

Perspective cold fusion reactions for synthesis of superheavy nuclei

J.Hong, G.G.Adamian, N.V.Antonenko,
P.Jachimowicz, M.Kowal

What interesting fusion reactions can still be done with targets $^{204,206,207,208}\text{Pb}$?

- 1. Study of xn (neutron) evaporation channels with $x > 1$**
- 2. Production of new neutron-deficient isotopes and study of fissility vs isospin of SHN (Factory of SHN)**
- 3. Study of fusion (quasi-fission) probability**

Dynamics of fusion in the dinuclear system model

Evaporation residue cross section for the production of superheavy nuclei:

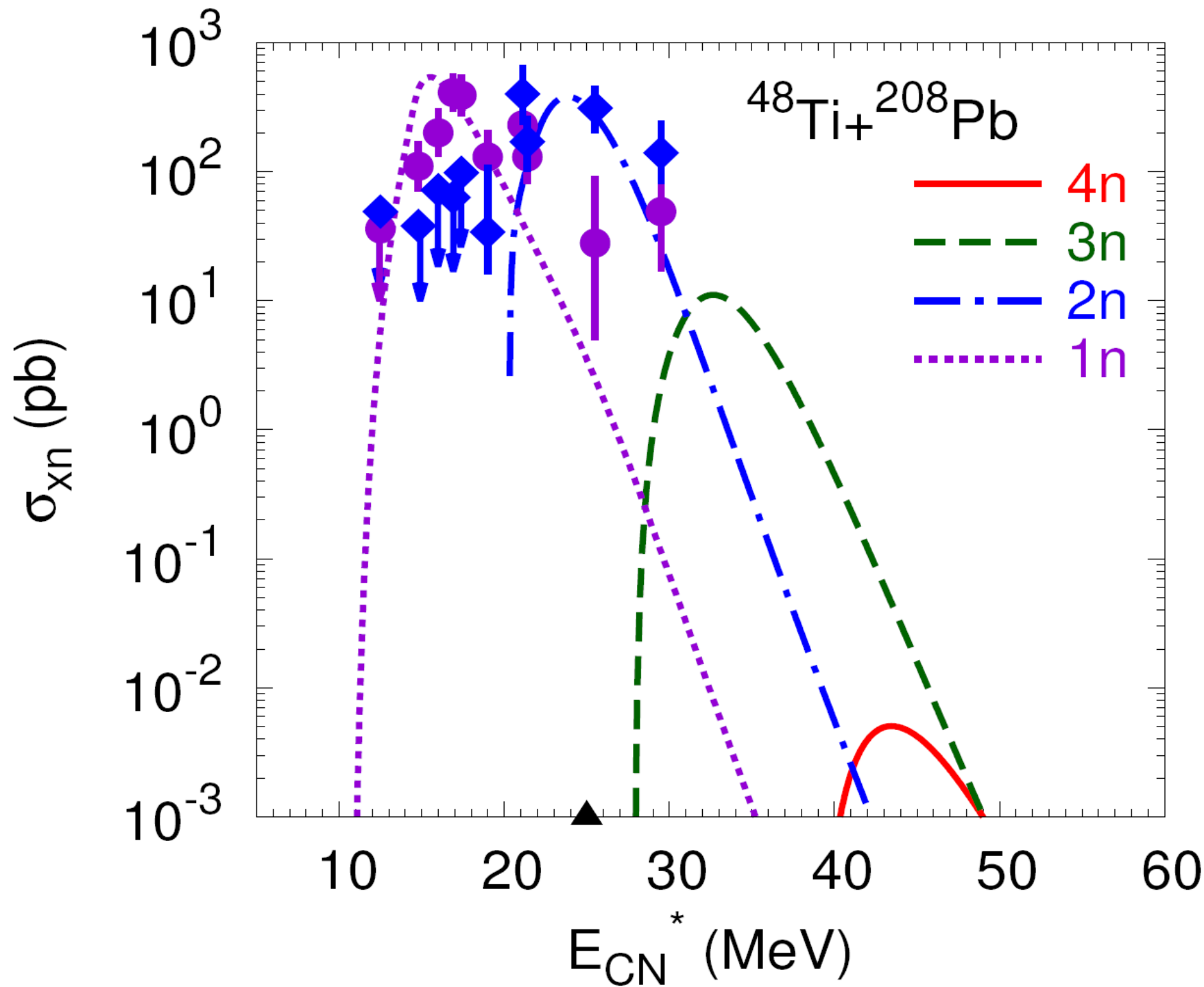
$$\sigma_{ER}^s(E_{c.m.}) = \sum_J \sigma_c(E_{c.m.}, J) P_{CN}(E_{c.m.}, J) W_{sur}^s(E_{c.m.}, J)$$

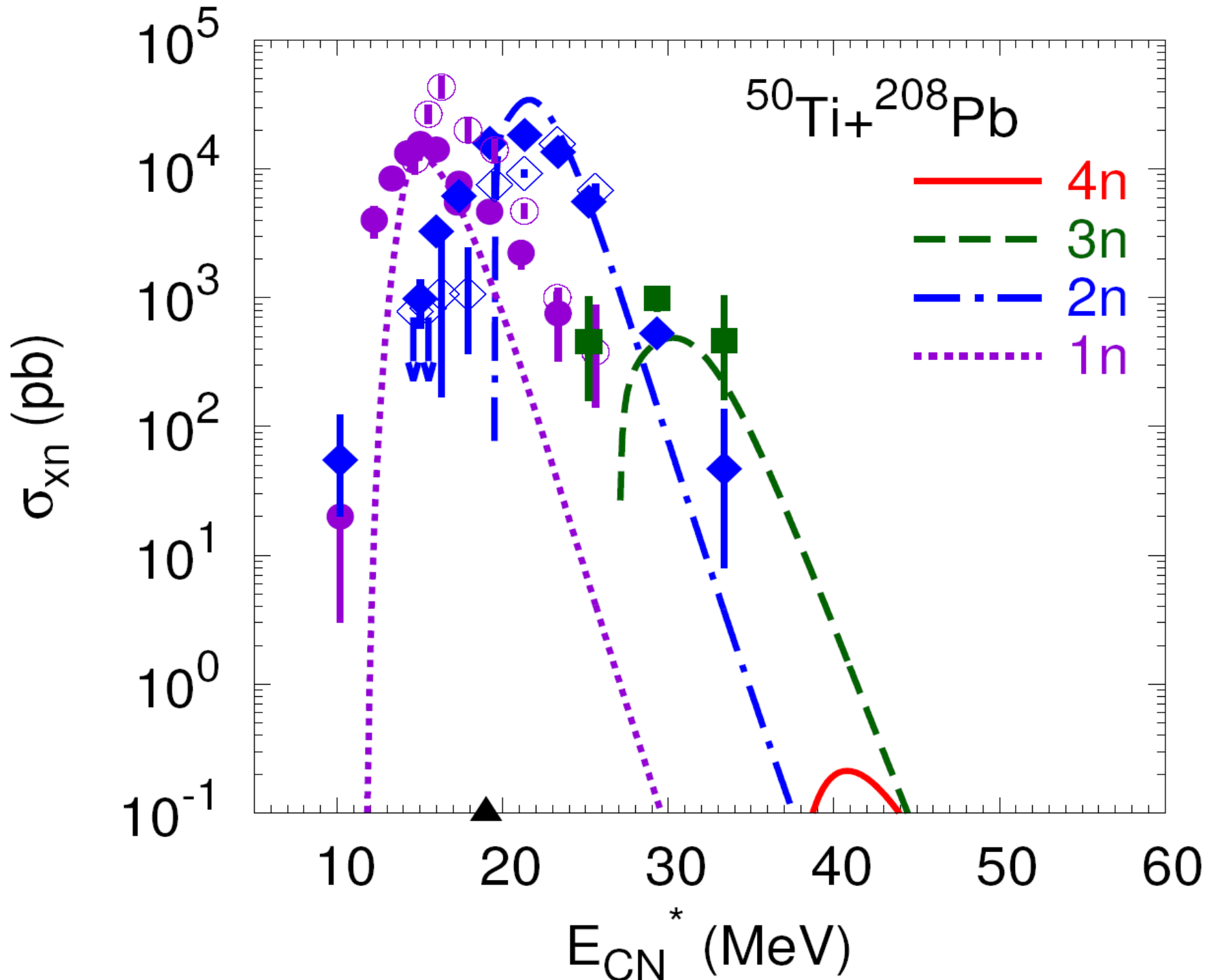
interplay of capture, fusion, survival probabilities

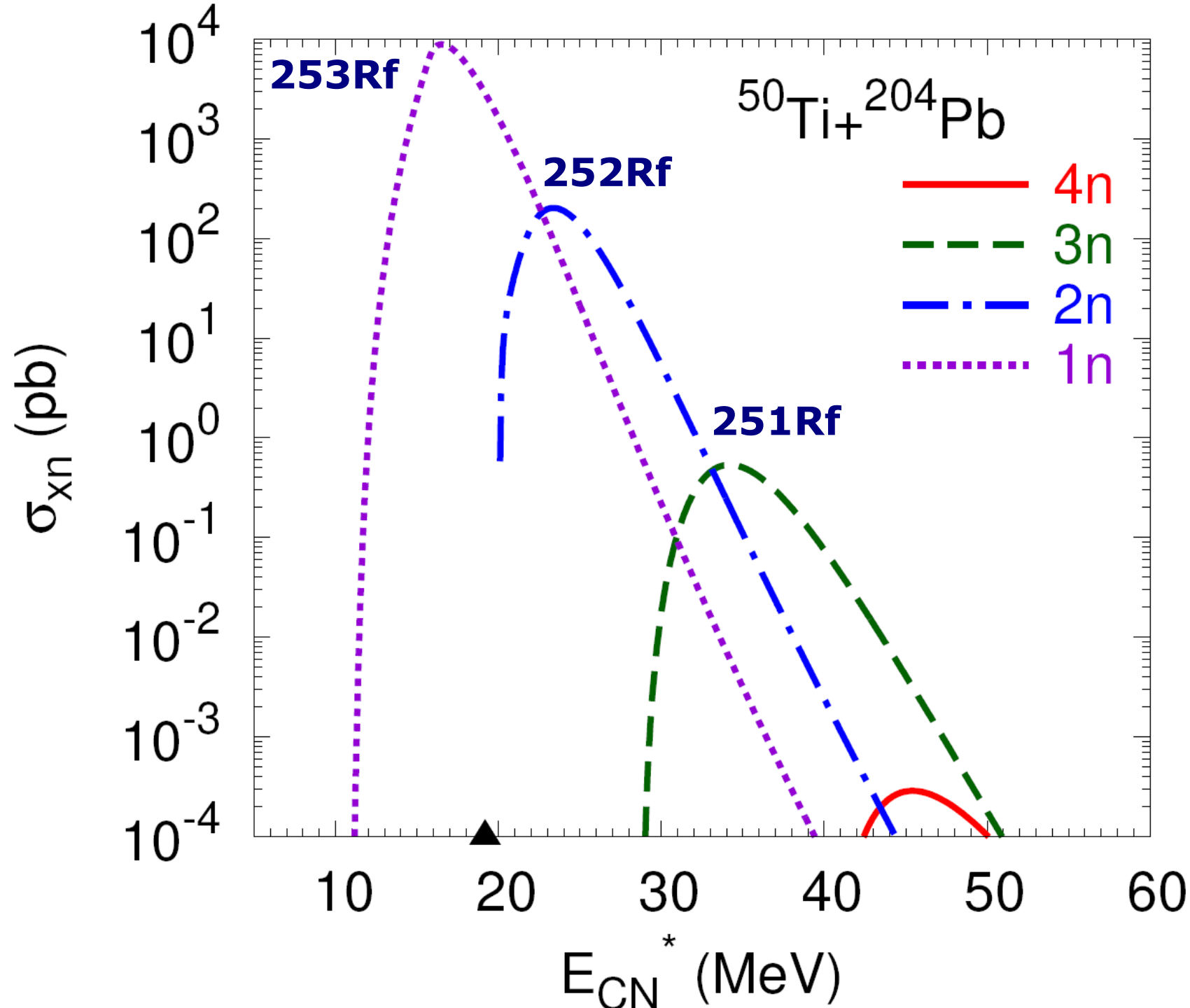
Q-value of fusion: 1n-channel is the sub-barrier in fusion with Ti, Cr, Fe

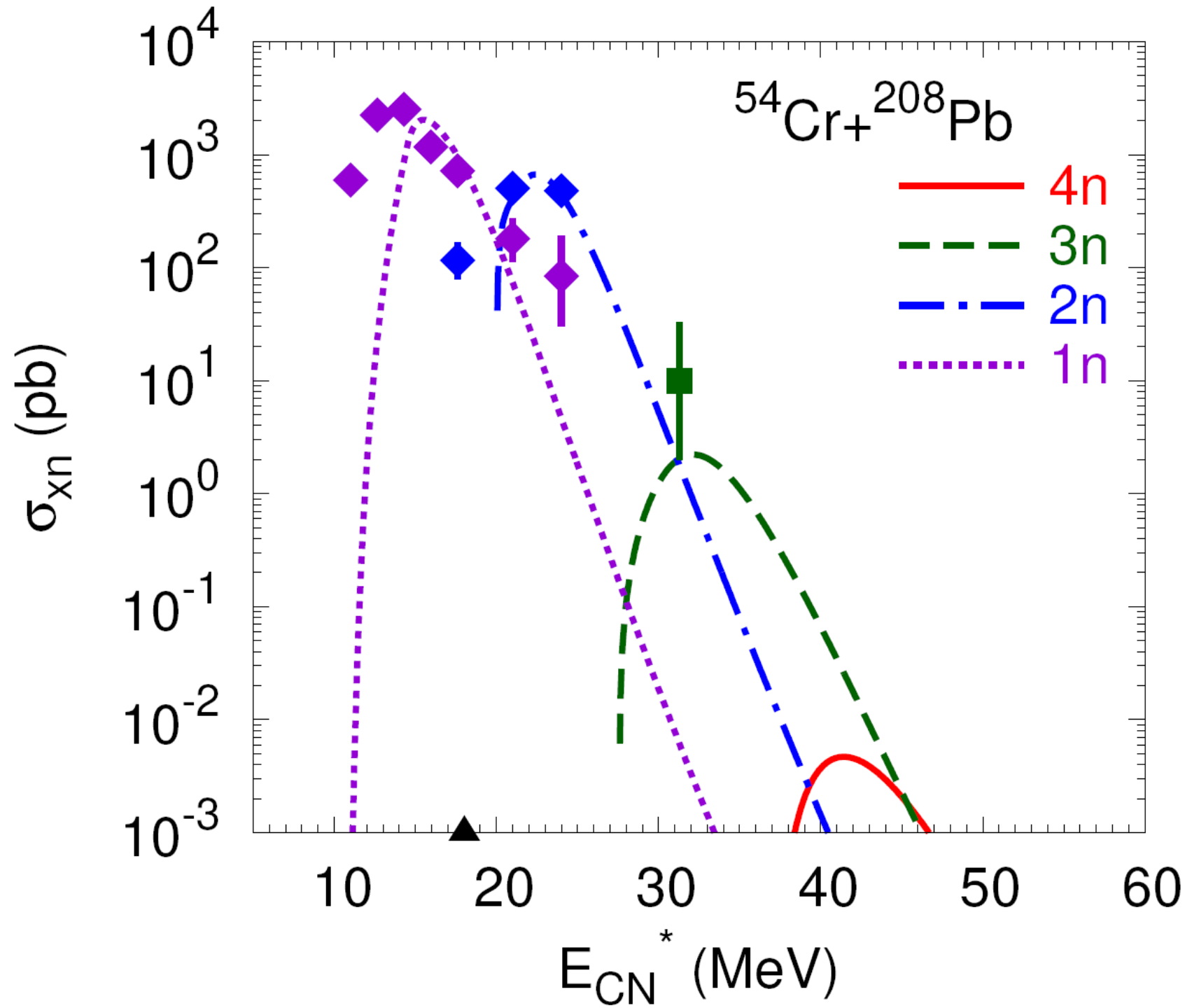
**Predictions of the properties of
heaviest nuclei are based on the
Macroscopic-Microscopic Model:**

