

Quantum spin dynamics in external classical fields

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Outline

- 1 Relativistic particles in curved spacetimes
 - Dynamics of spin and equivalence principle
 - Geometry of spacetime
- 2 Spin $\frac{1}{2}$ particle in curved spacetime
 - Dirac Hamiltonian for arbitrary metric
 - Electrodynamics in curved spacetime
 - Foldy-Wouthuysen Hamiltonian and equations of motion
- 3 Physical effects of spin dynamics
 - Manifestations in high energy physics experiments
 - Spin dynamics in a gravitational wave
 - Probing spacetime geometry
- 4 Conclusions and Outlook

Dynamics of spin and equivalence principle

- High-energy experiments take place in curved space or in noninertial frame (for example, on Earth)
- Equivalence principle (EP) - a cornerstone of gravity
- Newton's theory \Rightarrow Einstein's "falling elevator"
- Colella-Overhauser-Werner (1975) and Bonse-Wroblewski experiments - EP for quantum-mechanical systems:
- Measured phase shift due to inertial and gravitational force
- Gravity on *spin*: EP for relativistic particles?
- Classical theory of spin: Frenkel (1928), Mathisson (1937), Papapetrou (1951), Weyssenhoff-Raabe (1947)
- Compare classical rotator and quantum spin
- Measure spin effects to probe spacetime geometry!

Arbitrary Riemannian geometry in 4 dimensions

- Let t be time, x^a ($a = 1, 2, 3$) be spatial coordinates:

$$ds^2 = V^2 c^2 dt^2 - \delta_{\hat{a}\hat{b}} W^{\hat{a}}_c W^{\hat{b}}_d (dx^c - K^c c dt) (dx^d - K^d c dt)$$

V and K^a , and 3×3 matrix $W^{\hat{a}}_b$ depend arbitrarily on t, x^a .

- Their number $1 + 3 + 9 = 13$ but rotation $W^{\hat{a}}_b \rightarrow L^{\hat{a}}_{\hat{c}} W^{\hat{c}}_b$ is allowed with arbitrary $L^{\hat{a}}_{\hat{c}}(t, x) \in SO(3)$: $\implies 13 - 3 = 10$
- Coframe e_i^α with $g_{\alpha\beta} e_i^\alpha e_j^\beta = g_{ij}$, $g_{\alpha\beta} = \text{diag}(c^2, -1, -1, -1)$:

$$e_i^{\hat{0}} = V \delta_i^0, \quad e_i^{\hat{a}} = W^{\hat{a}}_b \left(\delta_i^b - c K^b \delta_i^0 \right), \quad a = 1, 2, 3$$

- Exact metric of flat spacetime in noninertial frame

$$V = 1 + \frac{\mathbf{a} \cdot \mathbf{r}}{c^2}, \quad W^{\hat{a}}_b = \delta_b^a, \quad K^a = -\frac{1}{c} (\boldsymbol{\omega} \times \mathbf{r})^a$$

Dirac particle in gravitational & electromagnetic field

- Fermion with moments (AMM $\mu' = \frac{(g-2)e\hbar}{4m}$ & EDM $\delta' = \frac{be\hbar}{2mc}$)

$$\left(i\hbar\gamma^\alpha D_\alpha - mc + \frac{\mu'}{2c}\sigma^{\alpha\beta}F_{\alpha\beta} + \frac{\delta'}{2}\sigma^{\alpha\beta}G_{\alpha\beta} \right)\psi = 0$$

- Spinor covariant derivative (with $\sigma_{\alpha\beta} = i\gamma_{[\alpha}\gamma_{\beta]}$)

$$D_\alpha = e^i_\alpha D_i, \quad D_i = \partial_i - \frac{ie}{\hbar}A_i + \frac{i}{4}\sigma^{\alpha\beta}\Gamma_{i\alpha\beta}$$

- Connection for general spacetime geometry

$$\Gamma_{i\hat{a}\hat{b}} = \frac{c^2}{V}W^b_{\hat{a}}\partial_b V e_i^{\hat{b}} - \frac{c}{V}Q_{(\hat{a}\hat{b})}e_i^{\hat{b}},$$

$$\Gamma_{i\hat{a}\hat{b}} = \frac{c}{V}Q_{[\hat{a}\hat{b}]}e_i^{\hat{b}} + (C_{\hat{a}\hat{b}\hat{c}} + C_{\hat{a}\hat{c}\hat{b}} + C_{\hat{c}\hat{b}\hat{a}})e_i^{\hat{c}}$$

