

Soft Physics for high energy Nucleus-Nucleus Collisions

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2nd China-Russia Joint Workshop on NICA Facility

Qing-Dao 2024年09月10-13日

Landscape of nuclear physics

degrees of freedom

**quarks
& gluons**



quarks, gluons

Energy
(MeV)

940
neutron mass



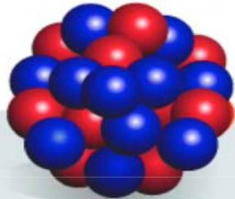
constituent quarks

hadrons



baryons, mesons

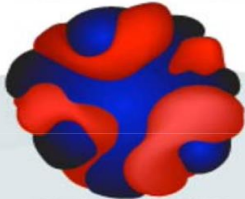
140
pion mass



protons, neutrons

8
proton separation
energy in lead

nuclei



nucleonic densities
and currents

1.32
vibrational
state in tin



collective coordinates

0.043
rotational
state in uranium

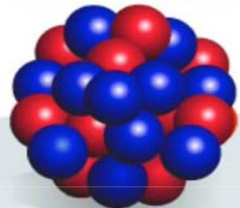
Landscape of nuclear physics

degrees of freedom

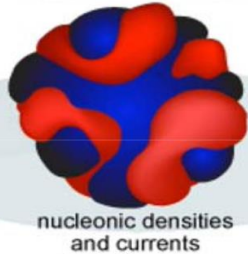
quarks & gluons



hadrons



nuclei



Energy (MeV)

940
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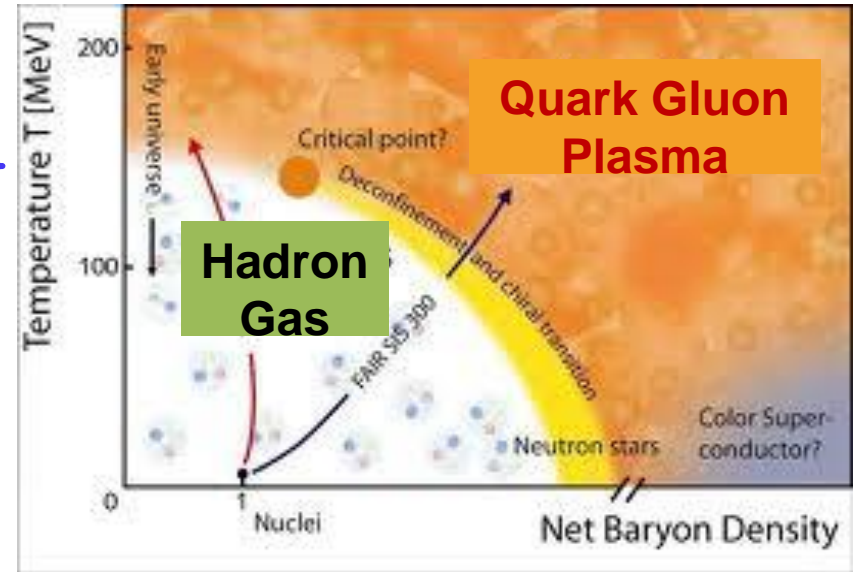
140
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proton separation energy in lead

