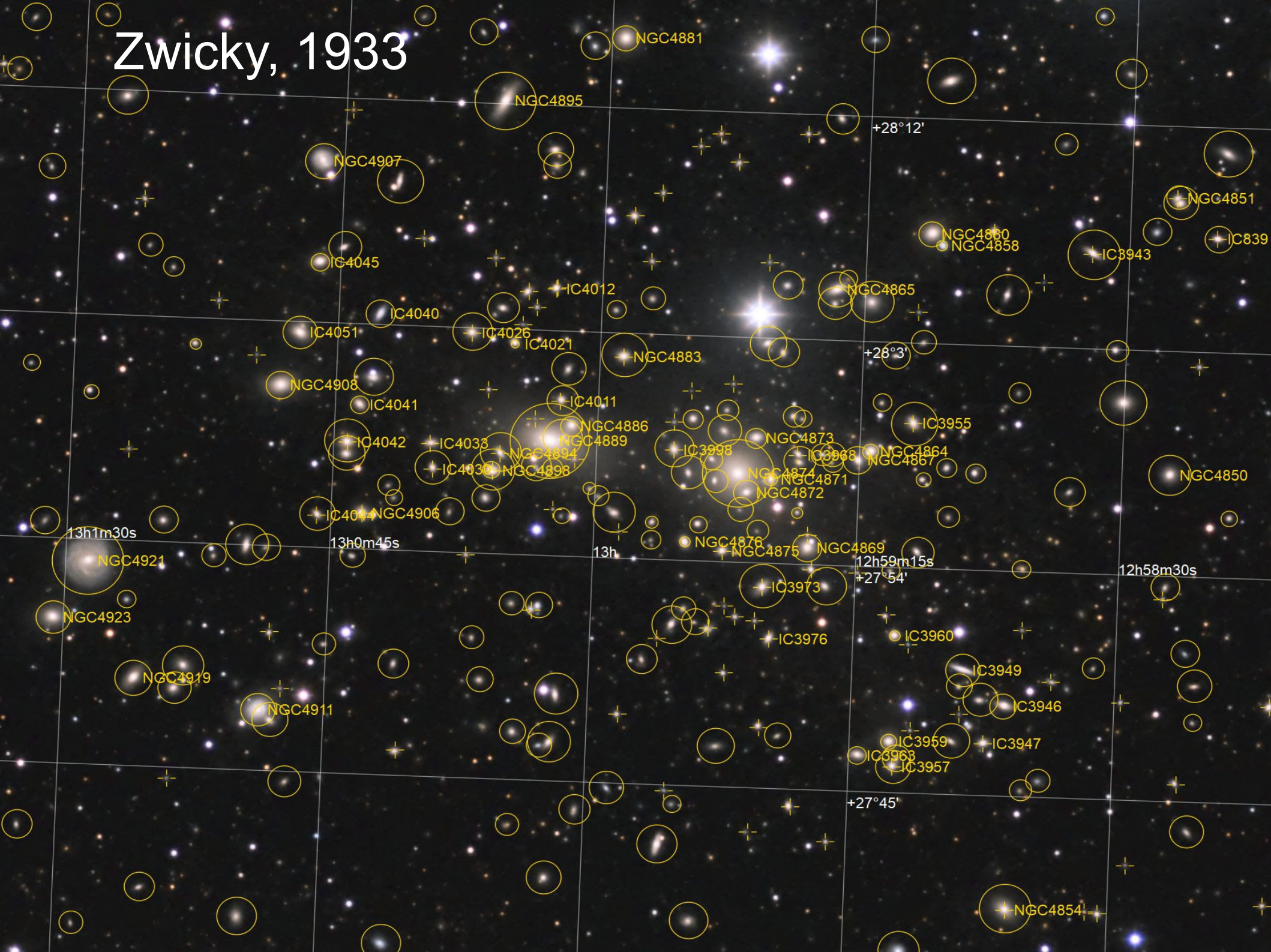


***The dark matter problem:
the past and the future***

*Anton N. Baushev,
Bogoliubov Laboratory of Theoretical Physics,
Joint Institute for Nuclear Research*

1) Why dark matter?

Zwicky, 1933



NGC4881

NGC4895

+28°12'

NGC4907

NGC4851

IC839

NGC4860
NGC4858

IC3943

IC4045

IC4012

NGC4865

IC4040

+28°3'

IC4051

IC4026
IC4021

NGC4883

NGC4908

IC4011

IC3955

IC4041

NGC4886
NGC4889

NGC4873

IC4042

IC4038
IC4030

NGC4894
NGC4898

IC3998

IC3968

NGC4864
NGC4867

NGC4850

IC4014
NGC4906

IC4004

NGC4876
NGC4875

NGC4869

12h59m15s
+27°54'

13h1m30s

13h0m45s

13h

12h58m30s

NGC4921

IC3973

NGC4923

IC3976

IC3960

NGC4919

IC3949

NGC4911

IC3946

IC3959
IC3963
IC3957

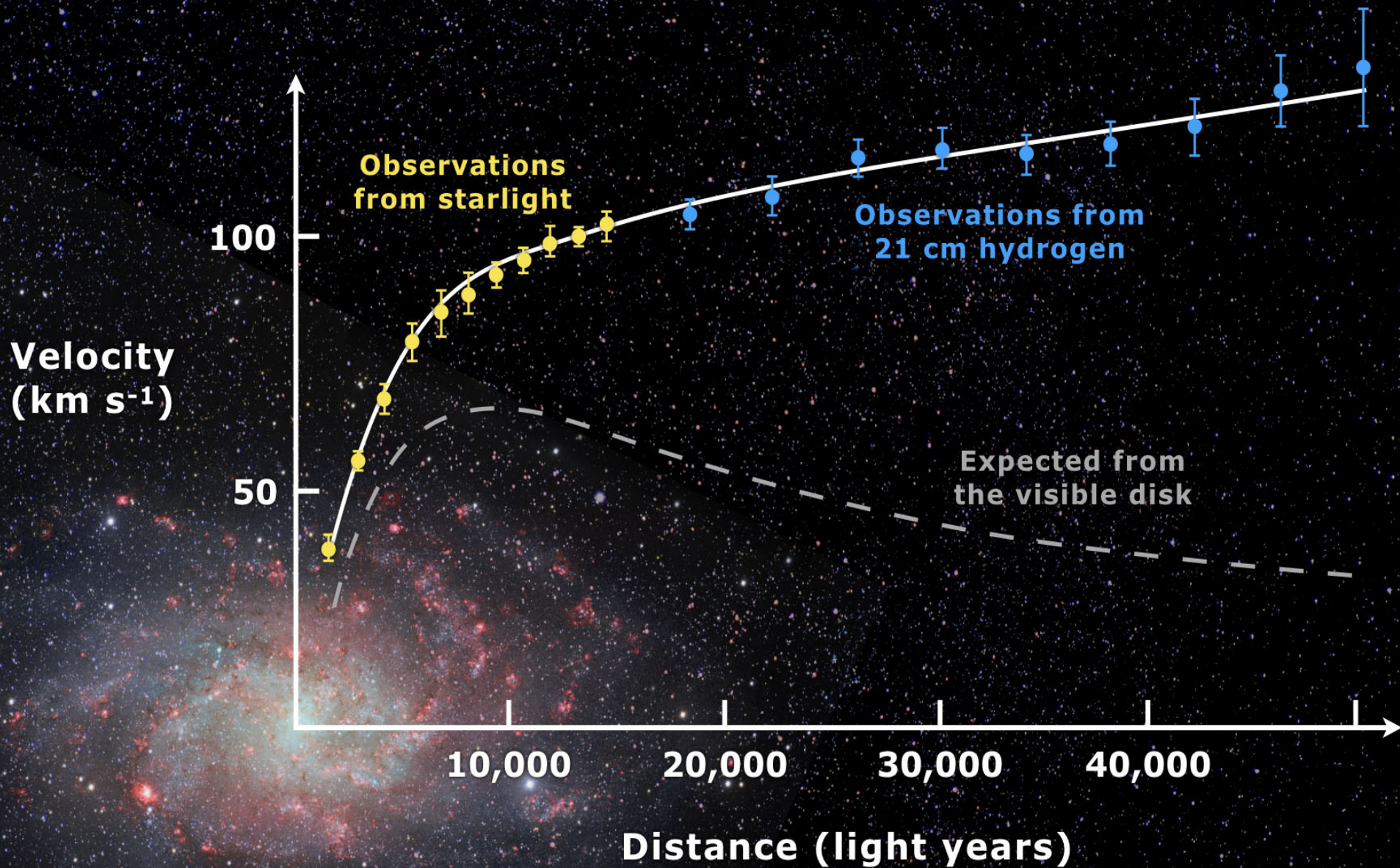
IC3947

+27°45'

NGC4854

Rotation curves of galaxies, 1977

Dark matter or modified gravity?



