



AYSS-2024

The acceptance and efficiency of the Highly Granular Neutron Detector prototype in the BM@N experiment

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27.10.2024

BM@N is the first operational experiment at the NICA complex Nuclotron-based Ion Collider fAcility SPD (Detector) BM@N (Detector) Extracted beam E-cooling MPD Collider Heavy Ion (Detector) Linac Nuclotron LU-20 Booster Baryonic Matter at Nuclotron

Introduction



- The Highly Granular Neutron Detector (HGND) at the BM@N experiment is under development for measuring the energy of neutrons up to 4 GeV produced in nucleus-nucleus collisions.
- Neutron measurements are necessary to obtain robust information on the symmetry energy of the Equation of State for high baryon density matter.
- A compact HGND prototype has already been designed and constructed to validate the concept of the full-scale HGND.
- For the first time, small prototype of the HGND was used in Xe+CsI at 3.8A GeV run at the BM@N.
- This work presents the results of the efficiency and geometric acceptance simulation of the HGND prototype for the detection of forward spectator neutrons from hadronic interactions and electromagnetic dissociation (EMD) of ¹²⁴Xe.



• Design of Highly Granular Neutron Detector prototype

HGND prototype in Xe+CsI@3.8A GeV run

UrQMD-AMC vs DCM-QGSM-SMM models

EMD in RELDIS model

HGND prototype efficiencies and acceptances

HGND prototype design



Scint. layer **Veto** 120x120x25 (MM)

1st (electromagnetic) part:

5 layers: Pb (8mm) + Scint. (25mm)

+ PCB + air

2nd (hadronic) part:

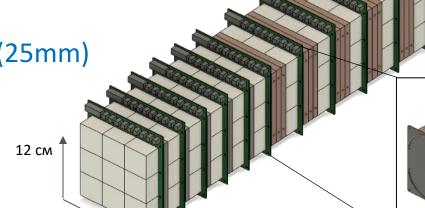
9 layers: Cu (30mm) + Scint. (25mm)

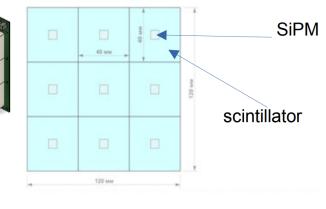
+ PCB + air

Scint. cell – 40 x 40 x 25 mm³ Total number of cells – 135

Total size – 12 x 12 x 82.5 cm³

Total length $\sim 2.5 \lambda_{int}$

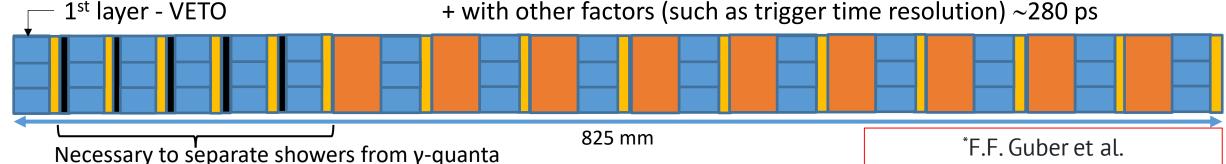




Hamamatsu S13360- 6050PE Photosensitive area – 6x6 mm² Number of pixels – 14400 Pixel size – 50 μm Gain $-1.7x10^6$ PDE – 40%

Time resolution of cell ~200 ps*,

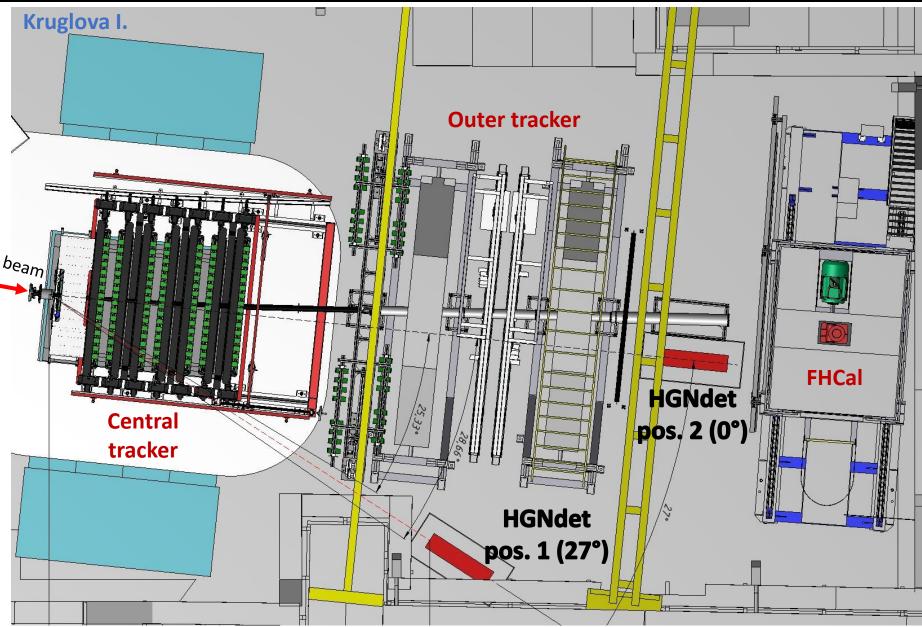
- + with light collection heterogeneity ~240 ps,
- + with other factors (such as trigger time resolution) ~280 ps



27.10.2024 A. Zubankov 10.1134/S0020441223030065

HGND prototype in the Xe+CsI@3.8A GeV run of BM@N





27° position:

Measurements of the neutron spectrum at ~ midrapidity.

0° position:

Test and calibration with known neutron energy (energy of a beam of spectator neutrons)

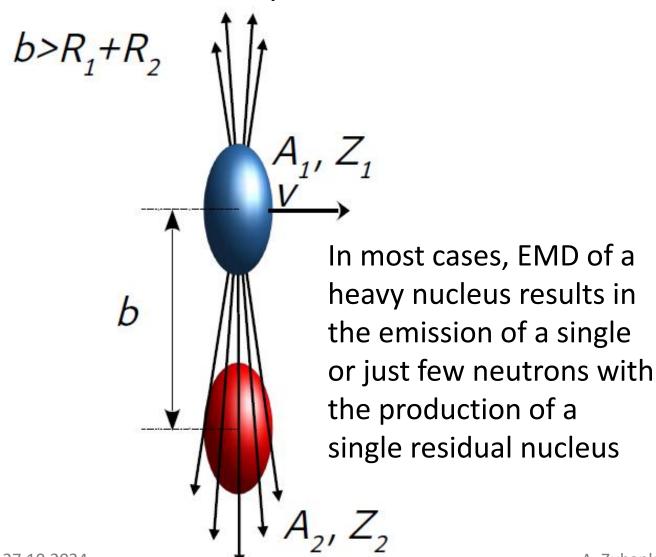


Interactions of nuclei



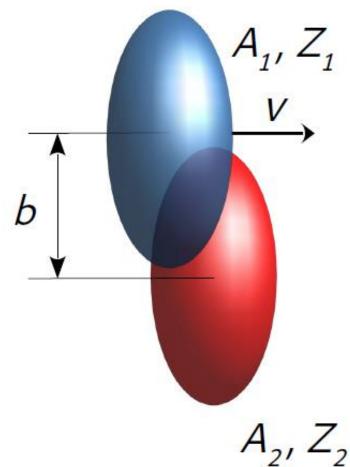
EMD:

