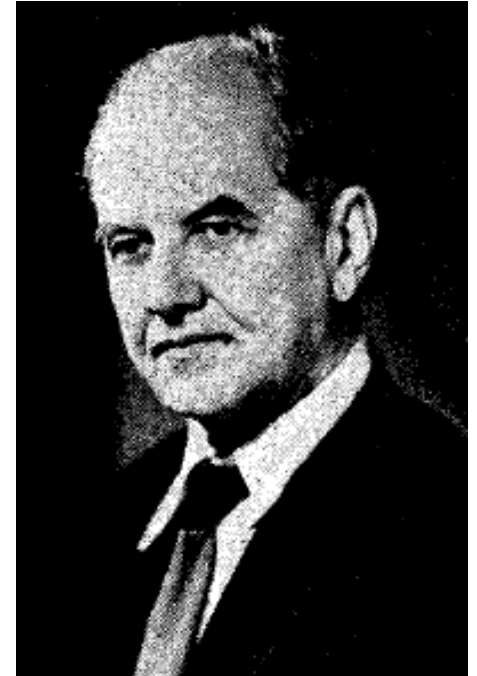
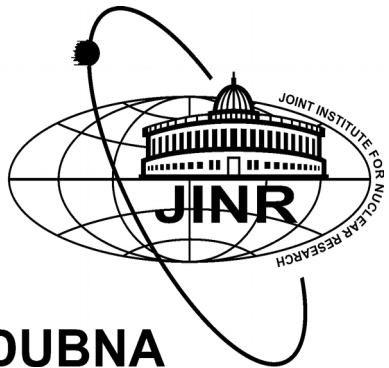


Nonequilibrium pion distribution in heavy-ion collisions from the Zubarev approach

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JINR Dubna; NRNU (MEPhI) Moscow



Colloquium Zubarev-100
Dubna 18./19.04.2018



Personal recollections – The academic family



N.N. Bogoliubov (1909 - 1992)



D.N. Zubarev (1917 – 1992)
PhD student of N.N.B.



G. Röpke (1941)
“Postdoc” of D.N.Z. (1968-70)

D. Blaschke (1959)



H. Reinholz (1959)



Recollections – The year 1959



Shirkov, [Zubarev](#), Logunov, ?, Galasiewicz, Tavkhelidze
Visit to Mzcheta (Georgia), 1959

Recollections – The year 1959



Shirkov, [Zubarev](#), Logunov, ?, Galasiewicz, Tavkhelidze
Visit to Mzcheta (Georgia), 1959



Recollections – The year 1959



Shirkov, [Zubarev](#), Logunov, ?, Galasiewicz, Tavkhelidze
Visit to Mzcheta (Georgia), 1959

Recollections – The year 1959



D.B. born 22.9.1959

Shirkov, [Zubarev](#), Logunov, ?, Galasiewicz, Tavkhelidze
Visit to Mzcheta (Georgia), 1959

Recollections – The year 1959



Shirkov with former vice Directors of BLTP (2009)

Shirkov, [Zubarev](#), Logunov, ?, Galasiewicz, Tavkhelidze
Visit to Mzcheta (Georgia), 1959

Recollections – The year 1959



Shirkov with former vice Directors of BLTP (2009)



Zubarev & Tokarchuk visit Rostock (1987)

Shirkov, [Zubarev](#), Logunov, ?, Galasiewicz, Tavkhelidze
Visit to Mzcheta (Georgia), 1959

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Shirkov with former vice Directors of BLTP (2009)



Zubarev & Tokarchuk visit Rostock (1987)



Galasiewicz at Wroclaw University (2008)

Shirkov, Zubarev, Logunov
Visit to Mzcheta (Georgia)







May 1984 – Inauguration of the Blokhintsev plate

Low-momentum pion enhancement – early stage

First heavy-ion collision experiments at CERN SPS with O-Au and S-S at 60 and 200 GeV/n
 → low momentum enhancement of negative particle (pions) yields relative to Bose-distribution