If a galaxy has a spherically symmetric density distribution $\rho(R)$, its mass depends on radius as $M(R) = \int_0^R \rho(r) \cdot 4\pi r^2 dr$. Speed V of the Keplerian rotation of a body on the circular orbit of radius R is given by Newton's law:

$$\frac{V^2}{R} = G \frac{M(R)}{R^2} \implies V = \sqrt{\frac{GM(R)}{R}}$$

Outside the galaxy, M stops growing with radius, and we get:

$$M = \text{const} \implies V = \sqrt{\frac{GM}{R}} \propto R^{-1/2}.$$

The flat tail $V = \text{const}$ corresponds to $M(R) \propto R$ and $\rho \propto R^{-2}$.

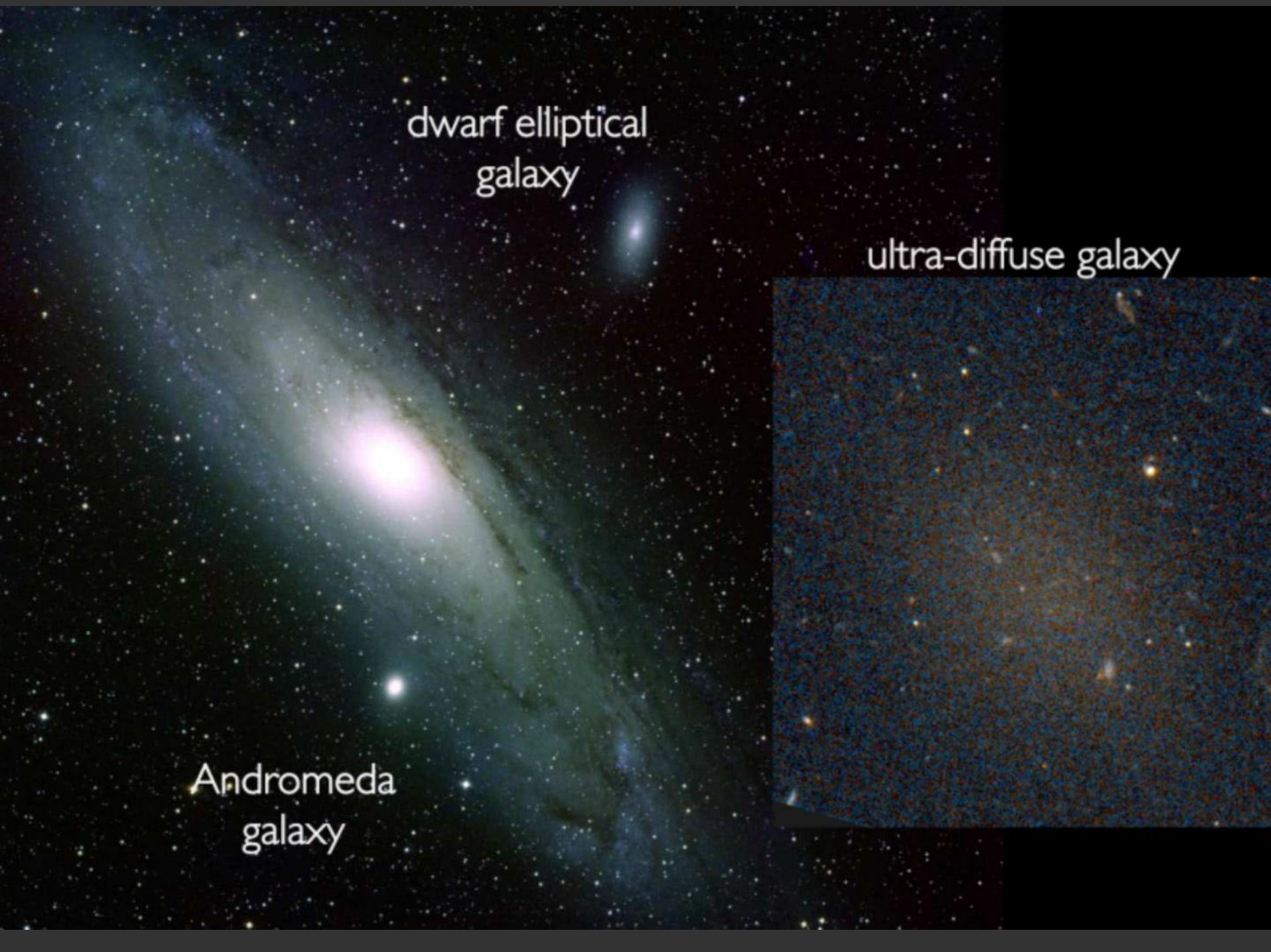