without overlap of nuclear densities



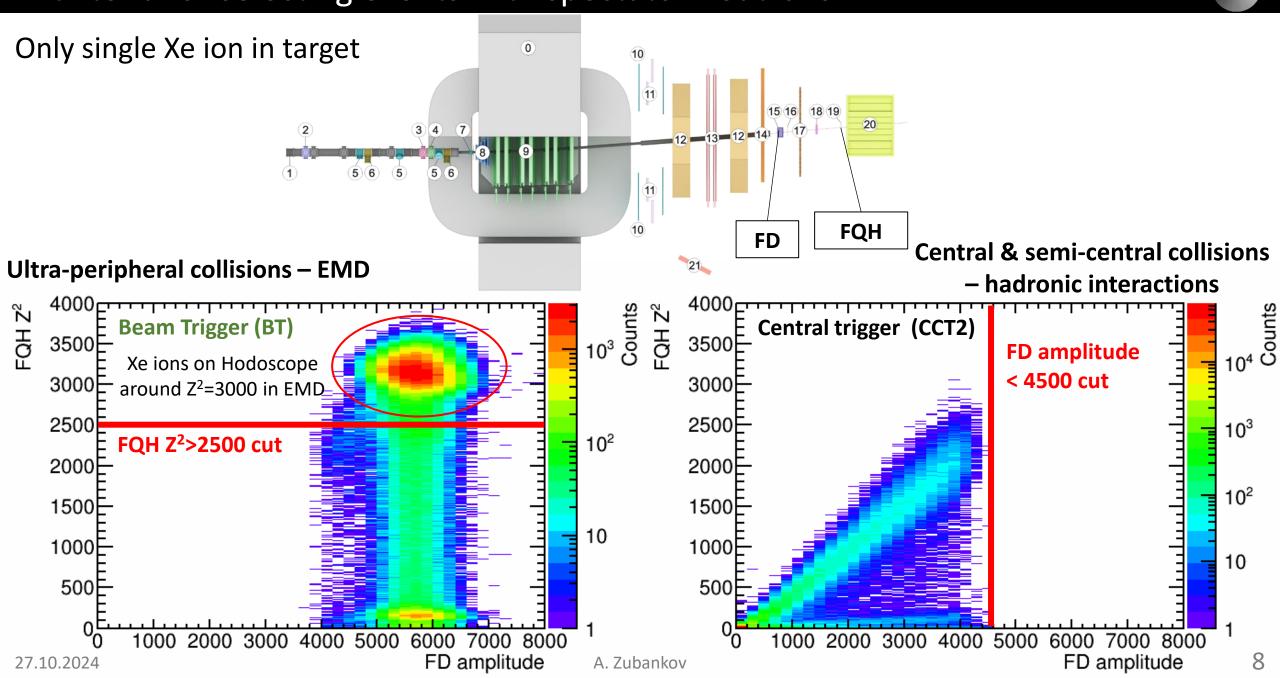
Hadronic interactions: with overlap of nuclear densities

$$b < R_1 + R_2$$



Criteria for selecting events with spectator neutrons

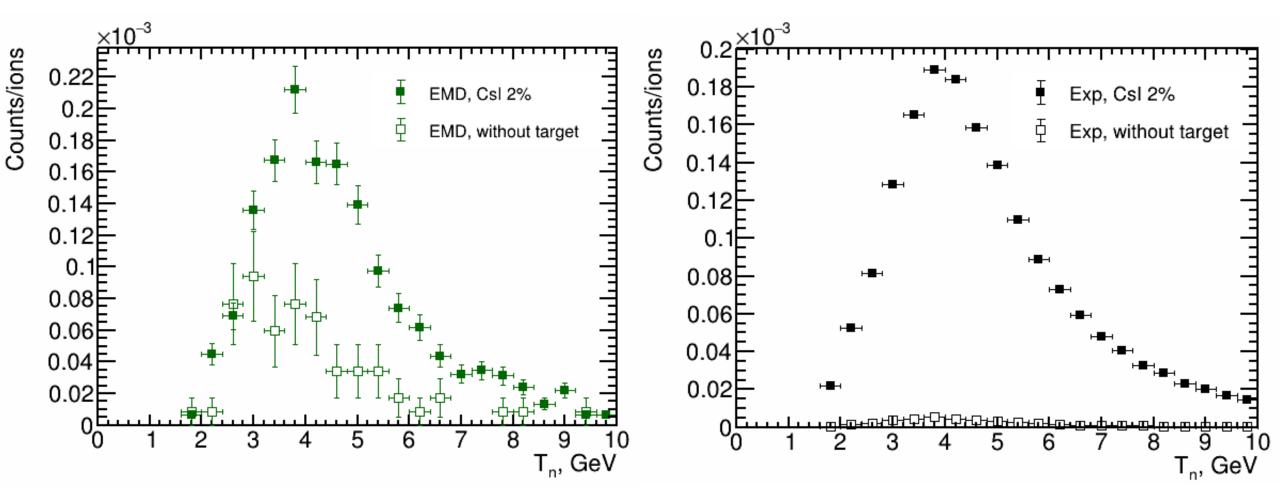




Criteria for selecting events with spectator neutrons

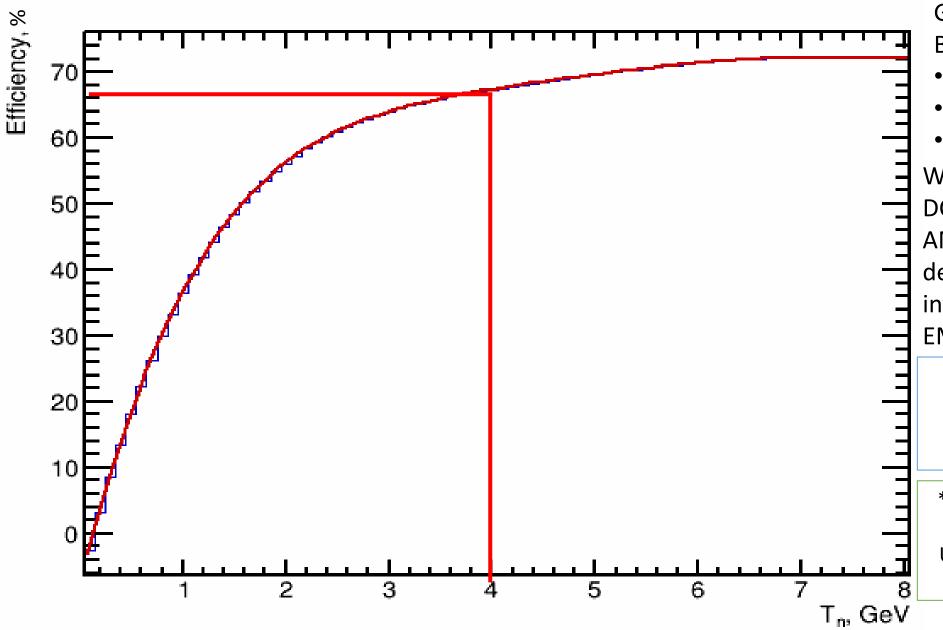


- Selection of events without charged particles, ToF cut, γ -cut (1.55 X_0 or 0.11 λ_{int})
 - Reconstruction of energy by maximum velocity
 - Scaled by incident ion beam rate



HGND prototype efficiency for neutrons





Geant4 simulation:

Box generator, only neutrons

- VETO-cut
- γ -cut
- ToF cut

We use the DCM-QGSM-SMM* and UrQMD-AMC models to estimate the detector efficiency for hadronic interactions, and RELDIS** for EMD

*M. Baznat et al., Monte-Carlo Generator of Heavy Ion Collisions DCM-SMM, *Phys. Part. Nucl. Lett.* **2020**, 17, 303.