Possible solution of the puzzle →

$$f(p, x) = \frac{1}{e^{(E-\mu)/T} \pm 1}, \quad E = \sqrt{p_T^2 + p_L^2 + m^2}$$

$$\frac{dN}{p_T dp_T dy} = (\pi R^2 \tau_{dec}) \frac{m_T}{(2\pi)^2}$$

$$\times \sum_{n=1}^{\infty} (\pm 1)^{n+1} e^{n\mu/T} K_1\left(n \frac{m_T}{T}\right),$$

Transverse mass $m_T = \sqrt{p_T^2 + m^2}$

Nonequilibrium chemical potential
 for pions $\mu = 126$ MeV; $T=167$ MeV

See also:

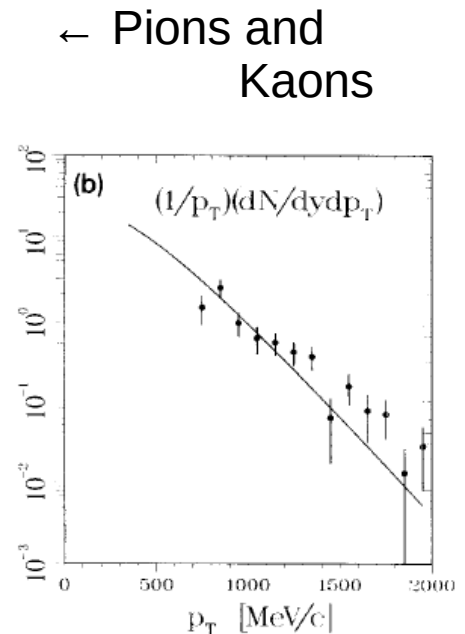
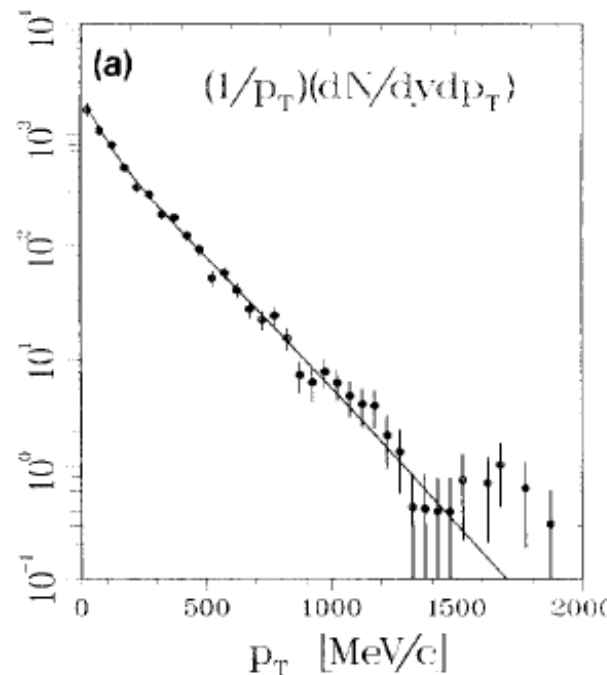
S. Gavin, “Pions in and out of equilibrium”
 NPA 544 (1992) 459; NPB 351 (1991) 561

Non-zero chemical potential
 and the shape of the p_T -distribution of hadrons
 in heavy-ion collisions

PLB 243 (1990) 181

M. Kataja and P.V. Ruuskanen

Department of Physics, University of Jyväskylä, Seminaarinkatu 15, SF-40100 Jyväskylä, Finland



← Pions and
 Kaons

Low-momentum pion enhancement – LHC revival

Explanation of hadron transverse-momentum spectra in heavy-ion collisions at $\sqrt{s_{NN}} = 2.76$ TeV within a chemical nonequilibrium statistical hadronization model

