

On a moduli space of the Wigner quasiprobability distributions

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The XXII International Scientific Conference of
Young Scientists and Specialists



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Context

Recently an ambiguity in specification of the Wigner quasiprobability distribution for a finite-dimensional quantum system has been studied. It was shown that for an N -level quantum system one can construct $N - 2$ parametric family of unitary non-equivalent Wigner quasiprobability distributions.

The main objective

In the report the moduli space of the Wigner quasiprobability distributions for N -dimensional quantum systems will be discussed and exemplified for low dimensional cases: for a single qubit, qutrit and quatrit.

Introduction

A quantum state is described by a density operator ϱ :

$$\varrho^\dagger = \varrho; \quad \text{tr}(\varrho) = 1; \quad \varrho \geq 0.$$

The Wigner function is constructed from the density matrix ϱ and the Stratonovich-Weyl kernel $\Delta(\Omega_N)$:

$$W_\varrho(\Omega) = \text{tr} (\varrho \Delta(\Omega)) .$$

Singular value decomposition: $\Delta(\Omega) = U(\Omega) P U^\dagger(\Omega)$,

where $P = \text{diag} ||\pi_1, \dots, \pi_N||$ and $\pi_1 \geq \pi_2 \geq \dots \geq \pi_N$.

Master equations:

$$\text{tr} (\Delta(\Omega)) = 1, \quad \text{tr} (\Delta(\Omega)^2) = N.$$

Parametrization of the moduli space

The Stratonovich-Weyl kernel

$$\Delta(\Omega|\nu) = \frac{1}{N} U(\Omega) \left[I + \kappa \sum_{\lambda \in H} \mu_s(\nu) \lambda_s \right] U(\Omega)^\dagger, \quad \kappa = \sqrt{N(N^2 - 1)/2},$$

where

- H is the **Cartan subalgebra** in $SU(N)$,
- parameter $\nu = (\nu_1, \dots, \nu_{N-2})$ labels members of the WF family,
- coefficients $\boxed{\sum_{s=2}^N \mu_{s^2-1}^2(\nu) = 1}$.

A density matrix of an N -dimensional quantum system

$$\varrho_\xi = \frac{1}{N} \left[I + \sqrt{\frac{N(N-1)}{2}} (\xi, \lambda) \right],$$

where

- ξ is an $(N^2 - 1)$ -dimensional Bloch vector,
- $\lambda = \{\lambda_1, \dots, \lambda_{N^2-1}\}$ is $\mathfrak{su}(N)$ algebra basis.

A family of the Wigner functions

$$W_{\xi}^{(\nu)}(\Omega_N) = \frac{1}{N} \left[1 + \frac{N^2 - 1}{\sqrt{N+1}} (\mathbf{n}, \boldsymbol{\xi}) \right],$$

where

- $\mathbf{n} = \mu_3 \mathbf{n}^{(3)} + \cdots + \mu_{N^2-1} \mathbf{n}^{(N^2-1)},$
- $\mathbf{n}^{(s^2-1)} = \frac{1}{2} \operatorname{tr} (U \lambda_{s^2-1} U^\dagger \lambda_\mu), \quad s = \overline{2, N}.$

The spectrum $\{\pi_1, \dots, \pi_N\}$ of the Stratonovich-Weyl kernel:

$$\pi_i = \frac{1}{N} \left(1 + \sqrt{2} \kappa \sum_{s=i+1}^N \frac{\mu_{s^2-1}}{\sqrt{s(s-1)}} - \kappa \sqrt{\frac{2(i-1)}{i}} \mu_{i^2-1} \right).$$

Constraints on the spherical angles

The spherical $(N - 2)$ angles:

$$\mu_3 = \sin \psi_1 \cdots \sin \psi_{N-2},$$

 \vdots

$$\mu_{i^2-1} = \sin \psi_1 \cdots \sin \psi_{N-i} \cos \psi_{N-i+1},$$

 \vdots

$$\mu_{N^2-1} = \cos \psi_1, \quad i = \overline{2, N}.$$

For decreasing order $\pi_1 \geq \cdots \geq \pi_N$

$$\boxed{\mu_3 \geq 0, \quad \mu_{(i+1)^2-1} \geq \sqrt{\frac{i-1}{i+1}} \mu_{i^2-1}, \quad i = \overline{2, N-1}.}$$

Examples: qubit, qutrit and quatrit

The Wigner function of a single qubit

A generic **qubit** quantum state is parameterized in a standard way

$$\varrho_{\text{qubit}} = \frac{1}{2} (I + \mathbf{r} \cdot \boldsymbol{\sigma})$$

by the Bloch vector $\mathbf{r} = (r \sin \psi \cos \phi, r \sin \phi \sin \phi, r \cos \psi)$.