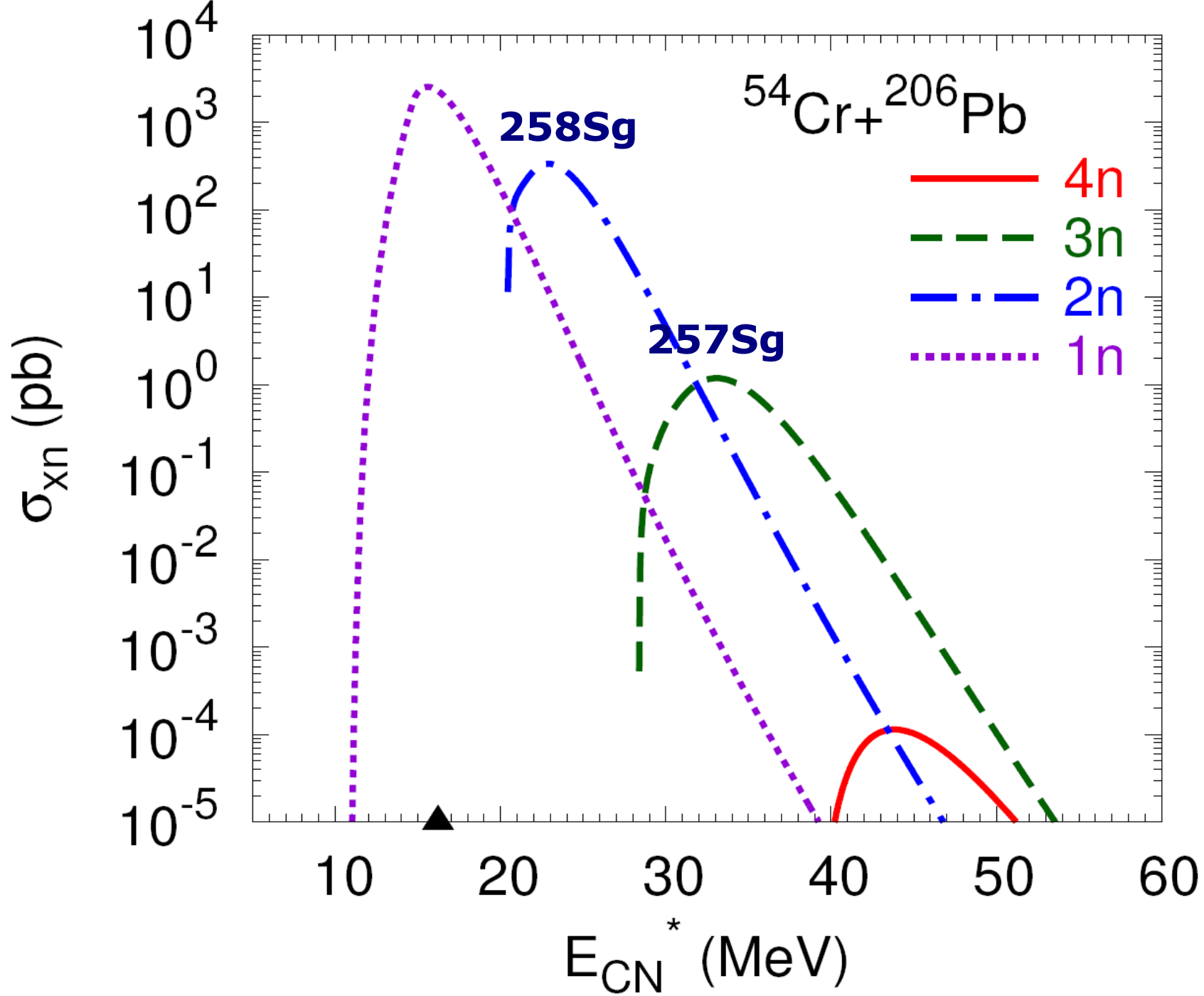
Mass Table by P.Jachimowicz, M.Kowal, J.Skalski,
At. Data Nucl. Data Tabl. **138** (2021) 101393

