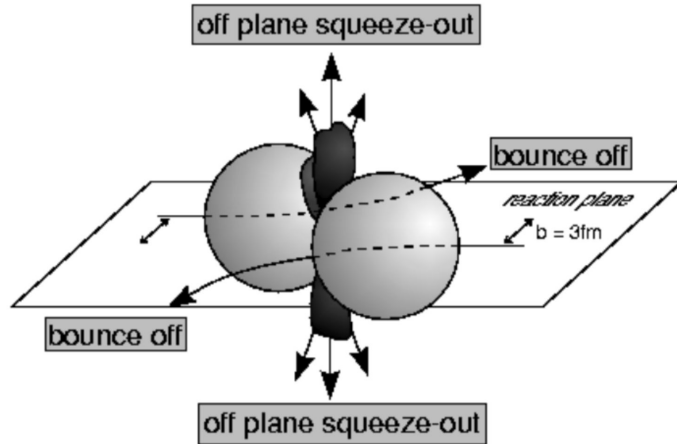


Performance study for the anisotropic flow measurements in MPD-FXT

P. Parfenov, M. Mamaev and A. Taranenko
(NRNU MEPhI, JINR)

Anisotropic flow & spectators



The azimuthal angle distribution is decomposed in a Fourier series relative to reaction plane angle:

$$\rho(\varphi - \Psi_{RP}) = \frac{1}{2\pi} \left(1 + 2 \sum_{n=1}^{\infty} v_n \cos n(\varphi - \Psi_{RP}) \right)$$

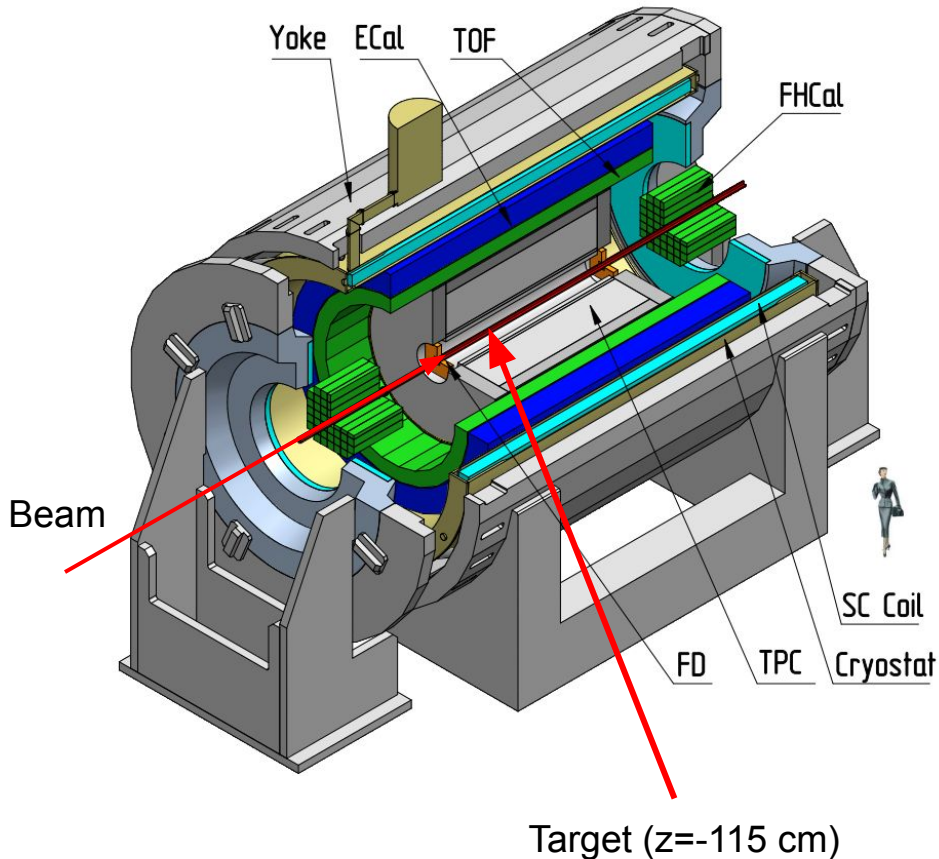
Anisotropic flow:

$$v_n = \langle \cos [n(\varphi - \Psi_{RP})] \rangle$$

Anisotropic flow is sensitive to:

- Time of the interaction between overlap region and spectators
- Compressibility of the created matter

MPD in Fixed-Target Mode (FXT)



- Model used: UrQMD mean-field
 - Bi+Bi, $E_{\text{kin}} = 1.45$ AGeV ($\sqrt{s_{\text{NN}}} = 2.5$ GeV)
 - Bi+Bi, $E_{\text{kin}} = 2.92$ AGeV ($\sqrt{s_{\text{NN}}} = 3.0$ GeV)
 - Bi+Bi, $E_{\text{kin}} = 4.65$ AGeV ($\sqrt{s_{\text{NN}}} = 3.5$ GeV)
- Point-like target
- GEANT4 transport
- Particle species selection via true-PDG code of the associated mc track

Flow vectors

From momentum of each measured particle define a u_n -vector in transverse plane:

$$u_n = e^{in\phi}$$

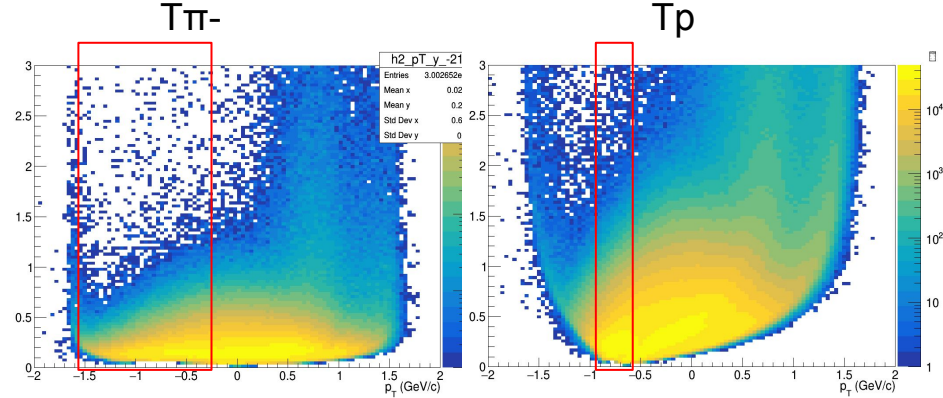
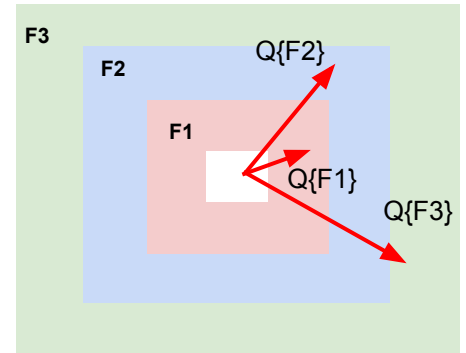
where ϕ is the azimuthal angle

Sum over a group of u_n -vectors in one event forms Q_n -vector:

$$Q_n = \frac{\sum_{k=1}^N w_n^k u_n^k}{\sum_{k=1}^N w_n^k} = |Q_n| e^{in\Psi_n^{EP}}$$

Ψ_n^{EP} is the event plane angle

Modules of FHCAL divided into 3 groups



Additional subevents from tracks not pointing at FHCAL:

Tp: p ; $-1.0 < y < -0.6$;

Tπ: π^- ; $-1.5 < y < -0.2$;

Flow methods for v_n calculation

Tested in HADES: M Mamaev et al 2020 PPNuclei 53, 277–281
M Mamaev et al 2020 J. Phys.: Conf. Ser. 1690 012122

Scalar product (SP) method:

$$v_1 = \frac{\langle u_1 Q_1^{F1} \rangle}{R_1^{F1}} \quad v_2 = \frac{\langle u_2 Q_1^{F1} Q_1^{F3} \rangle}{R_1^{F1} R_1^{F3}}$$

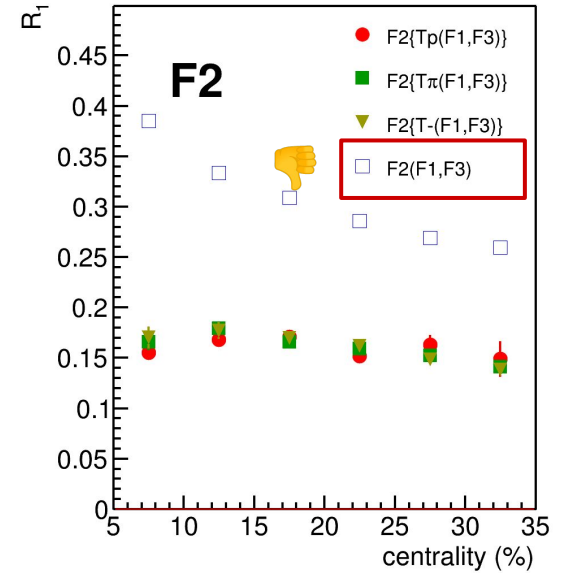
Where R_1 is the resolution correction factor

$$R_1^{F1} = \langle \cos(\Psi_1^{F1} - \Psi_1^{RP}) \rangle$$

Symbol “F2(F1,F3)” means R_1 calculated via
(3S resolution):

$$R_1^{F2(F1,F3)} = \frac{\sqrt{\langle Q_1^{F2} Q_1^{F1} \rangle \langle Q_1^{F2} Q_1^{F3} \rangle}}{\sqrt{\langle Q_1^{F1} Q_1^{F3} \rangle}}$$