Interacting galaxies

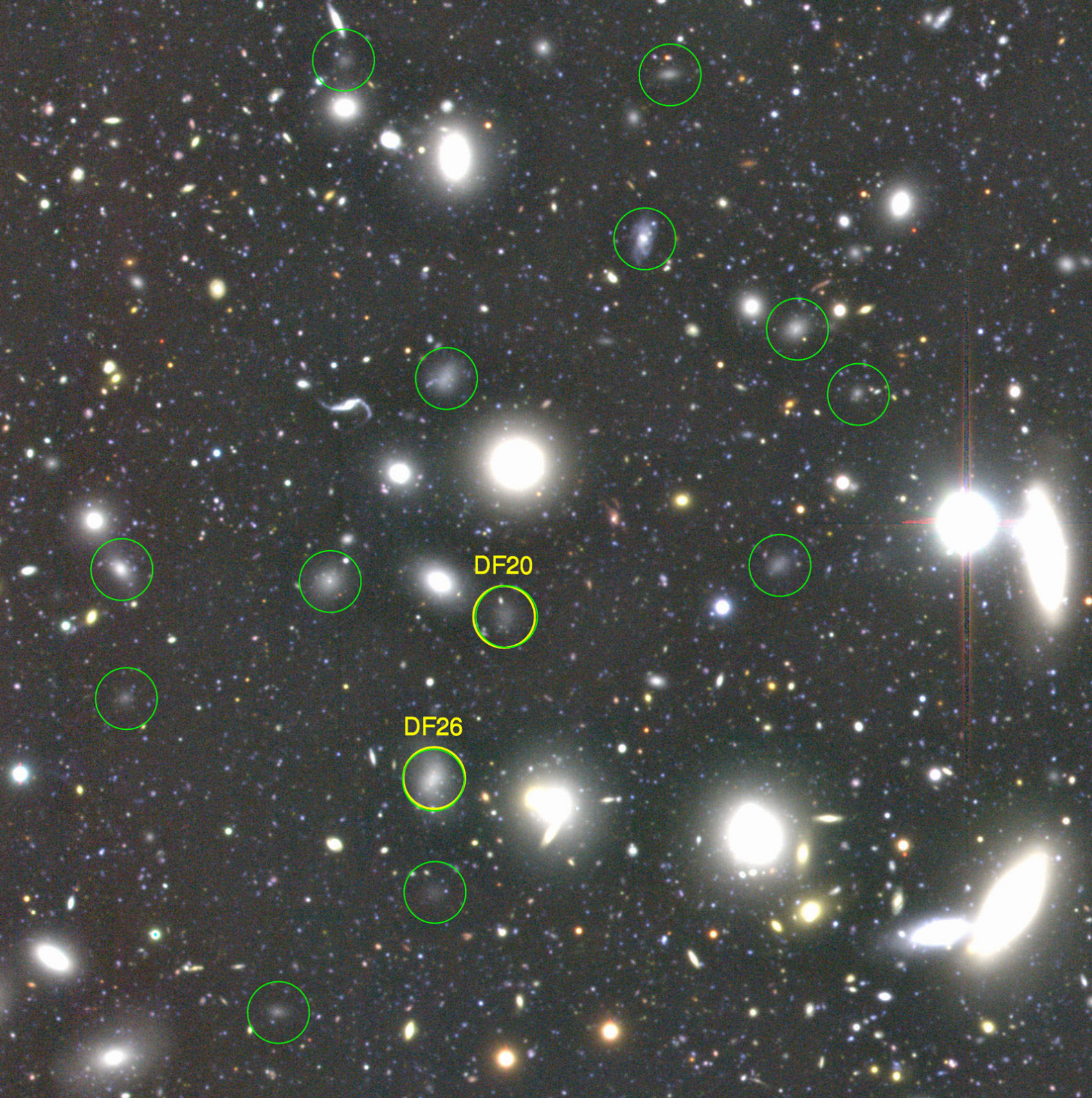


dwarf elliptical
galaxy

ultra-diffuse galaxy

Andromeda
galaxy





Ultra diffuse galaxies in the Coma cluster

Segue 1, a satellite of the Milky Way

$$M \sim 10^6 M_{\text{sun}}$$

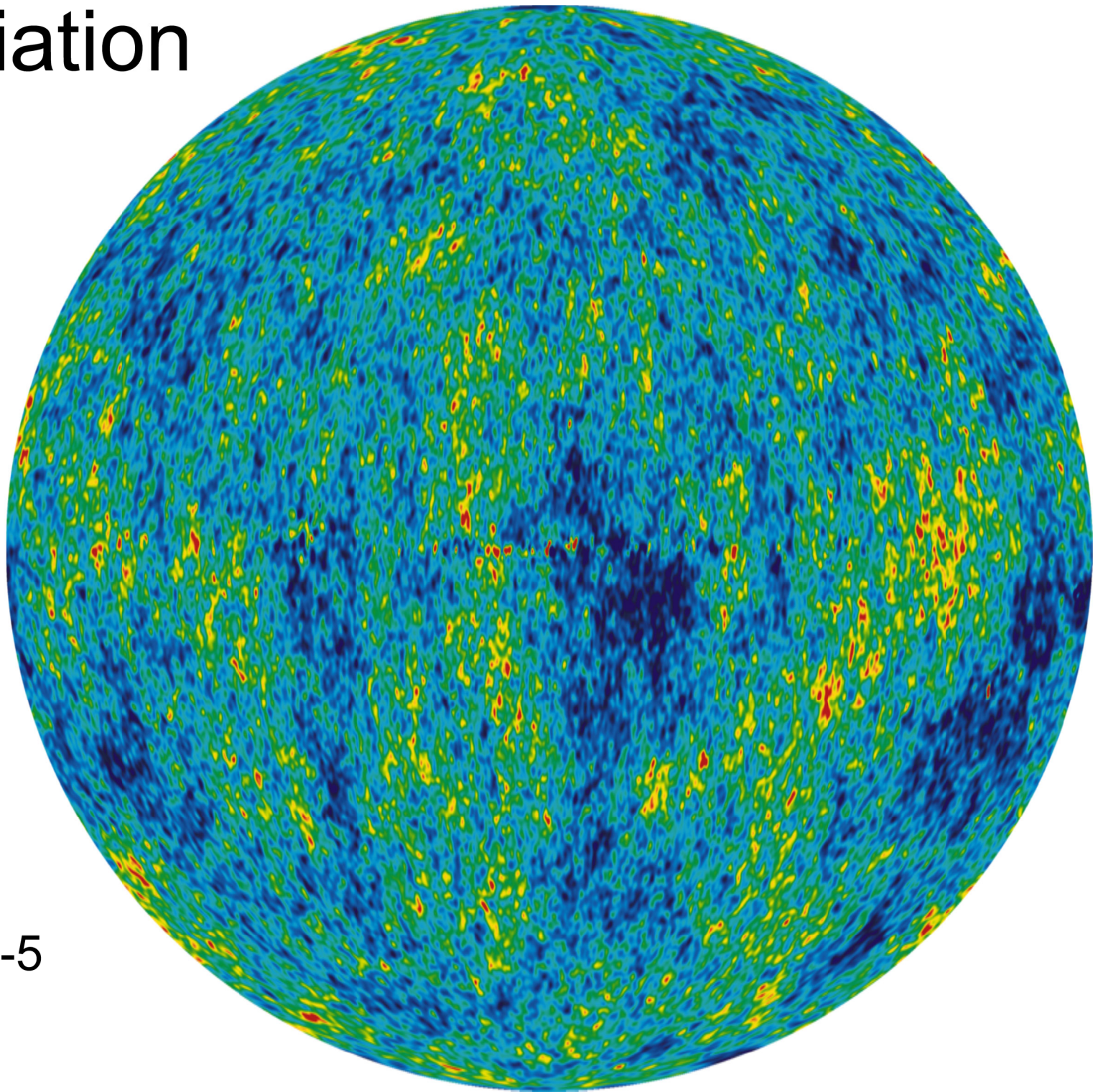


Relic radiation

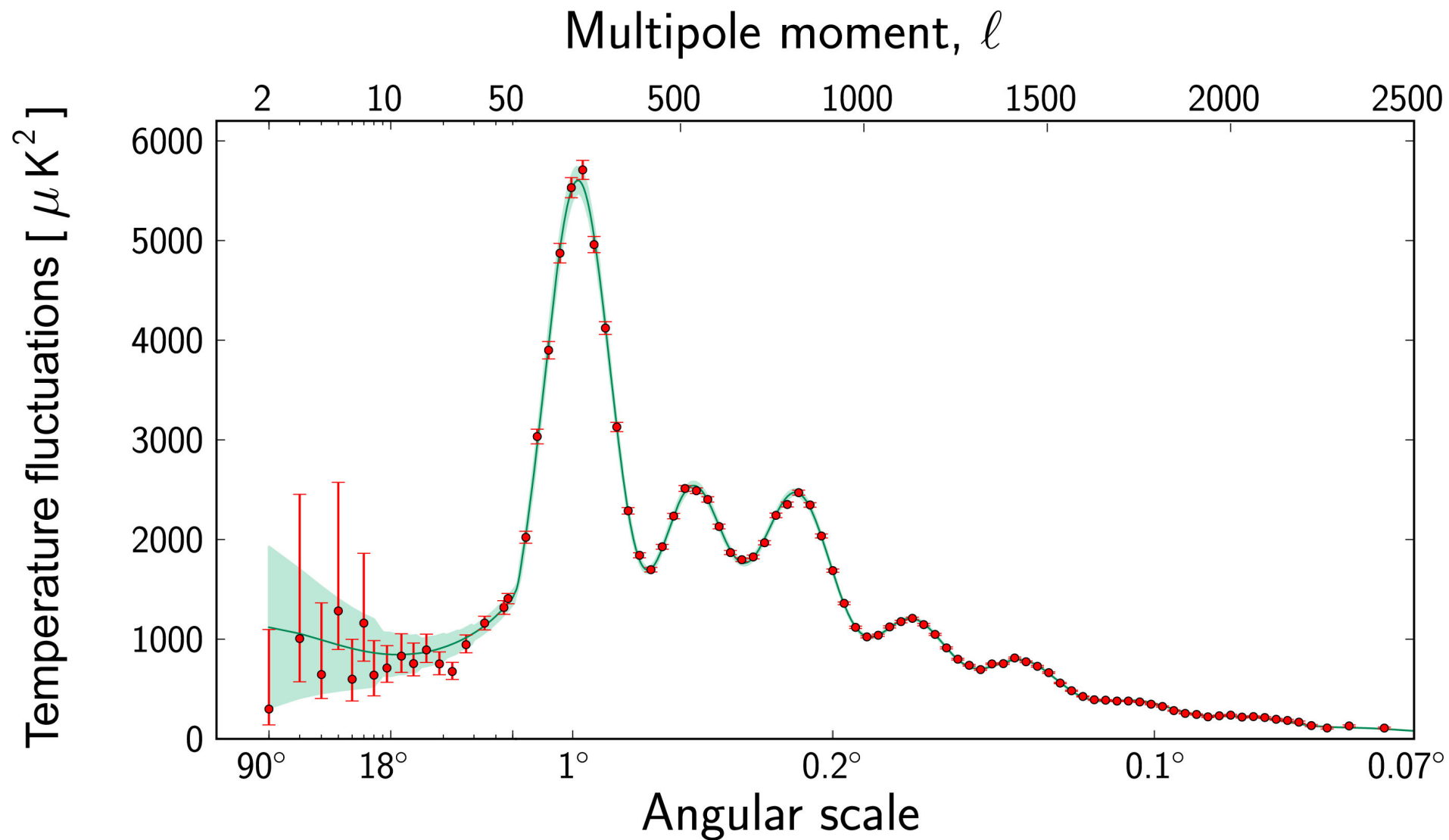
$$\Omega_m \approx 0.31$$

$$\Omega_b \approx 0.05$$

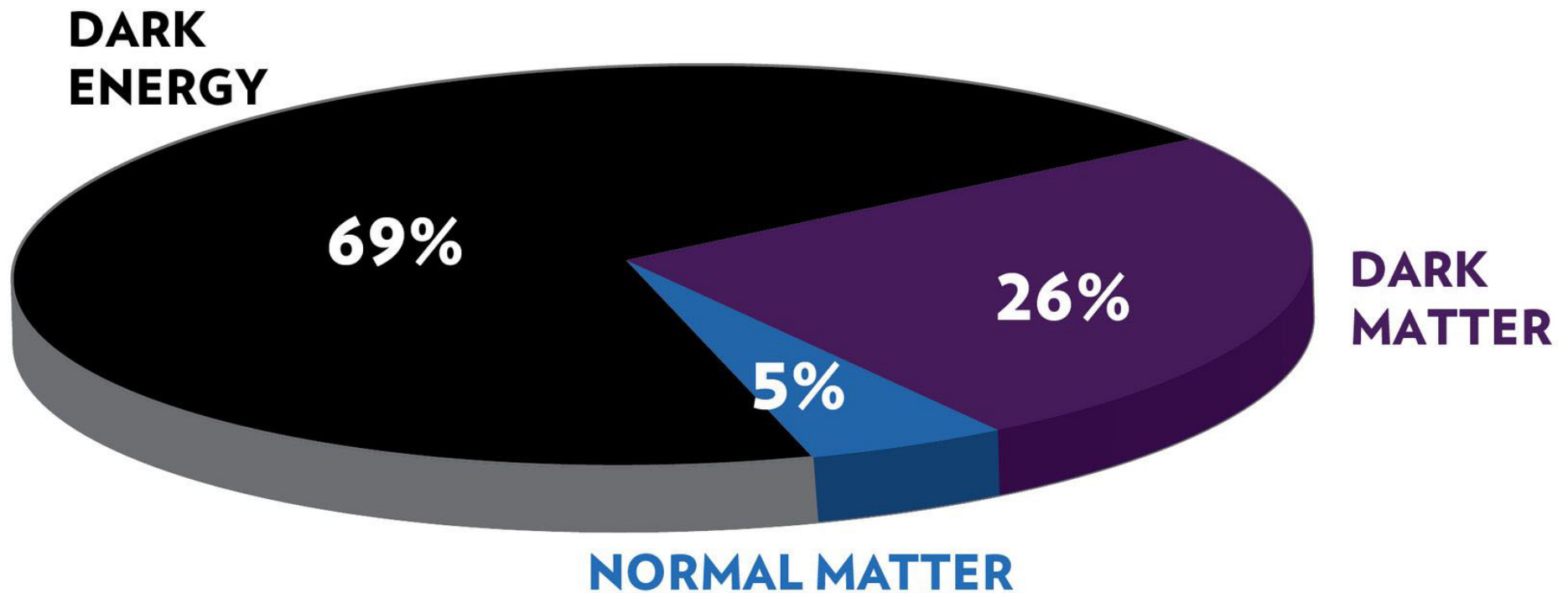
$$\Delta T/T \sim 10^{-5}$$



Spatial spectrum of relic radiation



ENERGY DISTRIBUTION OF THE UNIVERSE

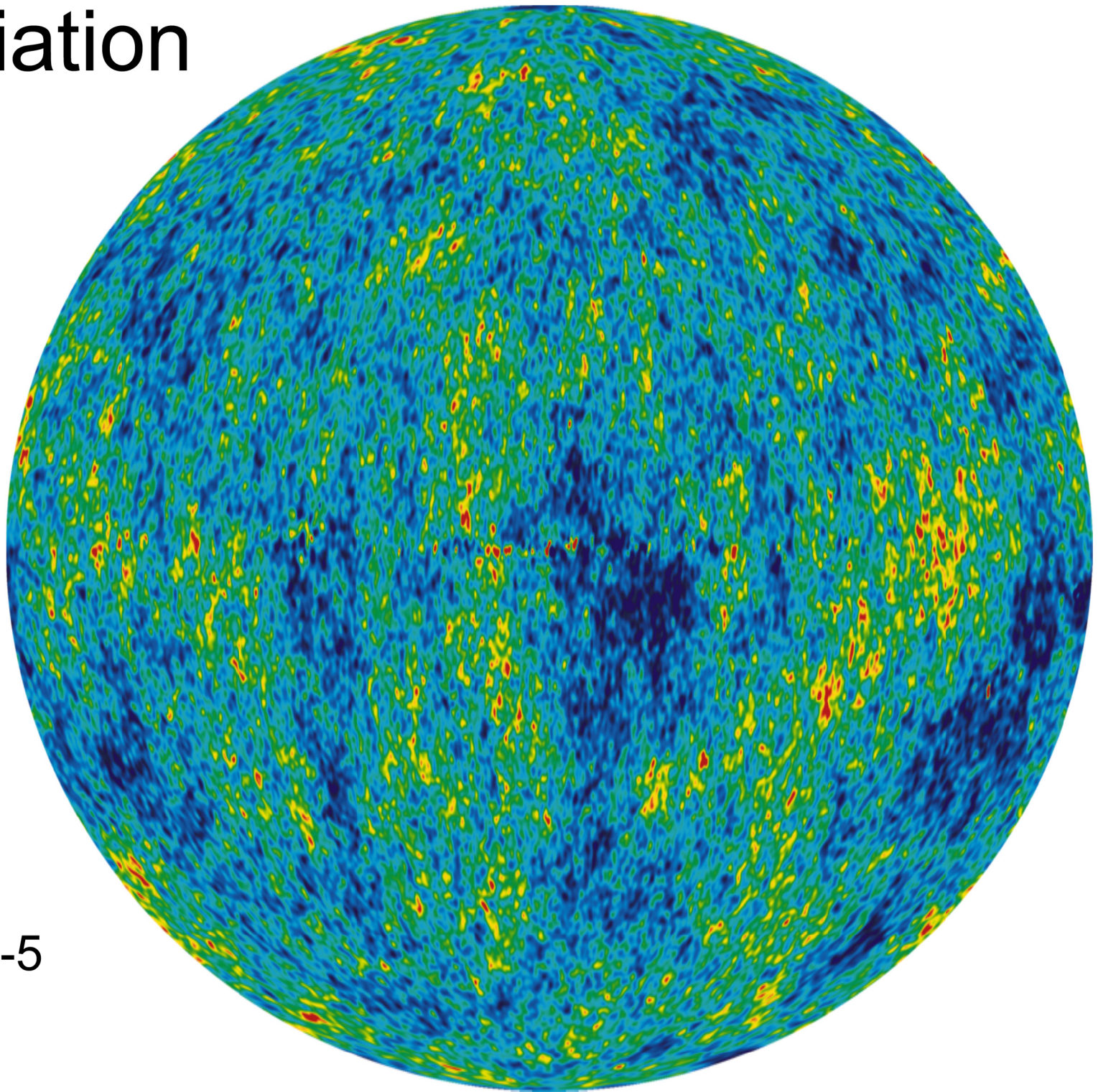