V. Begun, W. Florkowski, M. Rybczynski,
PRC 90, 014906 (2014)

Cooper-Frye formula

$$\frac{dN}{dyd^2p_T} = \int d\Sigma_\mu p^\mu f(p \cdot u),$$

Primordial distribution

$$f_i = \frac{g_i}{\Upsilon_i^{-1} \exp(\sqrt{p^2 + m_i^2}/T) \mp 1}.$$

Fugacity factor

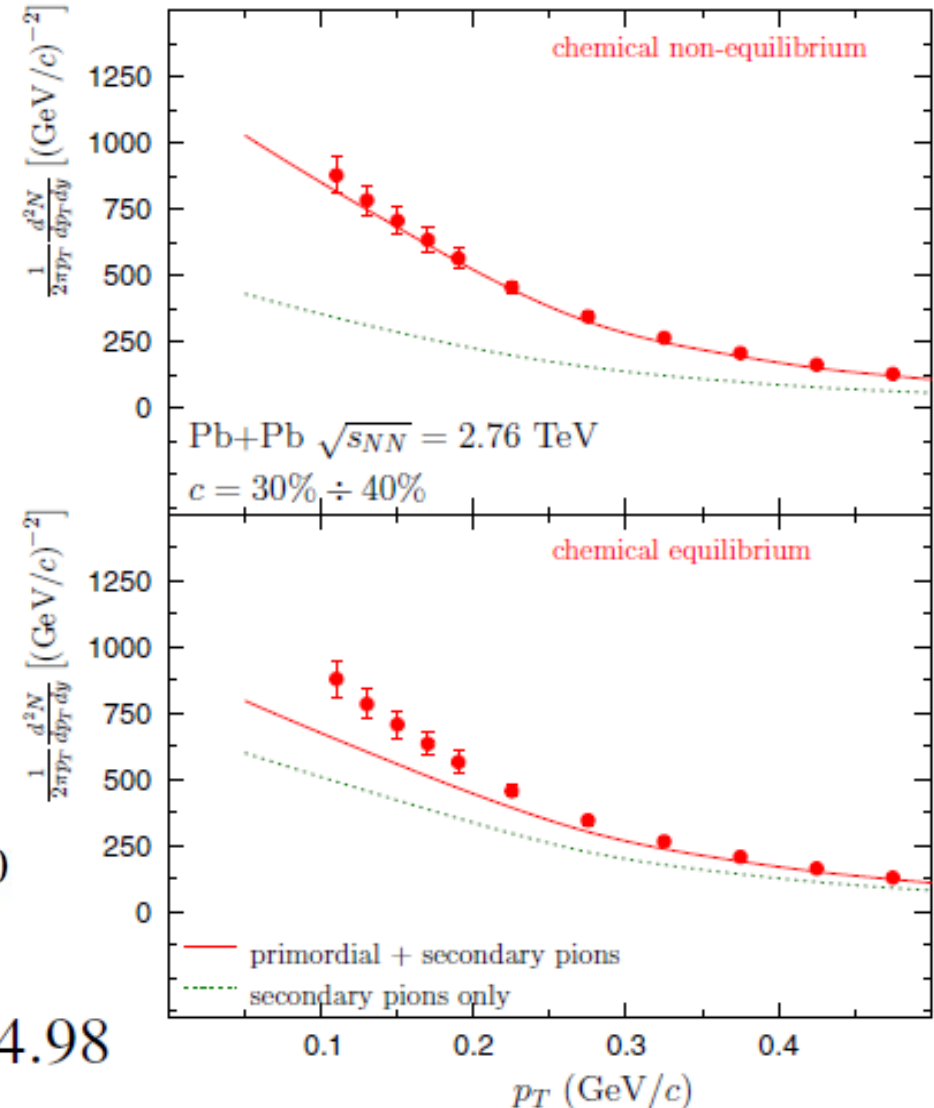
$$\begin{aligned} \Upsilon_i &= (\lambda_q \gamma_q)^{N_q^i} (\lambda_s \gamma_s)^{N_s^i} (\lambda_{\bar{q}} \gamma_{\bar{q}})^{N_{\bar{q}}^i} (\lambda_{\bar{s}} \gamma_{\bar{s}})^{N_{\bar{s}}^i} \\ &= \gamma_q^{N_q^i + N_{\bar{q}}^i} \gamma_s^{N_s^i + N_{\bar{s}}^i} \exp\left(\frac{\mu_B B_i + \mu_S S_i}{T}\right) \end{aligned}$$