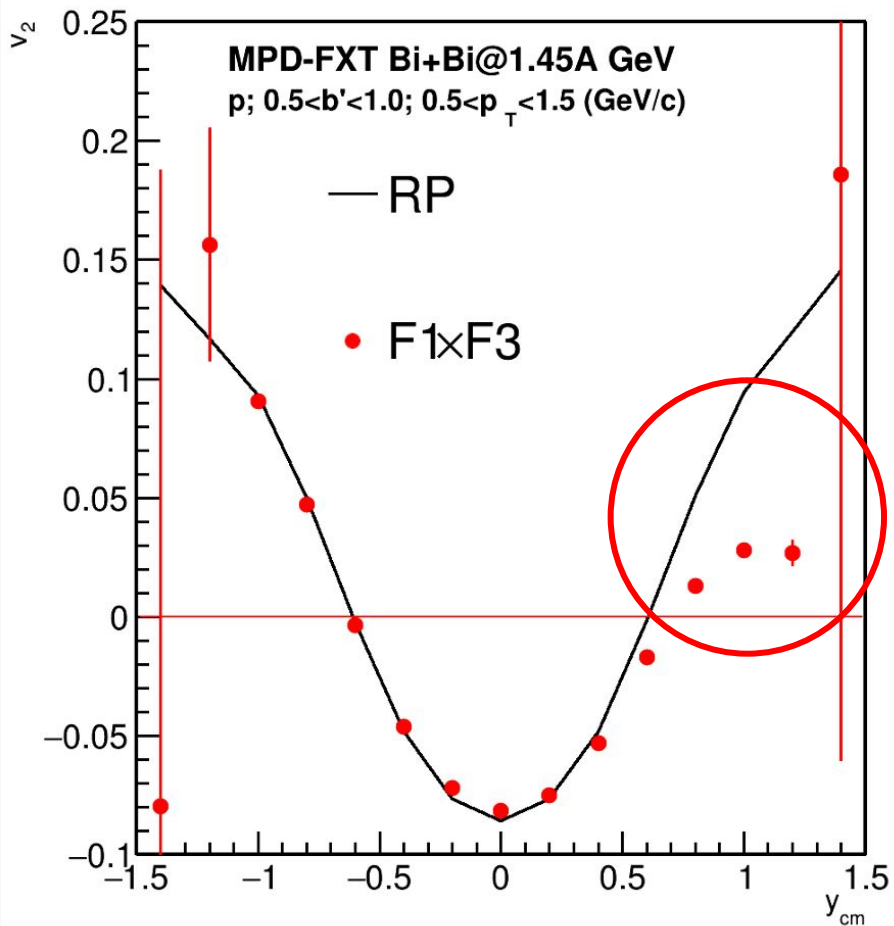
Method helps to eliminate non-flow
Using 2-subevents doesn't



Symbol “F2{Tp}(F1,F3)” means R_1
calculated via (4S resolution):

$$R_1^{F2\{Tp\}(F1,F3)} = \langle Q_1^{F2} Q_1^{Tp} \rangle \frac{\sqrt{\langle Q_1^{F1} Q_1^{F3} \rangle}}{\sqrt{\langle Q_1^{Tp} Q_1^{F1} \rangle \langle Q_1^{Tp} Q_1^{F3} \rangle}}$$

Previous results: main issue



Discrepancy between “reco” and “sim”:

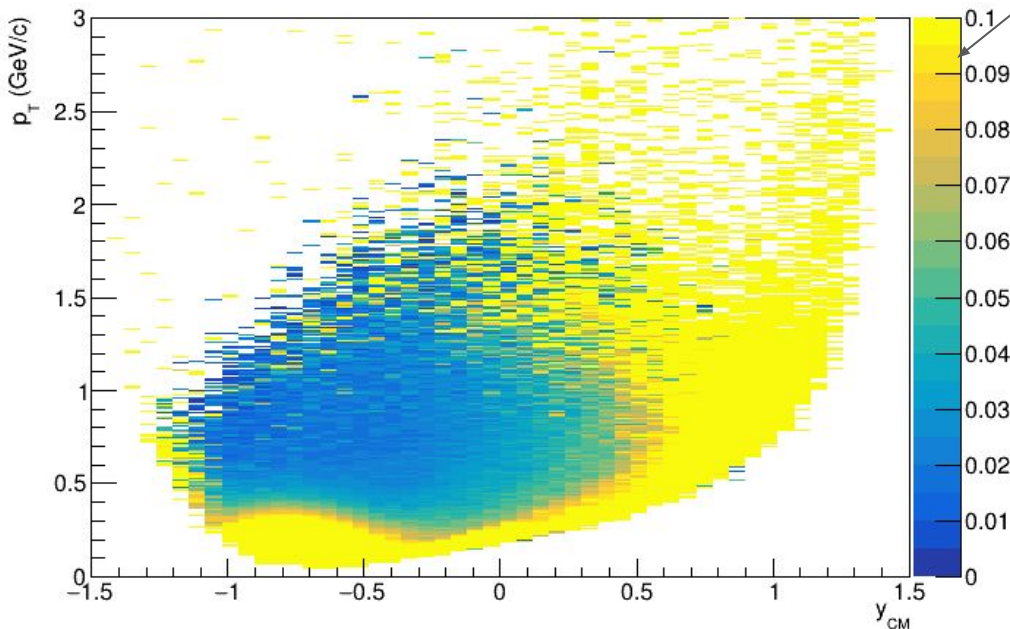
- Efficiency?
- Track quality?

Let's look at the track quality...

