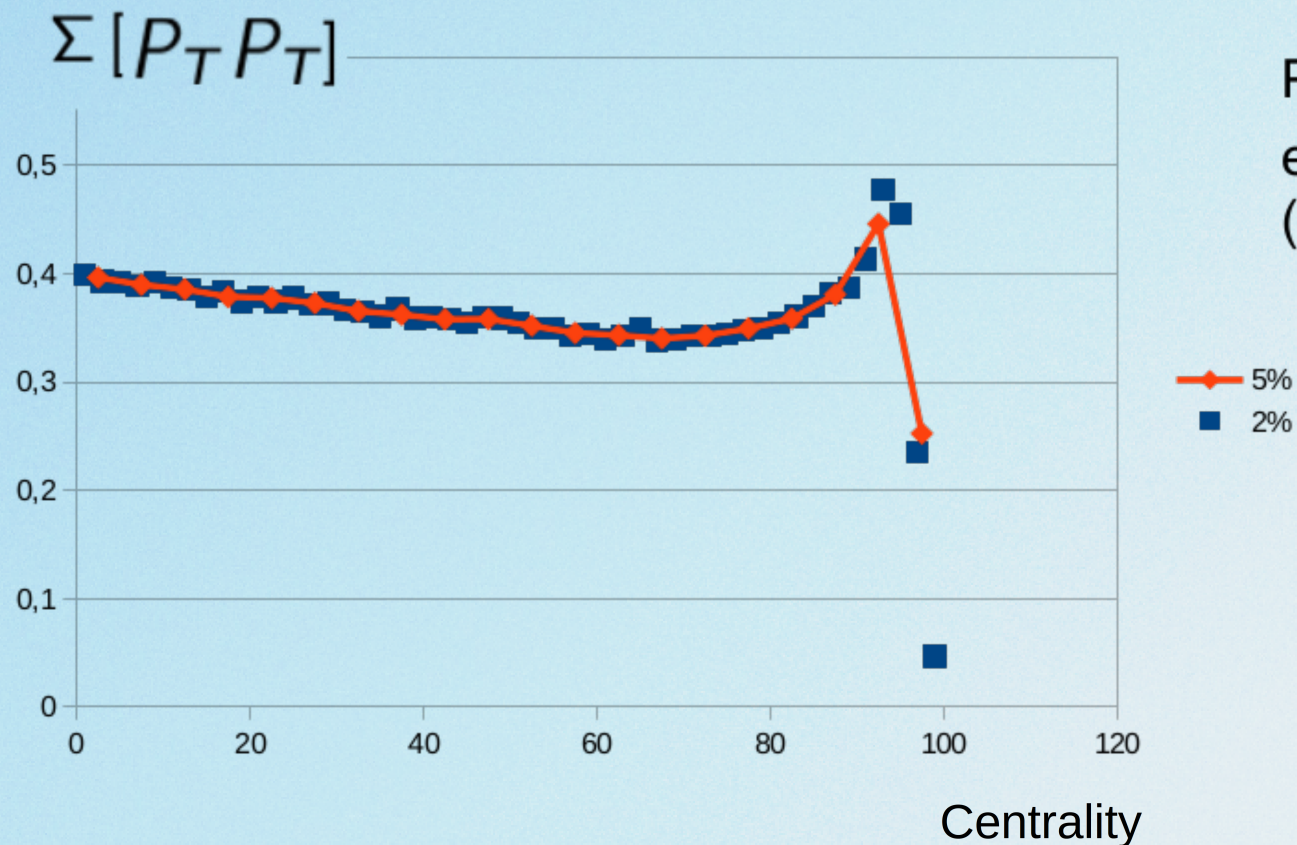


Pb-Pb, 2.76 TeV
eta window:
(-0.8, 0.8)

—◆— 5%
—■— 2%

Long-range SI pt-pt fluctuations

$$\Sigma[A, B] = \frac{1}{C_\Sigma} (\langle B \rangle_\omega [A] + \langle A \rangle_\omega [B] - 2(\langle AB \rangle - \langle A \rangle \langle B \rangle))$$



Pb-Pb, 2.76 TeV
eta window:
(-0.8, 0.8)