Relic radiation

$$\Omega_m \approx 0.31$$

$$\Omega_b \approx 0.05$$

$$\Delta T/T \sim 10^{-5}$$

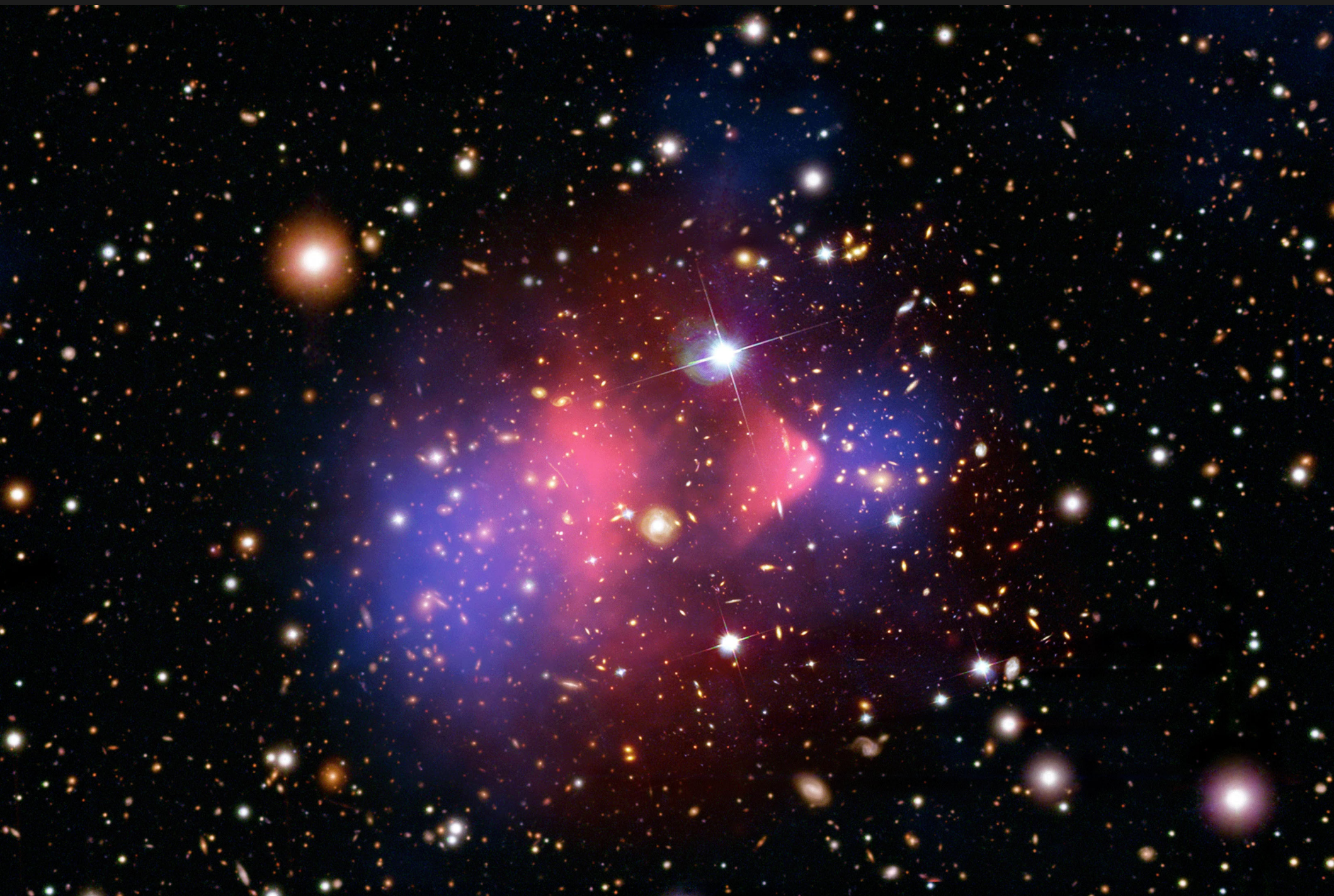


The metric of the Friedmann universe has the form $ds^2 = c^2 dt^2 - a^2(t) dl^2$, where $a(t)$ is scale factor. Instead of $a(t)$ or t , the redshift $z + 1 = a_0/a(t)$ is usually used. The radiation temperature in the Universe is $T \propto a^{-1}(t) \propto (z + 1)$.

For CMB ($z_r \sim 1200$) $\frac{\delta T}{T} \sim 10^{-5}$. Relative disturbance of density of the “ordinary” matter at this moment is $\frac{\delta T}{T} \sim \frac{\delta \rho_b}{\rho_b} < 10^{-4}$.

The density perturbations grow after the CMB formation as $\frac{\delta \rho}{\rho} \propto a(t)$ (the matter-dominated Universe). If there were no dark matter, the matter density perturbations would now have an amplitude $\frac{\delta \rho_b}{\rho_b} \sim z_r \cdot 10^{-4} \simeq 0.1$, and there would be no cosmic structures.

The bullet cluster, 1E 0657-56



2) *Candidates*

WIMP

Weakly Interacting Massive Particle

$$\Omega_{dm} \left(\frac{H}{100 \text{ km}/(\text{s Mpc})} \right)^2 \simeq (3 \cdot 10^{-27} \text{ cm}^3 \text{ s}^{-1} / \langle \sigma v \rangle)$$