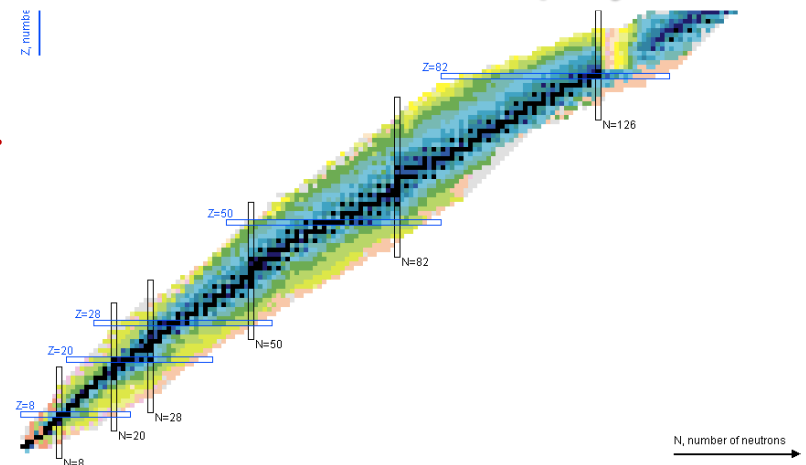
1.32
vibrational state in tin

0.043
rotational state in uranium

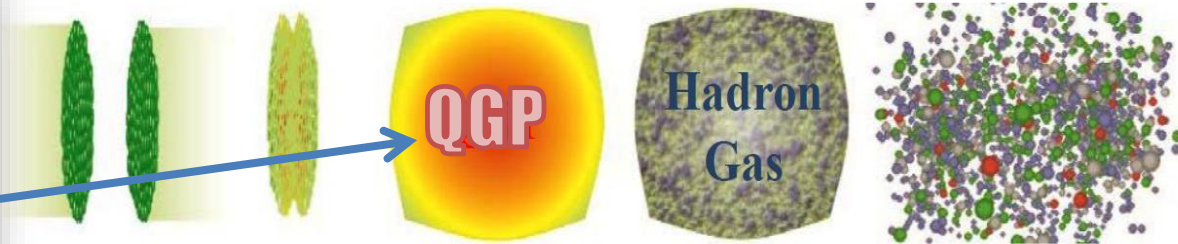
-intermediate and high energy nuclear physics