I. Pshenichnov, Electromagnetic Excitation and Fragmentation of Ultrarelativistic Nuclei. *Phys. Part. Nucl.* **2011, 42 (2), 215-250.

Ablation Monte Carlo: decay code from AAMCC



- The excited nuclear fragments are formed by means of MST-clusterization algorithm after UrQMD
 - A few excited nuclear prefragments can be formed, in contrast with DCM-QGSM-SMM, where all the spectator nucleon remain bound in one prefragment.
- Excitation energy of prefragment is calculated by hybrid approximation: a combination of Ericson formula for peripheral collisions and ALADIN approximation otherwise¹⁾
- Decays of prefragments are simulated as follows:
 - Fermi break-up model from Geant4 v9.2²⁾
 - Statistical Multifragmentation Model (SMM) from Geant4 v10.4²⁾
 - Weisskopf-Ewing evaporation model from Geant4 v10.4²⁾

- 1) R. Nepeivoda, et al., Particles 5 (2022) 40
- 2) J. Alison et al. Nucl. Inst. A 835 (2016) 186
- 3) 55th Geant4 Techical Forum

https://indico.cern.ch/event/1106118/contributions/4693132/

• They were validated and adjusted to describe the data³⁾.

Combining UrQMD and AAMCC



- AMC is developed to simulate secondary decays of spectator fragments created in other models, in particular UrQMD.
- It is assumed that spectator matter is formed out of nucleons that do not undergo any collisions.

UrQMD:

- Version 3.4
- Cascade mode in this work
- Offset radius 5 fm
- Evolution time 100 fm/c
- Other parameters are set to default values



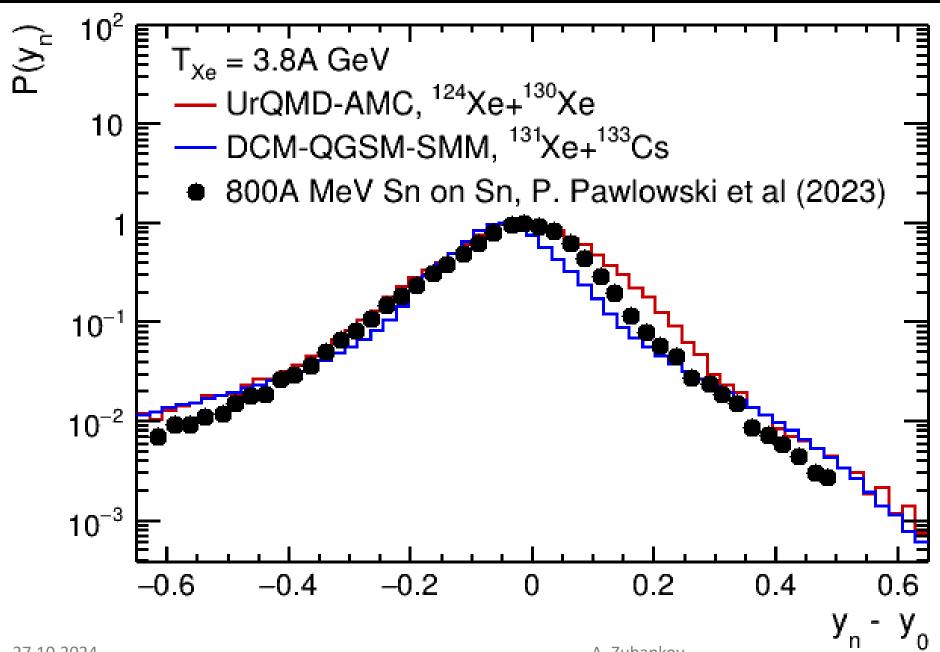
AMC:

- Find spectator nucleons
- Define prefragments via MSTclustering
- Constant d = 2.7 fm
- Model prefragments decays
- All the participant data remain intact



UrQMD-AMC vs DCM-QGSM-SMM





DCM-QGSM-SMM and UrQMD-AMC describe the experiment well in the rapidity region y_n - y_0 <0.

In the region y_n - y_o >0, DCM-QGSM-SMM underestimates the data whereas UrQMD-AMC overestimates.

For DCM-QGSM-SMM, there is a shift in the rapidity relative to the beam rapidity.