1) *Neutralino*

2) *The lightest Kaluza–Klein particle*

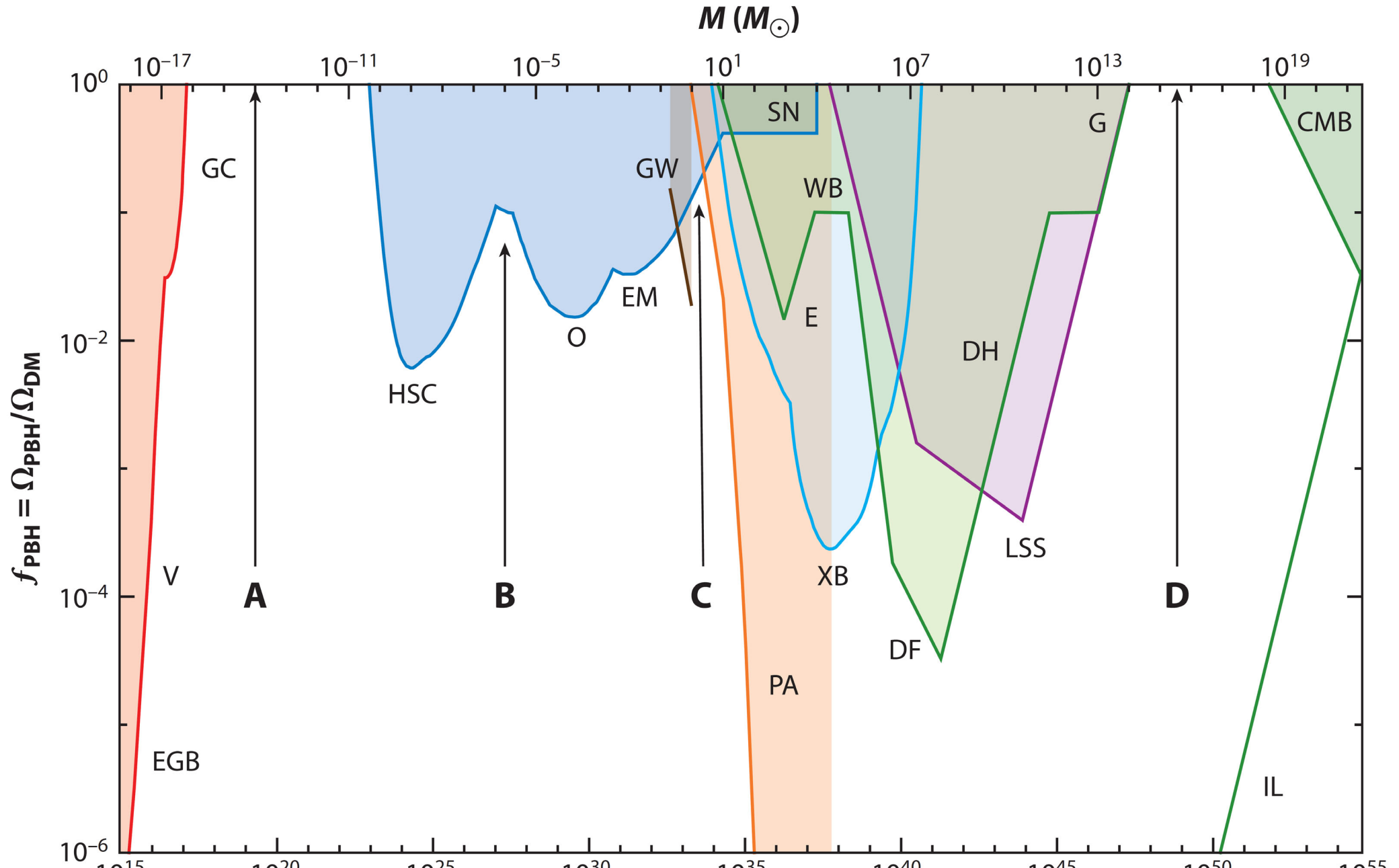
3)

Many other candidates

- 1) *axions*
- 2) *primordial black holes*
- 3) *light boson condensate*
- 4) *dark photons*
- 5) *sterile neutrinos*
- 6) *technicolor dark matter*
- 7) *many other candidates*
- 8) *something else..*

The dark matter can be cold or warm, but cannot be hot!

Primordial black holes (*Carr & Kühnel, 2020*)



3) How to search it?

The LZ Detector

7 tonne liquid xenon
time-projection
chamber

Instrumentation conduits

Existing
water tank

Liquid Xe
heat
exchanger

Gadolinium-loaded
liquid scintillator

120 outer
detector
PMTs

High voltage
feedthrough

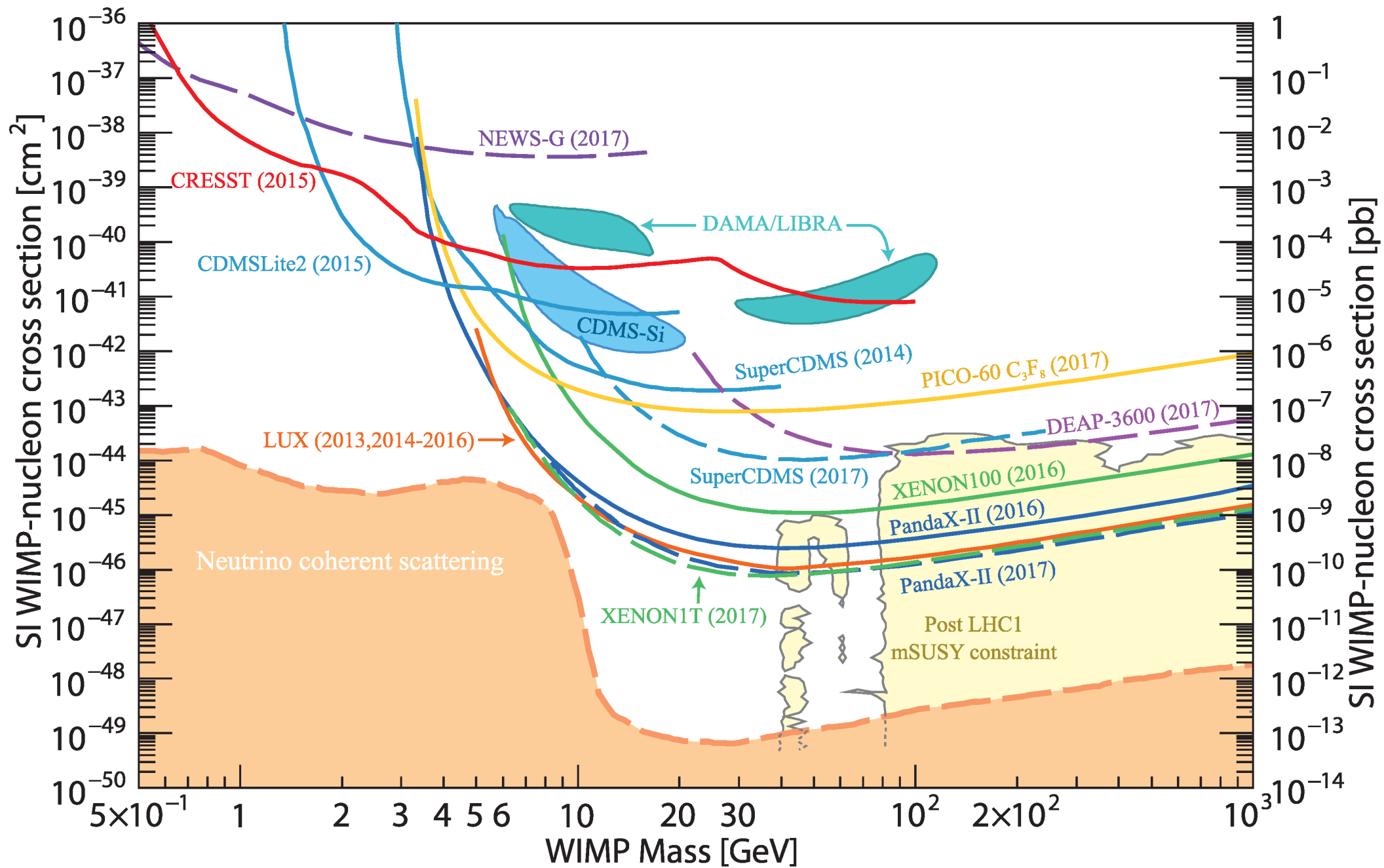
494 photomultiplier tubes (PMTs)