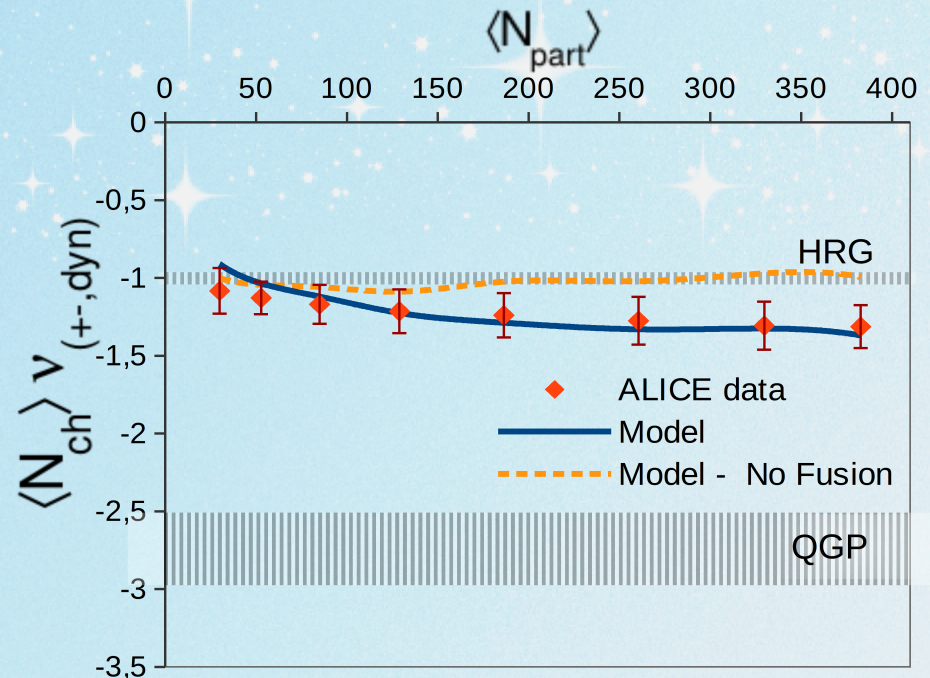
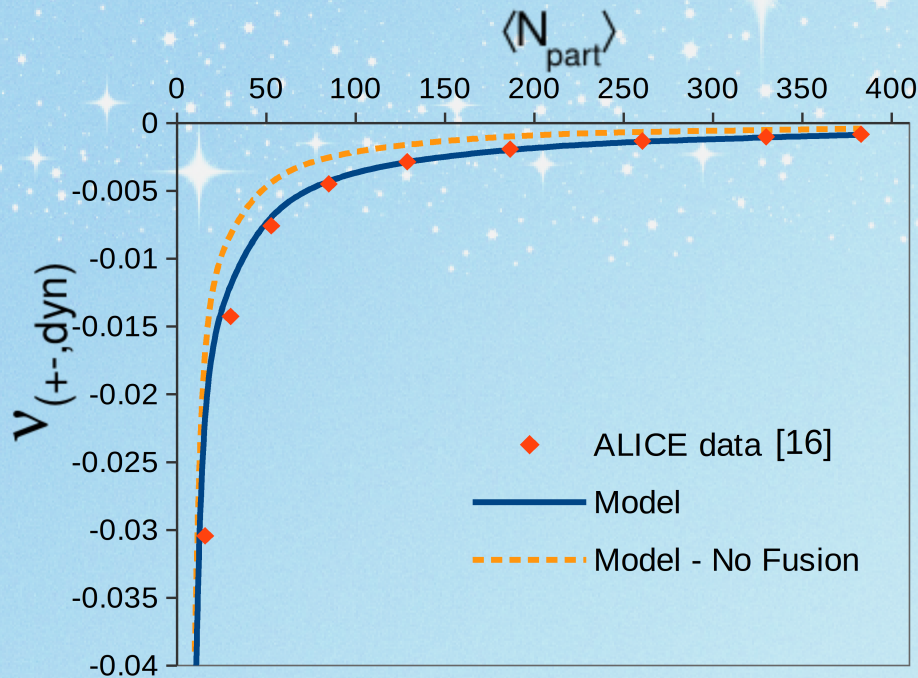
Charge fluctuations

- Expressed in terms of

$$v_{\text{dyn}} = \left\langle \left(\frac{N_+}{\langle N_+ \rangle} - \frac{N_-}{\langle N_- \rangle} \right)^2 \right\rangle - \left(\frac{1}{\langle N_+ \rangle} + \frac{1}{\langle N_- \rangle} \right)$$

- Can be related with degrees of freedom of medium
- $\langle N_{\text{ch}} \rangle v_{(+-, \text{dyn})} = -1$ for hadron gas
- $\langle N_{\text{ch}} \rangle v_{(+-, \text{dyn})} = -2.5 \dots -3$ for QGP