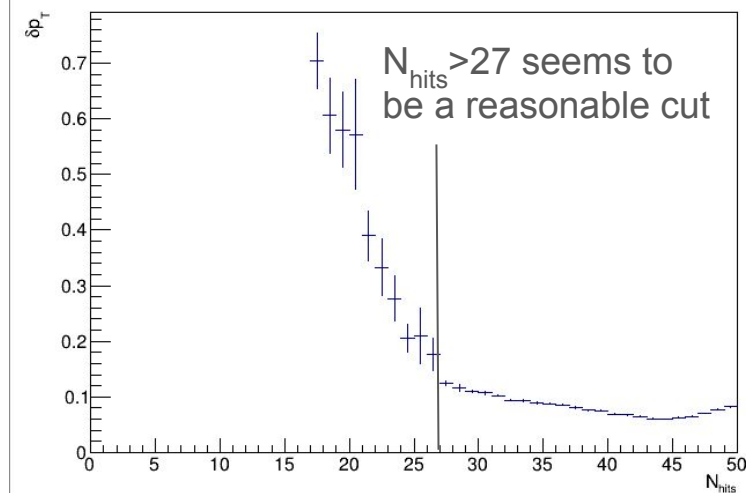
Basic track quality check: p_T

$$\Delta p_T = \frac{|p_T^{\text{reco}} - p_T^{\text{mc}}|}{p_T^{\text{mc}}}$$

Pt-resolution for reconstructed protons in Ycm-pT plane



Pt-resolution for reconstructed protons vs. Nhits

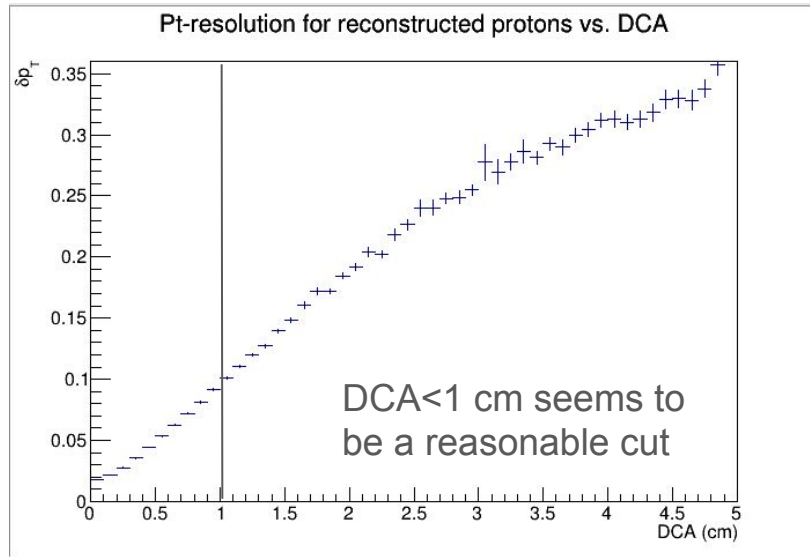
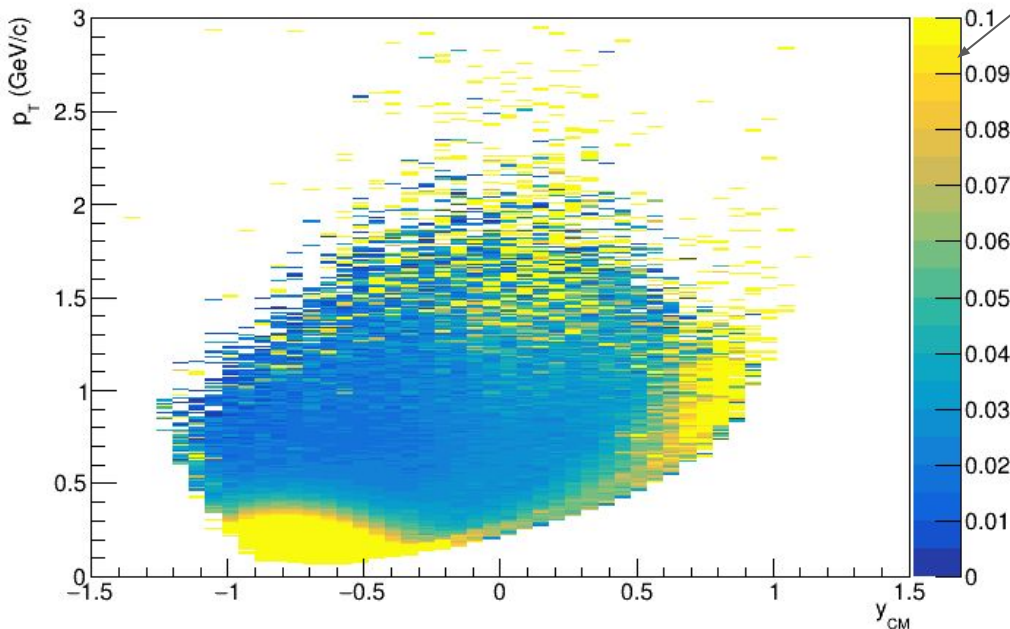


It seems the pt-resolution drops in the forward rapidity region ($y_{\text{CM}} > 0.5$)

Basic track quality check: p_T

$$\Delta p_T = \frac{|p_T^{\text{reco}} - p_T^{\text{mc}}|}{p_T^{\text{mc}}}$$

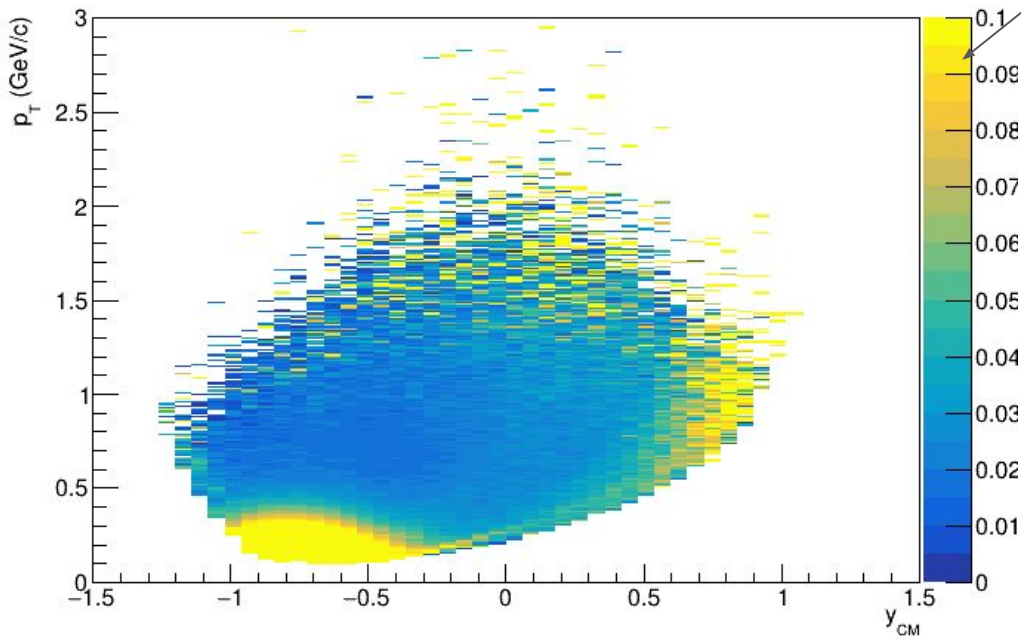
Pt-resolution for reconstructed protons in Ycm-pT plane



Cut $N_{\text{hits}} > 27$ seems to improve the situation

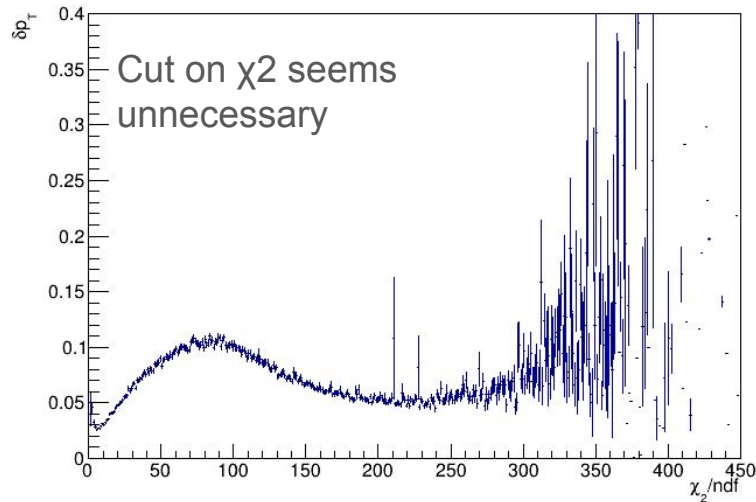
Basic track quality check: p_T

Pt-resolution for reconstructed protons in Ycm-pT plane



$$\Delta p_T = \frac{|p_T^{\text{reco}} - p_T^{\text{mc}}|}{p_T^{\text{mc}}}$$

Pt-resolution for reconstructed protons vs. Chi2

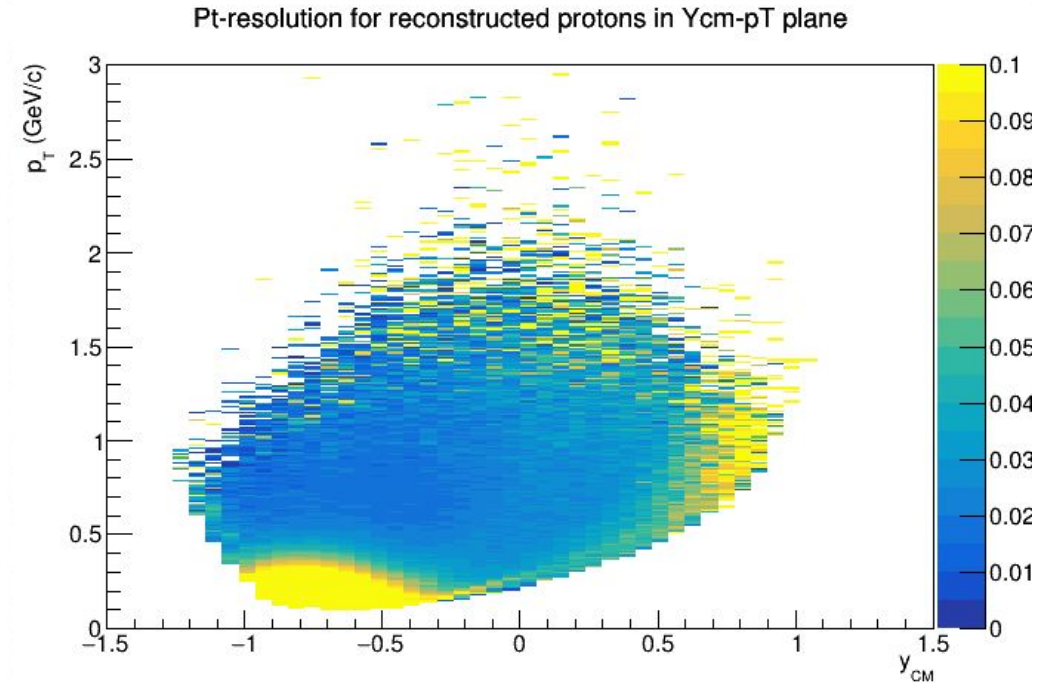


Cut DCA < 1 cm slightly improve the situation

Track cuts based on p_T -resolution check

Protons:

- $N_{\text{hits}} > 27$
- $\text{DCA} < 1 \text{ cm}$



Now let's look at the efficiency plots with the new cuts

(y-pt) distribution, efficiency and δp_T (protons)

$$\text{eff} = \frac{\frac{dN}{dydp_T}(\text{reco})}{\frac{dN}{dydp_T}(\text{sim})}$$

$$\Delta p_T = \frac{|p_T^{\text{reco}} - p_T^{\text{mc}}|}{p_T^{\text{mc}}}$$

Bi+Bi $\sqrt{s_{NN}}=2.5$ GeV

Cuts for reco tracks:

- Nhits>27
- DCA< 1 cm
- PID (pdg code)
- Primary (motherId)

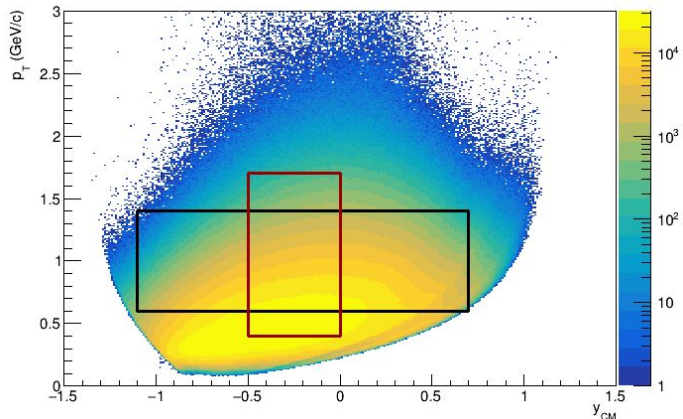
Cuts for sim particles:

- PID (pdg code)
- Primary (motherId)

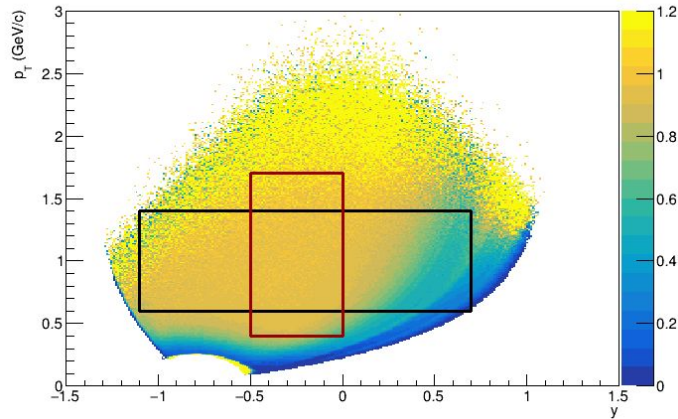
Black box: acceptance window for $v_n(y)$

Red box: acceptance window for $v_n(p_T)$

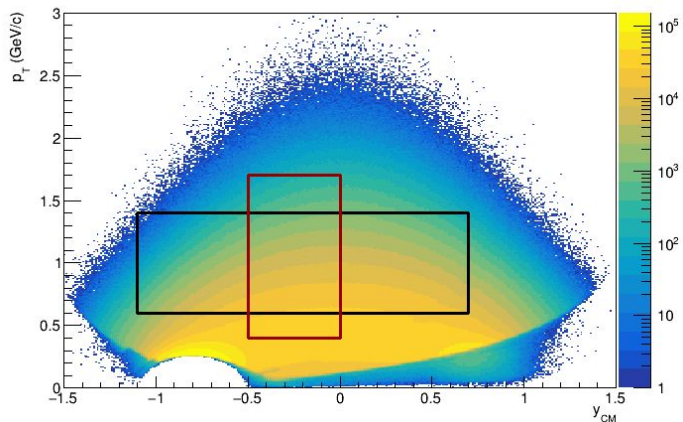
Reconstructed protons Ycm-pT



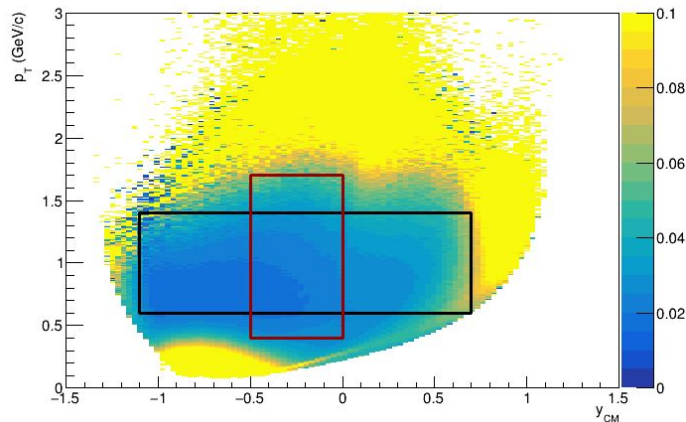
Efficiency (Y-pT) of primary protons



Simulated protons Ycm-pT



Pt-resolution for reconstructed protons in Ycm-pT plane



(y-pt) distribution, efficiency and δp_T (protons)

$$\text{eff} = \frac{\frac{dN}{dydp_T}(\text{reco})}{\frac{dN}{dydp_T}(\text{sim})}$$

$$\Delta p_T = \frac{|p_T^{\text{reco}} - p_T^{\text{mc}}|}{p_T^{\text{mc}}}$$

Bi+Bi $\sqrt{s_{NN}}=3.0$ GeV

Cuts for reco tracks:

- Nhits>27
- DCA< 1 cm
- PID (pdg code)
- Primary (motherId)

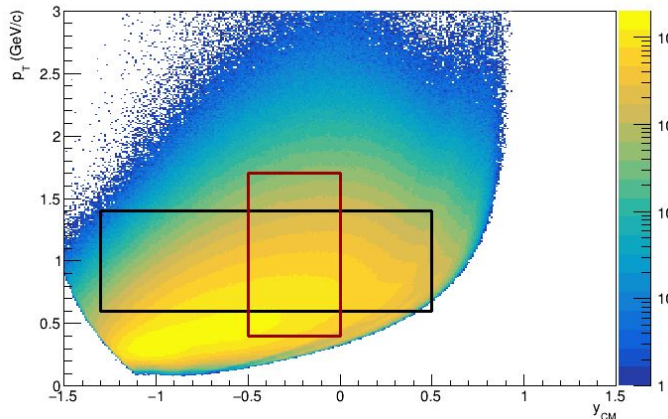
Cuts for sim particles:

- PID (pdg code)
- Primary (motherId)

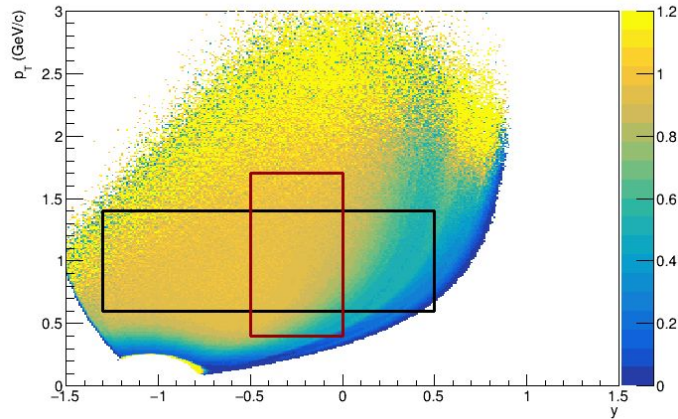
Black box: acceptance window for $v_n(y)$

Red box: acceptance window for $v_n(p_T)$

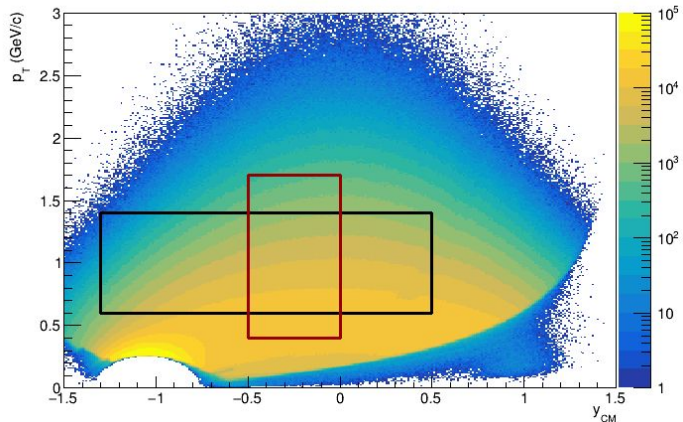
Reconstructed protons Ycm-pT



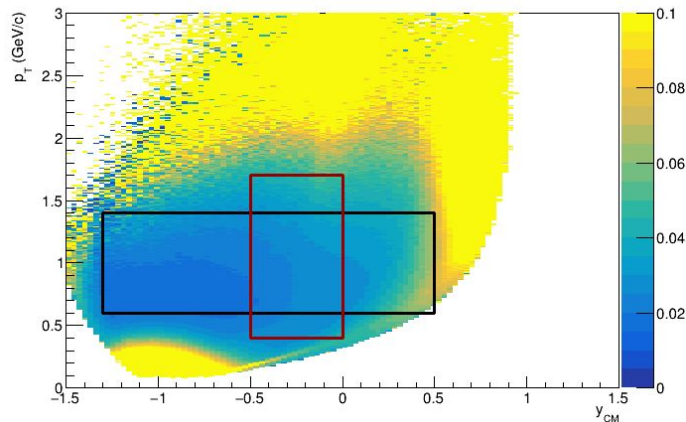
Efficiency (Y-pT) of primary protons



Simulated protons Ycm-pT



Pt-resolution for reconstructed protons in Ycm-pT plane



(y-pt) distribution, efficiency and δp_T (protons)

$$\text{eff} = \frac{\frac{dN}{dydp_T}(\text{reco})}{\frac{dN}{dydp_T}(\text{sim})}$$

$$\Delta p_T = \frac{|p_T^{\text{reco}} - p_T^{\text{mc}}|}{p_T^{\text{mc}}}$$

Bi+Bi $\sqrt{s_{NN}}=3.5$ GeV

Cuts for reco tracks:

- Nhits>27
- DCA< 1 cm
- PID (pdg code)
- Primary (motherId)

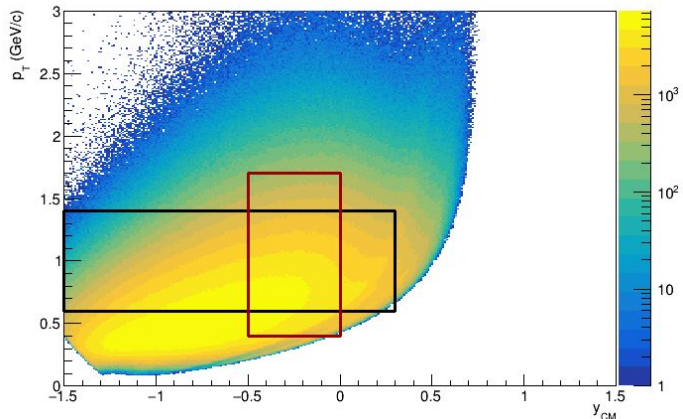
Cuts for sim particles:

- PID (pdg code)
- Primary (motherId)

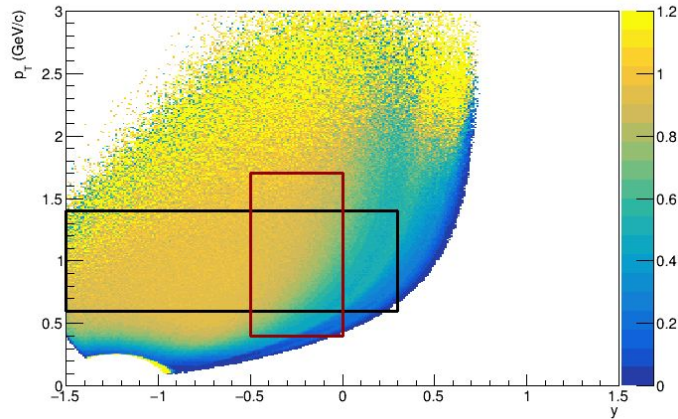
Black box: acceptance window for $v_n(y)$

Red box: acceptance window for $v_n(p_T)$

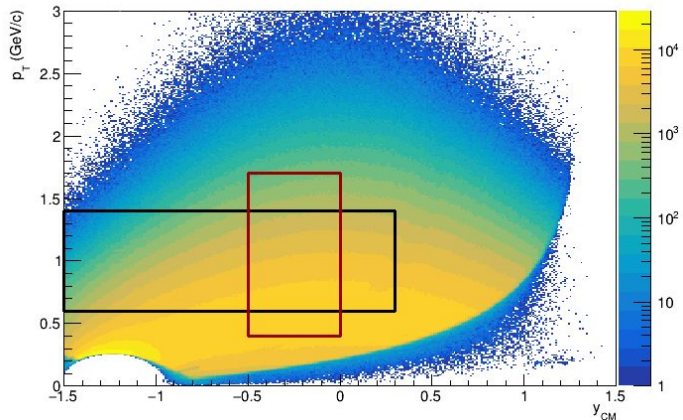
Reconstructed protons Ycm-pT



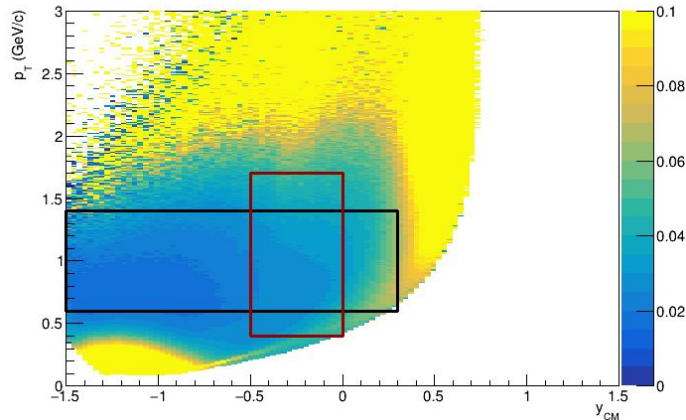
Efficiency (Y-pT) of primary protons



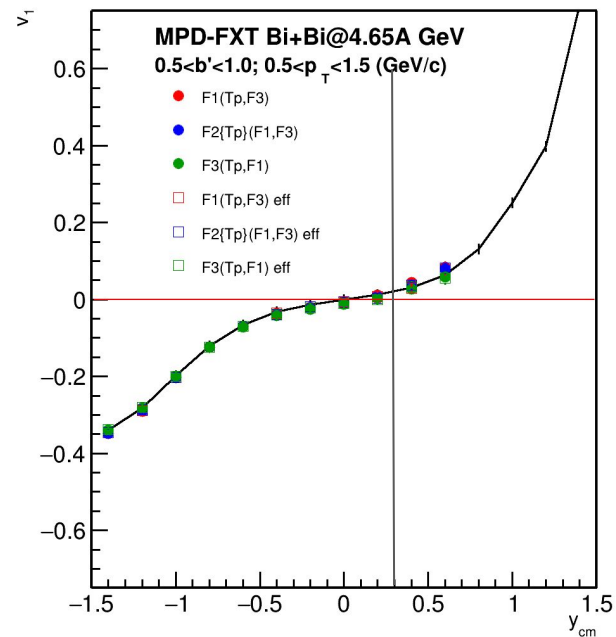
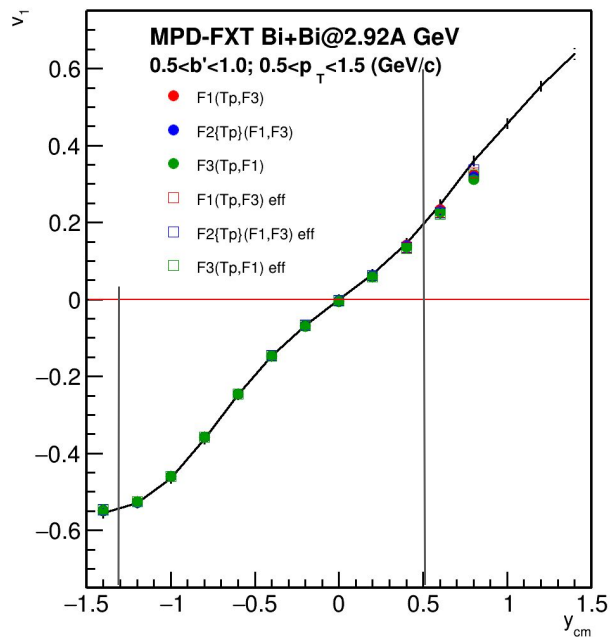
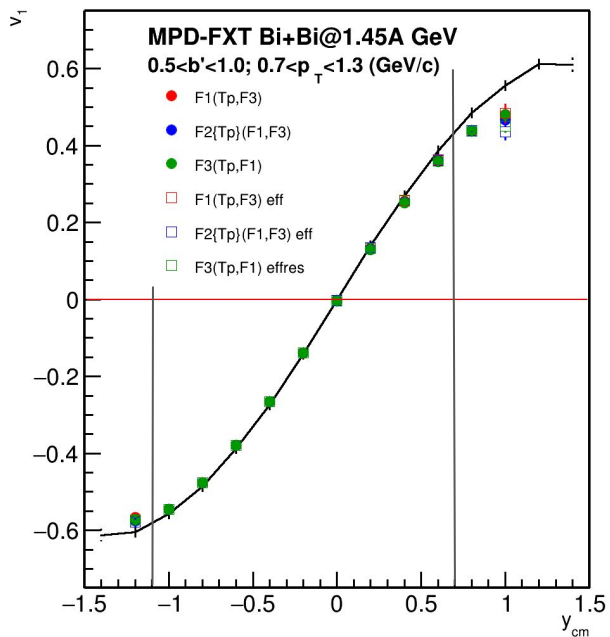
Simulated protons Ycm-pT