- Here anholonomy $C_{\hat{a}\hat{b}}^{\hat{c}} = W^d_{\hat{a}}W^e_{\hat{b}}\partial_{[d}W^{\hat{c}}_{e]}$ and

$$Q_{\hat{a}\hat{b}} = g_{\hat{a}\hat{c}}W^d_{\hat{b}}\left(\frac{1}{c}\dot{W}^{\hat{c}}_d + K^e\partial_e W^{\hat{c}}_d + W^{\hat{c}}_e\partial_d K^e\right)$$

Dirac Hamiltonian

- Naive Hamiltonian is not Hermitian. Rescale wave function $\psi \rightarrow \left(\sqrt{-g}e_0^0\right)^{\frac{1}{2}} \psi$ and recast Dirac wave equation into Schrodinger form $i\hbar\frac{\partial\psi}{\partial t} = \mathcal{H}\psi$

Dirac Hamiltonian (with $\mathcal{F}^b_a = VW^b_{\hat{a}}$ and $\pi = -i\hbar\nabla - eA$)

$$\begin{aligned} \mathcal{H} = & \beta mc^2 V + e\Phi + \frac{c}{2} (\pi_b \mathcal{F}^b_a \alpha^a + \alpha^a \mathcal{F}^b_a \pi_b) \\ & + \frac{c}{2} (\mathbf{K} \cdot \boldsymbol{\pi} + \boldsymbol{\pi} \cdot \mathbf{K}) + \frac{\hbar c}{4} (\boldsymbol{\Xi} \cdot \boldsymbol{\Sigma} - \Upsilon \gamma_5) \\ & - \beta V (\boldsymbol{\Sigma} \cdot \boldsymbol{\mathcal{M}} + i\boldsymbol{\alpha} \cdot \boldsymbol{\mathcal{P}}) \end{aligned}$$

- Here $\beta = \gamma^{\hat{0}}$, $\alpha^a = \gamma^{\hat{0}}\gamma^{\hat{a}}$, $\gamma_5 = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix}$, $\boldsymbol{\Sigma} = \begin{pmatrix} \boldsymbol{\sigma} & 0 \\ 0 & \boldsymbol{\sigma} \end{pmatrix}$,

$$\Upsilon = V \epsilon^{\hat{a}\hat{b}\hat{c}} \Gamma_{\hat{a}\hat{b}\hat{c}} = -V \epsilon^{\hat{a}\hat{b}\hat{c}} \mathcal{C}_{\hat{a}\hat{b}\hat{c}}, \quad \Xi_{\hat{a}} = \frac{V}{c} \epsilon_{\hat{a}\hat{b}\hat{c}} \Gamma_{\hat{0}}^{\hat{b}\hat{c}} = \epsilon_{\hat{a}\hat{b}\hat{c}} Q^{\hat{b}\hat{c}}$$

Electrodynamics in curved spacetime

- Gravity is universal: affects also electromagnetism. How?
- Basic objects: field strength F , excitation H and current J

Maxwell's theory – without coordinates and frames

$$dF = 0, \quad dH = J, \quad H = \lambda_0 \star F, \quad \lambda_0 = \sqrt{\epsilon_0/\mu_0}$$

- Coordinates x^i : $F = \frac{1}{2}F_{ij}dx^i \wedge dx^j$, $H = \frac{1}{2}H_{ij}dx^i \wedge dx^j$,
 and $J = \frac{1}{6}J_{ijk}dx^i \wedge dx^j \wedge dx^k$ are (1 + 3) decomposed:

$$\mathbf{E}_a = \{F_{10}, F_{20}, F_{30}\}, \quad \mathbf{B}^a = \{F_{23}, F_{31}, F_{12}\}$$

$$\mathbf{H}_a = \{H_{01}, H_{02}, H_{03}\}, \quad \mathbf{D}^a = \{H_{23}, H_{31}, H_{12}\}$$

- $\mathbf{J}^a = \{-J_{023}, -J_{031}, -J_{012}\}, \quad \rho = J_{123}$
- Maxwell equations are recast into standard form

$$\nabla \times \mathbf{E} + \dot{\mathbf{B}} = 0, \quad \nabla \cdot \mathbf{B} = 0,$$

$$\nabla \times \mathbf{H} - \dot{\mathbf{D}} = \mathbf{J}, \quad \nabla \cdot \mathbf{D} = \rho$$