UrQMD-AMC vs DCM-QGSM-SMM for BM@N



Counts/ions

10⁻⁵

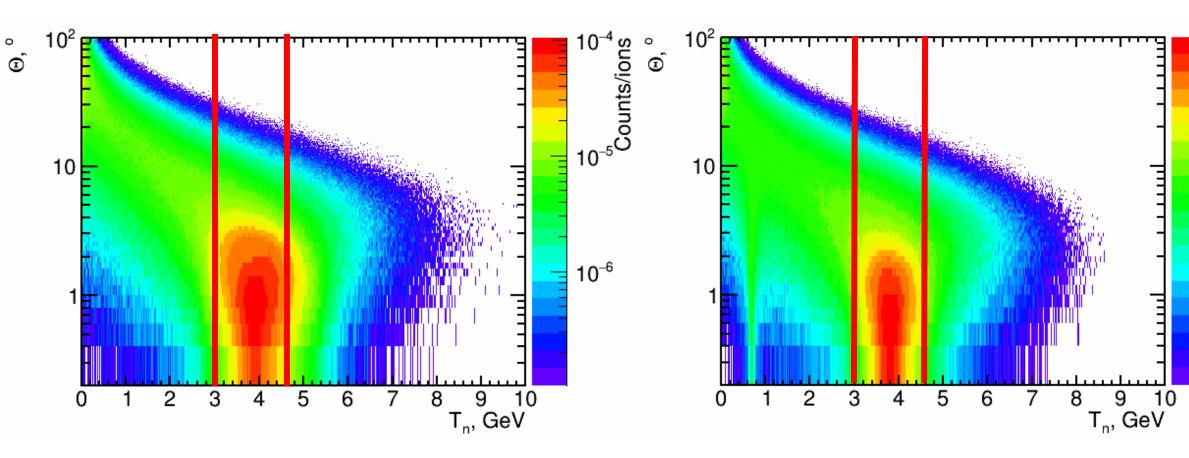
 10^{-6}

UrQMD-AMC

 $3.8A \text{ GeV} ^{124}\text{Xe} + ^{130}\text{Xe}$

DCM-QGSM-SMM

 $3.8A \text{ GeV} ^{131}\text{Xe} + ^{133}\text{Cs}$

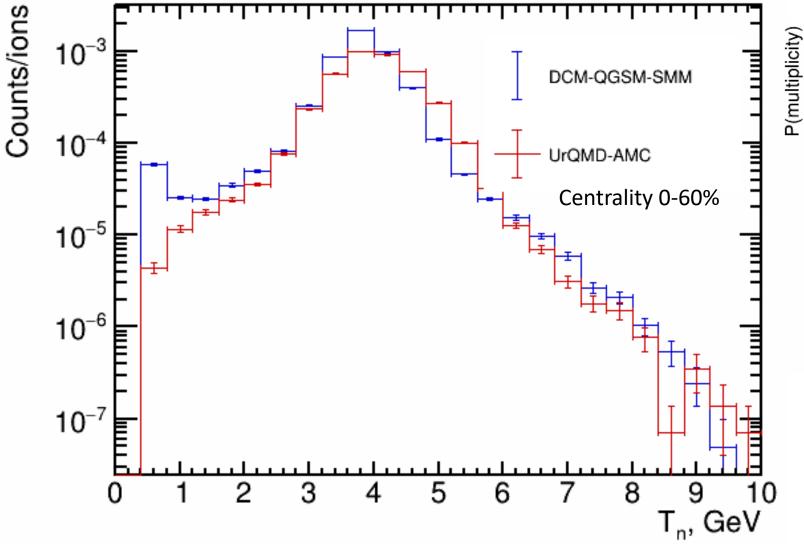


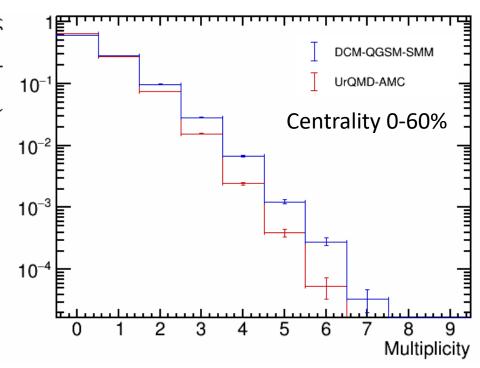
Spectator neutron multiplicity – 17.70

Spectator neutron multiplicity – 16.01

Spectator neutrons on the surface of the HGND prototype





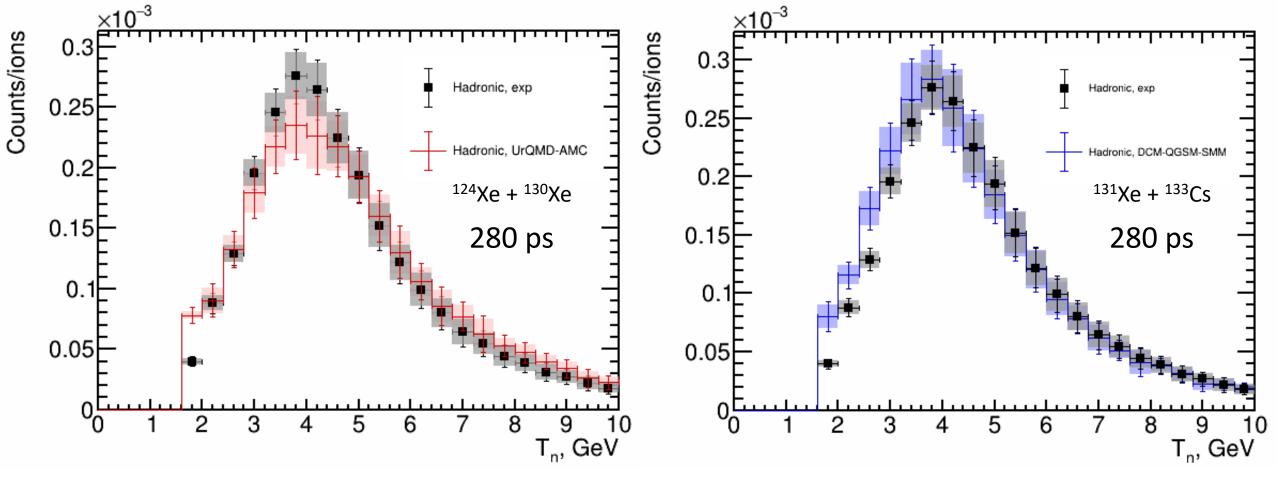


Neutron multiplicity on the surface

- 1.31 for UrQMD-AMC
- 1.44 for DCM-QGSM-SMM

Reconstructed neutron energy spectra for hadronic interactions





The difference in the shape and peak position of the reconstructed spectra of the models is noticeable, which is also due to the difference in the mean kinetic energy of neutrons and their multiplicity.

HGND prototype efficiencies for hadronic interactions



$$lpha cc = rac{N_{hit}}{N_{gen}} \qquad arepsilon = rac{N_{rec}}{N_{hit}}$$

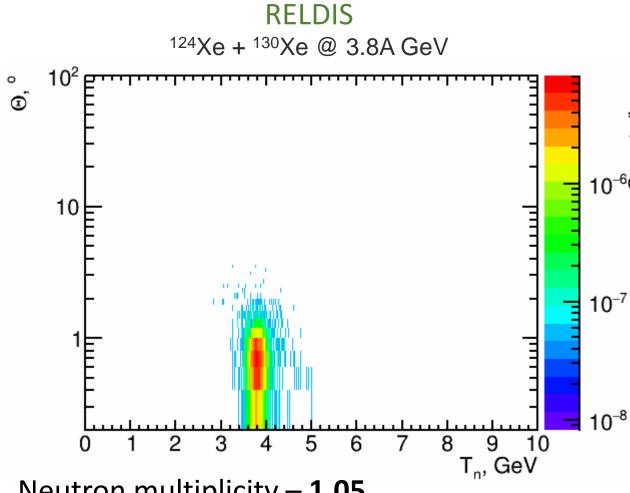
Model	асс, %	ε, %	acc x ε, %
DCM-QGSM-SMM	3.37 ± 0.02	39.41 ± 0.20	1.328 ± 0.007
UrQMD-AMC	2.50 ± 0.02	46.47 ± 0.29	1.162 ± 0.007
Combined	2.94±0.01±0.44 (stat) (sys)	42.94±0.18±3.53 (sys)	1.245±0.005±0.083 (stat) (sys)

The difference in *acc* is explained by the differences in angular distribution of primary neutrons (17.70 vs 16.01) and in average multiplicity of neutrons hitting the detector (1.31 vs 1.44).

The difference in ε is due to the difference in average multiplicity of neutrons hitting the detector (1.31 vs 1.44).

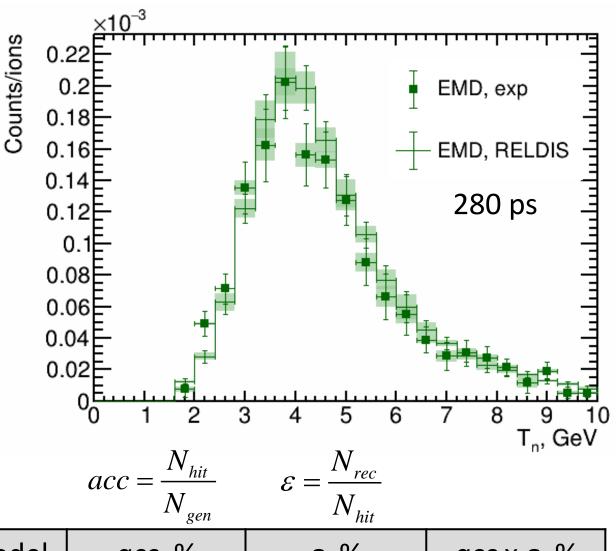
RELDIS – EMD in simulation





Neutron multiplicity – **1.05**

Neutron multiplicity on the surface – **1.02**



Model	асс, %	ε, %	acc x ε, %
RELDIS	34.31 ± 0.25	61.31 ± 0.45	21.04 ± 0.15

Conclusions



- The acceptance and efficiency of the HGND prototype in detecting projectile spectator neutrons from hadronic interactions were studied using UrQMD-AMC and DCM-QGSM-SMM models to generate primary collisions.
- The models were validated with GSI data on neutron production in 800A MeV ¹²⁴Sn + ¹²⁴Sn reaction.
- Some difference in the multiplicity of spectator neutrons and their energy spectra are found. This difference was used to estimate the systematic uncertainties.