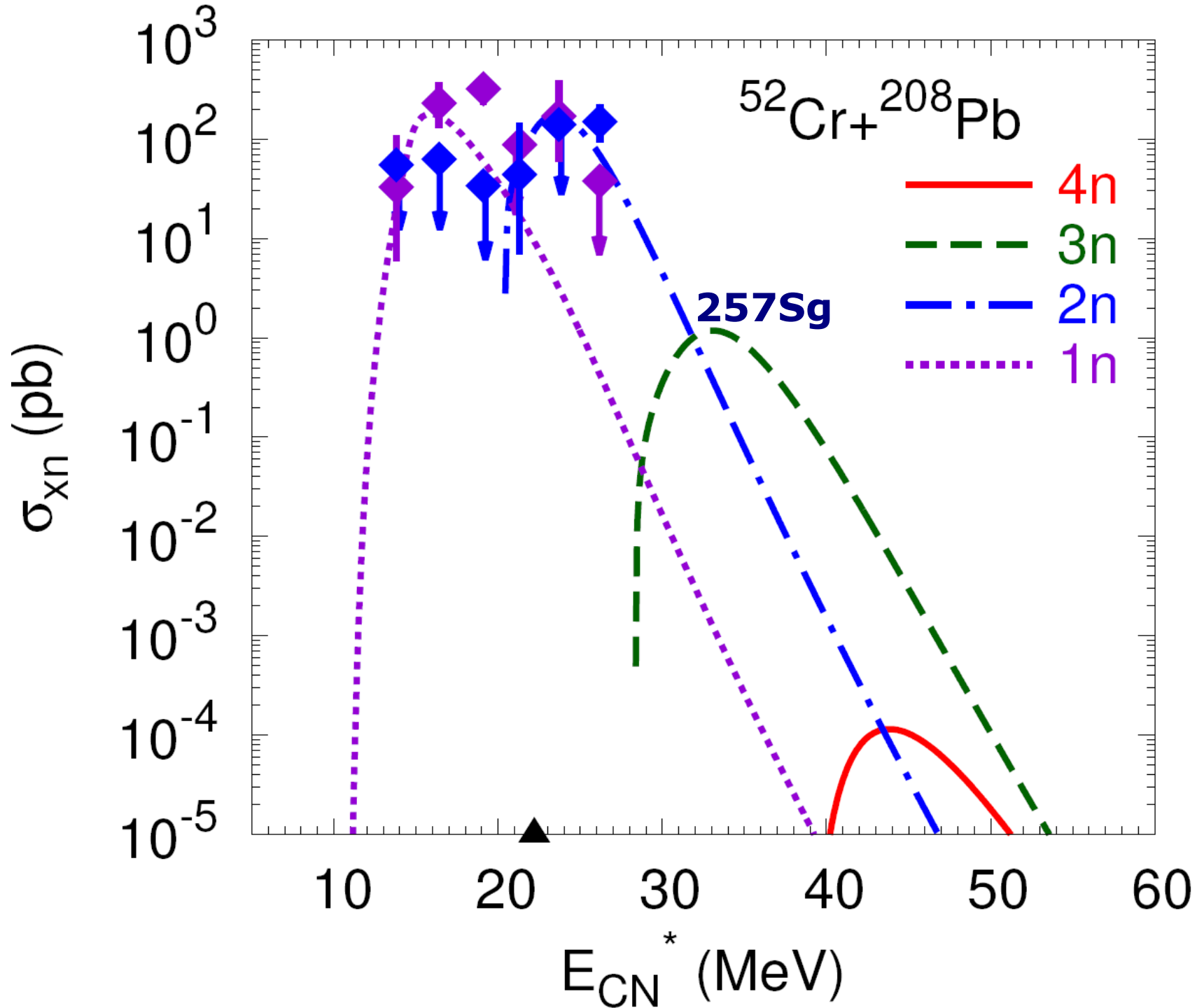


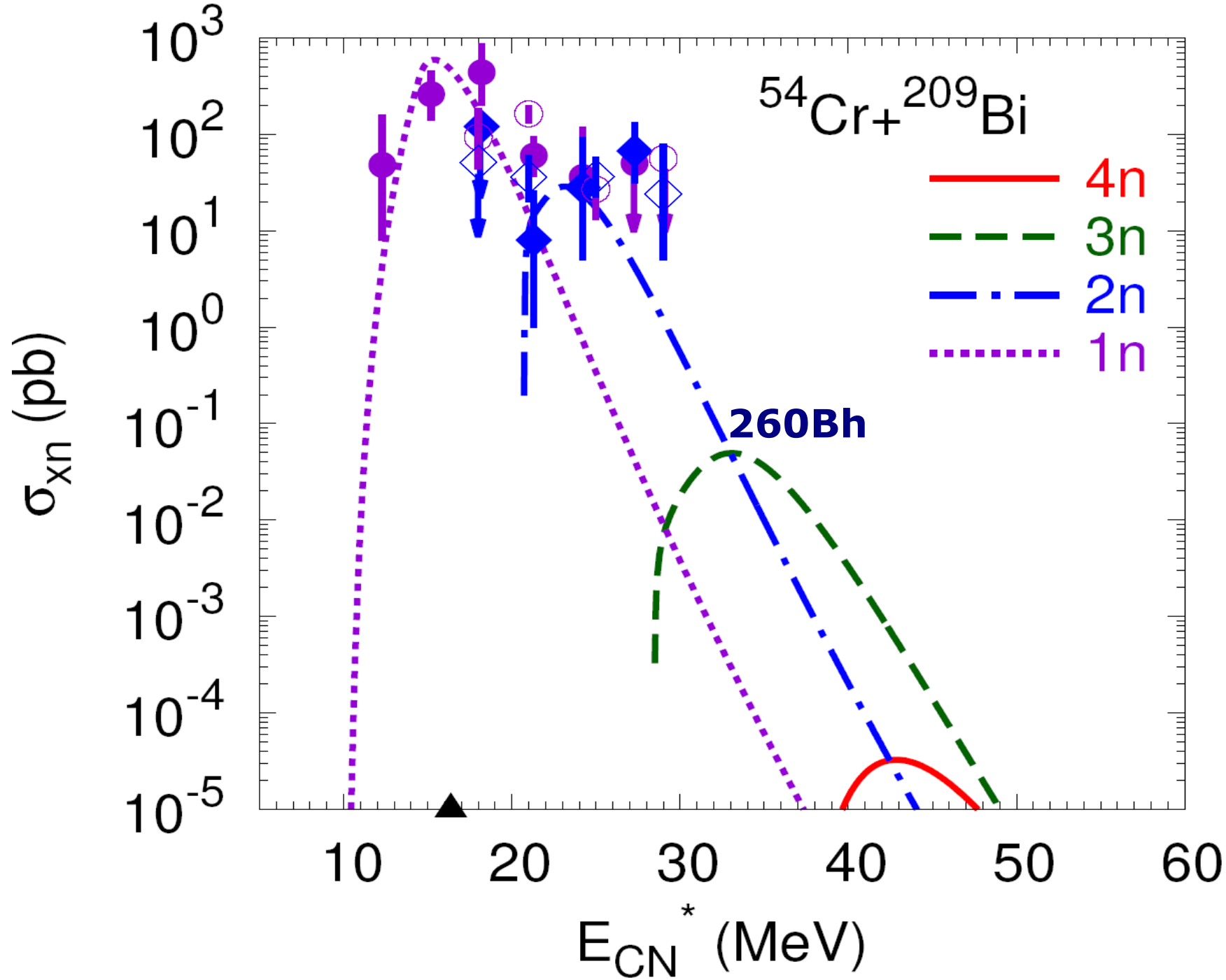


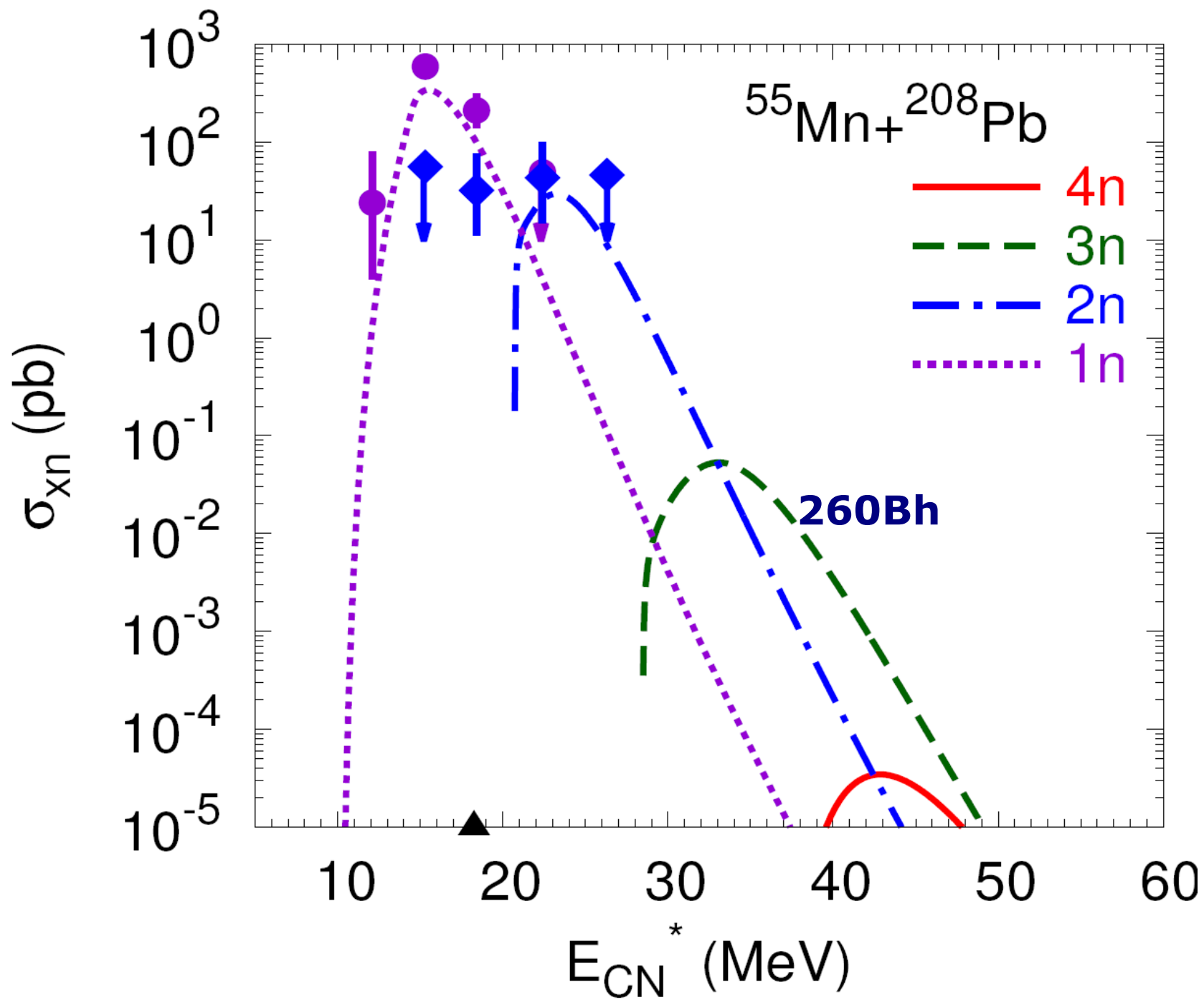




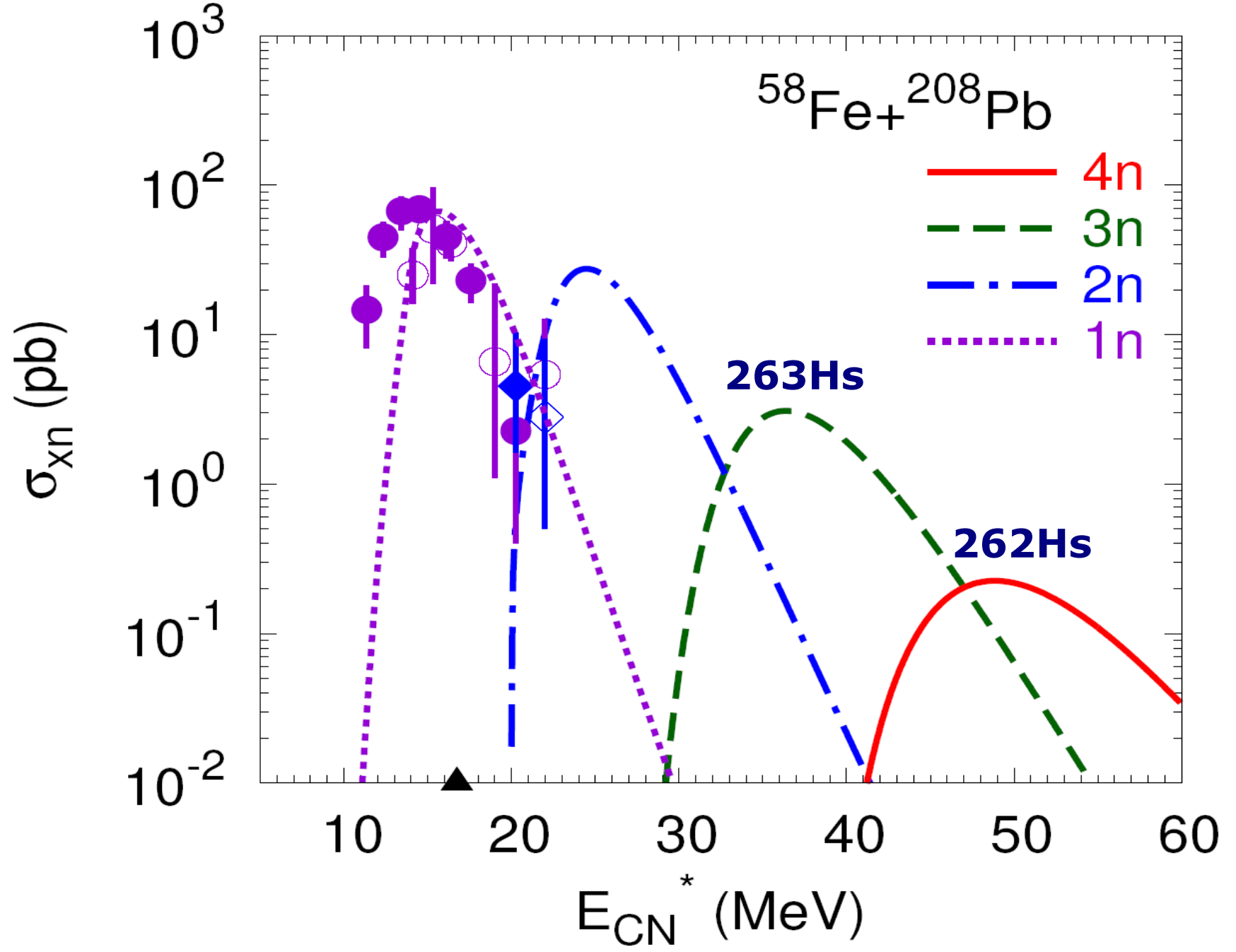


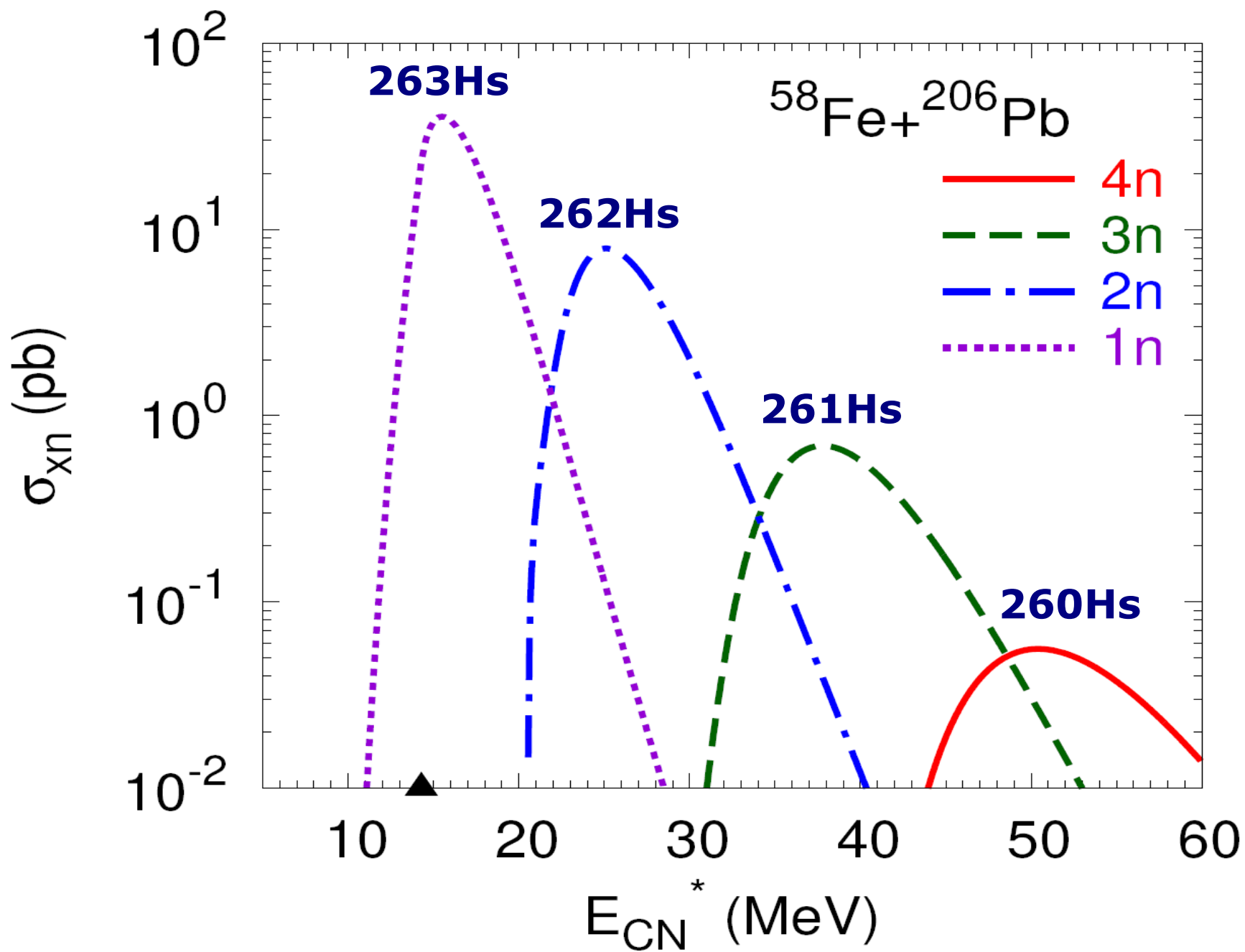


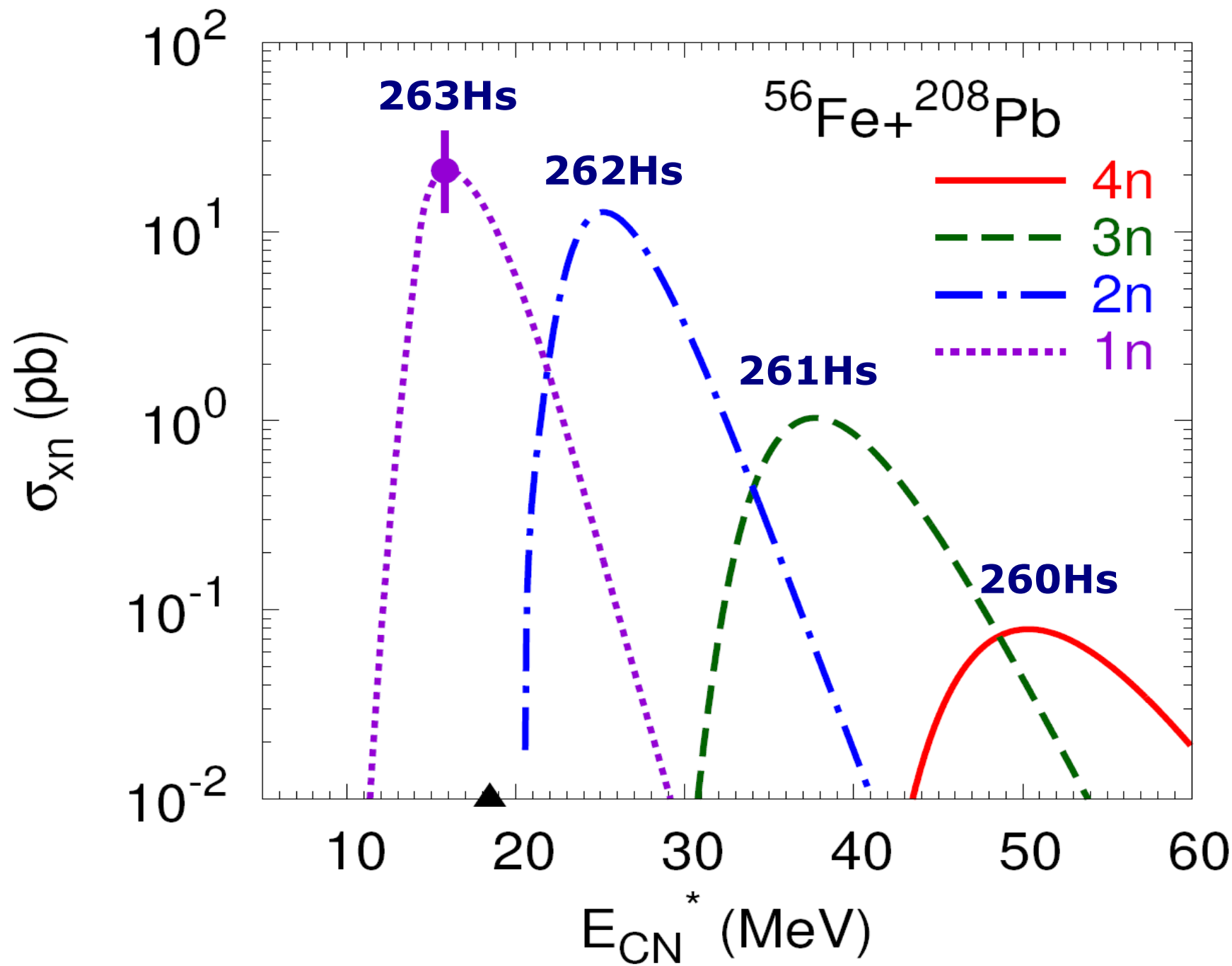


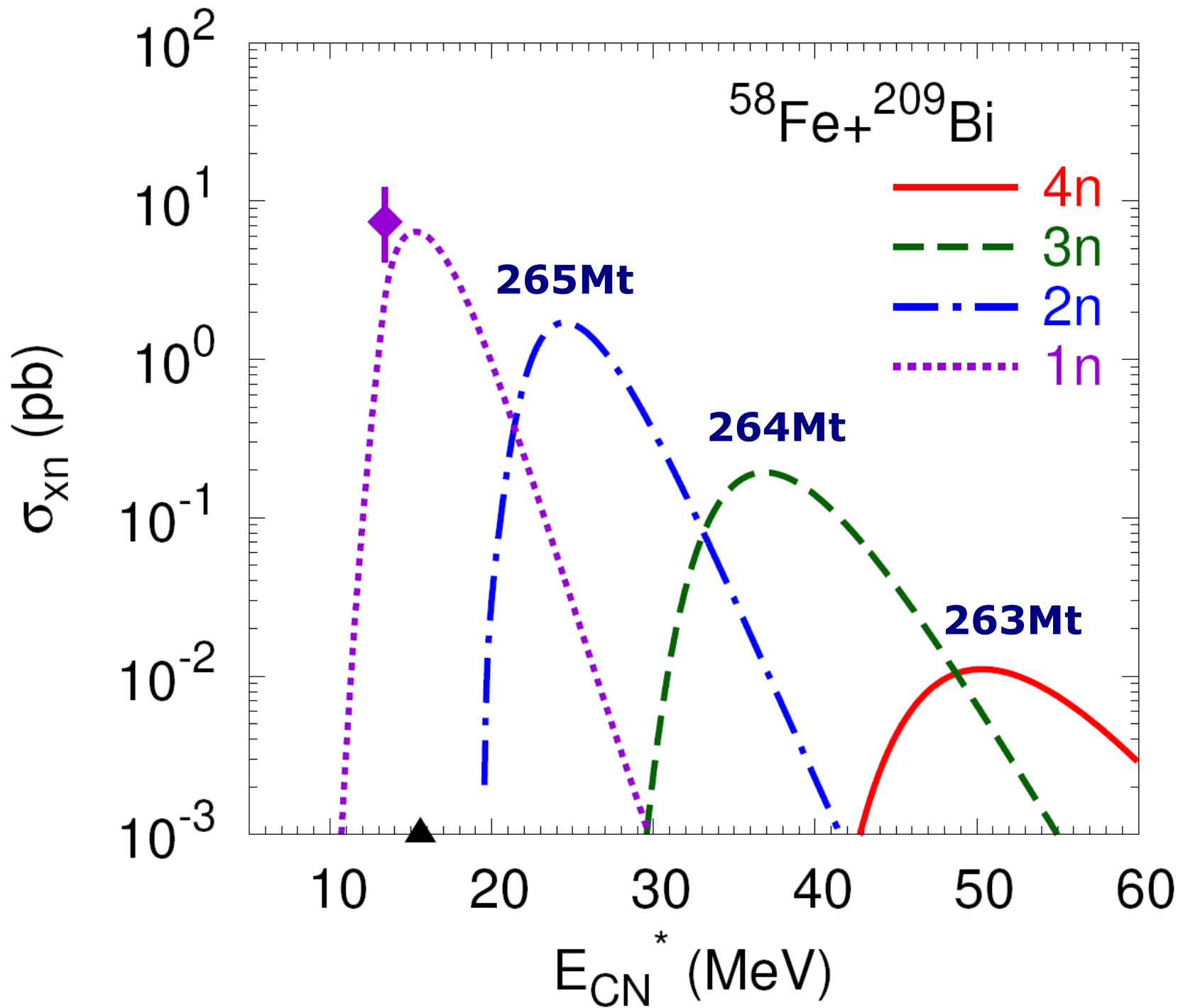


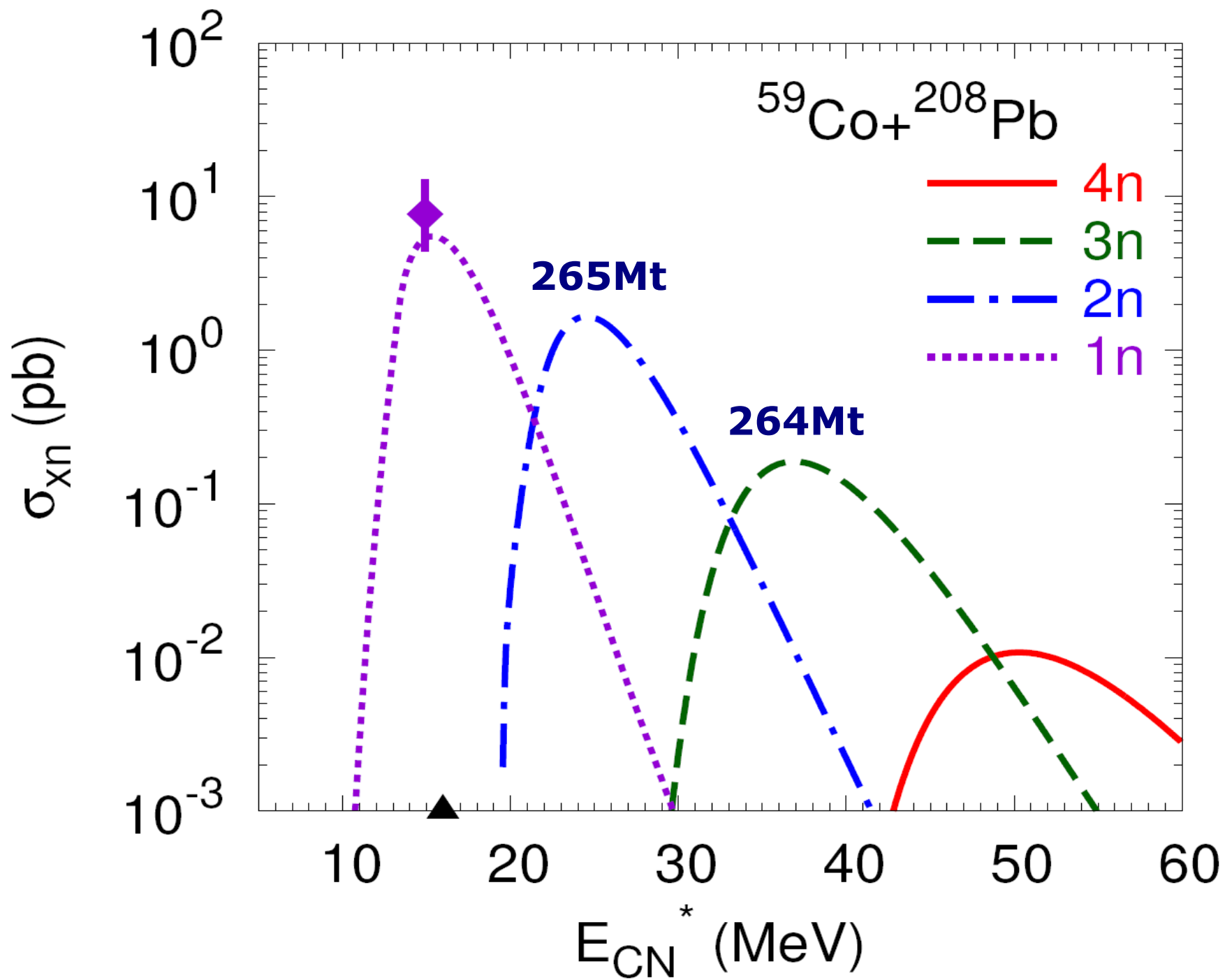