Additional 131 xenon "skin" PMTs

Neutron beampipes

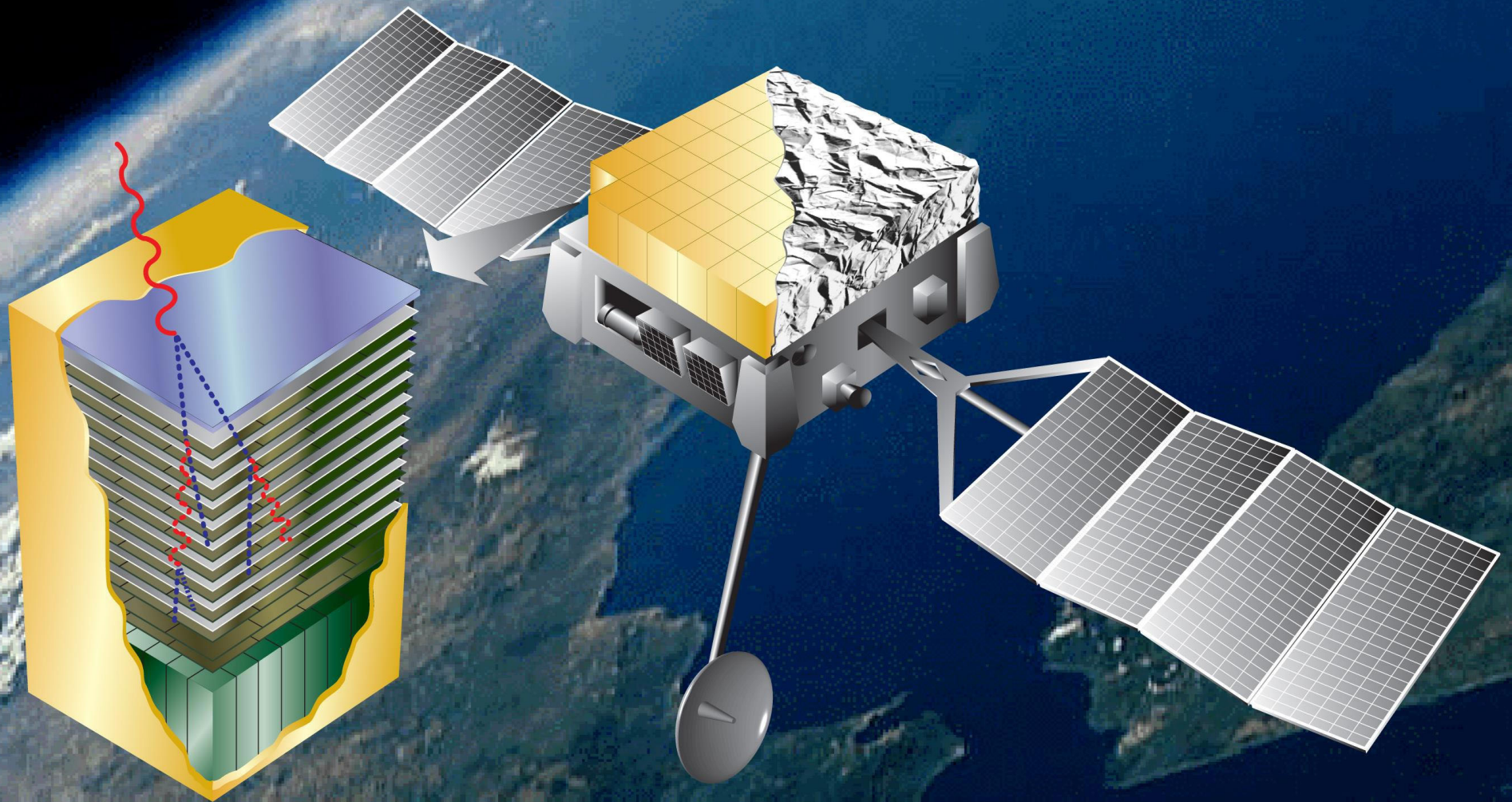


Results



Indirect detection

GAMMA-RAY LARGE AREA SPACE TELESCOPE



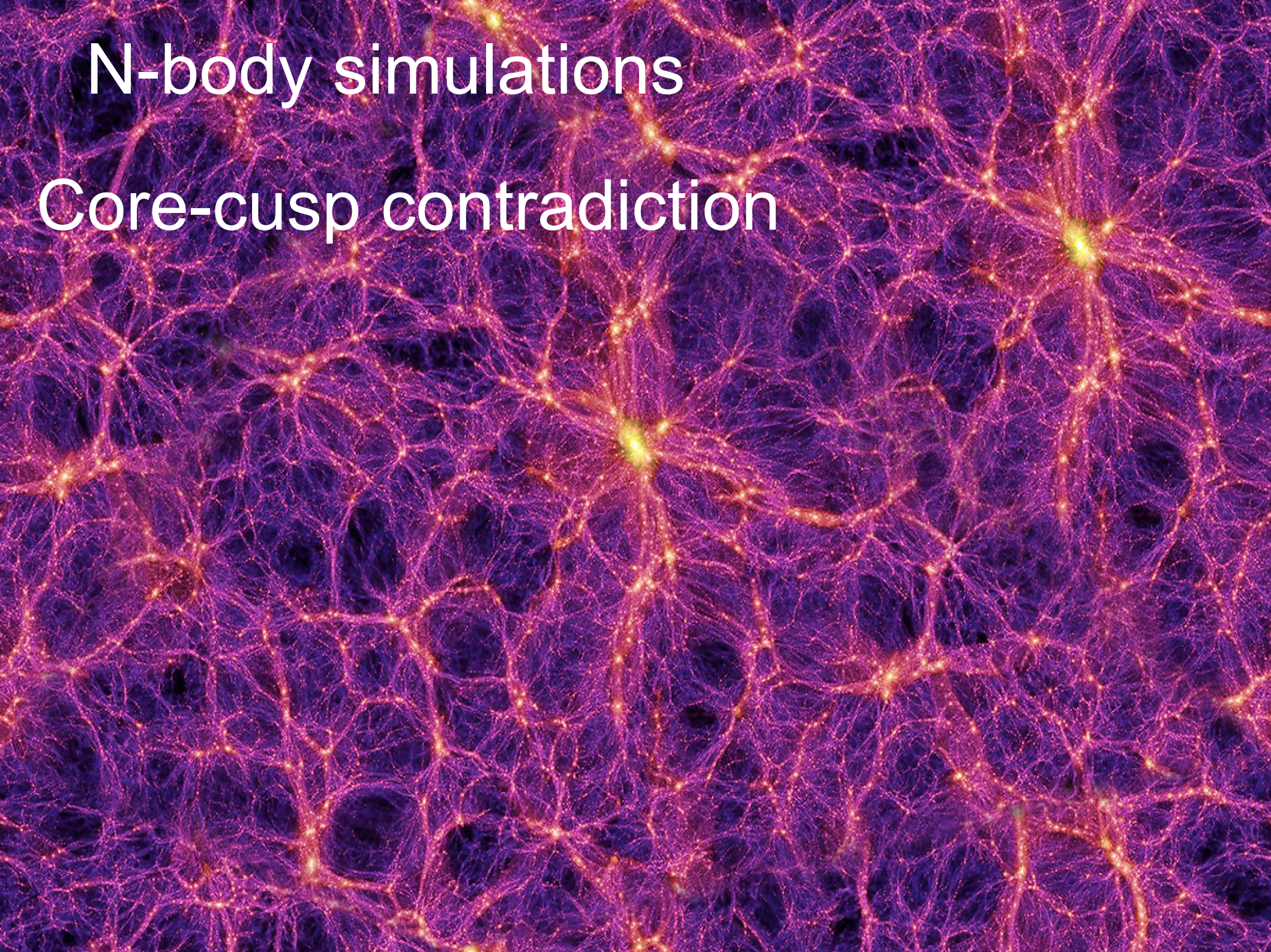
Exploded View:
One of Forty-nine Towers

- 10 Layers of 0.5 rad Length Converter (pb)
- 12 Layers of XY Silicon Strips
- Gamma Rays
- Positrons/Electrons