- Also, efficiency and acceptance have been investigated for neutrons from EMD using the RELDIS model.
- EMD in the BM@N experiment can be used as a source of high energy neutrons with multiplicity ≈1 per event.
- The estimated acceptance and efficiency of the HGND prototype are:

-
$$acc_{hadr} = 2.94 \pm 0.01_{(stat)} \pm 0.44_{(sys)}\%$$

-
$$\varepsilon_{hadr} = 42.94 \pm 0.18_{(stat)} \pm 3.53_{(sys)}\%$$

-
$$acc_{EMD} = 34.31 \pm 0.25\%$$

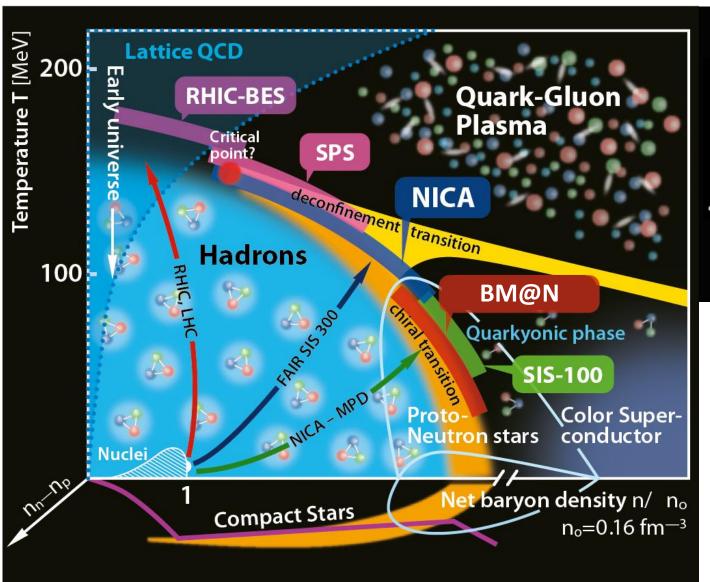
$$- \varepsilon_{FMD} = 61.31 \pm 0.45\%$$

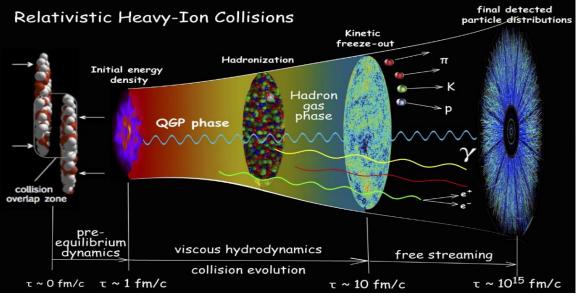


Backup

BM@N: studying the properties of dense baryonic matter



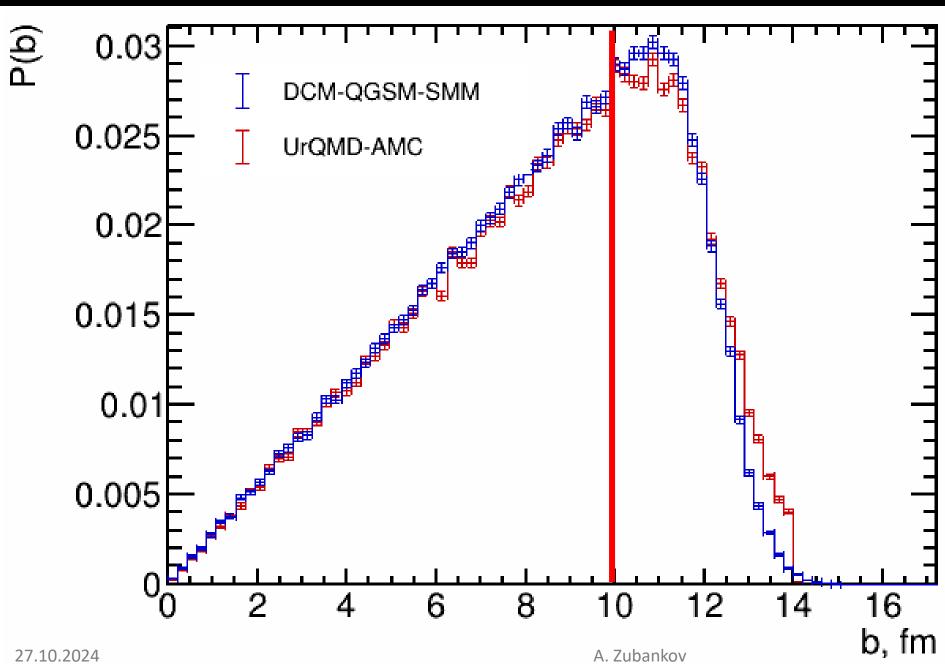




- Study of the QCD diagram at high baryon densities
- Study of the formation of multi-strange hyperons
- Search for hypernuclei in nucleus-nucleus collisions
- Study of the azimuthal asymmetry of charged particle yields in collisions of heavy nuclei.