The master equations determine the spectrum:

$$\text{spec}(P^{(2)}) = \left\{ \frac{1 + \sqrt{3}}{2}, \frac{1 - \sqrt{3}}{2} \right\}.$$

The **Wigner function** for a single qubit is

$$W_{\mathbf{r}}(\alpha, \beta) = \frac{1}{2} + \frac{\sqrt{3}}{2} (\mathbf{r}, \mathbf{n}),$$

where $\mathbf{n} = (-\cos \alpha \sin \beta, \sin \alpha \sin \beta, \cos \beta)$ is the unit 3-vector.

Qutrit kernel and its fundamental region

A generic **qutrit** state is given by the density matrix

$$\varrho_{\text{qutrit}} = \frac{1}{3} \left(I + \sqrt{3} \sum_{\nu=1}^8 \xi_{\nu} \lambda_{\nu} \right).$$

The **Stratonovich-Weyl kernel**

$$\Delta(\Omega_3) = U(\Omega_3) \frac{1}{3} \left[I + 2\sqrt{3} (\mu_3 \lambda_3 + \mu_8 \lambda_8) \right] U(\Omega_3)^{\dagger},$$

where the coefficients

$$\mu_3(\nu) = \frac{\sqrt{3}}{4} \sqrt{(1+\nu)(5-3\nu)}, \quad \mu_8(\nu) = \frac{1}{4}(1-3\nu)$$

are functions of the parameter $\nu = \frac{1}{3} - \frac{4}{3} \cos(\zeta)$ with $\zeta \in [0, \pi/3]$ being the moduli parameter of the unitary nonequivalent WF of a qutrit.

The Wigner function of a single qutrit

$$W_{\xi}^{(\nu)}(\Omega_3) = \frac{1}{3} + \frac{4}{3} [\mu_3(\mathbf{n}^{(3)}, \xi) + \mu_8(\mathbf{n}^{(8)}, \xi)],$$

with two orthogonal unit 8-vectors

$$n_{\nu}^{(3)} = \frac{1}{2} \text{tr} [U \lambda_3 U^{\dagger} \lambda_{\nu}] , \quad n_{\nu}^{(8)} = \frac{1}{2} \text{tr} [U \lambda_8 U^{\dagger} \lambda_{\nu}] .$$

The master equations

$$\text{tr} (\Delta(\Omega)) = 1 , \quad \text{tr} (\Delta(\Omega)^2) = 3$$

determine one-parametric family of kernels $P^{(3)}(\nu)$.

One-parametric $P^{(3)}(\nu)$ -family

- The spectrum of **generic** kernels:

$$\text{spec} \left(P^{(3)}(\nu) \right) = \left\{ \frac{1 - \nu + \delta}{2}, \frac{1 - \nu - \delta}{2}, \nu \right\},$$

where $\delta = \sqrt{(1 + \nu)(5 - 3\nu)}$ and $\nu \in (-1, -\frac{1}{3})$.

- Two **degenerate** kernels:

$$\text{spec} \left(P^{(3)}(-1) \right) = \{1, 1, -1\}, \quad \text{spec} \left(P^{(3)}(-1/3) \right) = \left\{ \frac{5}{3}, -\frac{1}{3}, -\frac{1}{3} \right\}.$$

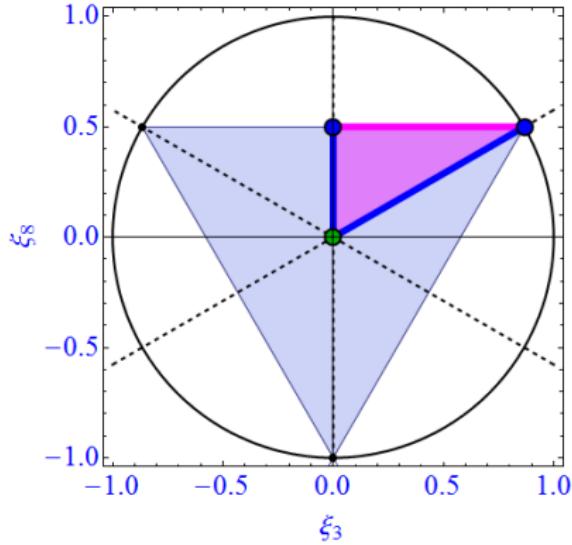
- The spectrum of **singular** kernel:

$$\text{spec} \left(P_{\det=0}^{(3)} \right) = \left\{ \frac{1 + \sqrt{5}}{2}, 0, \frac{1 - \sqrt{5}}{2} \right\}, \quad \text{tr} \left([P_{\det=0}^{(3)}]^m \right) = \mathcal{L}_m,$$

where the m -th **Lucas number** $\mathcal{L}_m = \phi^m + (-\phi)^{-m}$ and $\phi = \frac{1+\sqrt{5}}{2}$.

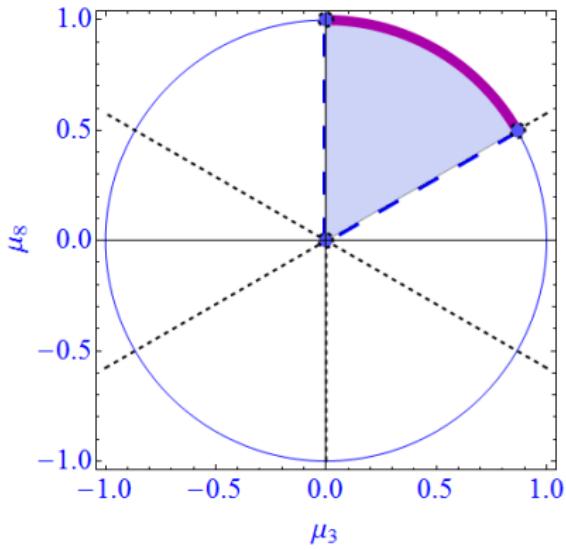
The ordering of the **density matrix** eigenvalues $1 \geq r_1 \geq r_2 \geq r_3 \geq 0$ and condition $\sum r_i = 1$ lead to

$$\xi_3 \geq 0, \quad \xi_8 \geq \frac{\xi_3}{\sqrt{3}}.$$



The ordering of the **SW kernel** eigenvalues $\pi_1 \geq \pi_2 \geq \pi_3$ and condition $\sum \mu_i^2 = 1$ lead to

$$\mu_3 = \sin \zeta, \quad \mu_8 = \cos \zeta, \quad 0 \leq \zeta \leq \frac{\pi}{3}.$$



Quatrit kernel and its fundamental region

A generic **quatrit** ($N = 4$) state is given by the density matrix

$$\varrho_{\text{quatrit}} = \frac{1}{4} \left(I + \sqrt{6} \sum_{\nu=1}^{15} \xi_{\nu} \lambda_{\nu} \right).$$

The **Stratonovich-Weyl kernel**

$$\Delta(\Omega_N|\nu) = U(\Omega_N) \frac{1}{4} \left[I + \sqrt{30} (\mu_3 \lambda_3 + \mu_8 \lambda_8 + \mu_{15} \lambda_{15}) \right] U(\Omega_N)^{\dagger}.$$

The **Wigner function** of a quatrit

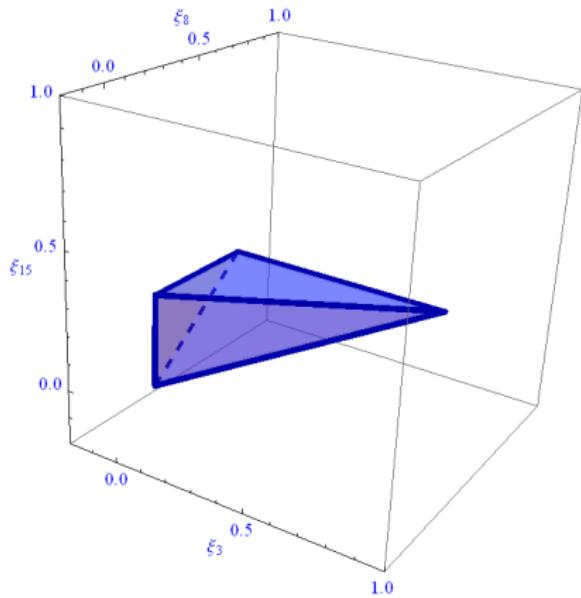
$$W_{\xi}^{(\nu)}(\Omega_4) = \frac{1}{4} + \frac{3\sqrt{5}}{4} \left[\mu_3(\mathbf{n}^{(3)}, \boldsymbol{\xi}) + \mu_8(\mathbf{n}^{(8)}, \boldsymbol{\xi}) + \mu_{15}(\mathbf{n}^{(15)}, \boldsymbol{\xi}) \right],$$

with

$$n_{\nu}^{(3,8,15)} = \frac{1}{2} \text{tr} \left[U \lambda_{3,8,15} U^{\dagger} \lambda_{\nu} \right].$$