$^{58}\text{Fe} + ^{208}\text{Pb}$

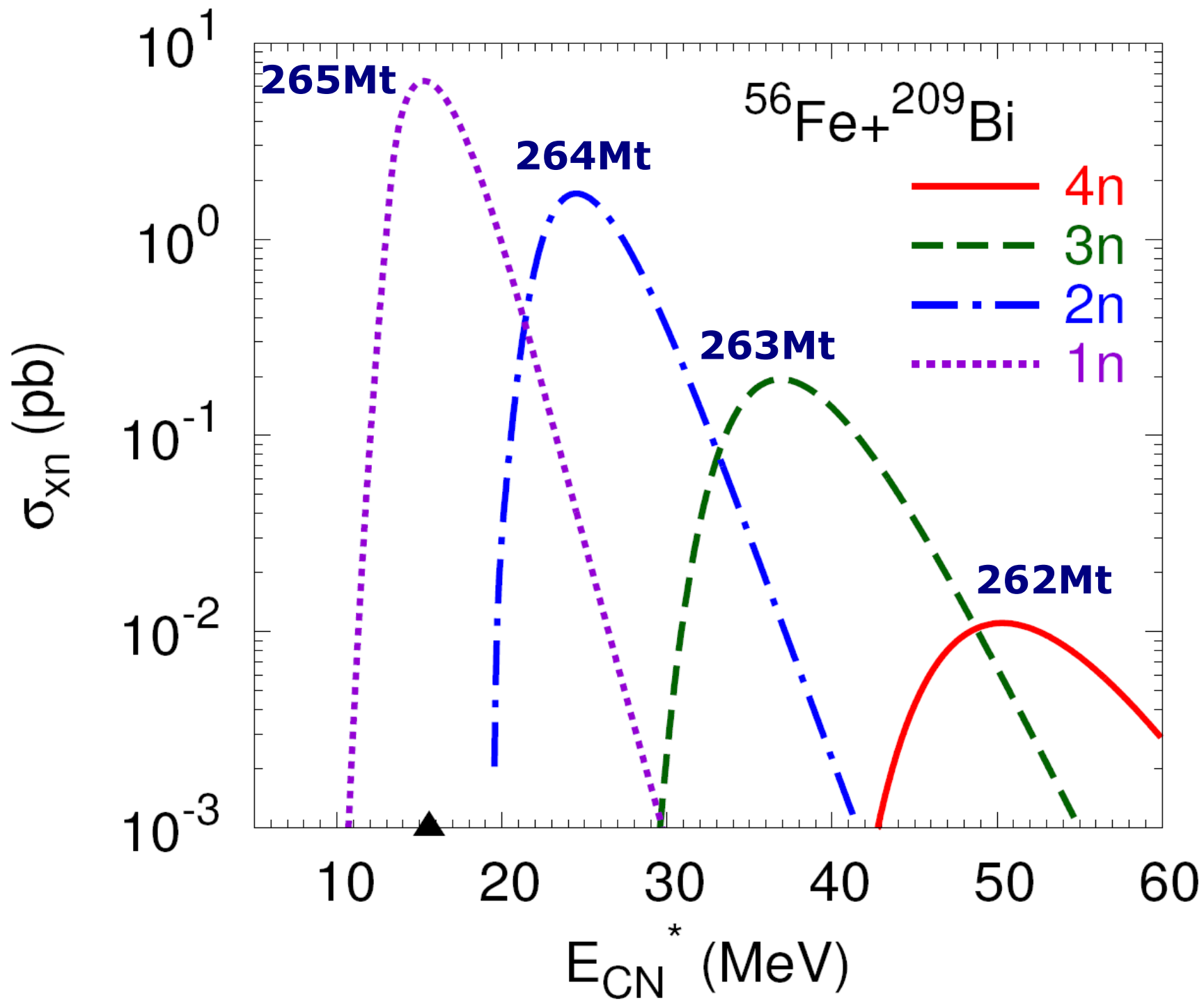




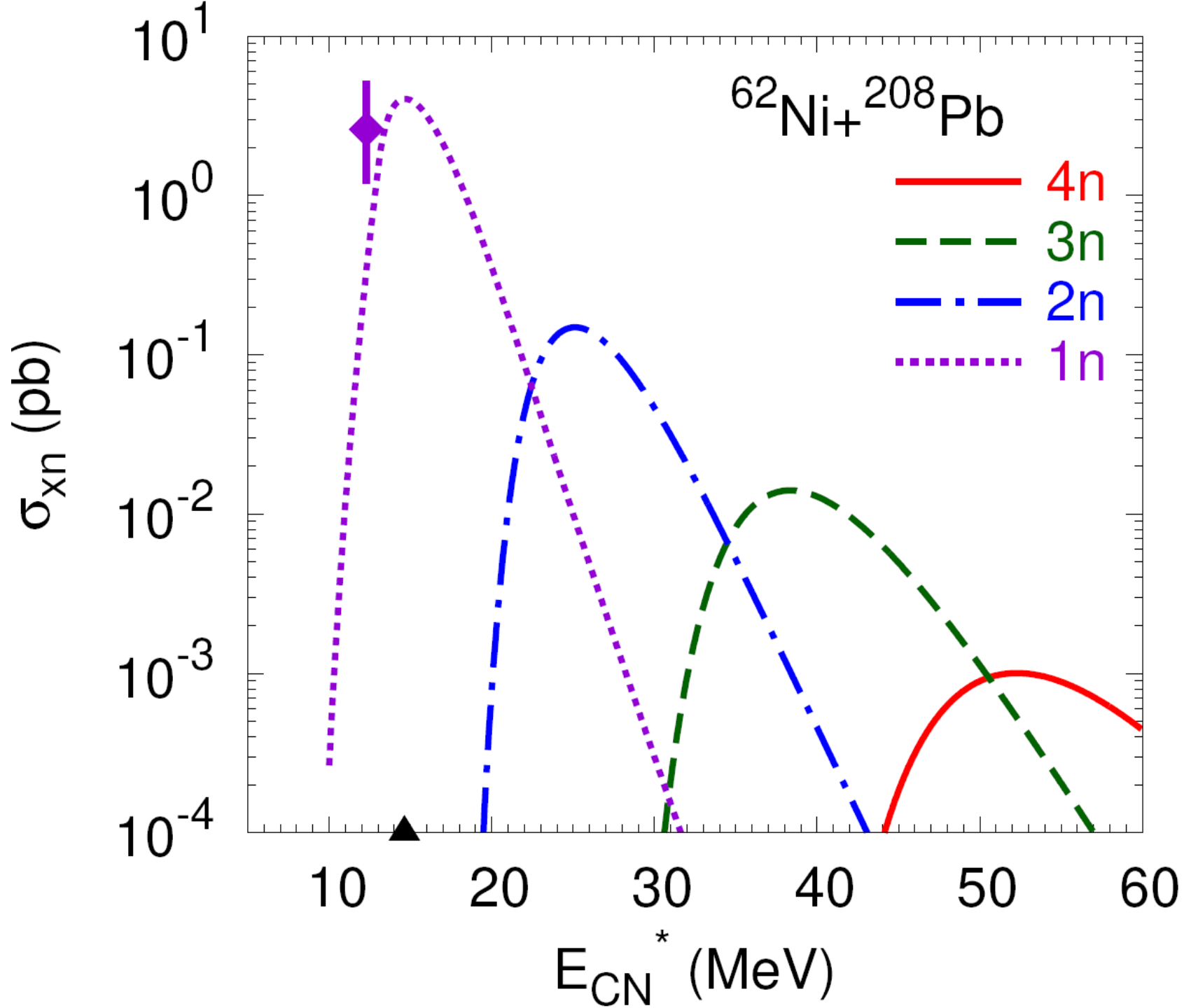


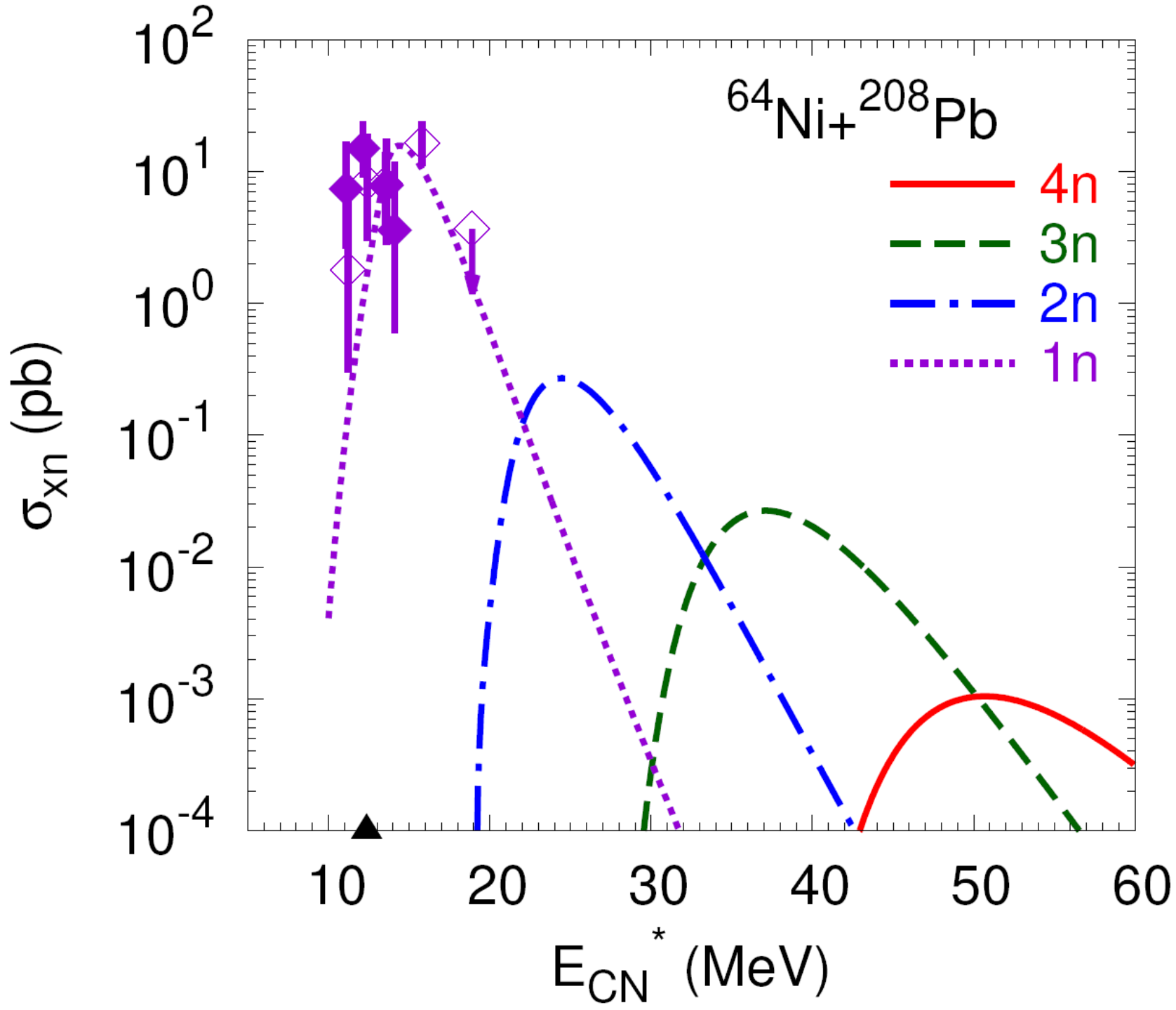




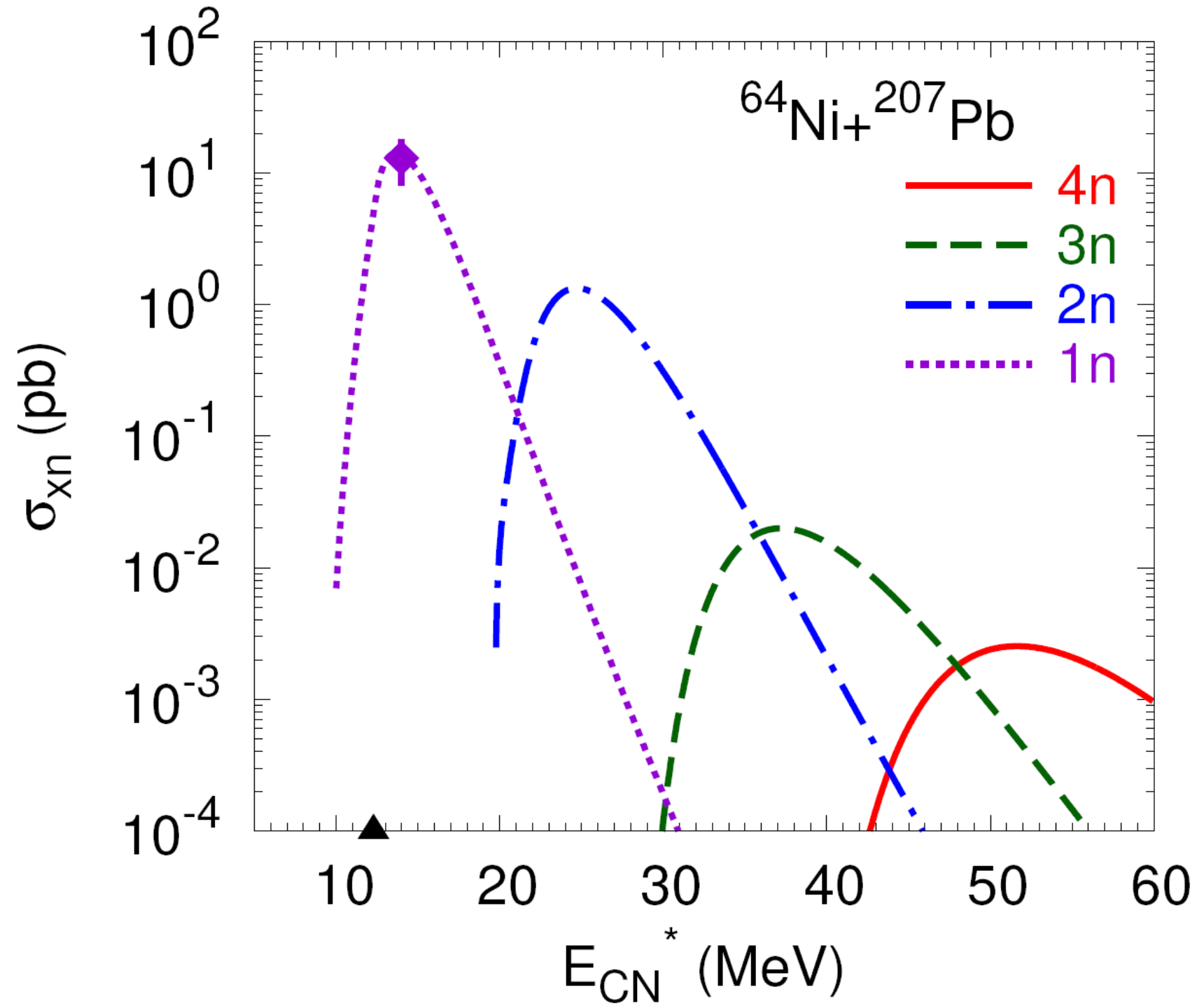


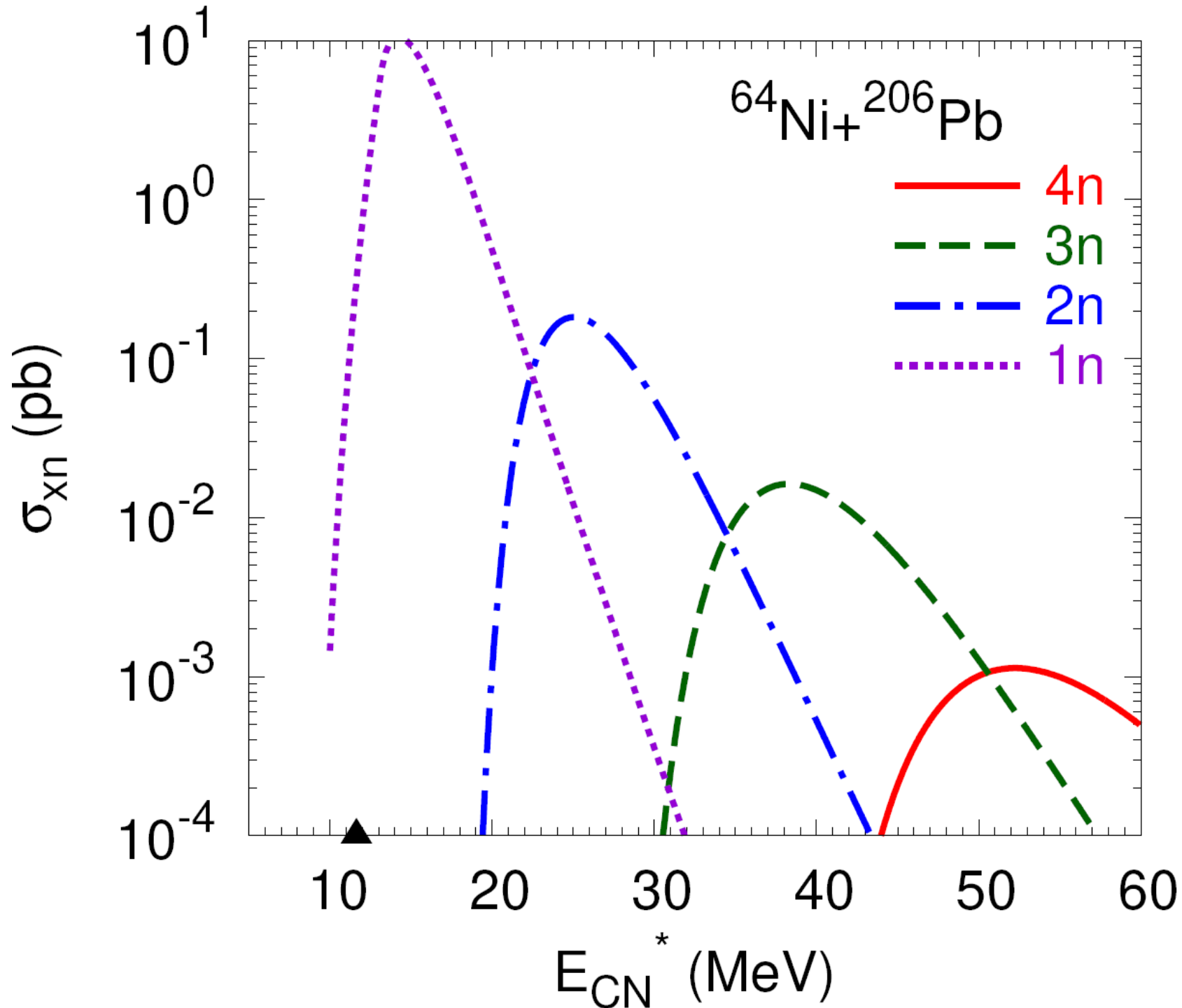
$^{62}\text{Ni} + ^{208}\text{Pb}$

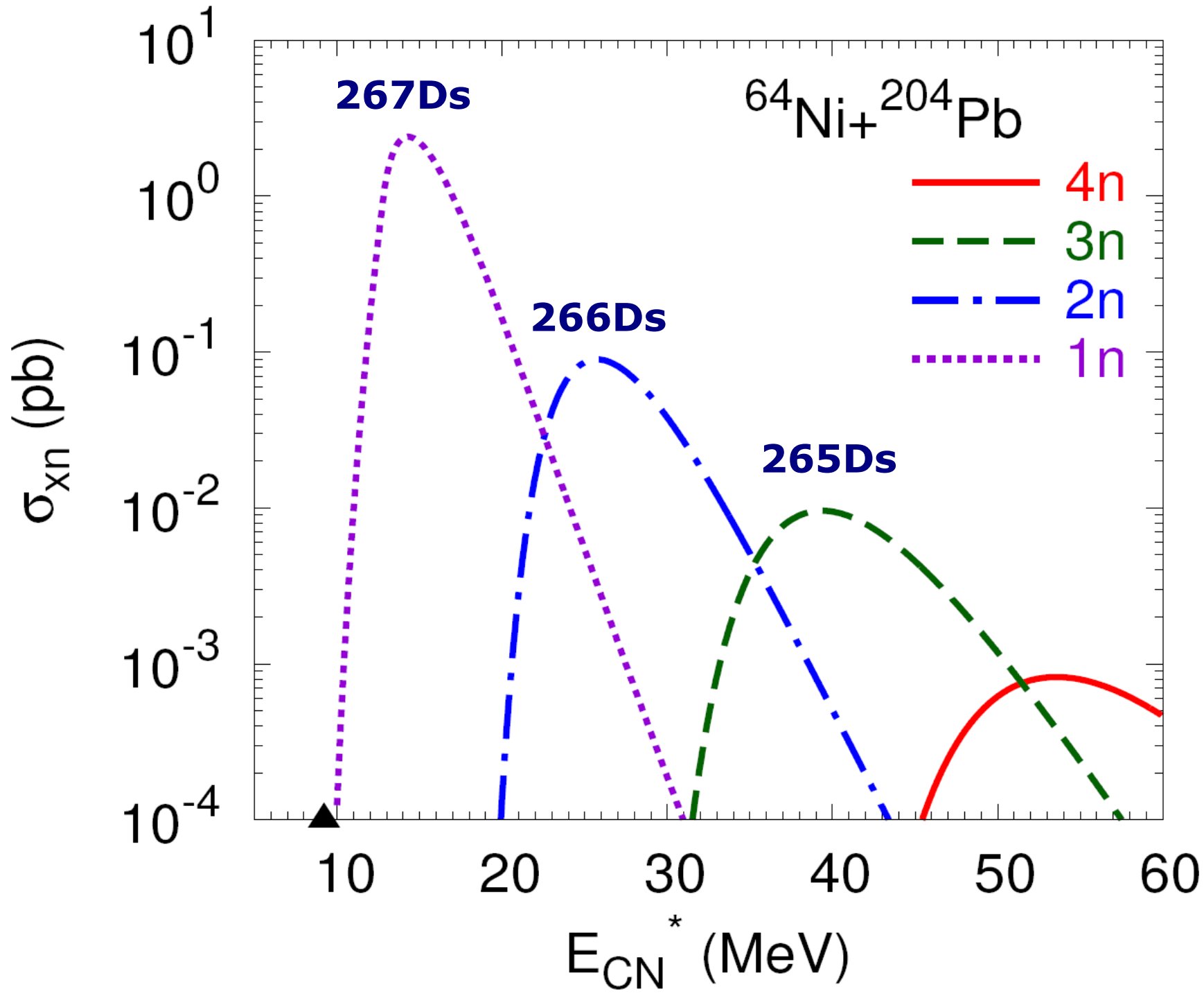




