

Structural reorganization in collapsed Langmuir monolayers

Структурная реорганизация ленгмюровских монослоев в области коллапса

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Методом рентгеновской дифракции в геометрии скользящего падения исследованы монослой арахидиновой кислоты на поверхности раствора нитрата церия. Обнаружено, что при поджатии в области коллапса монослой неожиданно сохраняет трансляционный порядок, формируя при этом гофрированную структуру. Предложен новый подход к численному моделированию двумерных дифракционных картин от монослоев с нарушенным ориентационным порядком. С использованием этого подхода описан механизм формирования необычных дифракционных карт, полученных от гофрированных монослоев в ходе синхротронных экспериментов.

Arachidic acid monolayers formed on an aqueous solution of cerium nitrate have been studied by grazing incidence X-ray diffraction. Despite compression beyond the collapse point, the monolayer unexpectedly retains its translational order and develops corrugation. This paper presents a method for simulating the two-dimensional diffraction patterns characteristic of monolayers with perturbed orientational order. The unusual diffraction data obtained from corrugated monolayers during synchrotron experiments are described using this approach.

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3 Introduction

4 Modern surface sensitive X-ray techniques have been recognized to of-
5 fer the immense opportunities for investigations of structural reorganization
6 in 2D nanosystems at liquid interfaces. Deeper understanding of molecular
7 mechanisms governing these complicated processes has significant implica-
8 tions both for life sciences and for nanotechnologies (ranging from industrial
9 applications to biomedicine). Nowadays, X-ray reflectometry, grazing inci-
10 dence diffraction, grazing incidence small-angle scattering, the X-ray waves
11 standing method, etc. are extensively used for characterization of organic
12 and bioorganic monolayers formed on liquid subphase. But the intricacy
13 of structural processes in such exceptionally challenging objects as molecular
14 layers on liquid surface requires the development of new adequate approaches
15 to analysis and interpretation of experimental data.

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16 Grazing incidence diffraction (GID) [1, 2] is one of the most informative
 17 methods for structural studies of organic monolayers. In this work we present
 18 the new theoretical formalism for quantitative analysis of experimental 2D
 19 diffraction maps from molecular monolayers with perturbed orientational or-
 20 der. The possibilities of the developed approach were demonstrated by inter-
 21 pretation of the GID data obtained for collapsed fatty acid monolayer formed
 22 on the water subphase in the presence of Ce(III) ions [3]. Cerium induced
 23 effects on the structural reorganization in collapsed phases have been eluci-
 24 dated. Unexpected trends in collapse behavior of fatty acid monolayer are
 25 discussed.

26 Results: experimental data and numeric simulations

27 The arachidic acid monolayers formed on a solution of cerium nitrate salt
 28 $\text{Ce}(\text{NO}_3)_3$ were studied. The monolayer has been compressed beyond the
 29 collapse point at surface pressure $\pi = 52$ mN/m, after which the compression
 30 was stopped. The GID measurements were carried out at the «Langmuir»
 31 station of the Kurchatov synchrotron radiation source. The photon energy
 32 of the incident beam was 13 keV. The GID measurements were performed at
 33 $\theta = 0.8\theta_c$ (θ_c is the critical angle of total external reflection for water) at a
 34 temperature of $T = 21^\circ\text{C}$. The experimental results are given in Fig. 1(a).

35 The unusual features of the 2D diffraction maps were observed, in par-
 36 ticular the rounding of the Bragg rods, which lie precisely on the circle in
 37 reciprocal space. This is unexpected for a two-dimensional polycrystalline
 38 structure. In a two-dimensional monolayer, the diffraction map typically
 39 exhibits a triplet of vertical Bragg rods: (02), (11) and (1 $\bar{1}$) (rectangular
 40 centered cell assumed). This features of diffraction maps are clearly seen
 41 in Fig. 1(a). The rounding of the Bragg rods was observed to lie exactly
 42 on a circle in reciprocal space with a center at the horizon $Q_z = 0 \text{ \AA}^{-1}$.
 43 This suggests the presence of an specific type of structural ordering in such
 44 monolayers.

45 To describe the obtained 2D maps, a theoretical diffraction model for dis-
 46 torted Langmuir monolayers was proposed. A stationary wave equation was
 47 used to describe the scattering processes. In the case of a plane continuous
 48 medium the solution $E_{i,f}(\mathbf{r}) = [R_{i,f}(z)e^{ik_z z} + T_{i,f}(z)e^{-ik_z z}] e^{-i\mathbf{k}_{\parallel}\cdot\mathbf{r}}$ describes
 49 the electric field of the incident and scattered waves in the scalar approxi-
 50 mation, where the indices i and f denote the incident and scattered waves,
 51 respectively; amplitudes $R_{i,f}(z)$ and $T_{i,f}(z)$ are reflection and transmission
 52 coefficients. The solution was derived from the Helmholtz equation [4]. How-
 53 ever, this was not enough to describe the scattering pattern, since the approx-
 54 imation of plane continuous medium does not describe the lateral structure
 55 of the monolayer.