Quatrit density matrix

In a quatrit case, there are 24 ways of the spec ($\varrho_{\text{quatrit}} = \{r_1, r_2, r_3, r_4\}$) ordering.



The fixed order of the eigenvalues

$$1 \geq r_1 \geq r_2 \geq r_3 \geq r_4 \geq 0, \\ 0 \leq r_i \leq 1, \quad \sum r_i = 1,$$

leads to

$$0 \leq \xi_3 \leq \sqrt{2/3},$$

$$\frac{\xi_3}{\sqrt{3}} \leq \xi_8 \leq \sqrt{2/3},$$

$$\frac{\xi_8}{\sqrt{2}} \leq \xi_{15} \leq 1/3.$$

The master equations

$$\text{tr}(\Delta(\Omega)) = 1, \quad \text{tr}(\Delta(\Omega)^2) = 4$$

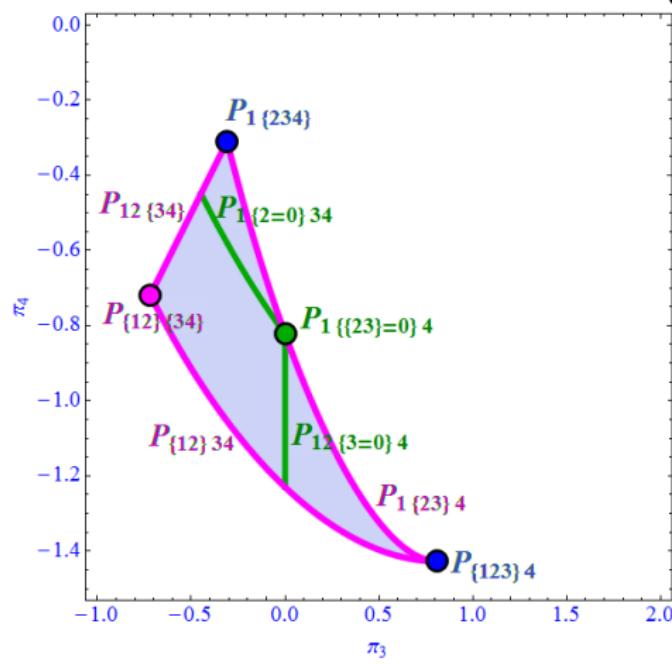
determine two-parametric family of kernels $P^{(4)}$ with $\pi_1 \geq \pi_2 \geq \pi_3 \geq \pi_4$:

- **Generic** kernel:

$$\text{spec}\left(P^{(4)}(\pi_3, \pi_4)\right) = \left\{ \frac{\gamma + \delta}{2}, \frac{\gamma - \delta}{2}, \pi_3, \pi_4 \right\},$$

where

$$\gamma = 1 - \pi_3 - \pi_4, \quad \delta = \sqrt{8 - 2(\pi_3^2 + \pi_4^2) - \gamma^2}.$$