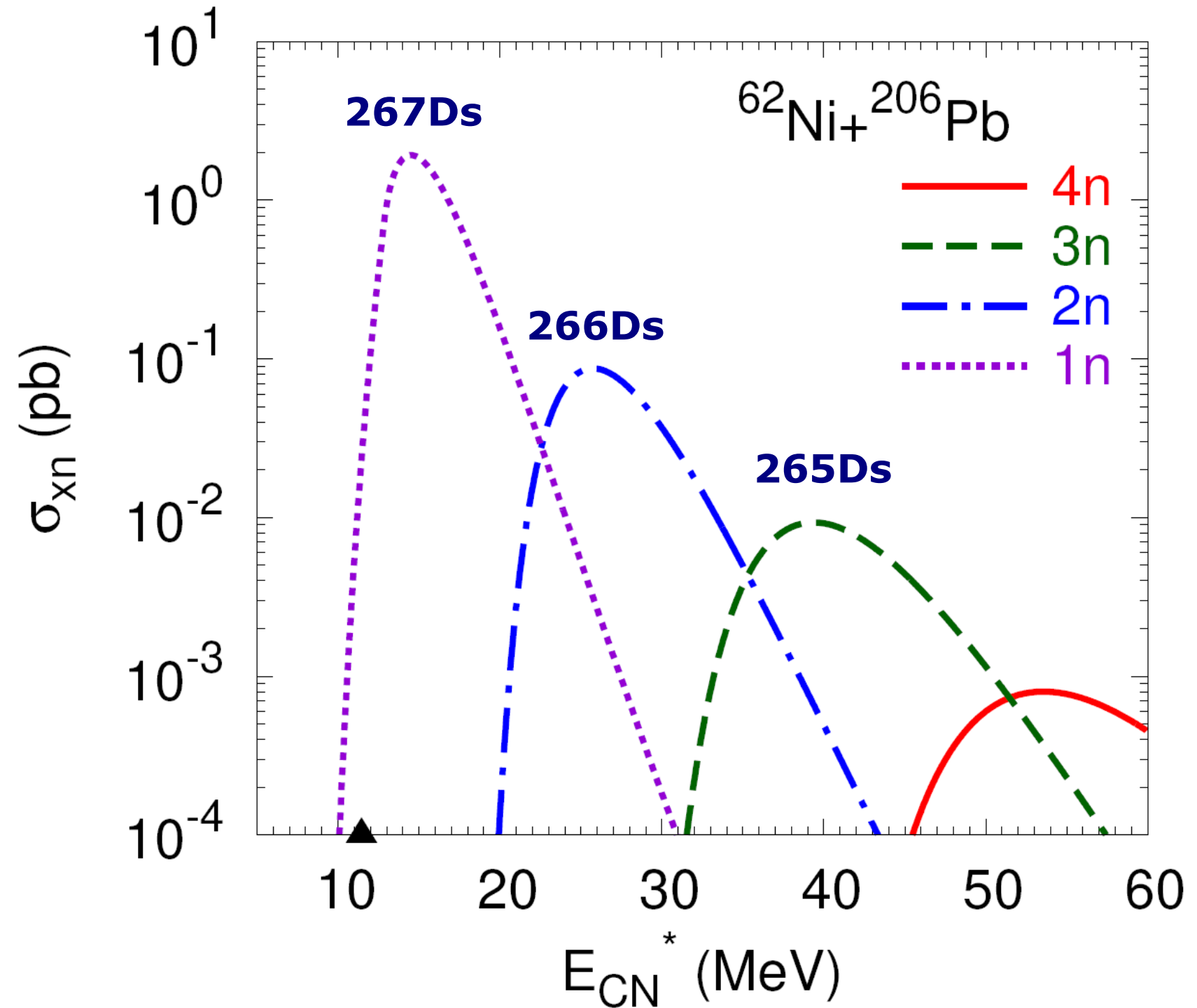
$^{64}\text{Ni} + ^{207}\text{Pb}$

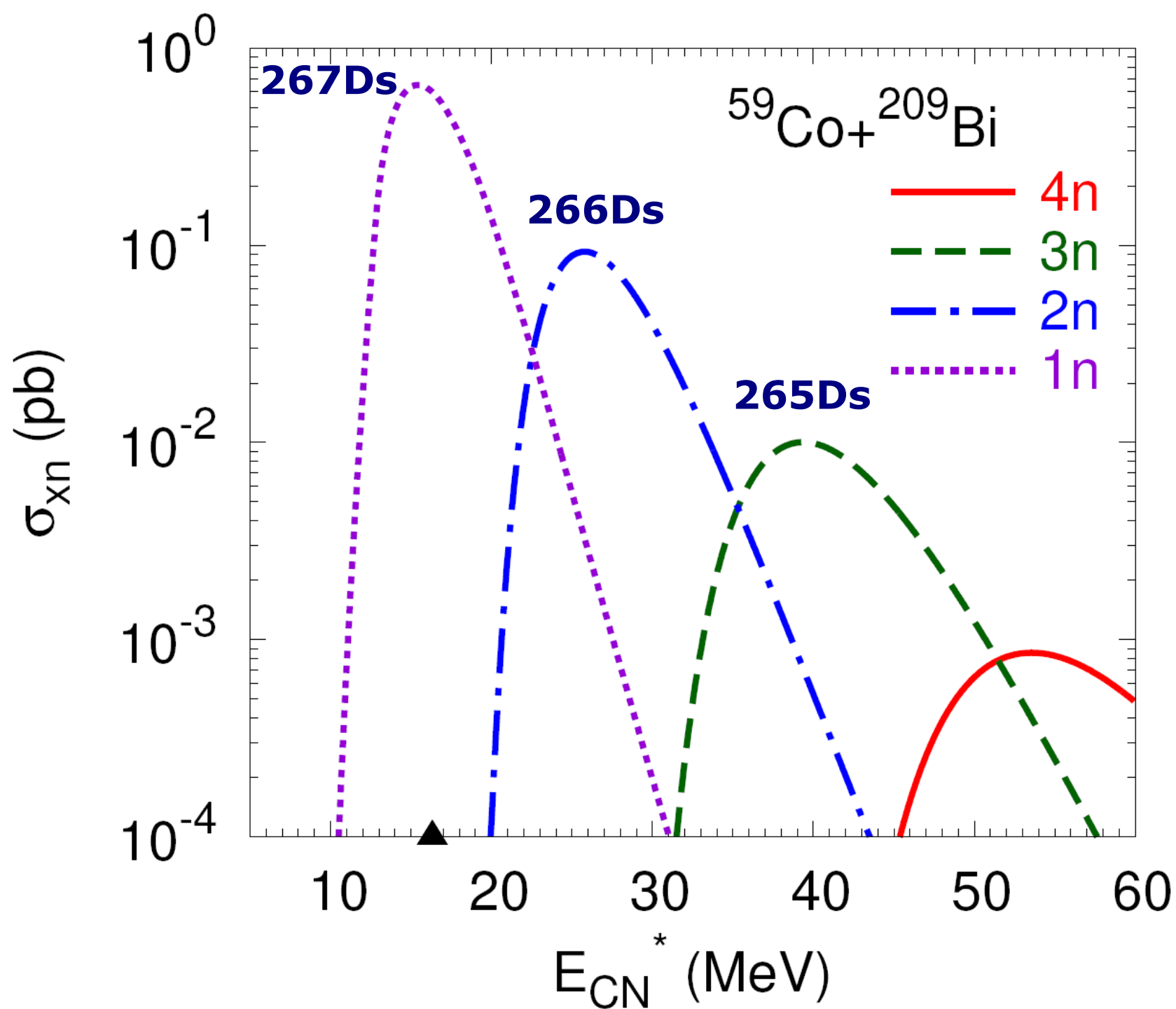


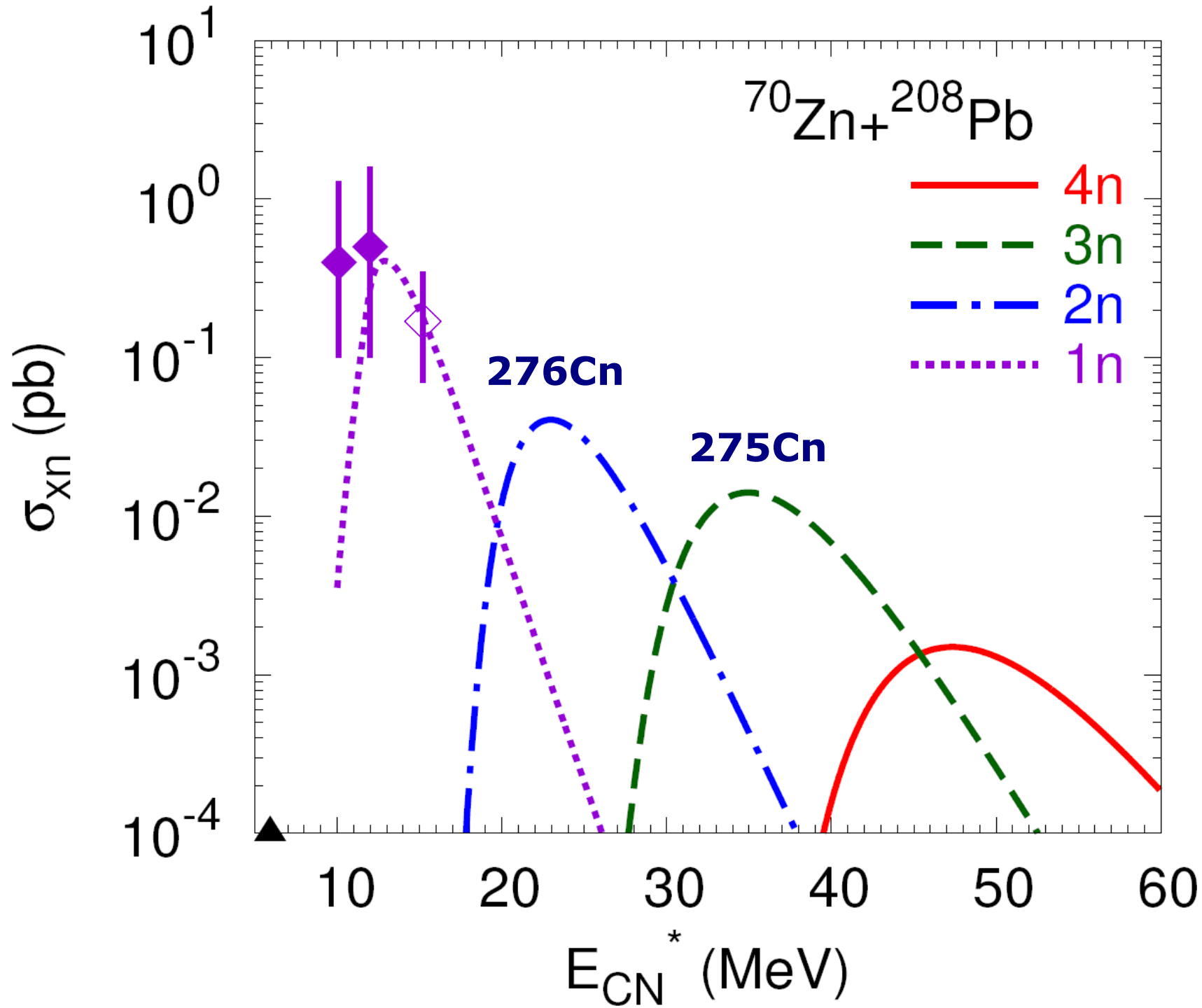


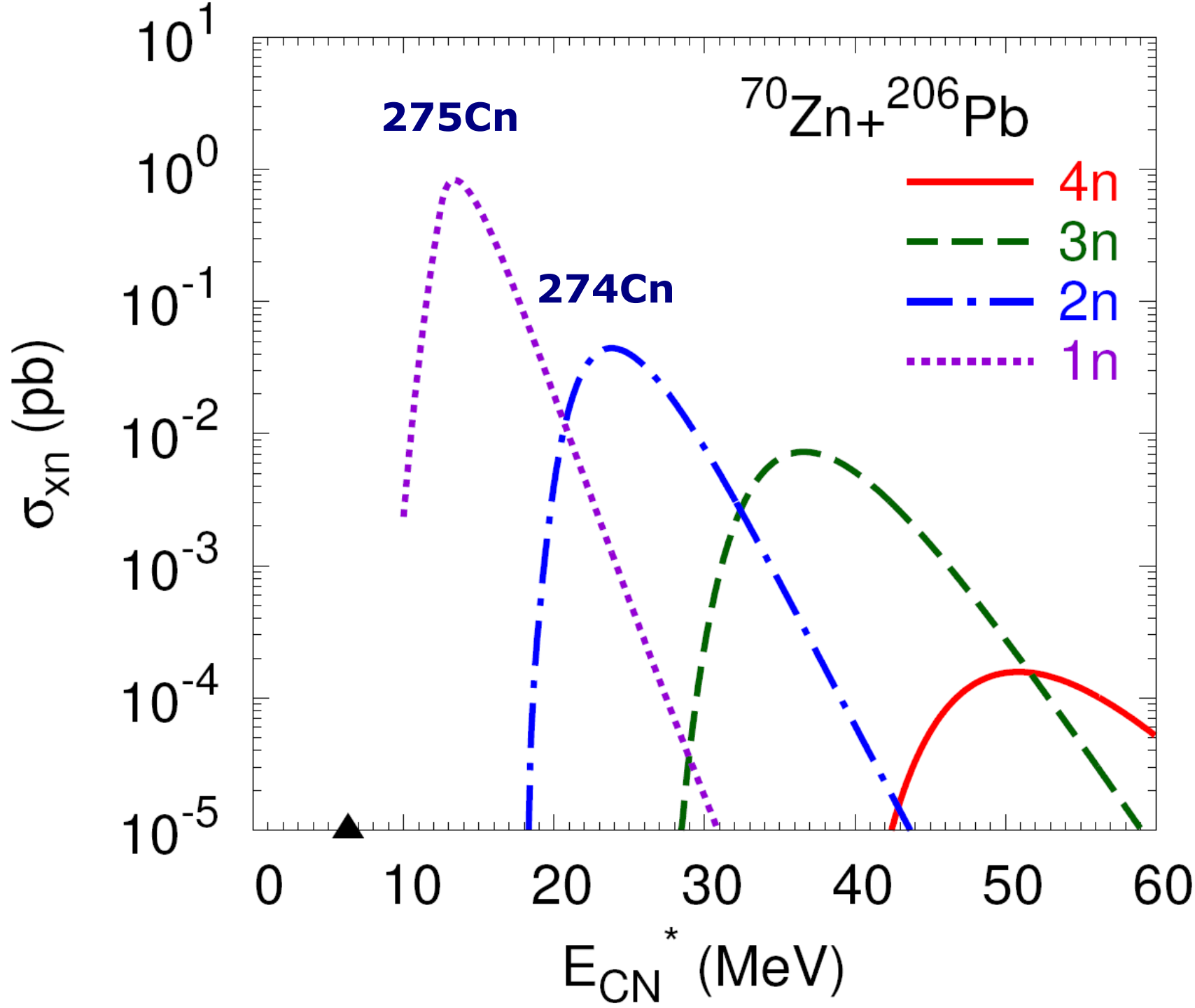


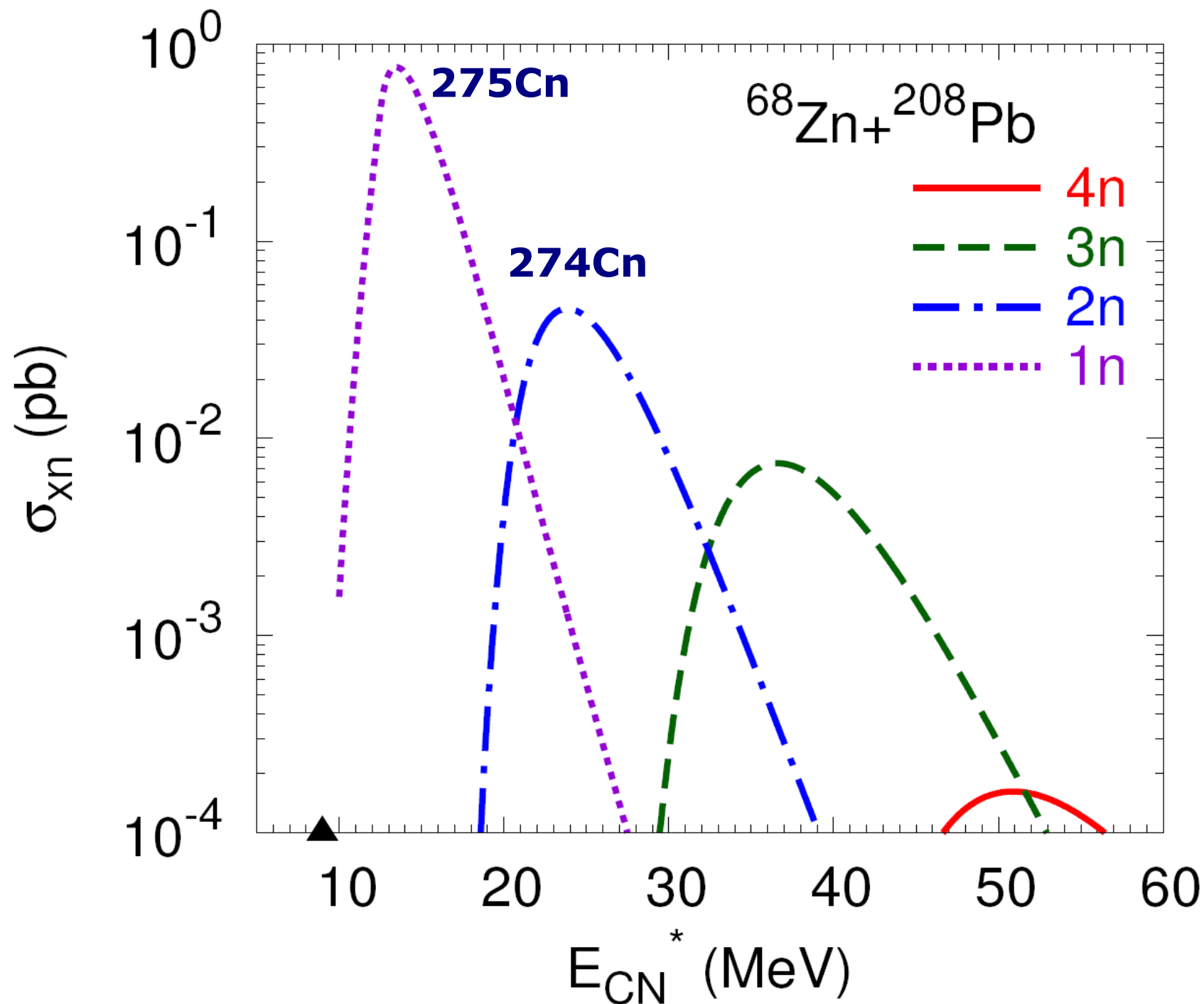
$^{62}\text{Ni} + ^{206}\text{Pb}$

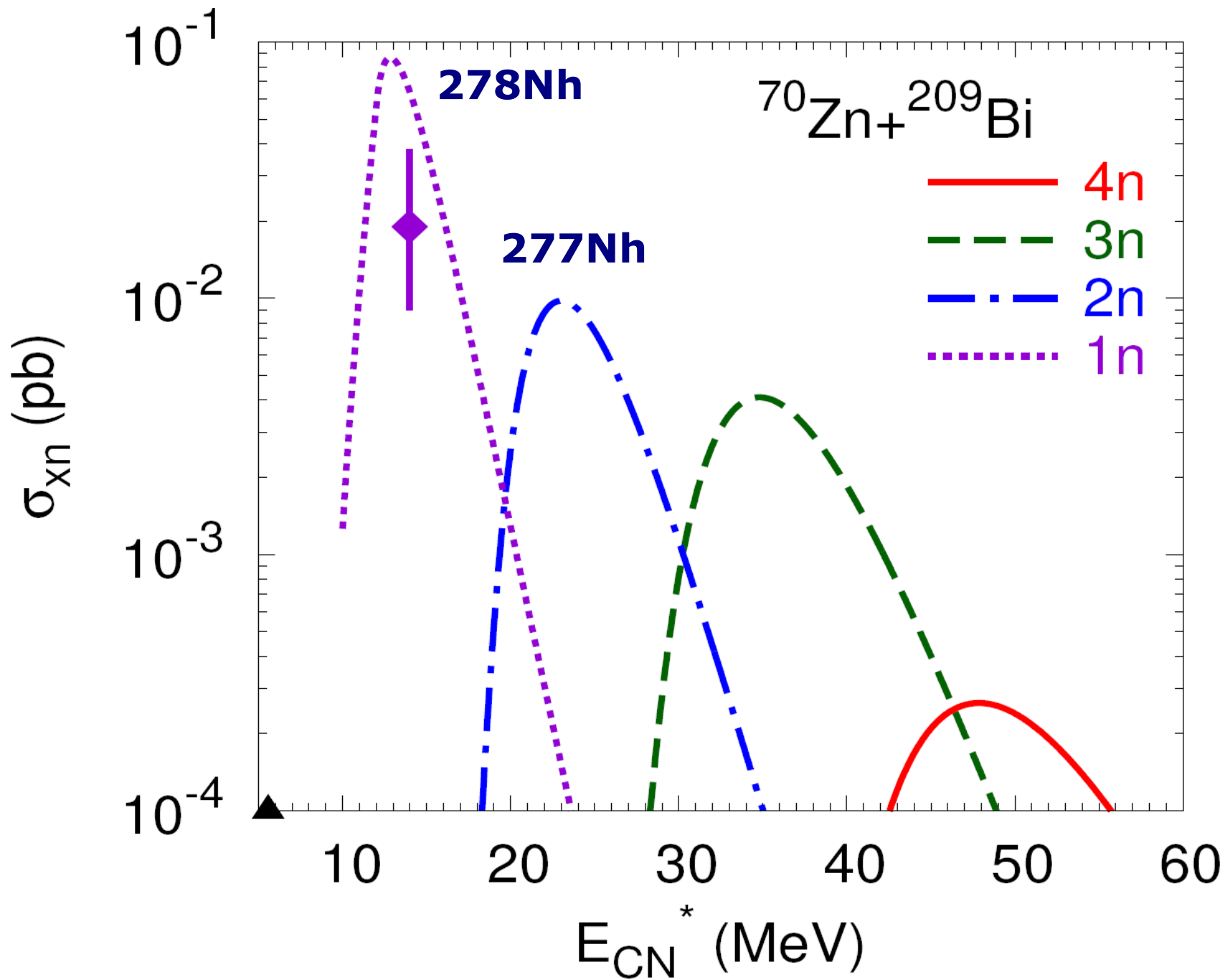