56 The Distorted wave Born approximation (DWBA) uses the solution $E_{i,f}(\mathbf{r})$
 57 for a planar structure, and the monolayer itself is perceived as a disturbance
 58 V_δ . As a disturbance V_δ we consider the addition $V_\delta(\mathbf{r})$ to the potential $V(z)$,

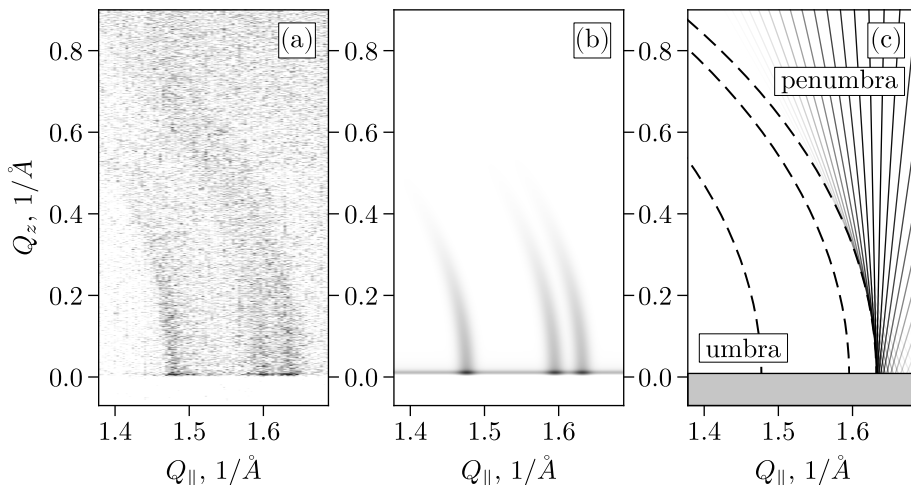


Fig. 1. (a) The diffraction map obtained from arachidic acid monolayer formed on cerium nitrate solution. In this context, Q_{\parallel} and Q_z represent the horizontal and vertical components of the scattering vector, respectively. (b) Corresponding numerical simulation of diffraction map. (c) A simulated schematic diffraction map, that demonstrate the mechanism of the formation of Bragg rods roundings. The picture consists of "umbra" and "penumbra", which are separated by rounded diffraction rods.

59 so that the resulting potential $V(\mathbf{r}) = V_{\delta}(\mathbf{r}) + V(z)$. The reciprocity principle
 60 is used to calculate the scattering amplitude f [5]:

$$f \propto \int_{\mathbb{R}^3} E_f V_{\delta} E_i d^3 r. \quad (1)$$

61 The diffraction intensity distribution is calculated by means of the well-
 62 known formula for the differential scattering cross section $d\sigma/d\Omega \propto ff^*$.
 63 The diffraction pattern from a 2D crystallite consists of rods in reciprocal
 64 space. These rods are perpendicular to the plane of the structure, which
 65 are observed as 2D diffraction pattern on the detector. Substitutions yield
 66 an expression for the differential scattering cross section, which consists of
 67 sixteen terms (scattering channels). In calculations, fifteen terms, containing
 68 the amplitudes $R_{i,f}$ from the solution $E_{i,f}(\mathbf{r})$, are frequently neglected. This
 69 results in the expression taking a simplified form only with coefficient $|T_i T_f|^2$:

$$\frac{d\sigma}{d\Omega} \propto |T_i T_f|^2 \langle \hat{V}_{\delta}(\mathbf{Q}) \hat{V}_{\delta}^*(\mathbf{Q}) \rangle, \quad (2)$$

70 where T_i and T_f are the transmission coefficients for the incident and scat-
 71 tered waves, respectively. The dynamic scattering effects of an evanescent
 72 wave were taken into account through the amplitude coefficients, and the
 73 multiplier $|\hat{V}_{\delta}|^2$ is responsible for diffraction. In this study, we considered the
 74 Langmuir monolayer within the two-dimensional polycrystal approximation.
 75 Subsequently, this multiplier is the modulus-square of the scattering poten-
 76 tial, which can be calculated in accordance with the tenets of the kinematic