-nuclear structure physics

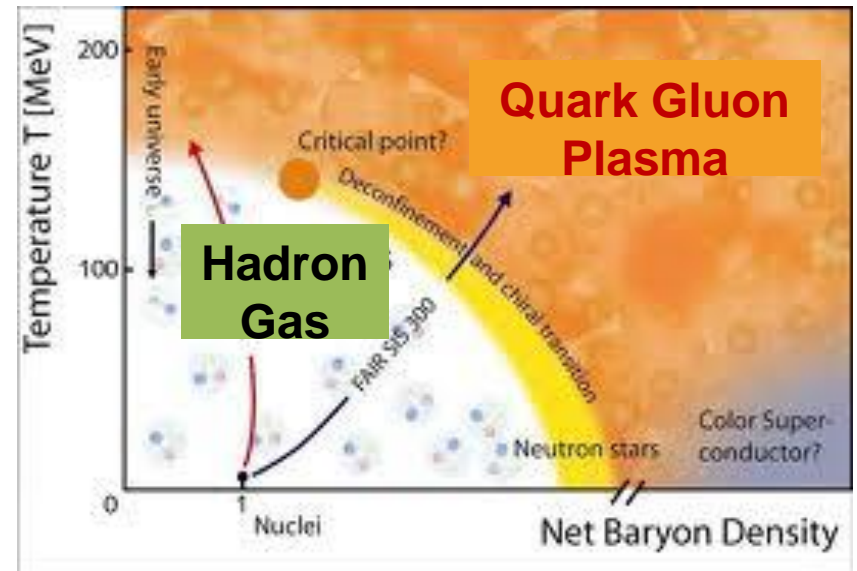


Relativistic heavy ion physics

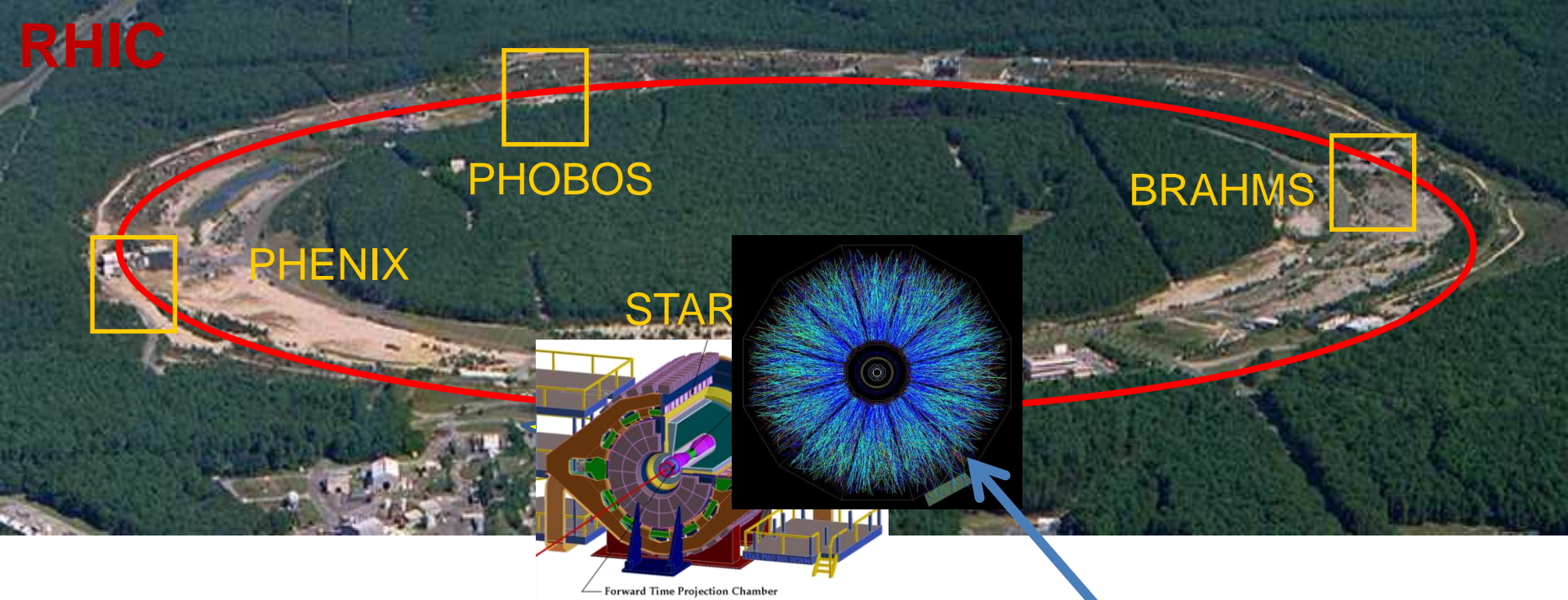


Relativistic heavy ion collisions

- create and study QGP
- the QCD phase diagram
- the deconfinement & chiral phase transition
- the QCD vacuum