- Gravity/inertia encoded in *constitutive relation* $H = H(F)$

$$D^a = \frac{\varepsilon_0 w}{V} \underline{g}^{ab} E_b - \lambda_0 \frac{w}{V} \underline{g}^{ad} \epsilon_{bcd} K^c B^b,$$

$$H_a = \frac{1}{\mu_0 w V} \{ (V^2 - K^2) \underline{g}_{ab} + K_a K_b \} B^b - \lambda_0 \frac{w}{V} \epsilon_{adc} K^c \underline{g}^{db} E_b$$

Here $K_a = \underline{g}_{ab} K^b$, $K^2 = \underline{g}_{ab} K^a K^b$ and $w = \det W^{\hat{c}}_d$.

- Frame e_i^α needed for fermions $\implies F_{\alpha\beta} = e_\alpha^i e_\beta^j F_{ij}$
- Components: $\mathfrak{E}_a = \{ \widehat{F}_{10}, \widehat{F}_{20}, \widehat{F}_{30} \}$ & $\mathfrak{B}^a = \{ \widehat{F}_{23}, \widehat{F}_{31}, \widehat{F}_{12} \}$
- Relation between holonomic and anholonomic fields

$$\mathfrak{E}_a = \frac{1}{V} W^b_{\hat{a}} (\mathbf{E} + c\mathbf{K} \times \mathbf{B})_b, \quad \mathfrak{B}^a = \frac{1}{w} W^{\hat{a}}_b \mathbf{B}^b$$

- Nonminimal coupling $-\beta V (\boldsymbol{\Sigma} \cdot \mathcal{M} + i\alpha \cdot \mathcal{P})$ governed by

$$\mathcal{M}^a = \mu' \mathfrak{B}^a + \delta' \mathfrak{E}^a, \quad \mathcal{P}_a = c\delta' \mathfrak{B}_a - \mu' \mathfrak{E}_a / c.$$

- Foldy-Wouthuysen transformation needed to reveal physics (uncouple positive and negative energy states)
- Recast generic Hamiltonian into

$$\mathcal{H} = \beta\mathfrak{M} + \mathcal{E} + \mathcal{O}, \quad \beta\mathfrak{M} = \mathfrak{M}\beta, \quad \beta\mathcal{E} = \mathcal{E}\beta, \quad \beta\mathcal{O} = -\mathcal{O}\beta$$

- Foldy-Wouthuysen unitary transformation

$$\psi_{FW} = U\psi, \quad \mathcal{H}_{FW} = U\mathcal{H}U^{-1} - i\hbar U\partial_t U^{-1}$$

- In arbitrary external fields (with $\epsilon = \sqrt{\mathfrak{M}^2 + \mathcal{O}^2}$)

$$U = \frac{\beta\epsilon + \beta\mathfrak{M} - \mathcal{O}}{\sqrt{(\beta\epsilon + \beta\mathfrak{M} - \mathcal{O})^2}} \beta$$

- For Dirac fermion we have explicitly: $\mathfrak{M} = mc^2V$ and

$$\begin{aligned} \mathcal{E} &= e\Phi + \frac{c}{2} (\mathbf{K} \cdot \boldsymbol{\pi} + \boldsymbol{\pi} \cdot \mathbf{K}) + \frac{\hbar c}{4} \boldsymbol{\Xi} \cdot \boldsymbol{\Sigma} - \beta V \boldsymbol{\Sigma} \cdot \boldsymbol{\mathcal{M}}, \\ \mathcal{O} &= \frac{c}{2} (\pi_b \mathcal{F}^b{}_a \alpha^a + \alpha^a \mathcal{F}^b{}_a \pi_b) - \frac{\hbar c}{4} \Upsilon \gamma_5 - i\beta V \boldsymbol{\alpha} \cdot \boldsymbol{\mathcal{P}} \end{aligned}$$