Pt-resolution for reconstructed protons in Ycm-pT plane

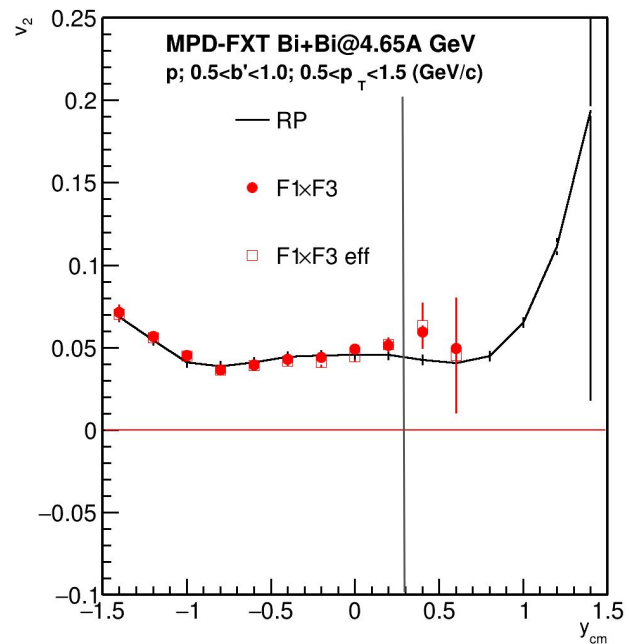
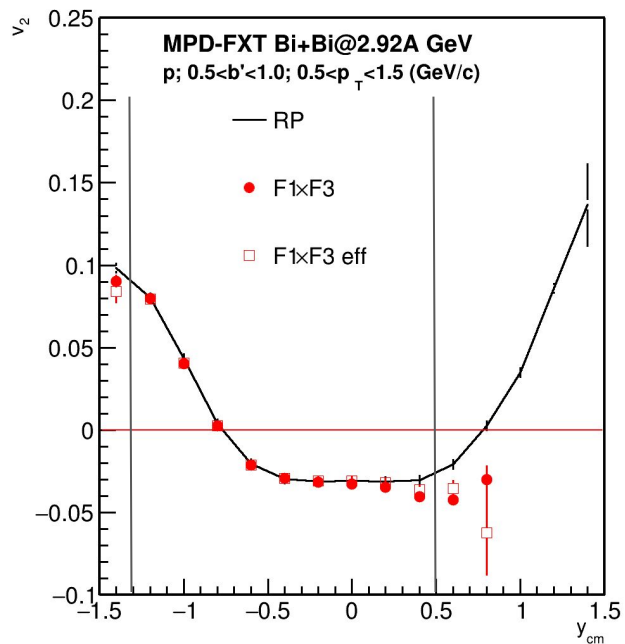
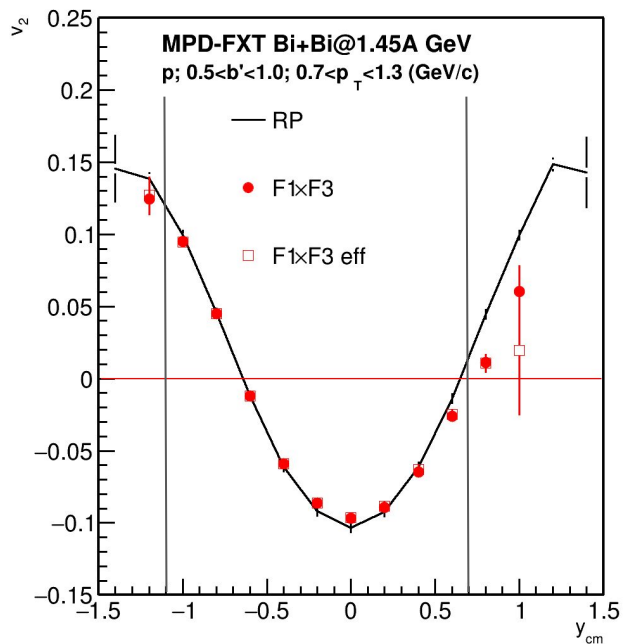


$v_1(y)$ of protons



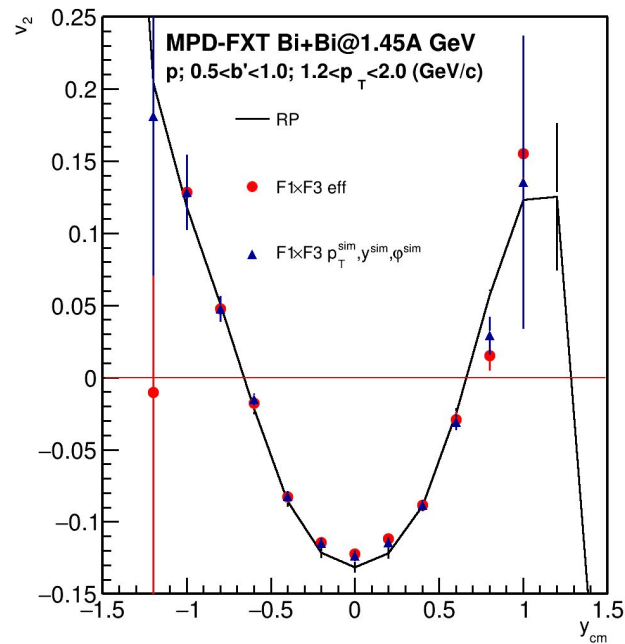
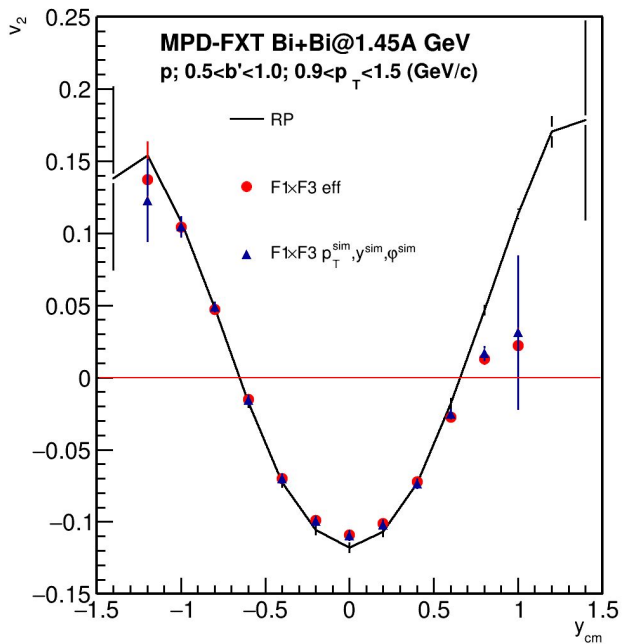
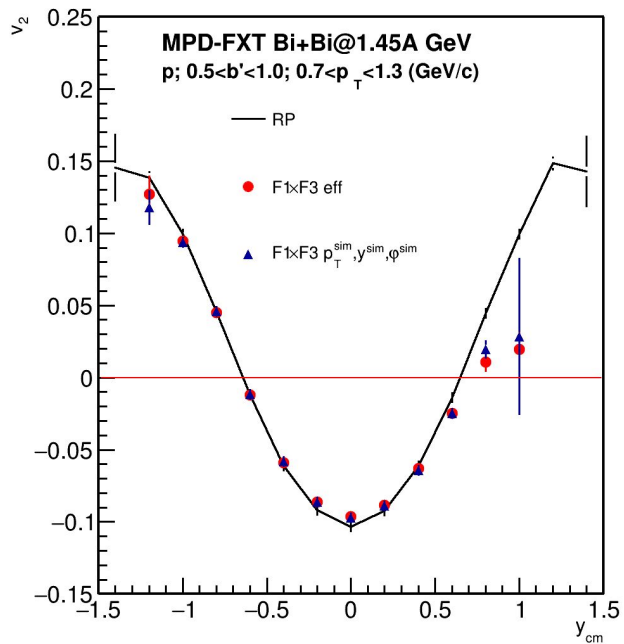
Efficiency corrections have no significant influence