Impact parameter distribution

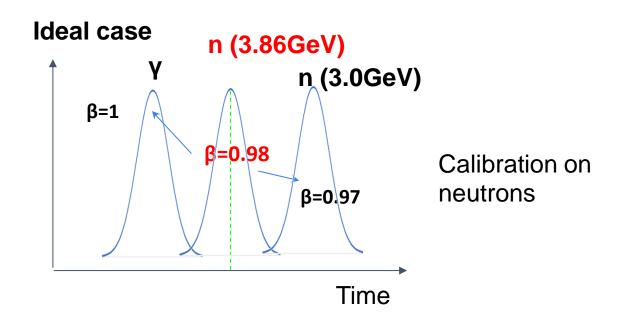


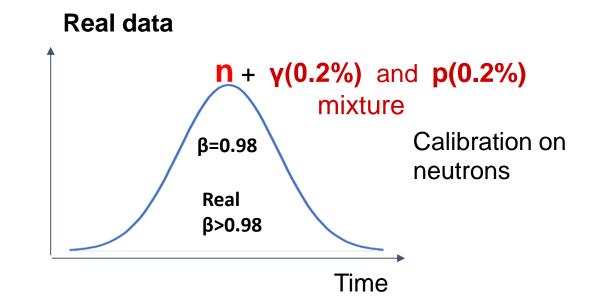


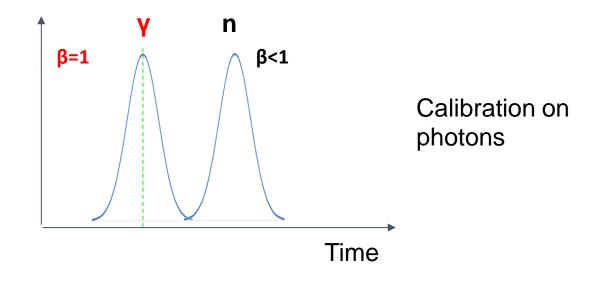
23

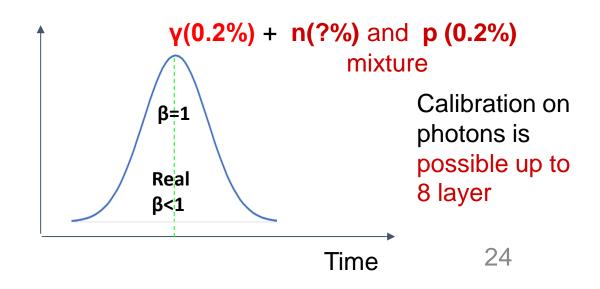
HGND calibration





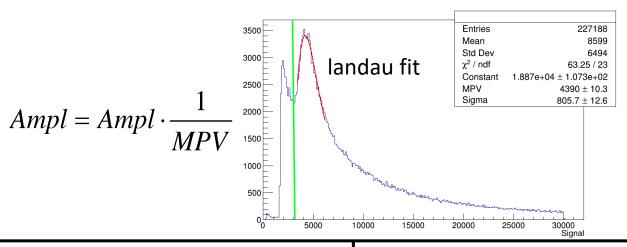




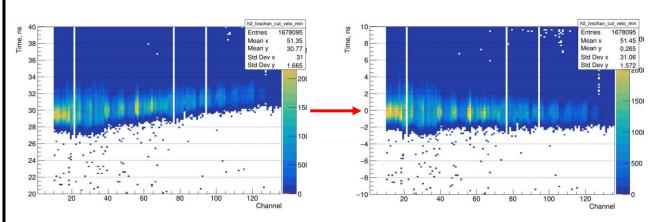


HGND calibration

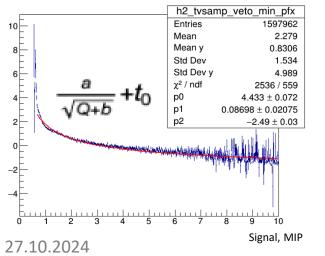
1. Amplitude normalization

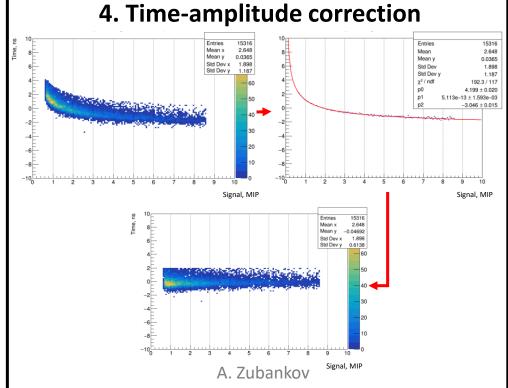


2. Time shift for all channels by the average fit value

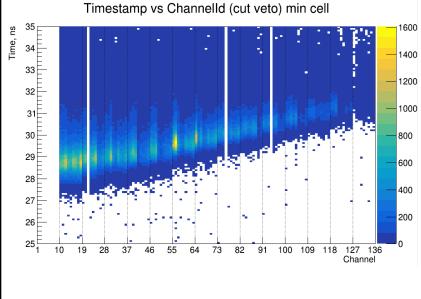


3. Determination of parameters of the approximating function for all channels & time limit





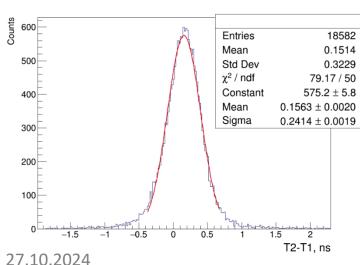
5. Time shift

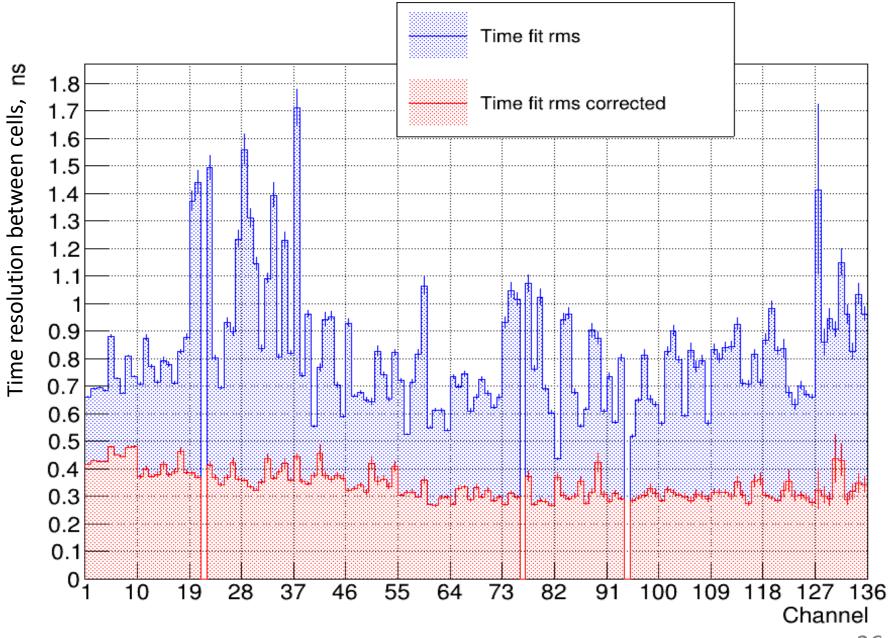


HGND calibration



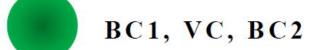
Time-amplitude correction of signals made it possible to get rid of the dependence of time on signal amplitude, which improved the time resolution by ~2.4 times.





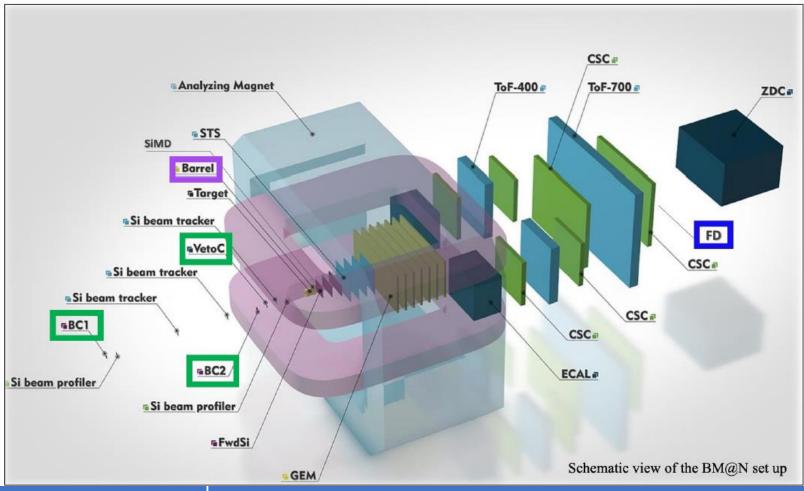
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26