RHIC

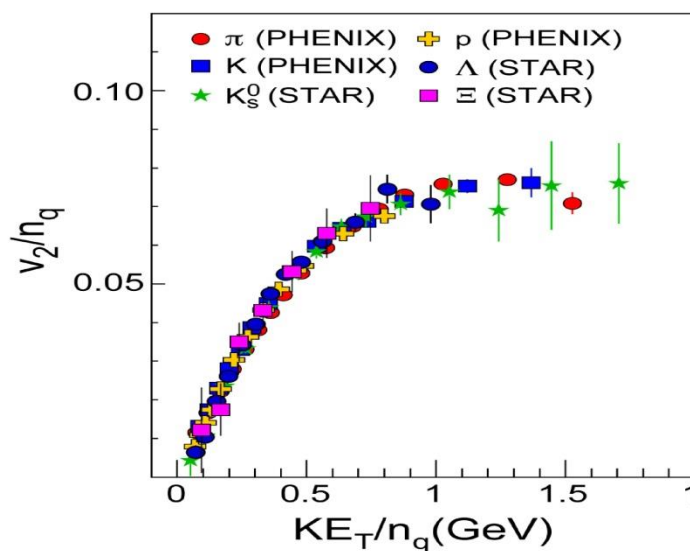
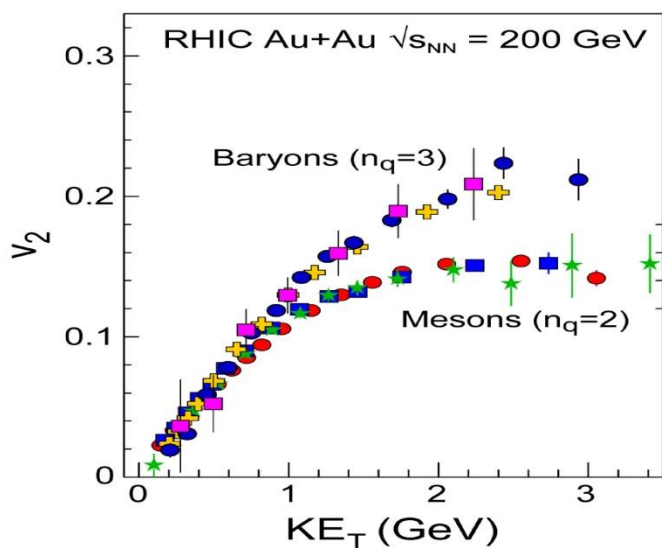
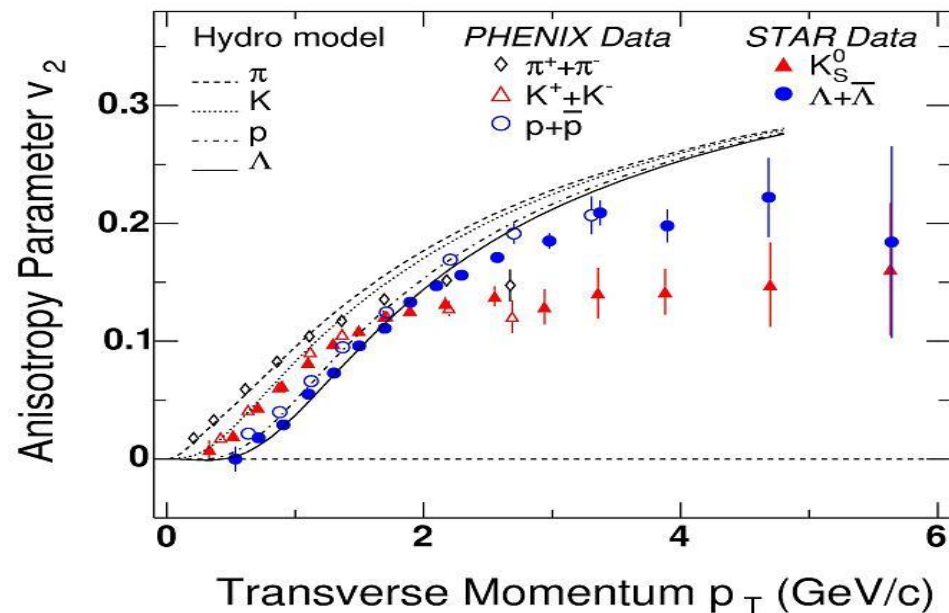
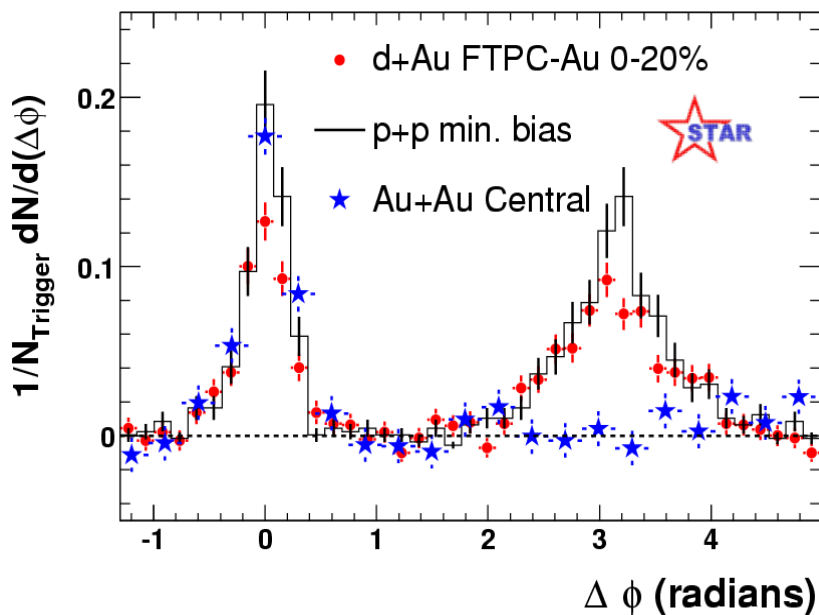