$v_2(y)$ of protons



Efficiency introduces no significant difference

$v_2(y)$ of protons: acceptance vs. resolution

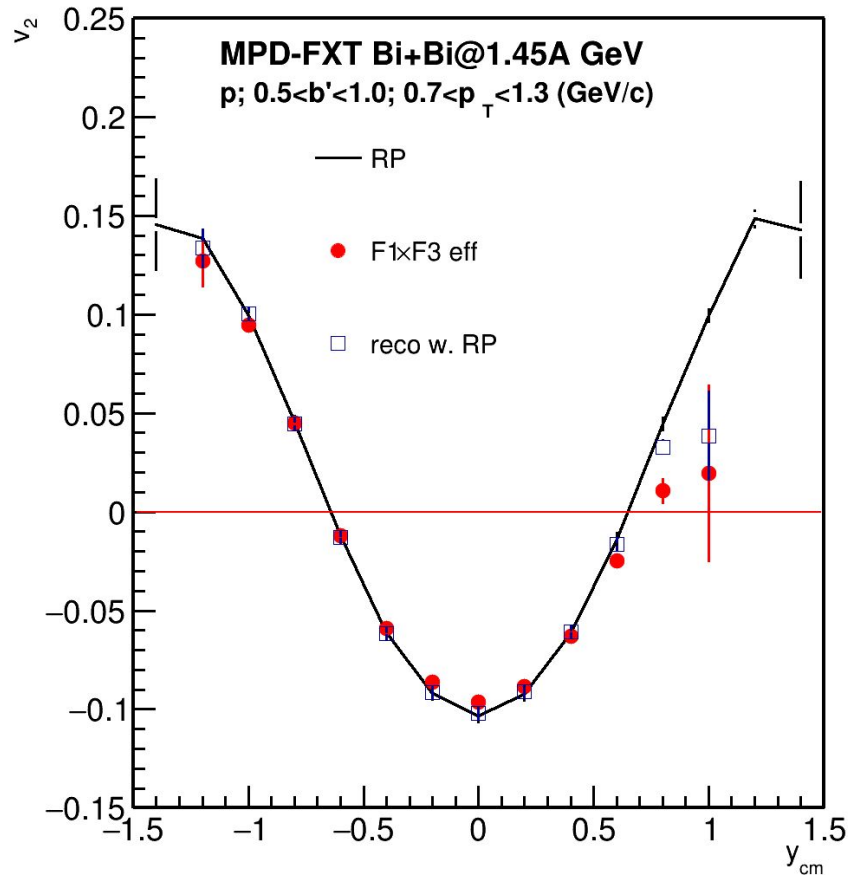


Need more statistics to check

Difference is mostly due to acceptance.

Effects related to the (p_T, y, ϕ) -resolution are small.

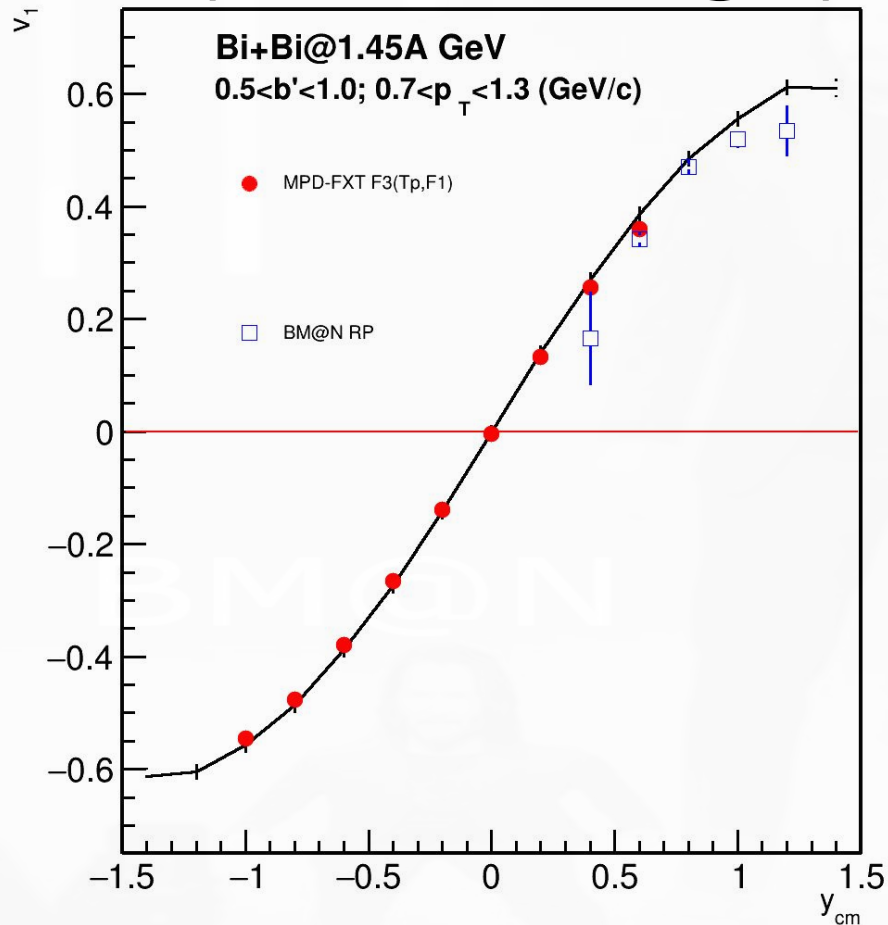
$v_2(y)$ of protons: effect of autocorrelation?



One additional source of the discrepancy might be due to autocorrelation caused by tracks that fall into FHCAL: they are both in u- and Q-vectors

Comparison with “reco” results w.r.t. RP suggests that might happen

Comparison with BM@N performance



BM@N TOF system (TOF-400 and TOF-700) has poor midrapidity coverage at $\sqrt{s_{NN}} = 2.5$ GeV

- One needs to check higher energies ($\sqrt{s_{NN}} = 3, 3.5$ GeV)
- More statistics are required due to the effects of magnetic field in BM@N:
 - Only “yy” component of $\langle uQ \rangle$ and $\langle QQ \rangle$ correlation can be used

Despite the challenges, both MPD-FXT and BM@N can be used in v_n measurements:

- To widen rapidity coverage
- To perform a cross-check in the future

Summary

- Optimal cuts for tracks were provided:
 - Protons
 - Nhits>27
 - DCA<1 cm
- Good agreement between “reco” and “mc” within corresponding acceptance window
- Discrepancy between “reco” and “mc” at forward rapidity:
 - Comparison with associated mc track shows that non-zero (p_T, y, ϕ) -resolution has small effect on resulted v_n - difference is due to acceptance
 - Possible contribution to the discrepancy from the tracks that fall into FHCAL - in progress
- Comparison with BM@N for Bi+Bi at $\sqrt{s_{NN}} = 2.5$ GeV:
 - BM@N TOF acceptance has poor midrapidity coverage at $\sqrt{s_{NN}} = 2.5$ GeV
 - Both MPD-FXT and BM@N can be useful for the future flow studies at Nuclotron-NICA

Backup

MPD

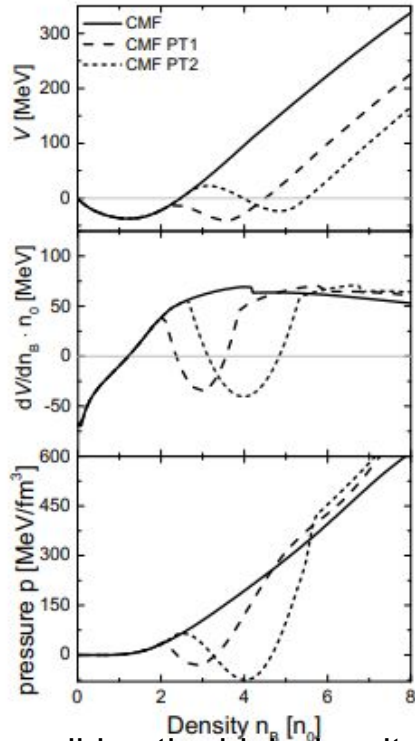
BM@N

v_n as a function of collision energy

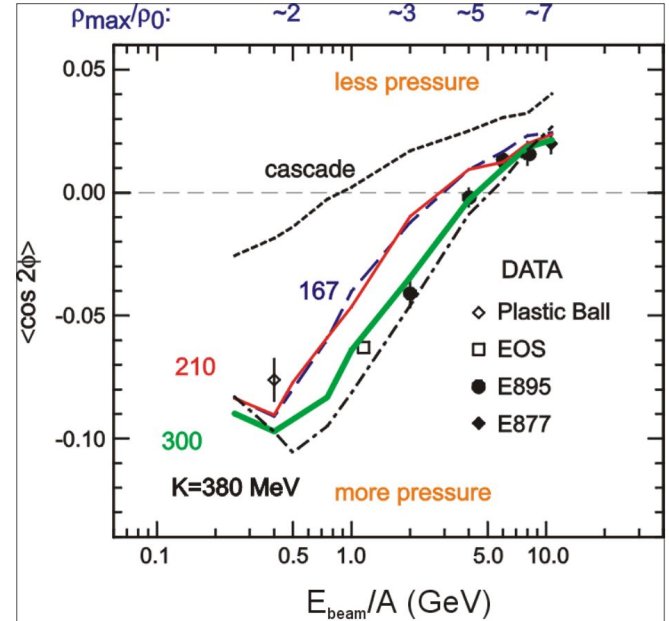
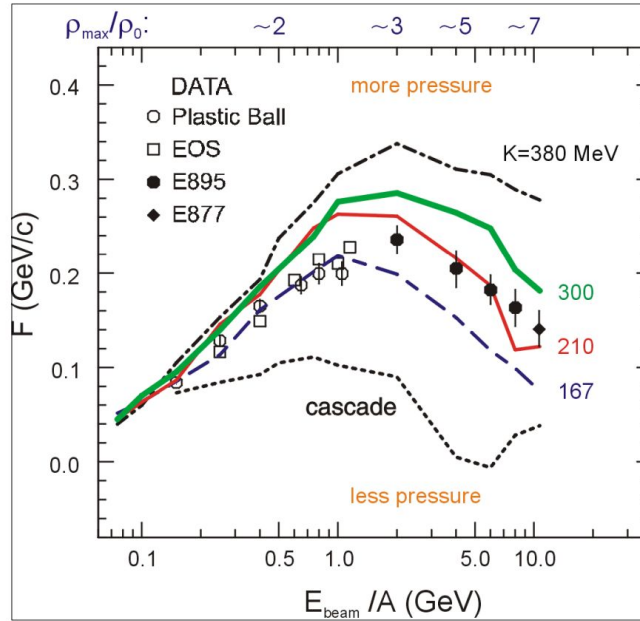
P. DANIELEWICZ, R. LACEY, W. LYNCH
[10.1126/science.1078070](https://doi.org/10.1126/science.1078070)

v_1 suggests softer EOS

v_2 suggests harder EOS



EPJ Web of Conferences 276, 01021 (2023)