Fit for semiperipheral collisions, $c=30\%-40\%$

$$T = 139.85 \text{ MeV}, \quad \gamma_q = 1.62, \quad \gamma_s = 2.0$$

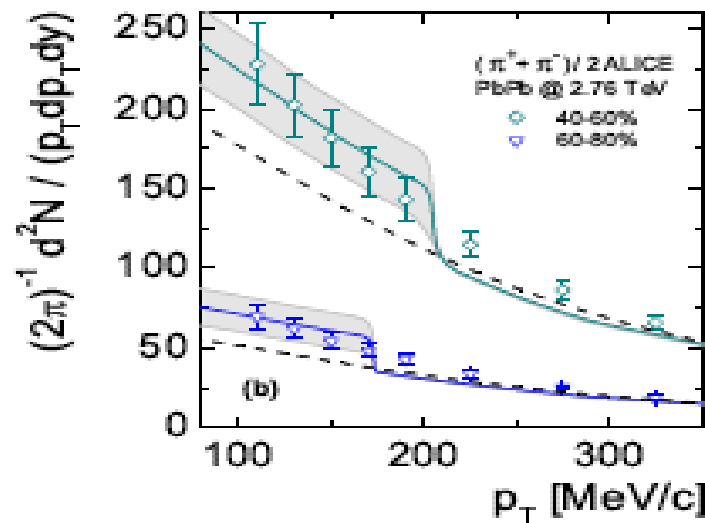
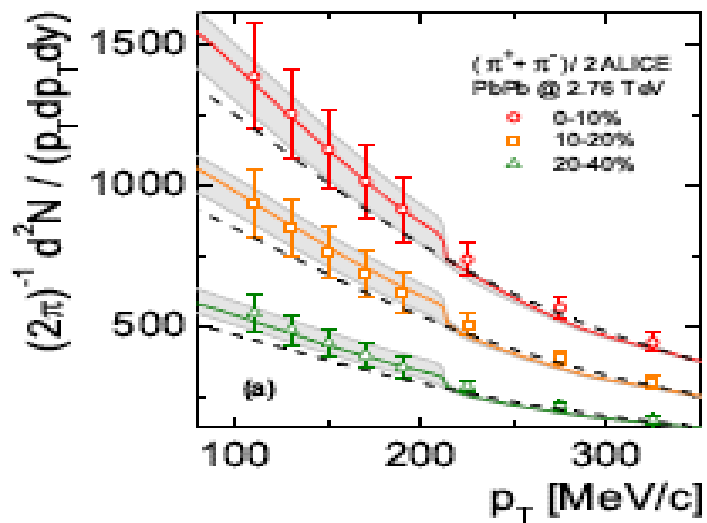
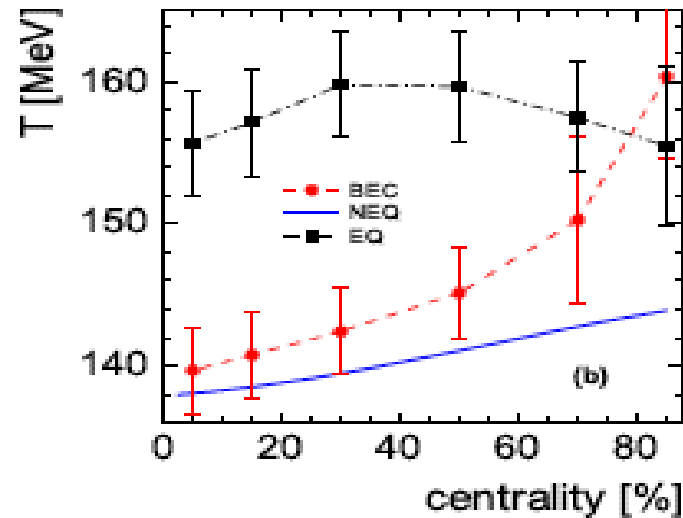
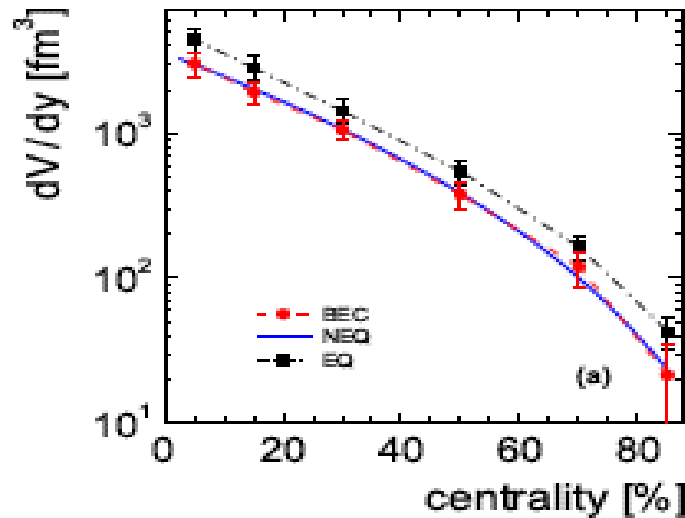
Low-pT Bose enhancement of pions