little bang: the different stage for a relativistic heavy ion collisions



The QGP was discovered

RHIC (2000--)



QGP-the most perfect fluid in the world

BNL News, 2005

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[New state of matter more remarkable than predicted -- raising many new questions](#)

April 18, 2005

TAMPA, FL -- The four detector groups conducting research at the [Relativistic Heavy Ion Collider](#) (RHIC) -- a giant atom "smasher" located at the U.S. Department of Energy's Brookhaven National Laboratory -- say they've created a new state of hot, dense matter out of the quarks and gluons that are the basic particles of atomic nuclei, but it is a state quite different and even more remarkable than had been predicted. In [peer-reviewed papers](#) summarizing the first three years of RHIC findings, the scientists say that instead of behaving like a gas of free quarks and gluons, as was expected, the matter created in RHIC's heavy ion collisions appears to be more like a *liquid*.

"Once again, the physics research sponsored by the Department of Energy is producing historic results," said Secretary of Energy Samuel Bodman, a trained chemical engineer. "The DOE is the principal federal funder of basic research in the physical sciences, including nuclear and high-energy physics. With today's announcement we see that investment paying off."

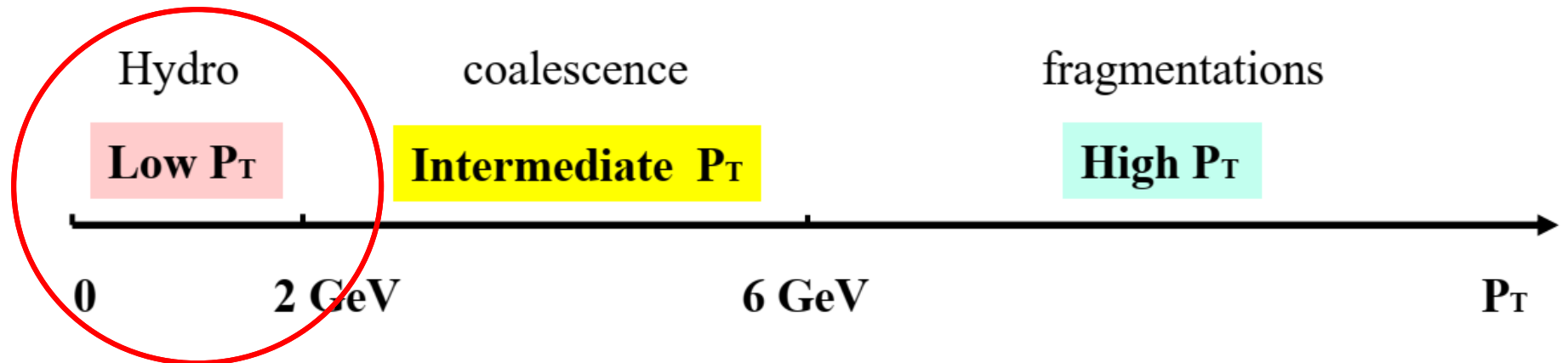
"The truly stunning finding at RHIC that the new state of matter created in the collisions of gold ions is more like a liquid than a gas gives us a profound insight into the earliest moments of the universe," said Dr. Raymond L. Orbach, Director of the DOE Office of Science.

