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Direct Photon Interferometry in Bi-Bi Collision at

$$\sqrt{s_{NN}} = 9.2 \text{ GeV}$$

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We consider Bose-Einstein correlations of direct photons in central Bi-Bi collisions at $\sqrt{s_{NN}} = 9.2$ GeV. Two-photon correlations are calculated in UrQMD hybrid model. Two equations of state, one with the first-order QCD phase transition and the other with the pure hadron gas, are considered to study the effects of

the formation of quark-gluon plasma (QGP) and mixed phase on the observables. 7 The extracted correlation radii in out-side-long parametrization demonstrate sig-nificant increasing in outward and longitudinal directions for the scenario including QGP, while correlations in side direction are not sensitive to the equation of the state.

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Introduction

At high energy density and temperature quantum chromodynamics (QCD) 11 on the lattice predicts a transition from hadronic matter to the state of de-12 confined quarks and gluons referred as quark-gluon plasma (QGP). The prop-13 erties of such a matter have been extensively studied experimentally in heavy 14 ion collisions at SPS, RHIC, LHC etc. at different regions of net baryon den-15 sity and temperature on the QCD phase diagram [1]. However, the critical 16 point of the QCD phase diagram, at which the first-order phase transition 17 line ends and transition becomes a crossover, has not been observed so far. 18 Therefore, one of the key point in experimental program of NICA [2] is the 19 exploration of the QCD phase diagram at the expected region of the critical 20 point, covering energies $\sqrt{s_{NN}} = 3-11$ GeV. 21

In this paper, we investigate the photon interferometry to study the prop-22 erties of QGP at NICA energies. Photon interferometry is based on the Bose-23 Einstein (BE) correlations, momentum correlations of particles that follow 24 BE statistics. The BE correlations known to be a powerful experimental 25 tool in heavy-ion physics. Hadronic BE correlations have been systemati-26 cally studied over the last decades and provided comprehensive results on 27 the space-time properties of the fireball created in a heavy-ion collision [3, 4]. 28 However, hadronic correlations reflect dimensions of the fireball at the mo-29 ment of its freeze-out, while photons, being color neutral, are mostly emitted 30 from the hottest parts of the fireball, probing the innermost hottest part of 31 the fireball and pre-equilibrium thermalization. 32

Model description

Direct photon BE correlations are calculated using UrQMD [5] hybrid 34 model [6]. The space-time evolution of Bi–Bi collision at $\sqrt{s_{NN}} = 9.2$ GeV is 35 modeled with ideal hydrodynamics [6]. We considered central Bi–Bi collision 36 with the impact parameter b = 0 fm. The fireball in the hydrodynamic model 37 is divided to $100 \times 100 \times 100$ cells with a cell size parameter dx = 0.2 fm. The 38 fluid dynamics is based on the transformation of energy, momentum and net 39 baryon density. The equation of state (EoS) is needed as an additional input 40 to calculate the pressure, temperature, chemical potential and QGP fraction 41 in a cell corresponding to the energy and the baryon number densities [6]. 42 Thus, temperature (T), chemical potential $(\mu_{\rm B})$, 4-velocity vector (u) and 43 fraction of QGP (f_{QGP}) are evaluated for each cell during hydrodynamical 44 evolution in UrQMD hybrid model. The system evolves with a time step 45 dt = 0.08 fm/c until the energy density in all cells decreases below $5\epsilon_0$, 46 where $\epsilon_0 = 146.5 \text{ MeV/fm}^3$ is the nuclear ground state energy density. In 47 UrQMD hybrid model the hydrodynamical evolution is calculated assuming 48 zero viscosity of the system, which equivalent to zero mean free path. 49

Thermal photon emission rate $R(T(x), \mu_{\rm B}(x), p_{\gamma} \cdot u(x))$ is computed for each cell (in a given 4-coordinates x). Only photons emitted at the rapidly y = 0 are considered, correspondingly, these photons in the laboratory system have 4-momentum p_{γ} . The full emission rate is decomposed to contributions from QGP and HG as following,

$$R(T(x), \mu_{\rm B}(x), p_{\gamma} \cdot u(x)) = f_{\rm QGP} \cdot R_{\rm QGP} + (1 - f_{\rm QGP}) \cdot R_{\rm HG}, \qquad (1)$$

where R_{QGP} is estimated according to [7] and R_{HG} is taken as the thermal photon production rate in hadronic reactions described in [8].