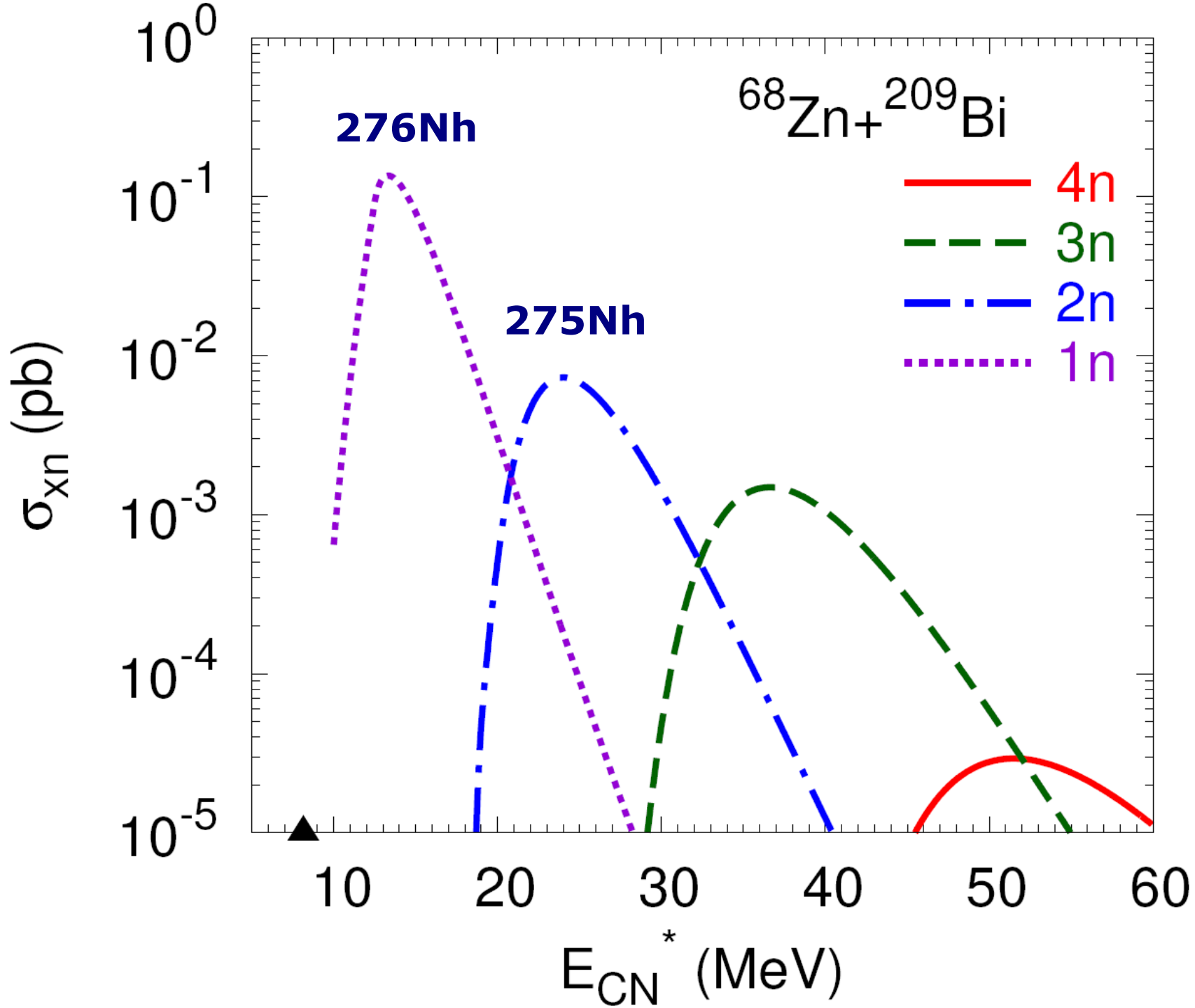


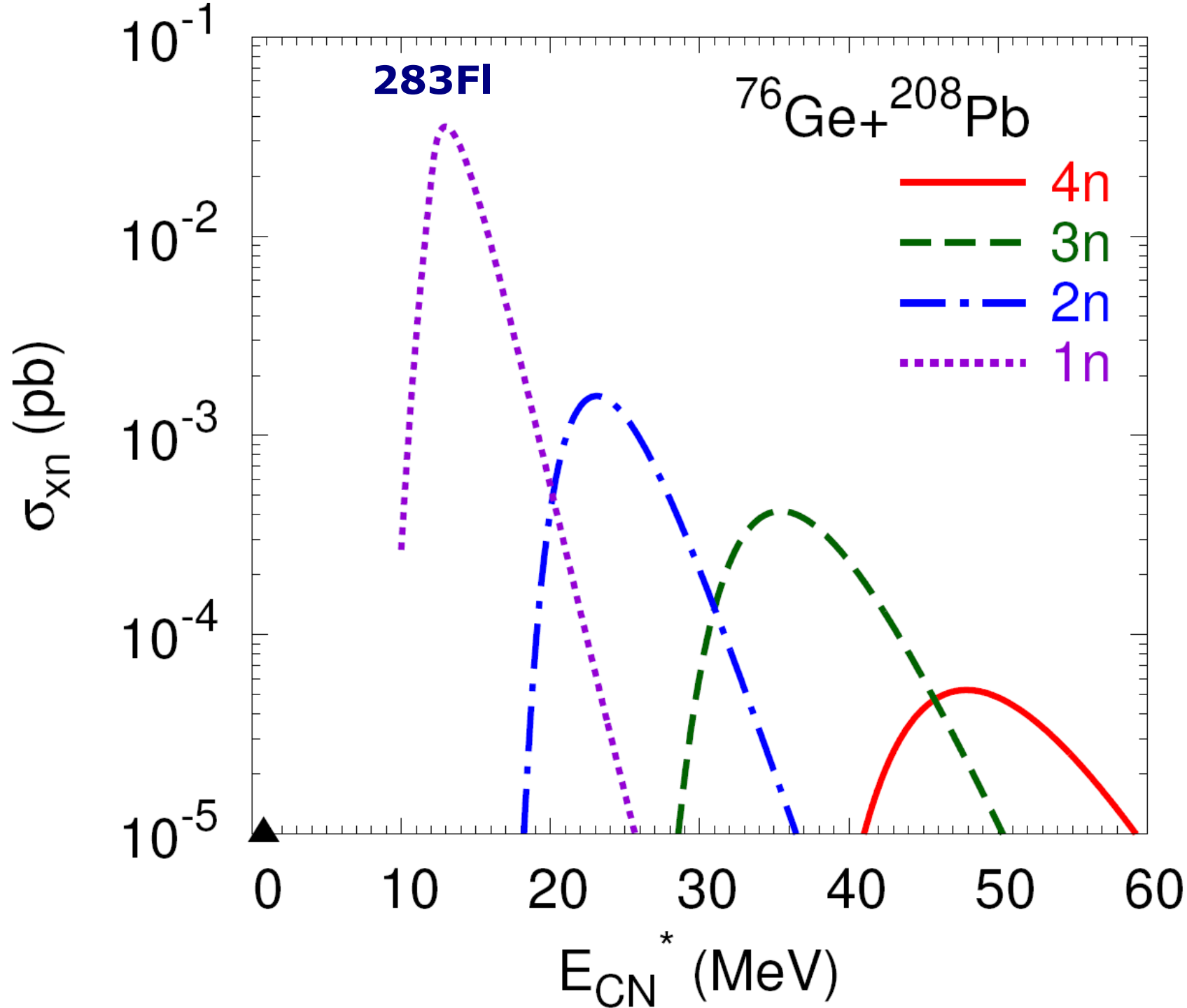






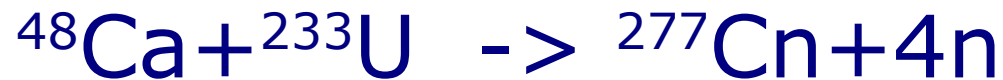








$$\frac{P_{CN}(1n)}{P_{CN}(4n)} \approx \frac{W_{sur}^{4n}}{W_{sur}^{1n}}$$



Summary

1. Using the cold fusion reactions in xn-channels, one can directly produce and study the neutron-deficient SHN in Factory