BD

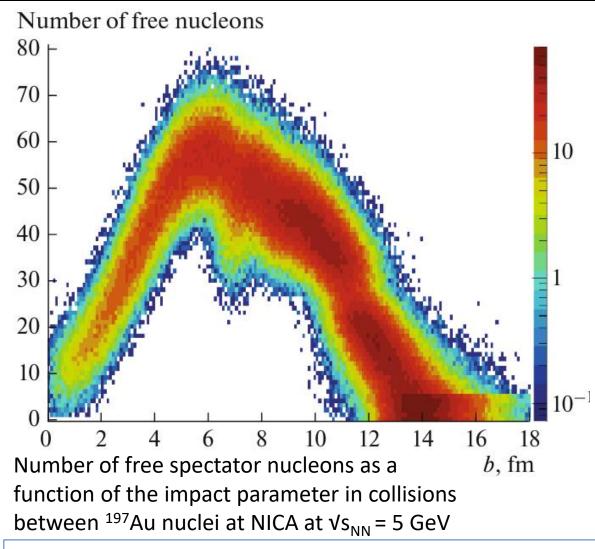
FD



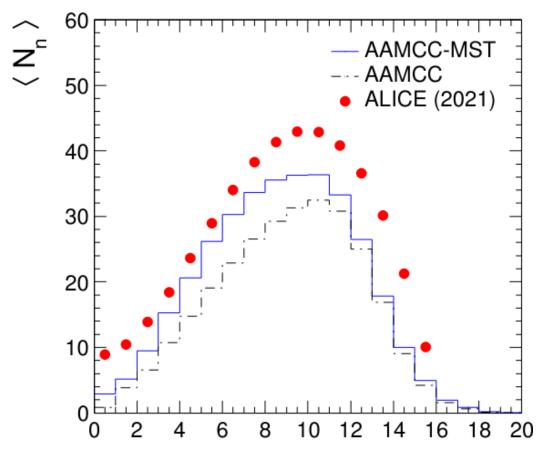
Trigger type	Trigger logic	
Beam Trigger (BT)	BT = BC1 * BC2 * !VC	
Min. Bias Trigger (MBT)	$\mathbf{MBT} = \mathbf{BT} * \mathbf{!FD}$	
Centrality Trigger 1 (CCT1)	CCT1 = BT * BD	
27.10.2024 Centrality Trigger 2 (CCT2)	A. Zubankov $\mathbf{CCT2} = \mathbf{MBT} * \mathbf{BD}$	

Nuclear interaction





A. Svetlichnyi & I. Pshenichnov, Formation of Free and Bound Spectator Nucleons in Hadronic Interactions between Relativistic Nuclei. *Bulletin of the Russian Academy of Sciences: Physics* **2020**, 84 (8), 911–916.



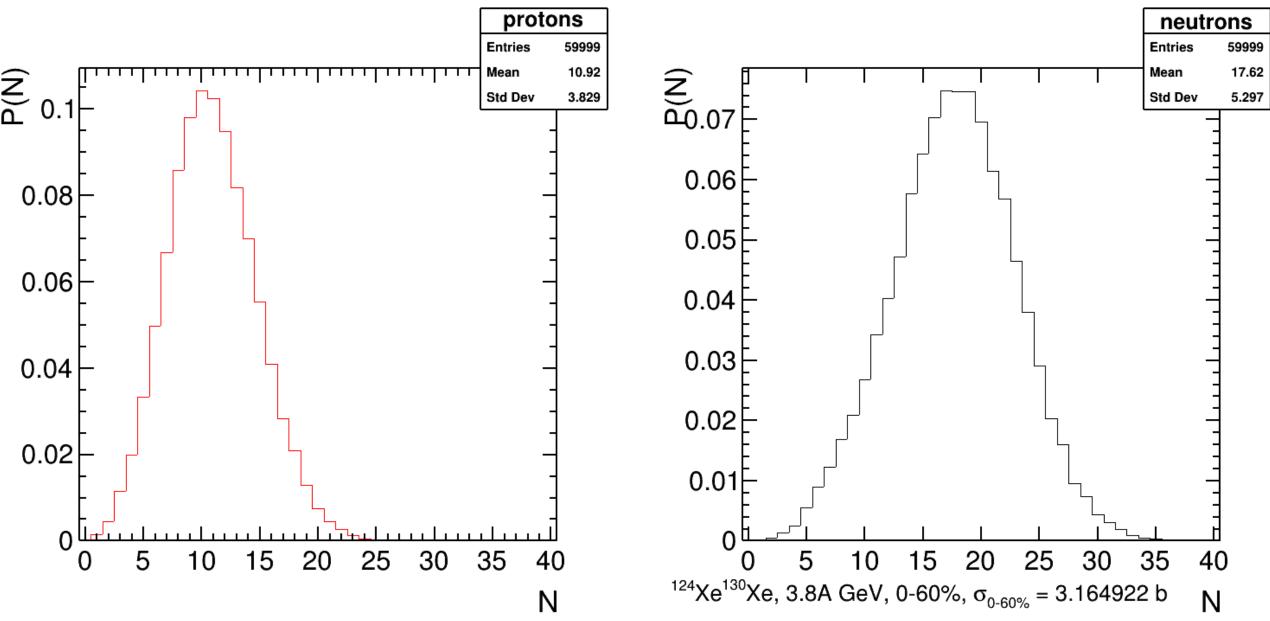
b, fm

Average multiplicities of neutrons in 208 Pb $^{-208}$ Pb collisions at $v_{NN} = 5.02$ TeV as functions of the collision impact parameter

Nepeivoda, R. et al., Pre-Equilibrium Clustering in Production of Spectator Fragments in Collisions of Relativistic Nuclei. *Particles* **2022**, 5, 40–51.