Charge fluctuations

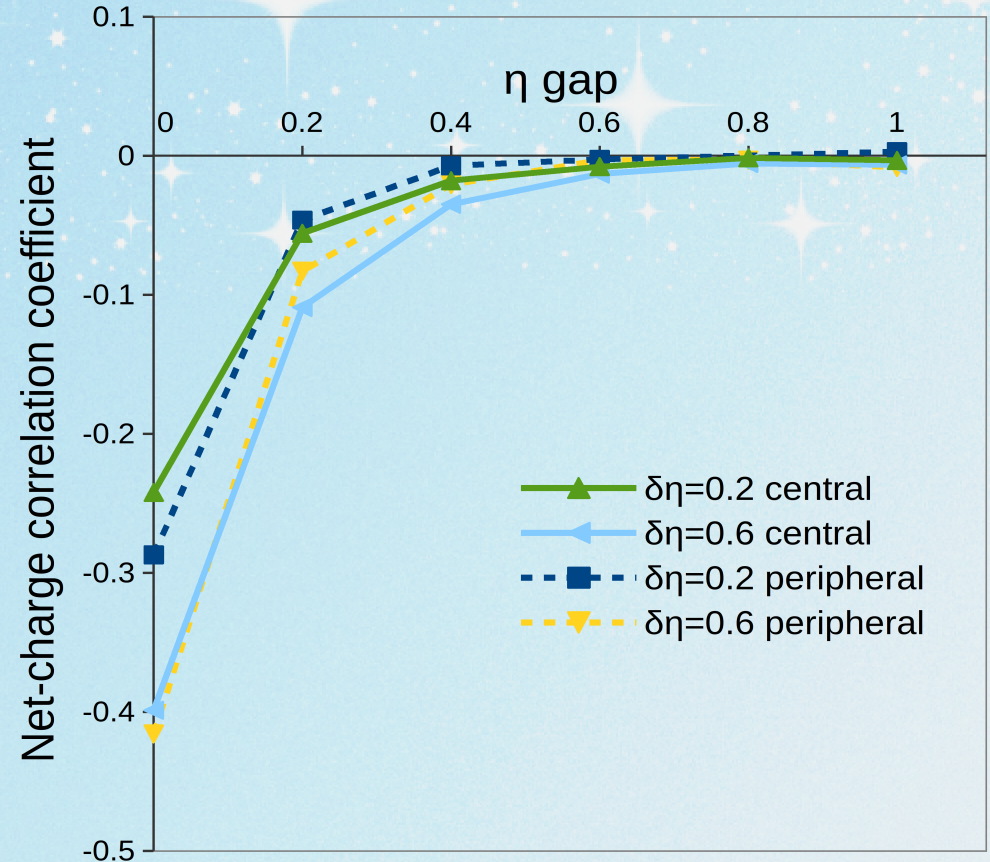
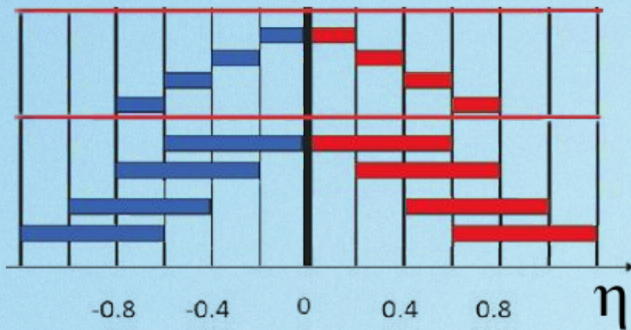


- MC model with string fusion successfully describes the centrality dependence of dynamical net-charge fluctuation.
- Scaling variable $\langle N_{ch} \rangle V_{+-,dyn}$ decreases with centrality towards the level of QGP estimation (which is in agreement with experiment)

In case of no fusion, it remains constant at the level of HRG

Net-charge correlations

$$b = \frac{\langle Q_F Q_B \rangle - \langle Q_F \rangle \langle Q_B \rangle}{\langle Q_F^2 \rangle - \langle Q_F \rangle^2}$$



* In small window can be related to balance function
(thanks to I. Altybeev for pointing it out)

- Absolute value of net-charge forward-backward correlation coefficient decreases with the gap.
- Weak dependence on the centrality
- A hint of a transition from negative to positive at large rapidity gap – to be checked at lower energies (RHIC, SPS, NICA)

Conclusions

- Centrality determination and a width of centrality class influences the value of fluctuational and correlational observables
- Strongly intensive observables and correlations between intensive observables show robust behaviour against the volume fluctuations and the centrality determination methods
- String fusion model explains non-monotonic behaviour of Pt-pt forward-backward correlations
- String fusion model explains the centrality behaviour of the net charge fluctuations in Pb-Pb collisions at LHC energy (string fusion improves the agreement with the experiment).
- The predictions for net-charge forward-backward correlation coefficient have been obtained.
- The observables proposed provide reference to the properties of the initial state of AA collisions and are useful to study both at LHC and NICA energies