2. For future, population of the yrast rotational band of SHN produced in cold fusion ($^{50}\text{Ti} + ^{208}\text{Pb} \rightarrow ^{256}\text{Rf} + 2\text{n}$)

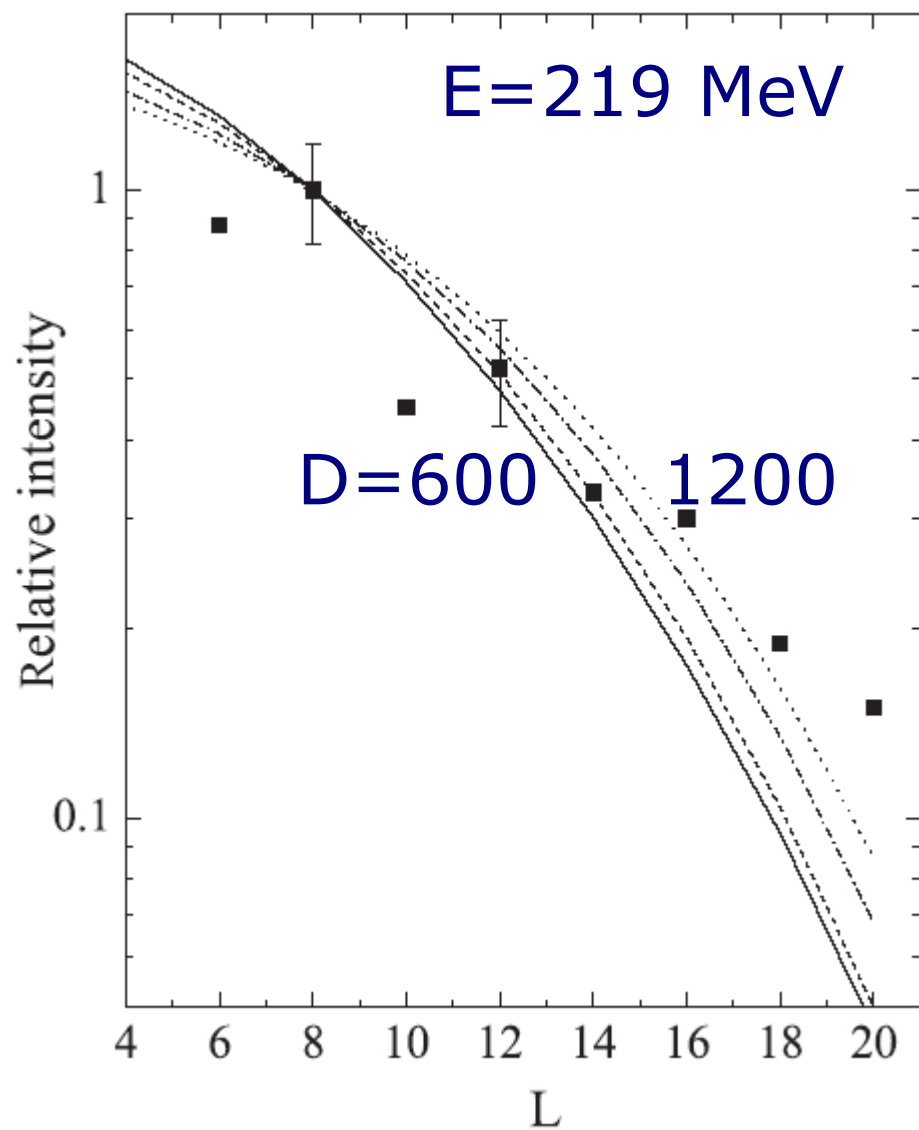
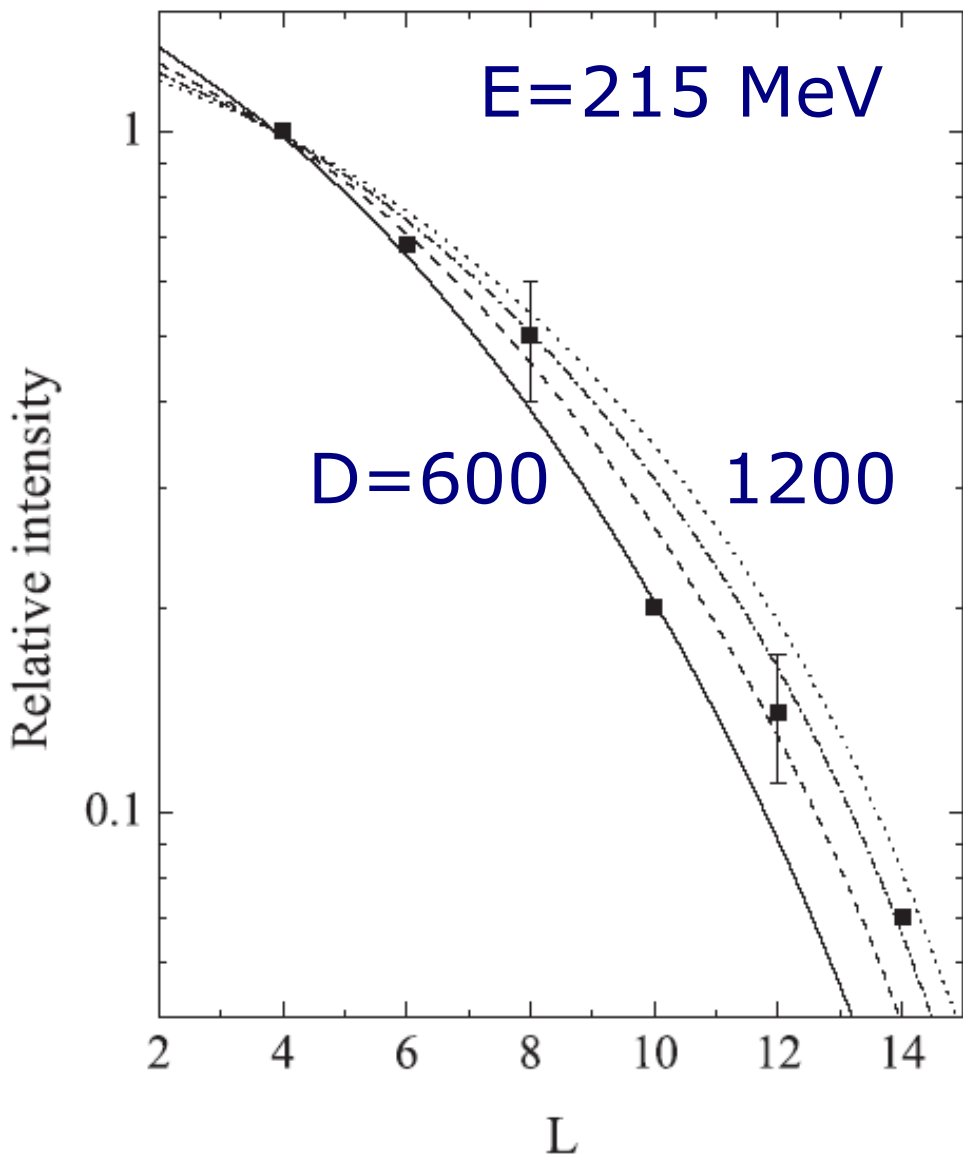
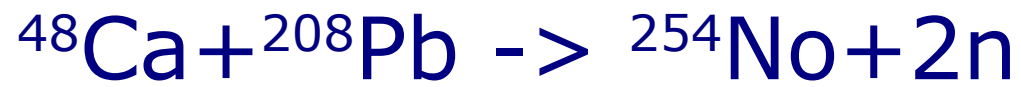
Dependence of fission barrier on spin

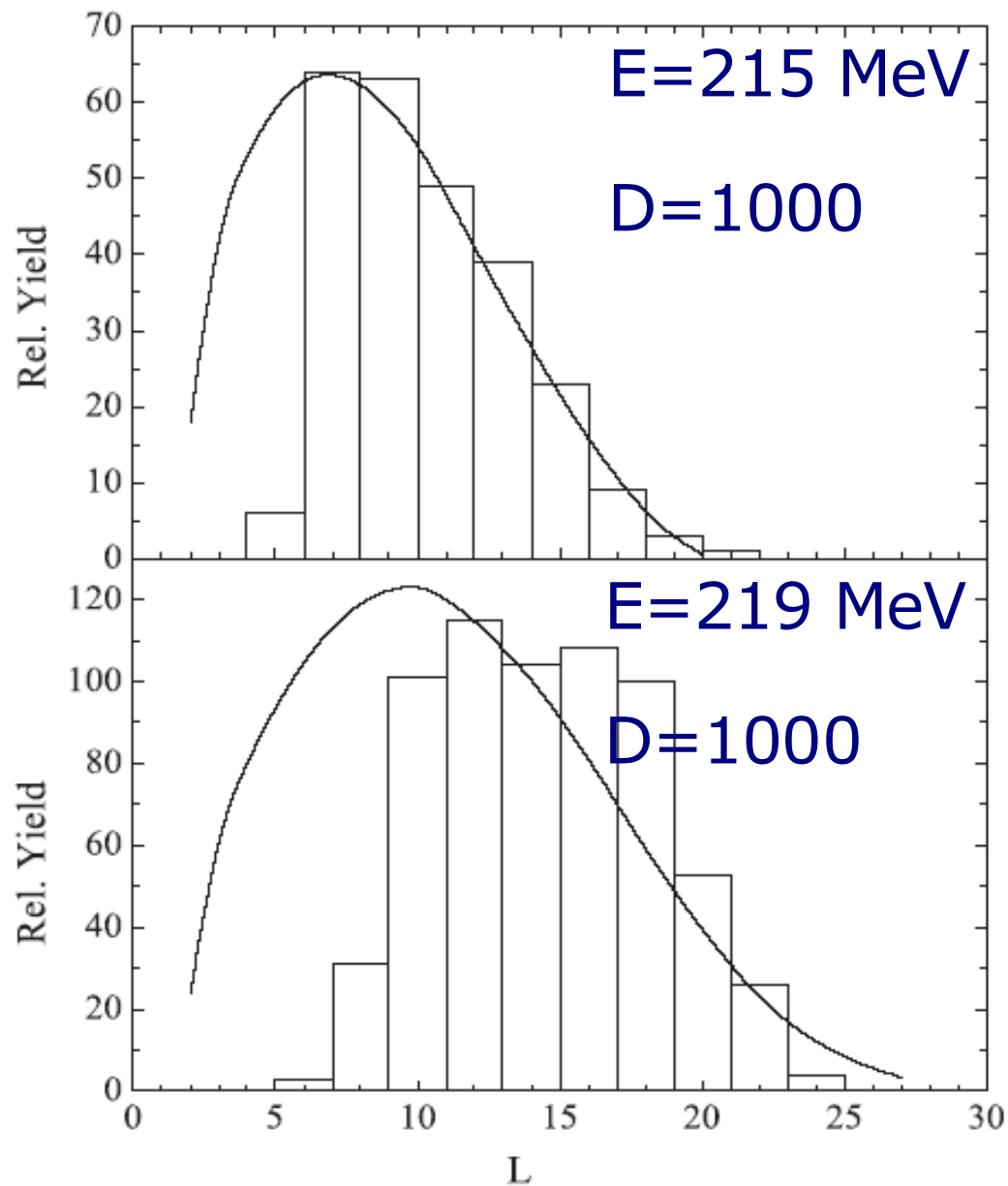
$$B_f(E_{CN}^*, J) = B_f^{LD}(J) + B_f^M(E_{CN}^* = 0) \\ \times \exp[-E_{CN}^*(J)/E_D] \exp[-J(J+1)/D]$$

Damping parameter **D from exper. data**

$$E_D = \alpha_0 A^{4/3} / a,$$

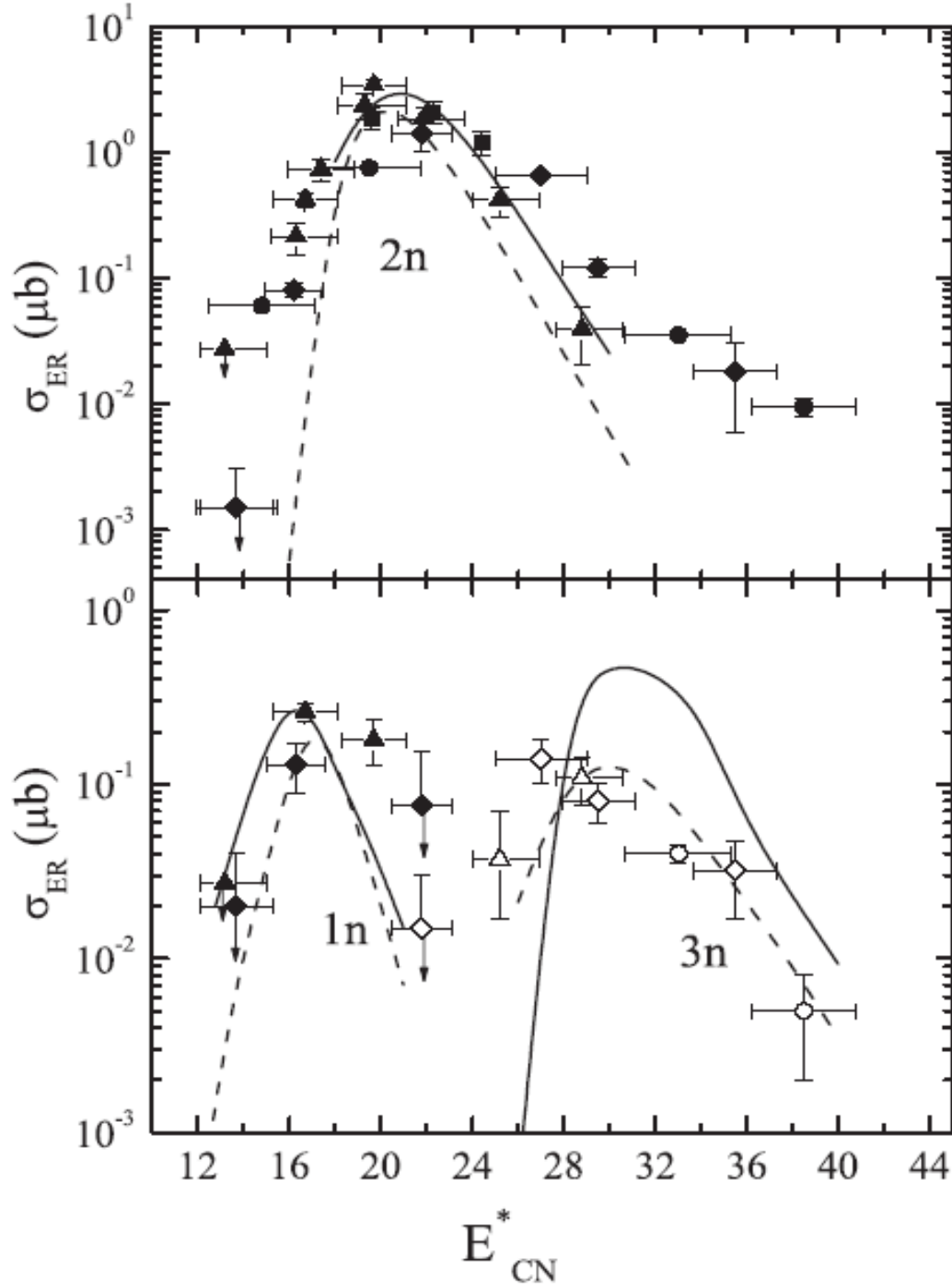
where $\alpha_0 = 0.4$.