- FW Hamiltonian $\mathcal{H}_{FW} = \mathcal{H}_{FW}^{(1)} + \mathcal{H}_{FW}^{(2)} + \mathcal{H}_{FW}^{(3)} + \mathcal{H}_{FW}^{(4)}$:

$$\mathcal{H}_{FW}^{(1)} = \beta\epsilon' + \frac{\hbar c^2}{16} \left\{ \frac{1}{\epsilon'}, \left(2\epsilon^{cae} \Pi_e \{ \pi_b, \mathcal{F}^d{}_c \partial_d \mathcal{F}^b{}_a \} + \Pi^a \{ \pi_b, \mathcal{F}^b{}_a \Upsilon \} \right) \right\} \\ + \frac{\hbar m c^4}{4} \epsilon^{cae} \Pi_e \left\{ \frac{1}{\mathcal{T}}, \left\{ \pi_d, \mathcal{F}^d{}_c \mathcal{F}^b{}_a \partial_b V \right\} \right\},$$

$$\mathcal{H}_{FW}^{(2)} = \frac{c}{2} (K^a \pi_a + \pi_a K^a) + \frac{\hbar c}{4} \Sigma_a \Xi^a + \frac{\hbar c^2}{16} \left\{ \frac{1}{\mathcal{T}}, \left\{ \Sigma_a \{ \pi_e, \mathcal{F}^e{}_b \}, \left\{ \pi_f, \left[\epsilon^{abc} \times \right. \right. \right. \right. \\ \left. \left. \left. \left. \times \left(\frac{1}{c} \dot{\mathcal{J}}^f{}_c - \mathcal{F}^d{}_c \partial_d K^f + K^d \partial_d \mathcal{F}^f{}_c \right) - \frac{1}{2} \mathcal{F}^f{}_d \left(\delta^{db} \Xi^a - \delta^{da} \Xi^b \right) \right] \right\} \right\} \right\},$$

$$\mathcal{H}_{FW}^{(3)} = e\Phi - \frac{e\hbar c^2}{4} \left\{ \frac{1}{\epsilon'}, V^2 \Pi^a \mathfrak{B}_a \right\} - \frac{e\hbar c^2}{8} \left\{ \frac{1}{\mathcal{T}}, \left[2\hbar \mathcal{F}^b{}_a \partial_b (V^2 \mathfrak{E}^a) \right. \right. \\ \left. \left. - \Sigma_a \epsilon^{abc} \left(\{ \mathcal{F}^d{}_b, \pi_d \} V^2 \mathfrak{E}_c - V^2 \mathfrak{E}_b \{ \mathcal{F}^d{}_c, \pi_d \} \right) \right] \right\},$$

$$\mathcal{H}_{FW}^{(4)} = -\frac{c}{8} \left\{ \frac{1}{\epsilon'}, \left[\Sigma_a \epsilon^{abc} \left(\{ \mathcal{F}^d{}_b, \pi_d \} V \mathcal{P}_c - V \mathcal{P}_b \{ \mathcal{F}^d{}_c, \pi_d \} \right) - 2\hbar \mathcal{F}^b{}_a \partial_b (V \mathcal{P}^a) \right] \right\} \\ - V \Pi^a \mathcal{M}_a + \frac{c^2}{4} \left\{ \frac{1}{\mathcal{T}}, \left(\Pi^a \{ \{ \mathcal{F}^c{}_a \mathcal{F}^d{}_b V \mathcal{M}^b, \pi_c \}, \pi_d \} + \beta \hbar \{ \mathcal{F}^b{}_a [\mathcal{J}^a + K^c \partial_c (V \mathcal{P}^a)], \pi_b \} \right) \right\}$$

- Here $\{ , \}$ anticommutators, $\mathcal{T} = 2\epsilon'^2 + \{ \epsilon', m c^2 V \}$, $\mathbf{\Pi} = \beta \mathbf{\Sigma}$,

$$\epsilon' = \sqrt{m^2 c^4 V^2 + \frac{c^2}{4} \delta^{ac} \{ \pi_b, \mathcal{F}^b{}_a \} \{ \pi_d, \mathcal{F}^d{}_c \}}, \quad \mathcal{J}^a = \epsilon^{abc} \mathcal{F}^d{}_b \partial_d (V \mathcal{M}_c) + \frac{\partial \mathcal{P}^a}{c \partial t}$$

- This result is exact – no (weak field etc) approximations for $V, W_{\underline{b}}, K_{\underline{a}}$.