$$\mu_\pi = 2T \ln \gamma_q \simeq 134.9 \text{ MeV} \leftrightarrow m_{\pi^0} \simeq 134.98$$



Low-momentum pion enhancement at LHC - Onset of Bose-Einstein Condensation of pions ?

$$n = \int d^3p \frac{1}{(2\pi)^3} \frac{g}{\exp\left(\frac{\sqrt{p^2+m^2}-\mu}{T}\right) - 1} \left[1 + \frac{(2\pi)^3}{V} \delta(p_x) \delta(p_y) \delta(p_z) \right]$$



Nonequilibrium statistical approach

nonequilibrium statistical operator $\rho(t)$

$$\rho(t) = \lim_{\varepsilon \rightarrow 0} \varepsilon \int_{-\infty}^t dt' e^{\varepsilon(t-t')} U(t, t') \rho_{\text{rel}}(t') U^\dagger(t, t') \quad U(t, t') = \exp[(i/\hbar)H(t - t')]$$

limit $\varepsilon \rightarrow 0$ has to be taken after the thermodynamic limit

relevant statistical operator $\rho_{\text{rel}}(t)$ according to Zubarev

The missing irrelevant correlations $\rho_{\text{irrel}}(t) = \rho(t) - \rho_{\text{rel}}(t)$

are assumed to be formed dynamically by the action of the time evolution operator $U(t, t')$

The relevant information is given by the averages of the relevant observables B_n

The maximum

of the information entropy at given averages is the generalized Gibbs distribution $\langle B_n \rangle^t = \text{Tr}\{\rho(t) B_n\}$

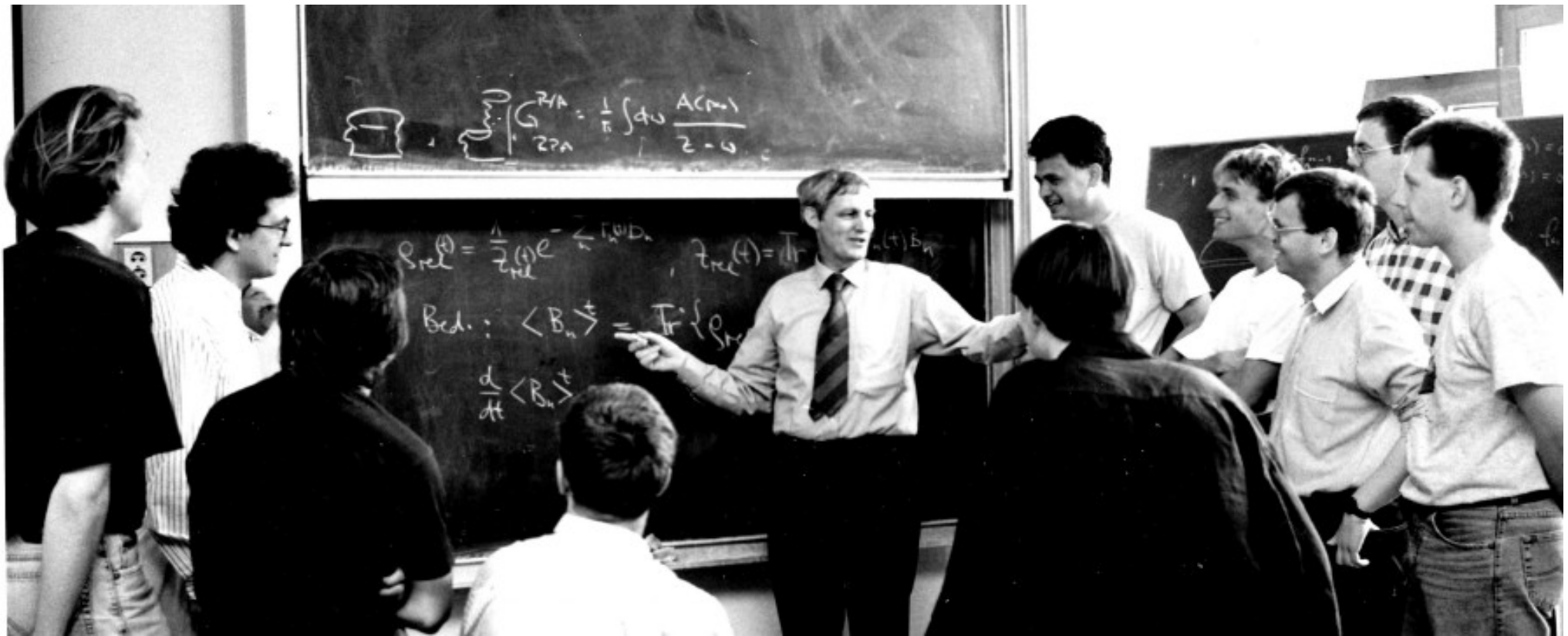
$$\rho_{\text{rel}}(t) = \frac{1}{Z_{\text{rel}}(t)} e^{-\sum_n F_n(t) B_n}; \quad Z_{\text{rel}}(t) = \text{Tr} e^{-\sum_n F_n(t) B_n}$$