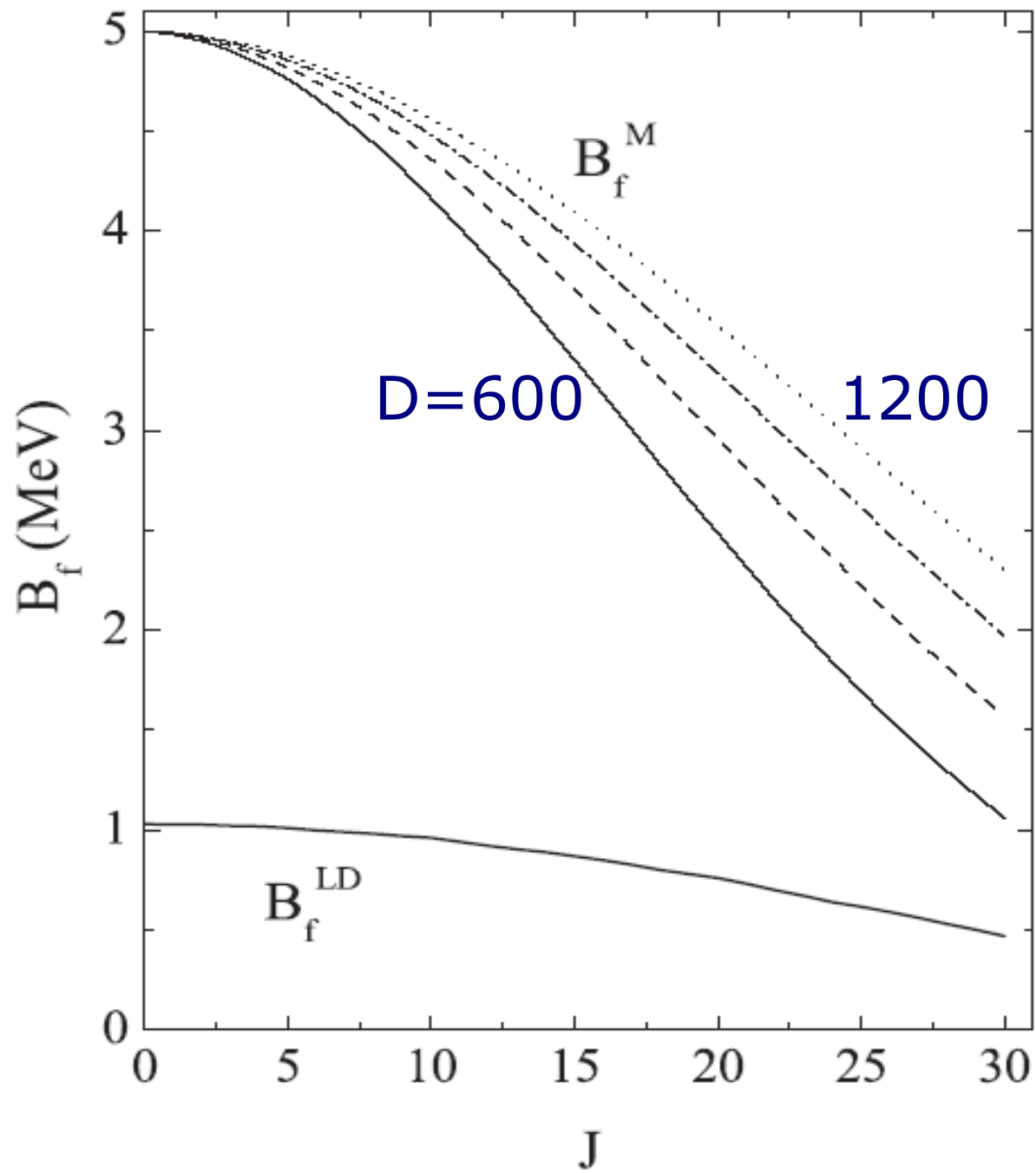




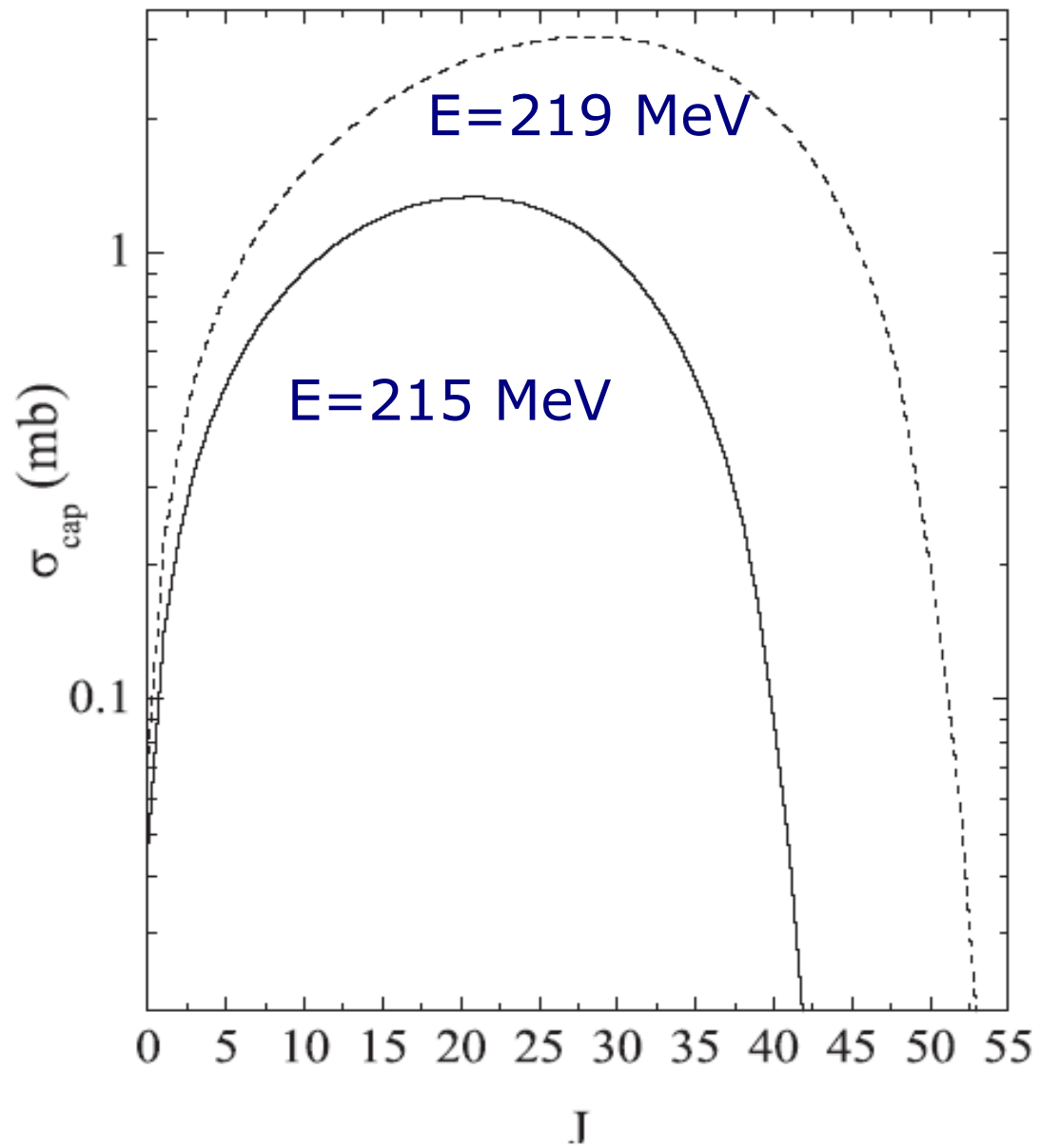
$^{48}\text{Ca} + ^{208}\text{Pb}$

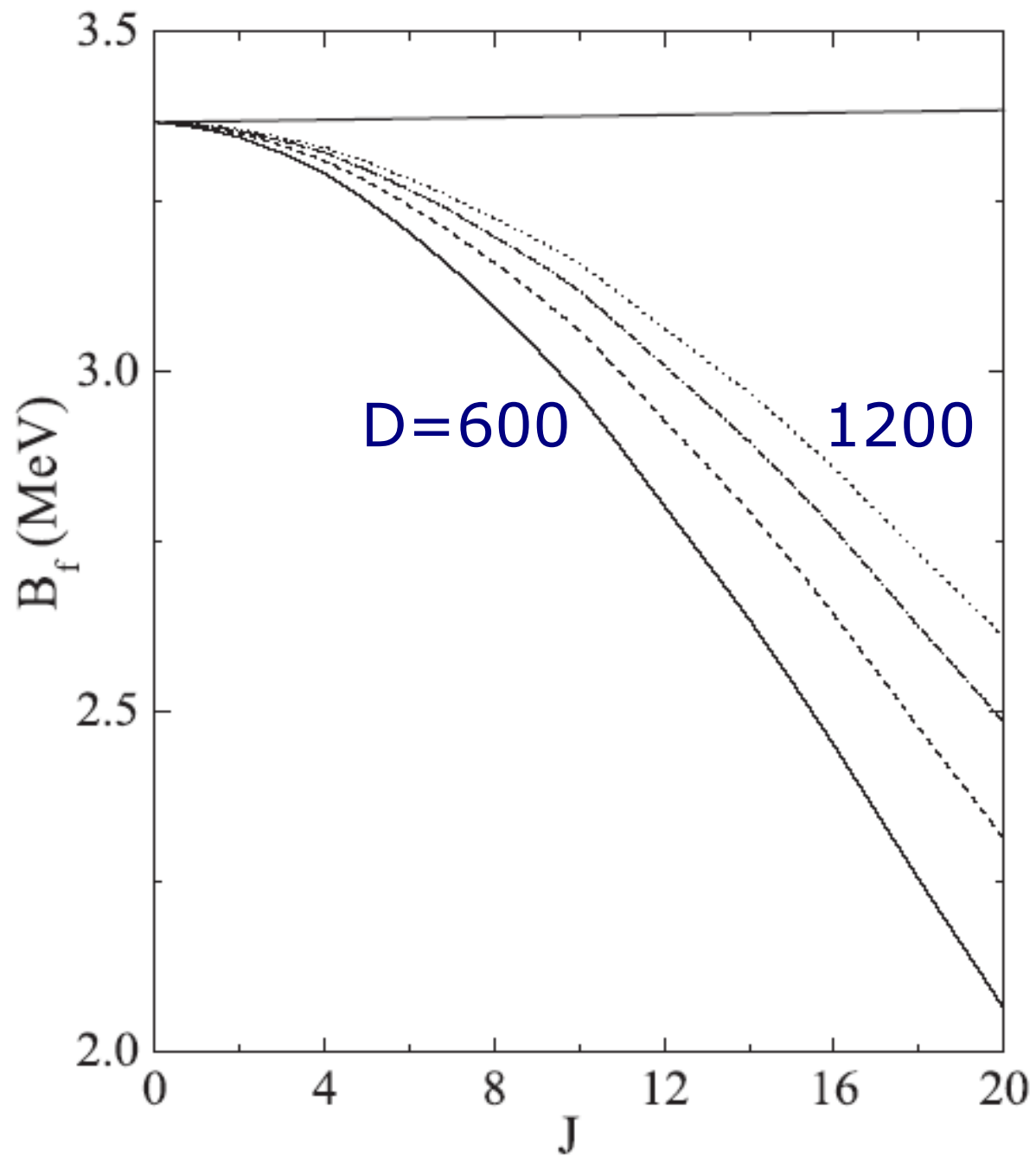
$D=1000$

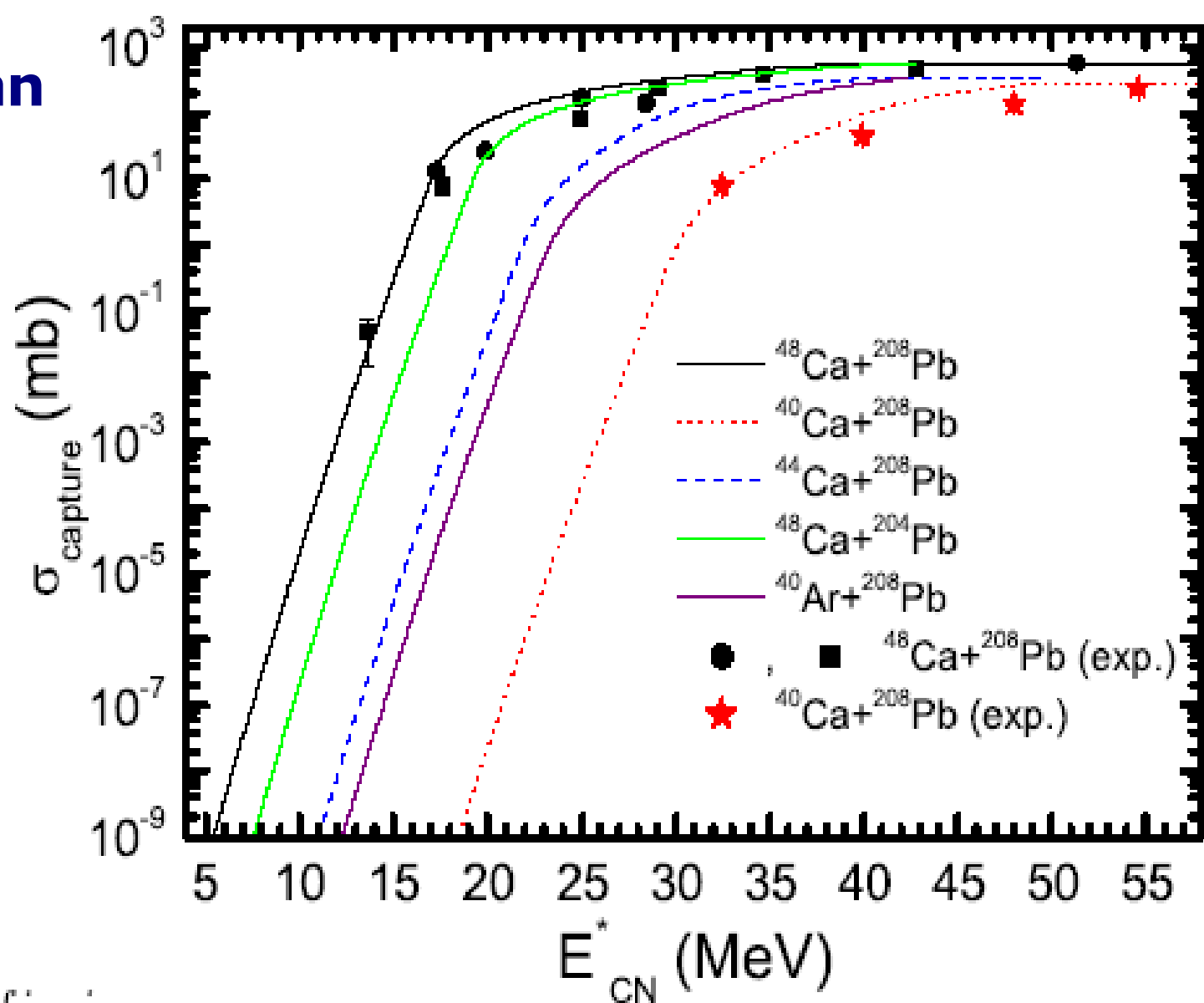




$^{48}\text{Ca} + ^{208}\text{Pb}$







$$P_{\text{CN}} \approx \frac{\sigma_{\text{fission}}}{\sigma_{\text{capture}}}$$

This allows us to obtain an access to the isotopes which are unreachable in other reactions due to the lack of proper projectile-target combinations

A weak drop of the cross section is due to

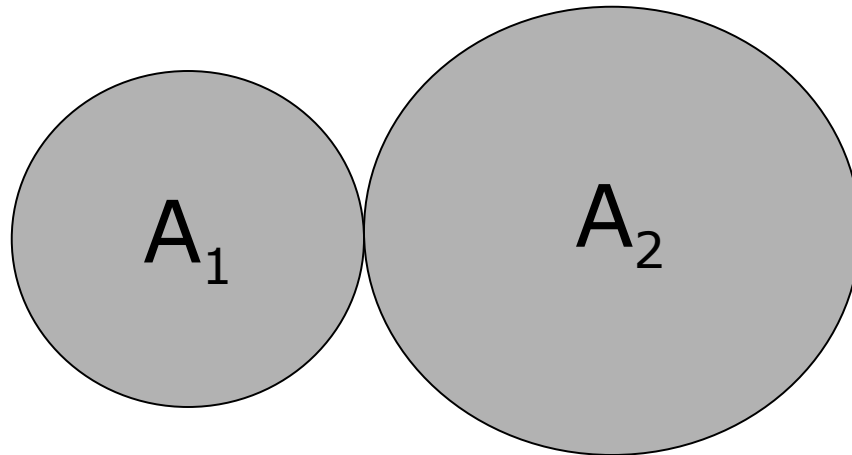
- 1. the interplay of fusion and survival probabilities**
- 2. a weak change of the difference between the fission barrier height and neutron binding energy at 1-4 steps of n-evaporation**

Energies of the maximum of cross section in **1n-,2n-channels** are considerably smaller than the Coulomb barrier height for the sphere-side orientation plus **Q**-value: $V_b + Q$.

The larger the value of $V_b + Q$, the smaller the cross sections are for **1n-,2n-channels**

Mass asymmetry coordinate

$$\eta = \frac{A_1 - A_2}{A_1 + A_2}$$



$\eta = 0$ for $A_1 = A_2$, $\eta = \pm 1$ for A_1 or $A_2 = 0$