Quantum dynamics of spinning particle

- Evolution of spin (polarization operator $\mathbf{\Pi} = \beta \mathbf{\Sigma}$)

$$\frac{d\mathbf{\Pi}}{dt} = \frac{i}{\hbar} [\mathcal{H}_{FW}, \mathbf{\Pi}] = \mathbf{\Omega}_{(1)} \times \mathbf{\Sigma} + \mathbf{\Omega}_{(2)} \times \mathbf{\Pi}$$

- Semiclassical precession velocity of spin

$$\Omega_{(1)}^a = \frac{c^2}{\epsilon} \mathcal{F}^d{}_c \pi_d \left(\frac{1}{2} \Upsilon \delta^{ac} - \epsilon^{akl} V C_{kl}{}^c + \frac{\epsilon}{\epsilon + mc^2 V} \epsilon^{abc} W^k{}_{\hat{b}} \partial_d V \right. \\ \left. + \frac{eV^2}{\epsilon' + mc^2 V} \epsilon^{acb} \mathfrak{E}_b - \frac{2V}{c\hbar} \epsilon^{acb} \mathcal{P}_b \right)$$

$$\Omega_{(2)}^a = \frac{c}{2} \Xi^a - \frac{c^3}{\epsilon(\epsilon + mc^2 V)} \epsilon^{abc} Q_{(bd)} \delta^{dn} \mathcal{F}^k{}_n \pi_k \mathcal{F}^l{}_c \pi_l \\ - \frac{ec^2 V^2}{\epsilon} \mathfrak{B}^a + \frac{2V}{\hbar} \left(-\mathcal{M}^a + \frac{c^2}{\epsilon(\epsilon + mc^2 V)} \delta^{an} \mathcal{F}^c{}_n \pi_c \mathcal{F}^d{}_b \pi_d \mathcal{M}^b \right)$$

- Here $\epsilon = \sqrt{m^2 c^4 V^2 + c^2 \delta^{cd} \mathcal{F}^a{}_c \mathcal{F}^b{}_d \pi_a \pi_b}$

Semiclassical FW Hamiltonian

$$\mathcal{H}_{FW} = \beta \epsilon + e\Phi + \frac{c}{2} (\mathbf{K} \cdot \boldsymbol{\pi} + \boldsymbol{\pi} \cdot \mathbf{K}) + \frac{\hbar}{2} \mathbf{\Pi} \cdot \mathbf{\Omega}_{(1)} + \frac{\hbar}{2} \mathbf{\Sigma} \cdot \mathbf{\Omega}_{(2)}$$

- Physical spin \mathbf{s} precesses wrt rest frame: $\frac{d\mathbf{s}}{dt} = \boldsymbol{\Omega} \times \mathbf{s}$

Spin dynamics on Earth (with $\mathbf{g} = -\frac{GM}{r^3} \mathbf{r}$, $\gamma = 1/\sqrt{1-v^2/c^2}$)

$$\begin{aligned} \boldsymbol{\Omega} = & \frac{e}{m} \left\{ -\frac{1}{\gamma} \mathfrak{B} + \frac{1}{\gamma+1} \frac{\mathbf{v} \times \boldsymbol{\mathcal{E}}}{c^2} \right\} - \boldsymbol{\omega} + \frac{2\gamma+1}{\gamma+1} \frac{\mathbf{v} \times \mathbf{g}}{c^2} \\ & - \frac{2\mu'}{\hbar} \left\{ \mathfrak{B} - \frac{\mathbf{v} \times \boldsymbol{\mathcal{E}}}{c^2} - \frac{\gamma}{\gamma+1} \mathbf{v} \frac{\mathfrak{B} \cdot \mathbf{v}}{c^2} \right\} \\ & - \frac{2\delta'}{\hbar} \left\{ \boldsymbol{\mathcal{E}} + \mathbf{v} \times \mathfrak{B} - \frac{\gamma}{\gamma+1} \mathbf{v} \frac{\boldsymbol{\mathcal{E}} \cdot \mathbf{v}}{c^2} \right\} \end{aligned}$$

- Analysis of manifestations of terrestrial rotation and gravity in precision high-energy physics: *influence not negligible*
- E.g.: Earth's gravity produces same effect as deuteron's EDM of $\delta' = 1.5 \times 10^{-29}$ e·cm in planned dEDM experiment with magnetic focusing (AGS proposal EDM Collaboration)

Spin dynamics in a gravitational wave

- Gravitational and electromagnetic fields not superimposed but their action on fermion is combined in a nontrivial way \implies new prospects for detection of grav. wave effects?
- In coordinates (t, x, y, z) , weak gravitational wave is

$$V = 1, \quad \mathbf{K} = 0, \quad W^{\hat{a}}_{\hat{b}} = \begin{pmatrix} 1 + w_{\oplus} & w_{\otimes} & 0 \\ w_{\otimes} & 1 - w_{\oplus} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Functions $w_{\otimes}(\varphi)$, $w_{\oplus}(\varphi)$ of phase $\varphi = \omega(t - \frac{z}{c})$, describe 2 polarizations of a plane wave with frequency ω along z