- **Degenerate kernels:**

- **Triple degenerate**

$$P_{\{123\}4}^{(4)} : \pi_1 = \pi_2 = \pi_3 \neq \pi_4 ,$$

$$P_{1\{234\}}^{(4)} : \pi_1 \neq \pi_2 = \pi_3 = \pi_4 .$$

- **Double degenerate**

$$P_{\{12\}\{34\}} : \pi_1 = \pi_2 \neq \pi_3 = \pi_4 ,$$

$$P_{\{12\}34} : \pi_1 = \pi_2 \neq \pi_3 \neq \pi_4 ,$$

$$P_{1\{23\}4} : \pi_1 \neq \pi_2 = \pi_3 \neq \pi_4 ,$$

$$P_{12\{34\}} : \pi_1 \neq \pi_2 \neq \pi_3 = \pi_4 .$$

- **Singular kernels**

$$P_{1\{2=0\}34} : \pi_1 \neq \pi_2 = 0 \neq \pi_3 \neq \pi_4 ,$$

$$P_{12\{3=0\}4} : \pi_1 \neq \pi_2 = 0 \neq \pi_3 \neq \pi_4 ,$$

$$P_{1\{\{23\}=0\}4} : \pi_1 \neq \pi_2 \neq \pi_3 = 0 \neq \pi_4 ,$$

with $\text{tr} \left(P_{1\{\{23\}=0\}4}^m \right) = \mathcal{R}_m .$

Note that

$$\mathcal{R}_m = \mathcal{R}_{m-1} + \frac{3}{2} \mathcal{R}_{m-2}, \quad \mathcal{R}_1 = 1, \quad \mathcal{R}_2 = 4;$$

$$\mathcal{L}_m = \mathcal{L}_{m-1} + \mathcal{L}_{m-2}, \quad \mathcal{L}_1 = 2, \quad \mathcal{L}_2 = 1.$$

Parameterizing μ by two spherical coordinates

$$\mu_3 = \sin \psi_1 \sin \psi_2, \quad \mu_8 = \sin \psi_1 \cos \psi_2, \quad \mu_{15} = \cos \psi_1$$

and using the constraints coming from the requirement of a decreasing order of the SW kernel's eigenvalues

$$\mu_3 \geq 0, \quad \mu_8 \geq \frac{\mu_3}{\sqrt{3}}, \quad \mu_{15} \geq \frac{\mu_8}{\sqrt{2}},$$

we have:

$$\left[\begin{array}{l} \left\{ \begin{array}{l} \psi_2 \in \left(0, \frac{\pi}{3}\right], \\ 0 < \psi_1 \leq \operatorname{arccot}(\cos \psi_2 / \sqrt{2}) ; \end{array} \right. \\ \\ \left\{ \begin{array}{l} \psi_2 = 0, \\ 0 < \psi_1 \leq \operatorname{arccot}(1/\sqrt{2}); \end{array} \right. \quad \text{(See Figure 1)} \\ \\ \psi_1 = 0. \end{array} \right.$$

Girard's theorem: the spherical excess of a triangle determines the solid angle

$$\pi/2 + \pi/3 + \pi/3 - \pi = 4\pi/24.$$

Any fixed order of eigenvalues corresponds to one of 24 possible ways to tessellate a sphere.

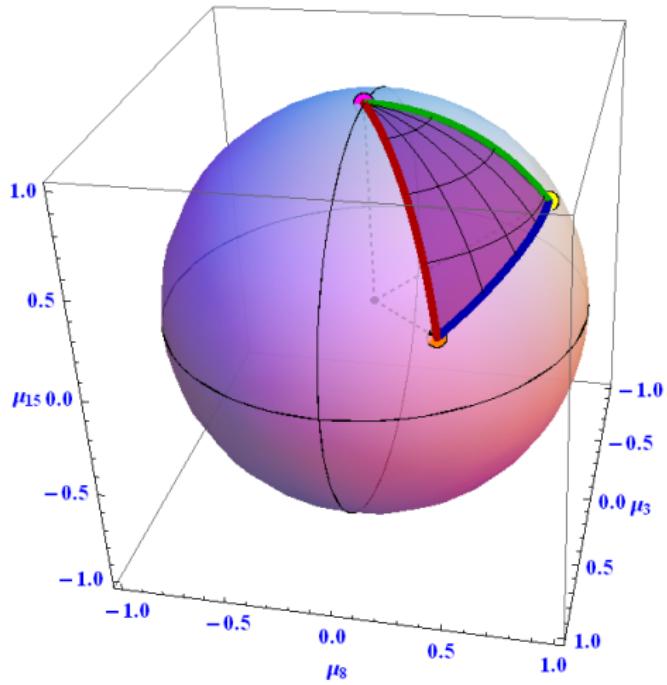


Figure 1: Möbius (2,3,3) triangle with $(\pi/2, \pi/3, \pi/3)$ angles.

Conclusions

An ambiguity in the master equation's solution for Stratonovich-Weyl kernel is analyzed and the corresponding moduli spaces of the Wigner QPFD is determined for $N = 3, 4$ quantum systems:

- for the qutrit the moduli space is the $\frac{\pi}{3}$ arc of the unit circle,
- for the quatrit the moduli space is (2, 3, 3) Möbius triangle.

The basic goal of our further studies is
understanding of a physical meaning of the Wigner function moduli space.

Thank you for attention