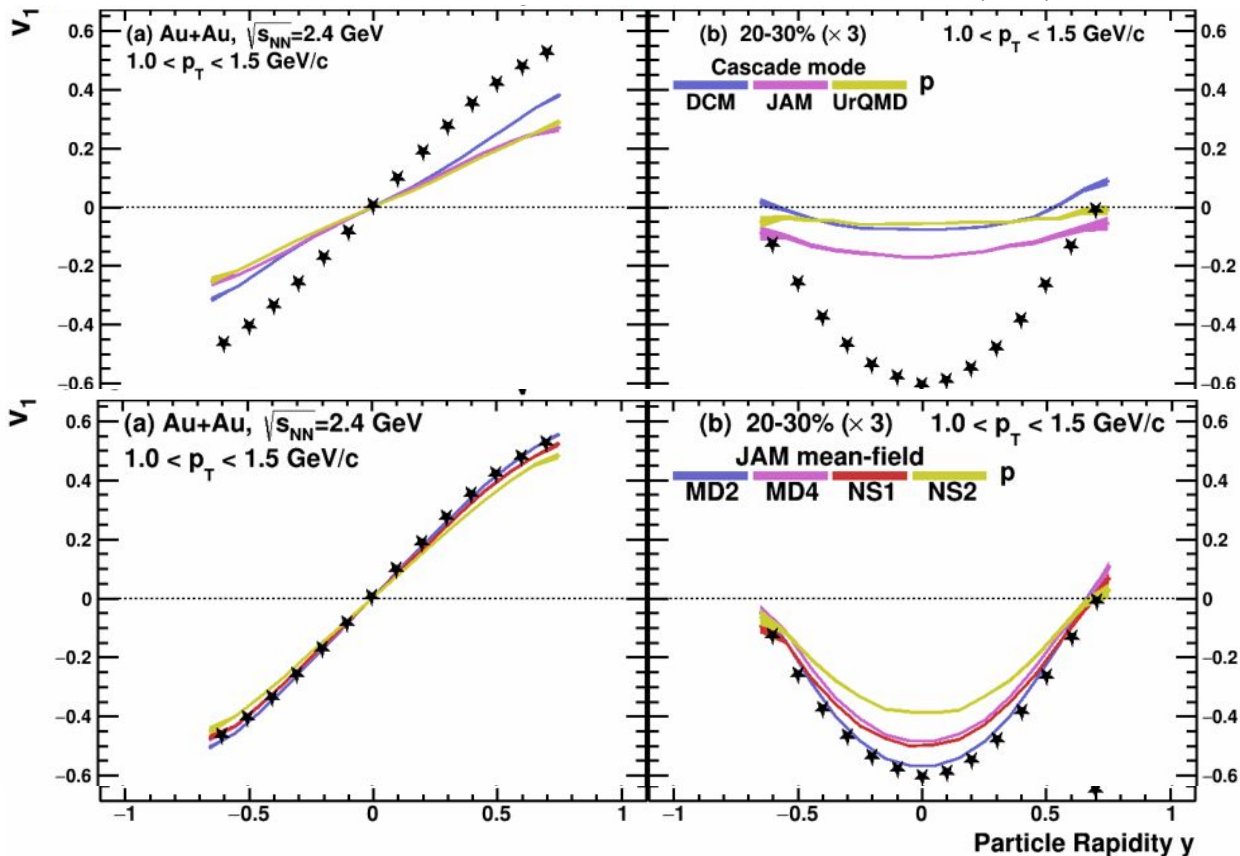


Describing the high-density matter using the mean field
 Flow measurements constrain the mean field

Discrepancy is probably due to non-flow correlations

Selecting the model

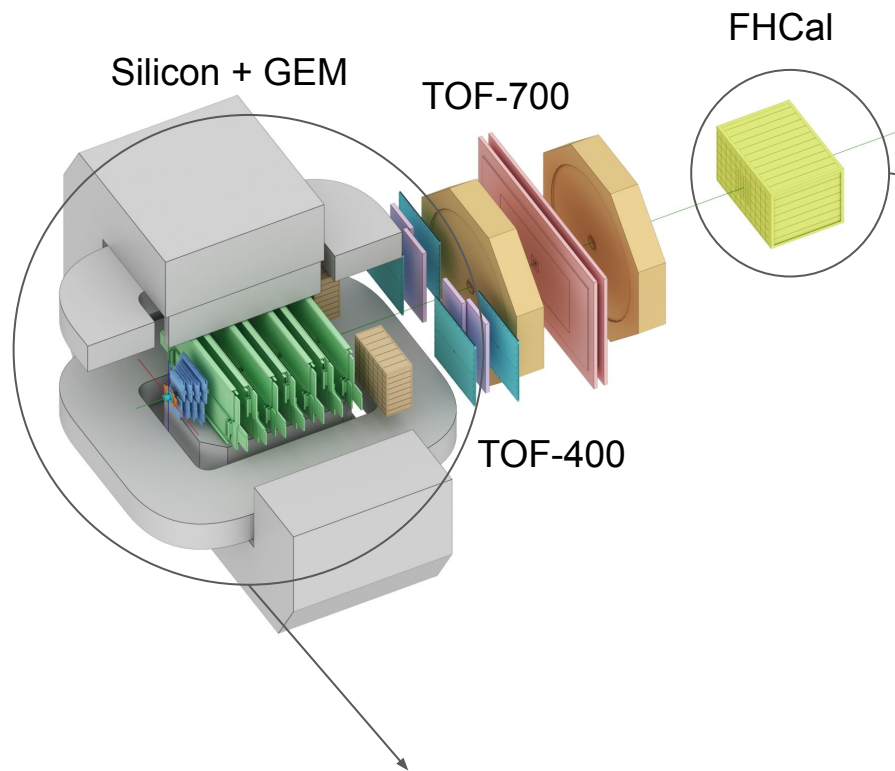
P.Parfenov Particles 5 (2022) 4, 561-579



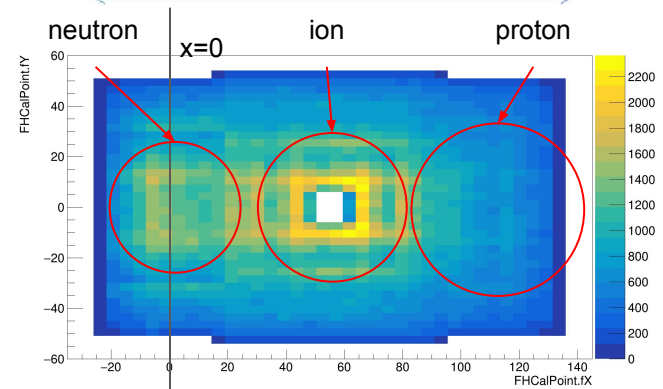
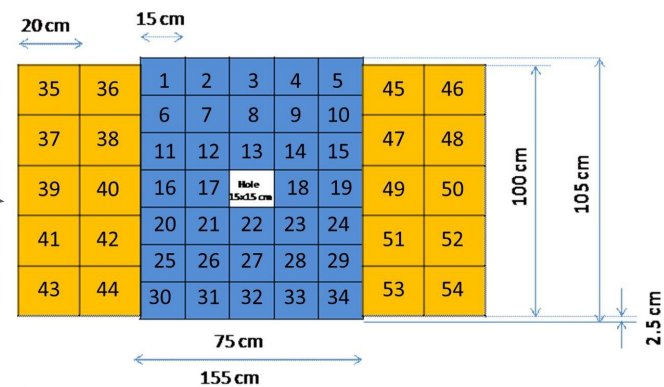
Cascade models fail to reproduce v_n at low-energy heavy-ion collision

Mean field models reproduce the v_n rather well

The BM@N experiment (GEANT4 simulation for RUN8)



Square-like tracking system within the magnetic field deflecting particles along X-axis



Charge splitting on the surface of the FHCAL is observed due to magnetic field