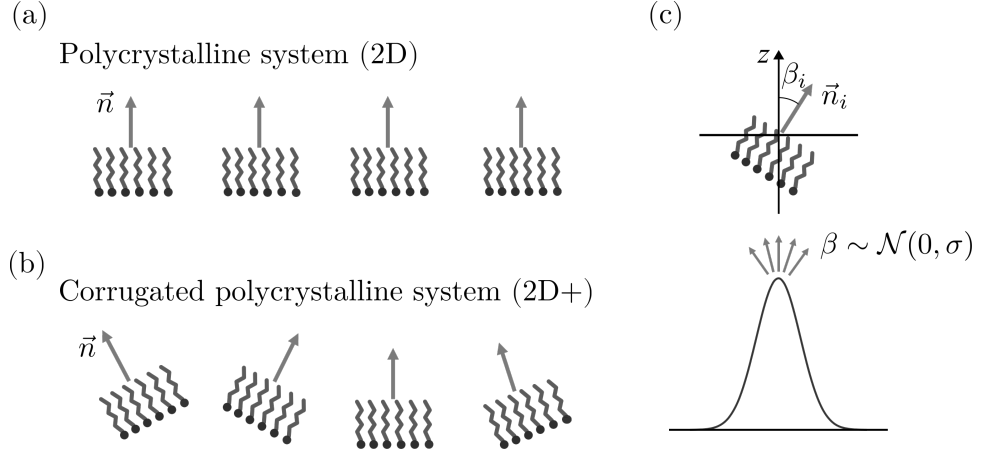


Fig. 2. Schematic representation of a two-dimensional monolayer (a) evolving into a corrugated film with crystallite deviation relative to the horizon (b) distributed as a normal stochastic variable (c).

77 theory of diffraction:

$$\frac{d\sigma}{d\Omega} \propto |T_i T_f|^2 \langle |S(\mathbf{Q})|^2 |F(\mathbf{Q})|^2 J(\mathbf{Q}) \rangle, \quad (3)$$

78 where $S(\mathbf{Q})$ is a structural factor of the plane ordering and $F(\mathbf{Q})$ is a form
 79 factor of a molecule. $J(\mathbf{Q})$ represents the correlation function that accounts
 80 for the angular deviations of the crystallites, while also encompassing the
 81 effects of interference from coherently scattering molecules within each ordered
 82 crystallite. It is important to note that each ordered crystallite has a
 83 finite correlation length, meaning that the summation in $J(\mathbf{Q})$ occurs for all
 84 molecules within one coherently scattering domain. In such a case, the reduc-
 85 tion in peak width can be explained by the increase in the size of coherently
 86 scattering domains.

87 It should be noted that if the system were three-dimensional, classical
 88 Debye-Scherrer rings would be observed. However, the observed sections of
 89 the rings are incomplete, indicating that the system is not of the three di-
 90 mensional nature, but represent some type of distorted 2D lattice. A natural
 91 explanation is that the pieces of the plane lattice (2D crystallites) are inclined
 92 in the direction of the monolayer normal. This leads to a tilt of the crystal-
 93 lites relative to the horizon, as illustrated in Fig. 2(a, b). Furthermore, we
 94 assumed that the tilt of the plane crystallites follows a normal distribution
 95 [Fig. 2(c)]. In such a way, the tilt of the crystallites was incorporated into
 96 the model through the application of the probability density function of the
 97 normal distribution, resulting in the formation of rounded rods through nu-
 98 merical integration. Using the aforementioned corrugated monolayer model,
 99 we conducted numerical modeling [Fig. 1(b)] and reconstruct the parameters
 100 associated with the distortion of the plane order in a polycrystalline mono-
 101 layer. The numerical modeling can be implemented using a Python software.

102 Fig. 1(c) demonstrates the mechanism of the formation of Bragg rods
 103 curvature. A simulated schematic diffraction map is shown in the Fig. 1(c).

104 For clarity, the simulation was performed only for the rightmost curved Bragg
105 rod. The three circles on which the rounded Bragg rods lie are marked
106 with dotted lines. The scattering picture consists of a "umbra" (shadow)
107 and "penumbra" (partial shade), which are separated by a set of diffraction
108 rods. The boundary between umbra and penumbra runs exactly along the
109 diffraction ring (dotted line), and the sum from the set of diffraction rods
110 gives the maximum intensity along the ring. Bragg rods are represented by
111 solid lines (the colors of which are represented in shades of gray in accordance
112 with the normal distribution). The increased intersection of diffraction lines
113 near the ring, as opposed to the region further away, results in the formation
114 of a caustic, which manifests as a rounded diffraction rod.

115 In conclusion, the theoretical model for the analyzing of diffraction from
116 two-dimensional molecular ensembles, taking into account their corrugated
117 structure, has been proposed. This approach allows the calculation of 2D
118 diffraction maps obtained by the grazing incidence X-ray diffraction method.
119 This analysis provides the possibility to obtain quantitative information about
120 the structure and organizational processes in Langmuir monolayers with un-
121 usual types of ordering, in particular in corrugated monolayers.

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