where the Lagrange multipliers $F_n(t)$ are determined by the self-consistency conditions

$$\langle B_n \rangle_{\text{rel}}^t = \langle B_n \rangle^t$$

with $\langle B_n \rangle_{\text{rel}}^t \equiv \text{Tr}\{\rho_{\text{rel}}(t) B_n\}$.

Nonequilibrium statistical operator formalism



Gerd Roepke explains the Zubarev approach to nonequilibrium in the Max-Planck group at Rostock University (1995)

Nonequilibrium approach to pion enhancement effect

Hamiltonian with instantaneous and retarded part

$$\mathcal{H} = \mathcal{H}^{(0)}\delta(t - t') + \mathcal{H}^{(r)}(t - t')$$

Occupation numbers $N_c = \sum_{\mathbf{p}} n_{\mathbf{p},c}$

Hamiltonian $H = \sum_c H_c + \frac{1}{2} \sum_{c,c'} V_{c,c'}^{\text{interaction}}$

Relevant statistical operator $\rho_{\text{rel}}(t) = \frac{1}{Z_{\text{rel}}(t)} e^{-\sum_c \beta_c(t)(H_c - \mu_c(t)N_c)}$

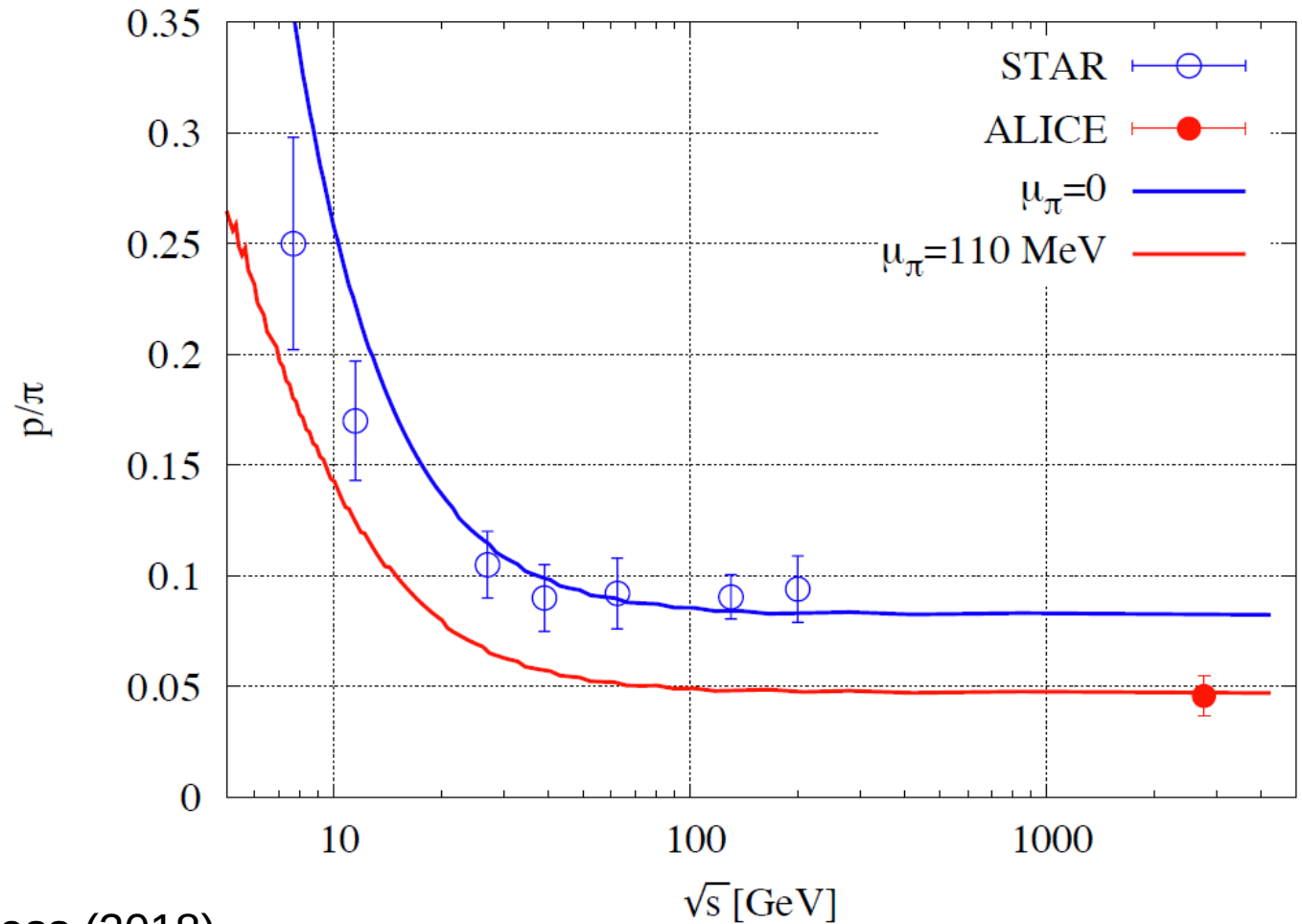
Occupation numbers $\langle n_{\mathbf{p},c} \rangle_{\text{rel}}^t = \frac{1}{e^{\beta_c(t)[E_c^{\text{mf}}(\mathbf{p}) - \mu_c(t)]} \pm 1}$