Fermi Gamma-ray Space Telescope

N-body simulations

Core-cusp contradiction



Isothermal profile

$$\rho \sim r^{-2}$$

Navarro-Frenk-White profile

$$\rho_{NFW} = \frac{\rho_s}{(r/r_s)(1 + r/r_s)^2}$$

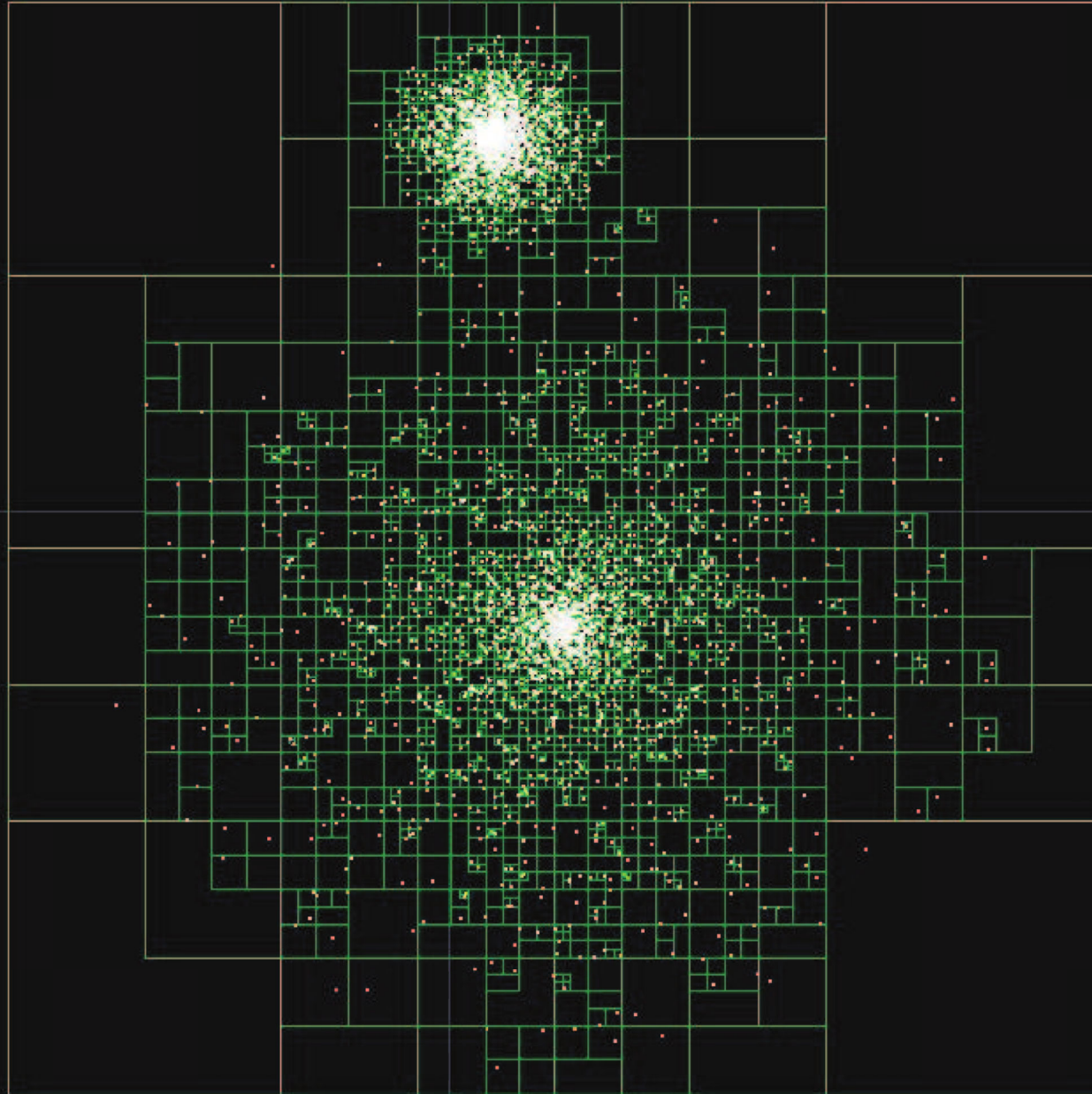
Einasto profile

$$\rho_{Ei} = \rho_s \exp \left\{ -2n \left[\left(\frac{r}{r_s} \right)^{\frac{1}{n}} - 1 \right] \right\}$$

Pseudo-isothermal profile

$$\rho = \frac{\rho_0 r_0^2}{r^2 + r_0^2}$$

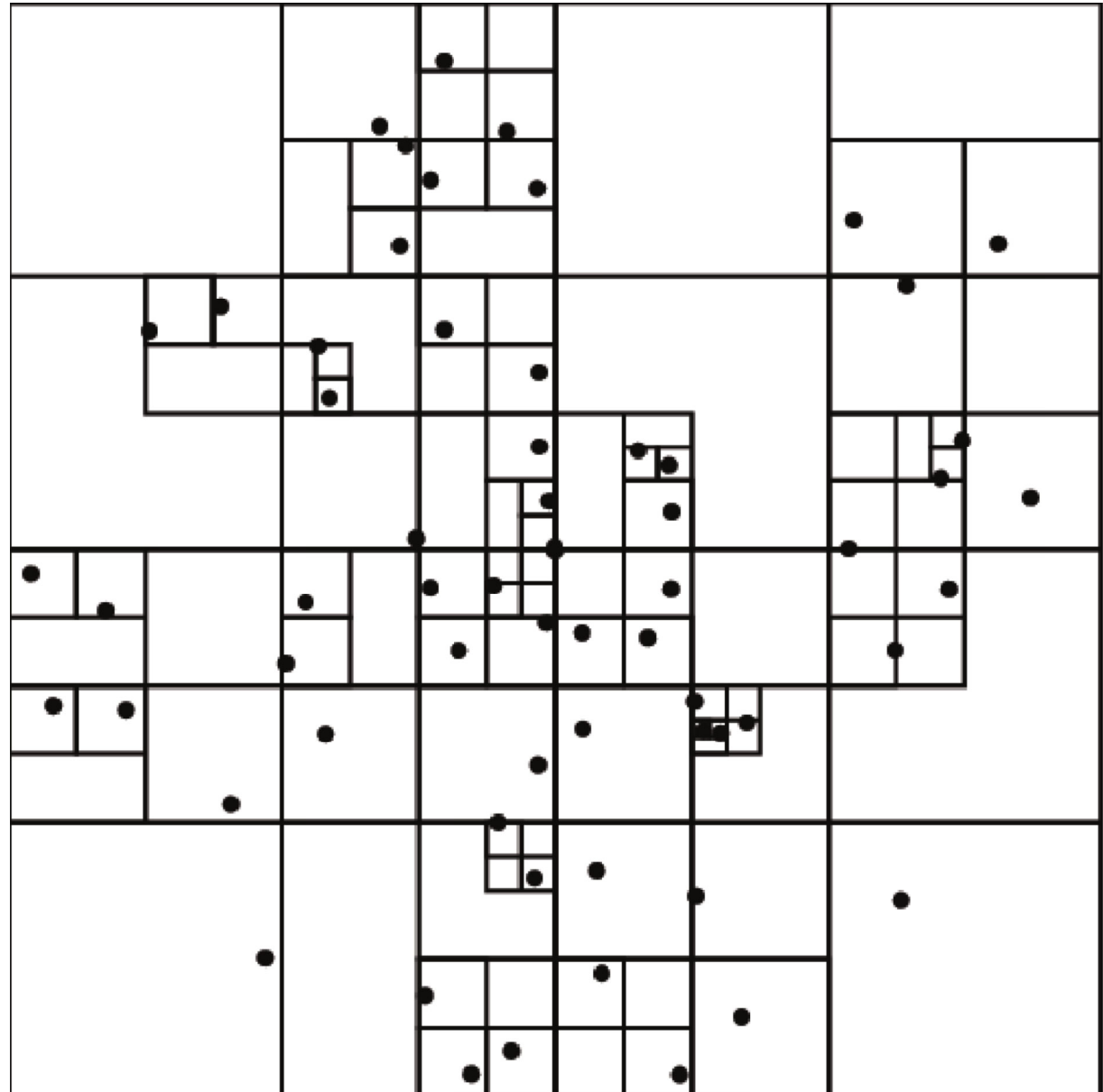
Tree algorithm



Tree algorithm

Cell opening

$$\frac{GM_{cell}}{r^2} \left(\frac{l}{r} \right)^2 \leq f_{acc} |\vec{a}|$$



Trajectory calculation algorithms