Nuclear interaction

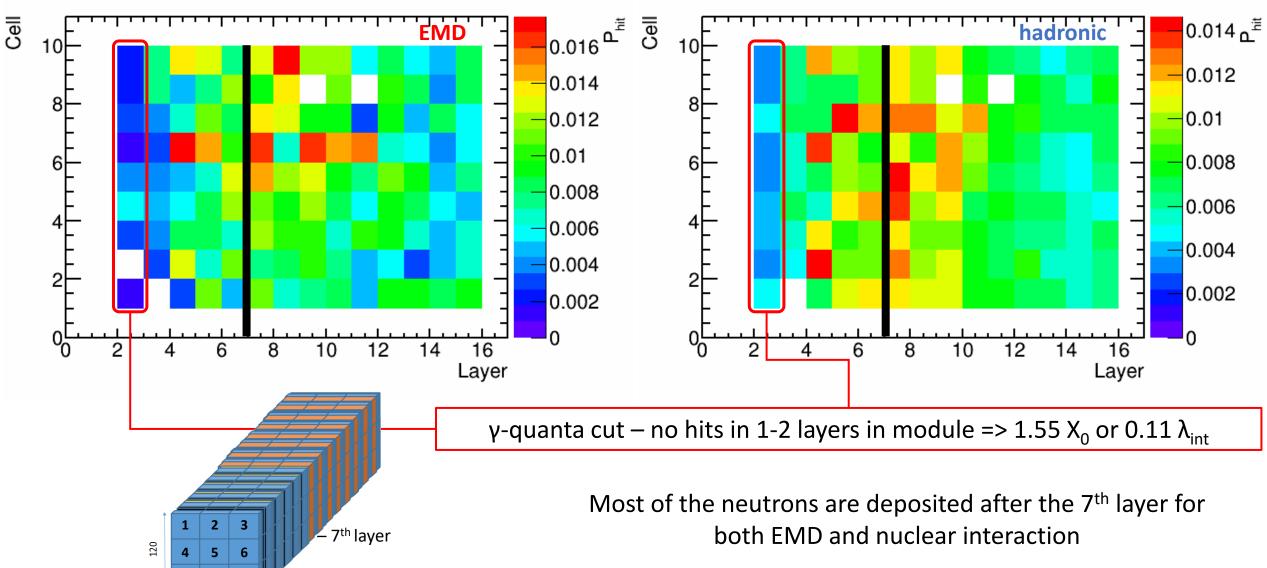




Fastest cells for EMD vs hadronic interactions



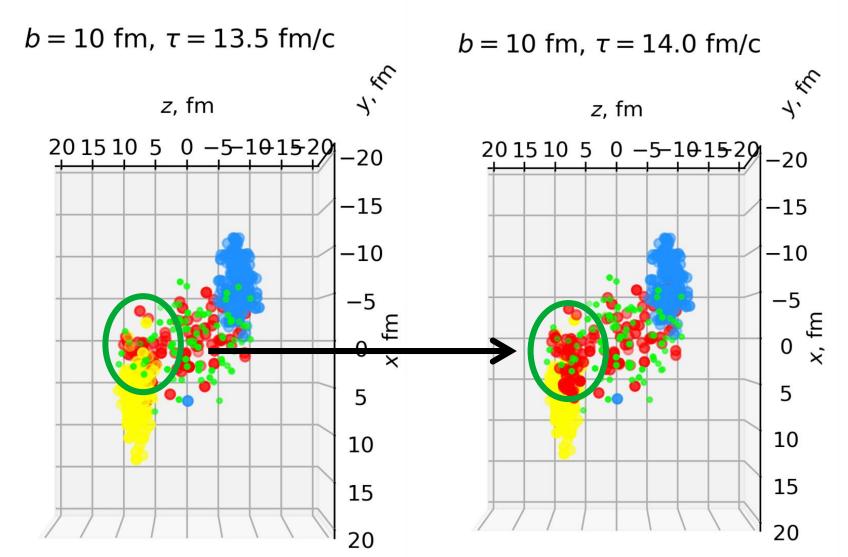
Comparison of hadronic interactions (CCT2) with electromagnetic dissociation (BT) Run 8281 (BT) vs 8300 (CCT2) 3.8 AGeV



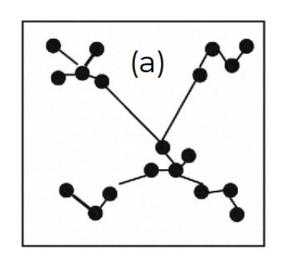
30

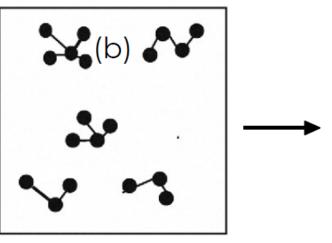
120

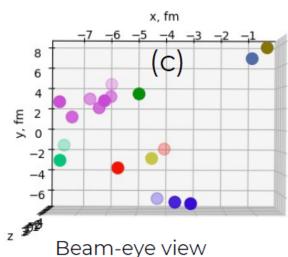
Knocking out some spectator nucleons by mesons



MST-clustering

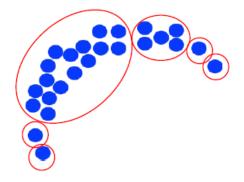






Clusters representation on the Side A

- Graph vertexes nucleons, edges weights Cartesian distances between them.
- (a) The minimum spanning tree is selected from the complete graph
- ullet (b) All edges with a weight greater than d are removed. d is the clustering parameter depending on the excitation energy
- (c) Connectivity components are separate (pre-)fragments

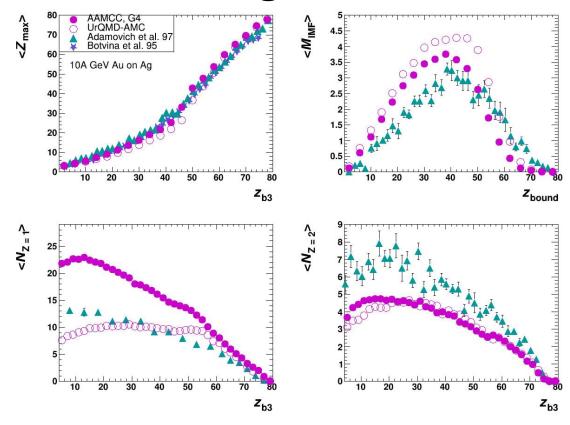


Prefragments in a central collision



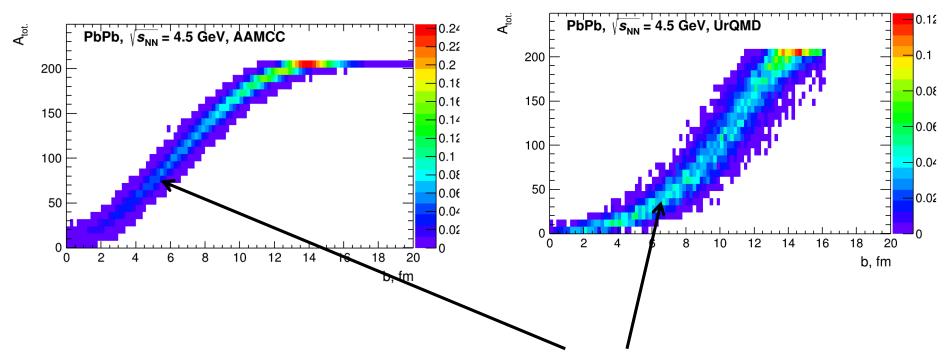
The prefragment is dynamically divided into several prefragments until thermodynamic equilibrium is reached.

¹⁹⁷Au fragmentation



- \cdot UrQMD-AMC and AAMCC describe Z_{max} . Models give similar numbers of He
- •UrQMD-AMC is systematically lower than AAMCC for Z_{bound} < 50. This is due to a smaller spectator volume in UrQMD.
- •AAMCC is closer to data on M_{IMF} , while UrQMD-AMC overestimates M_{IMF} in semi-central collisions. This is because of higher excitation energy of prefragments since more nucleons are removed.
- •The difference in H fragments can be attributed to the different number of participants, because of a larger contribution of protons from MST-clustering

Spectator matter volume as a function of impact parameter



UrQMD gives less spectators than AAMCC for all b

Abrasion-Ablation Monte Carlo for Colliders

- Abbreviated as AAMCC or A²MC²
- Nucleus-nucleus collisions are simulated by means of the Glauber Monte Carlo model
 Non-participated nucleons form spectator matter (prefragment)
- Excitation energy of prefragment can be calculated via three options:
 - Ericson formula based on the particle-hole model²⁾
 - parabolic ALADIN approximation³⁾ adjusted to describe the data for light and heavy nuclei
 - Hybrid approximation: a combination of Ericson formula for peripheral collisions and ALADIN approximation otherwise
- Deexcitation is simulated via MST-clusterisation⁴⁾ accomplished with decay models from Geant4⁵⁾

¹⁾ C. Loizides, J.Kamin, D.d'Enterria Phys. Rev. C 97 (2018) 054910

²⁾ T. Ericson Adv. In Phys. 9 (1960) 737

³⁾ A. Botvina et al. NPA 584

⁴⁾ R. Nepeivoda, et al., Particles 5 (2022) 40

⁵⁾ J. Alison et al. Nucl. Inst. A 835 (2016) 186