Equations of motion $\frac{d}{dt} \langle N_c \rangle^t = \frac{i}{\hbar} \langle [H, N_c] \rangle_{\text{rel}}^t - \frac{1}{\hbar^2} \int_{-\infty}^0 d\tau e^{\epsilon\tau} \text{Tr}\{\rho_{\text{rel}}(t)[H(\tau), [H, N_c]]\}$

$$\frac{d}{dt} \langle H_c^{\text{mf}} \rangle^t = \frac{i}{\hbar} \langle [H, H_c^{\text{mf}}] \rangle_{\text{rel}}^t - \frac{1}{\hbar^2} \int_{-\infty}^0 d\tau e^{\epsilon\tau} \text{Tr}\{\rho_{\text{rel}}(t)[H(\tau), [H, H_c^{\text{mf}}]]\}$$

Nonequilibrium approach to pion enhancement effect

LHC pion and proton puzzle related



G. Roepke et al., in progress (2018)

Conclusions

- Zubarev approach to nonequilibrium gives a sound conceptual basis to describe deviations from equilibrium distributions from particles produced In heavy-ion collisions, under nonequilibrium conditions
- perfect example for applications is the low-pT pion enhancement as described with a pion chemical potential that appears as a Lagrange multiplier of a generalized nonequilibrium statistical operator
- Extension to chemical nonequilibrium situations (e.g., strangeness production) shall be straightforward
- A microscopic process for the overpopulation of the pion phase space need to be worked out, e.g., a gluon-to-pion conversion process