The total thermal photon yield is obtained by integrating over space-time volume of the fireball

$$E_{\gamma} \frac{\mathrm{d}^3 N}{\mathrm{d} p_{\gamma}^3} = \int \mathrm{d}^4 x R(T(x), \mu_{\mathrm{B}}(x), p_{\gamma} \cdot u(x)) = \int \mathrm{d}^4 x R(x, p_{\gamma}), \qquad (2)$$

It has been shown (see for example [3]) that the correlation function of two-photon BE correlations (C_2) can be expressed assuming smoothness approximation as the following,

$$C_2(K,q) = 1 + \frac{1}{2} \cdot \frac{\left| \int d^4 x R(x,K) e^{iqx} \right|^2}{\left| \int d^4 x R(x,K) \right|^2},$$
(3)

where $K = (p_1 + p_2)/2$ is the mean pair 4-momentum and $q = p_1 - p_2$ is the relative pair 4-momentum. Factor 1/2 appears because of the photon spin.

In this work, we use out-side-long parametrization of $q = (q^0, q_o, q_s, q_l)$ meaning that outward $q_o = \mathbf{q}_T \cdot \mathbf{K}_T / K_T$ is collinear with transverse mean pair momentum \mathbf{K}_T , longitudinal direction is $q_l = q_z$, side q_s is perpendicular to q_o and q_l , $q^0 = \mathbf{q} \cdot \mathbf{K} / K$. In this parameterization, R_l and R_o are sensitive to the source lifetime, while R_s contains only spacial information.

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To investigate dependence of the photon BE correlation on the QCD 69 phase transition, two EoS are considered. The first option, Bag model EoS 70 [9], models the first-order QCD phase transition between HG and QGP. A 71 system with Bag model EoS consists of QGP, mixed phase and HG, both 72 R_{QGP} and R_{HG} contribute to the total emission rate as it is written in formula 73 (1). A system with Hadron Resonance gas EoS [10], in its turn, consists of 74 pure HG ($f_{\text{QGP}} = 0$) and only R_{HG} contributes to the emission rate. These 75 two EoS are standard incorporations in UrQMD hybrid model [6]. 76

Results

⁷⁸ Correlation functions C_2 versus q_o , q_s and q_l at $K_T = 0.38 \text{ GeV}/c$ are ⁷⁹ presented in Fig. 1 for Bag model EoS (left) and HG EoS (right). Such ⁸⁰ correlation functions are build for a range $0 < K_T < 1.4 \text{ GeV}/c$.



Fig. 1. Photon BE correlation functions dependending on q_o , q_s and q_l at $K_{\rm T} = 0.38 \text{ GeV}/c$ for systems with Bag model EoS (left) and HG EoS (right) The correlation functions in Fig. 1 are parametrized as

 $C_2 = 1 + \lambda \exp(-q^2 R^2), \qquad (4)$

assuming a Gaussian emitting source. From formula (4) a correlation radius 82 R for each direction can be extracted. In a high-multiplicity environment, 83 λ is proportional to the fraction of direct photons over inclusive photons 84 $((N_{\gamma}^{\rm dir}/N_{\gamma}^{\rm inc})^2 \sim 10^{-3})$ which provides possibility to measure direct photon 85 invariant yields [11]. Note, that because of the large lifetime of hadrons 86 decaying into final photons, decay photons do not participate in BE corre-87 lations and result only in reduction of the correlation strength. However, in 88 this paper we focus on calculation of correlation radii and assume $\lambda = 1/2$ 89 (no decay photon contribution) for the simplicity. 90

The obtained correlation radii in out-side-long parametrization for the system with the first-order QCD phase transition (Bag model EoS) and for the pure HG system are shown in Fig. 2.

The characteristic decreasing of correlation radii with $K_{\rm T}$ is caused by transverse flow developed due to expansion of a system. According to Fig. 2, R_s does not change with the EoS as it depends only on spatial size of the fireball and does not convolute time component of the system. On the other

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Fig. 2. Correlation radii in out-side-long parametrization as a function of $K_{\rm T}$ for systems with Bag model EoS (left) and HG EoS (right)

hand, R_o and R_l are larger in the scenario, including QGP. This is consistent with expectations of longer lifetime of the fireball in the case of EoS with first order phase transition, as R_o and R_l convolute spatial and time components of the system.

Summary

The BE correlations of direct photons in central (b = 0 fm) Bi–Bi col-103 lisions at $\sqrt{s_{NN}} = 9.2$ GeV are calculated in hydrodynamic approach using 104 UrQMD. To study the effect of QGP formation, two scenarios, the first-order 105 QCD phase transition EoS (Bag model EoS) and pure HG EoS are consid-106 ered. The obtained correlation radii are decreasing with $K_{\rm T}$ in both scenarios 107 according to the developing of the transverse flow in the expanding system. 108 The side correlation radius R_s is found to be insensitive on the EoS, while 109 R_o and R_l are increased for the case of QGP formation. 110

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