- Hamiltonian for fermion's spin in this spacetime reduces to

$$\mathcal{H}_{FW} = -(\mu_0 + \mu') \boldsymbol{\Pi} \cdot \boldsymbol{\mathfrak{B}}$$

Here Bohr's magneton $\mu_0 = \frac{e\hbar}{2m}$. Important observation: Anholonomic field $\boldsymbol{\mathfrak{B}}^a = W^{\hat{a}}_{\hat{b}} \mathbf{B}^b$ bears "imprint" of the gravitational wave on applied magnetic field \mathbf{B} !

- Recall particle with magnetic moment in flat space (no gravity) in constant homogeneous magnetic field: spin polarized along/against applied field. Additional rotating (alternating) field in plane perpendicular to original field \implies spin flip: magnetic resonance phenomenon occurs
- Suppose $\mathbf{B} = (B_0, 0, 0)$ with $B_0 = \text{const}$, and $w_{\oplus} = 0$,

$$w_{\otimes} = g_0 \cos \varphi = g_0 \cos(\omega t - \omega z/c)$$

describes wave with frequency ω and amplitude g_0 along z

- $\implies \mathfrak{B} = (B_0, B_0 w_{\otimes}, 0)$, ie *magnetic resonance conditions*
- Probability to get at t spin oriented oppositely to initial at t_0

$$P_{-\frac{1}{2}} = \frac{\sin^2 \{ \omega_0 g_0 (t - t_0) \Lambda / 4 \}}{\Lambda^2}$$

- Here Larmor frequency $\omega_0 = 2(\mu_0 + \mu')B_0/\hbar$, and

$$\Lambda^2 = 1 + \frac{4(1 - \xi)^2}{g_0^2}, \quad \xi = \frac{\omega}{\omega_0} \left(1 - \frac{g_0^2}{16\xi^2} \right)$$

Experimental bounds on torsion

- To probe spacetime geometry: dynamics of spin

$$\frac{d\Pi}{dt} = \frac{i}{\hbar} [\mathcal{H}_{FW}, \Pi] = \Omega \times \Pi$$

- *Theory*: spin precession to probe torsion: Adamowicz (1975), Rumpf (1980), Audretsch (1981), Lämmerzahl (1997); review W.T.Ni, Rep.Prog.Phys. 73 (2010) 056901
- *Experiment*: effect of Earth's gravity on nuclear spins Hg
- Spin Hamiltonian (torsion $\check{T}^\alpha = \frac{1}{2}\eta^{\mu\nu\lambda\alpha}T_{\mu\nu\lambda}$, $\check{\mathbf{T}} = \{\check{T}^a\}$)

$$\mathcal{H}_{FW} = -g_N \mu_N \mathbf{B} \cdot \Pi - \frac{\hbar}{2} \boldsymbol{\omega} \cdot \Sigma - \frac{\hbar c}{4} \check{\mathbf{T}} \cdot \Sigma.$$

- B.J. Venema et al, Phys. Rev. Lett. **68** (1992) 135

Limits on torsion from Zeeman energy levels measurements

- $|\check{\mathbf{T}}| < 4.3 \times 10^{-14} \text{m}^{-1}$
- Recent: C. Gemmel et al, Phys. Rev. **D82** (2010) 111901

Conclusions and Outlook

- Searches for spin effects in gravity is fundamental issue. Overview of relevant laboratory experiments: Wei-Tou Ni, Rep. Prog. Phys. **73** (2010) 056901
- Theoretical framework of fermion spin dynamics developed [based on: Obukhov, Silenko, Teryaev, Phys. Rev. **D90** (2014) 124068; Phys. Rev. **D94** (2016) 044019; Phys. Rev. **D96** (2017) 105005] applicable to arbitrary strong and time-dependent gravitational, inertial *and* electromagnetic fields
- Exact Foldy-Wouthuysen transformation constructed
- Effects of terrestrial gravity and rotation non-negligible
- Influence of gravitational wave on spin possibly detectable in the framework of a magnetic resonance type setup
- Probing spacetime geometry: from nuclear spin dynamics obtained new limits on spacetime torsion $T < 10^{-14} \frac{1}{\text{m}}$

Thanks !