Also of great interest to many following progress at RHIC is the emerging connection between the collider's results and calculations using the methods of string theory, an approach that attempts to explain

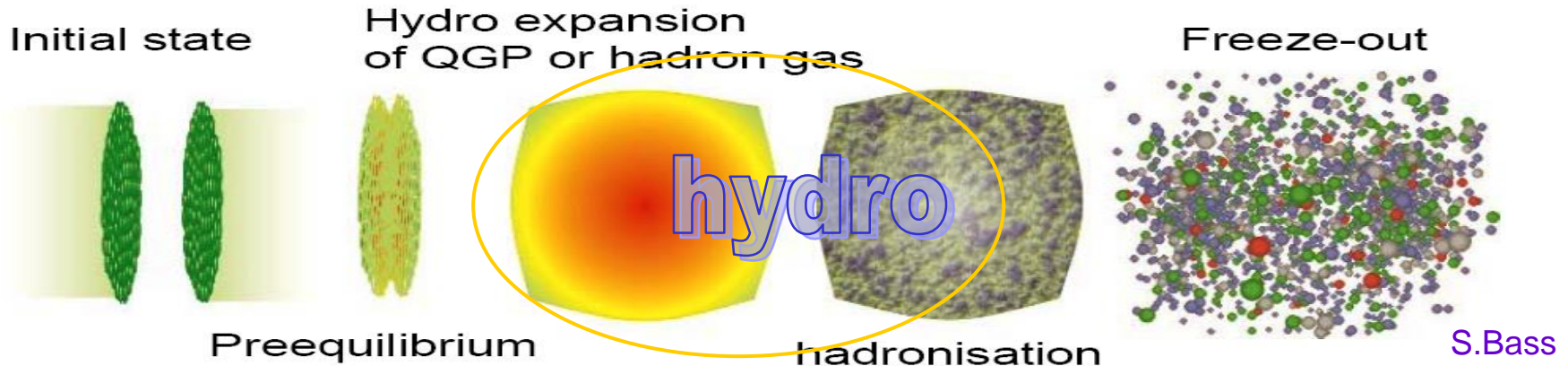


Secretary of Energy
Samuel Bodman

Hydrodynamics & collective flow in large systems (Au+Au & Pb+Pb collisions)



Viscous hydrodynamics



Conservation laws:

$$\partial_\mu T^{\mu\nu}(x) = 0, \quad \partial_\mu N_i^\mu(x) = 0,$$

2nd order I-S equ:

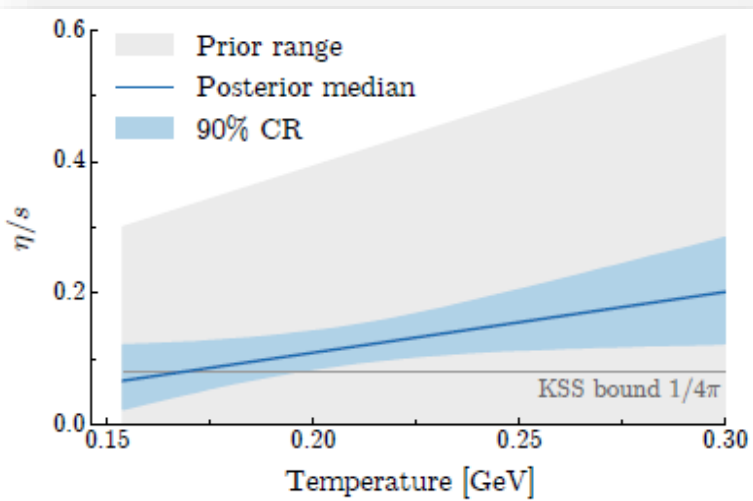
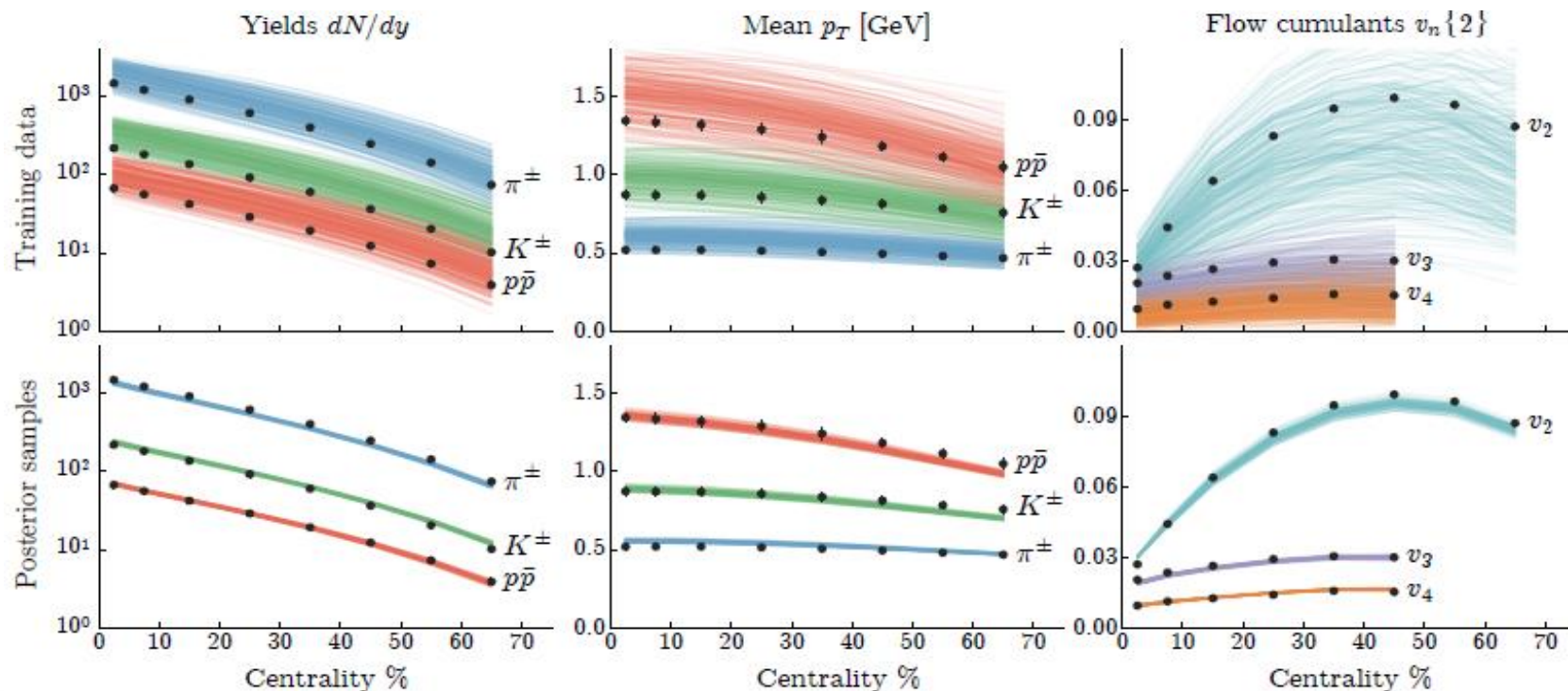
$$\dot{\Pi} = -\frac{1}{\tau_\Pi} \left[\Pi + \zeta\theta - l_{\Pi q} \nabla_\mu q^\mu + \Pi \zeta T \partial_\mu \left(\frac{\tau_\Pi u^\mu}{2\zeta T} \right) \right],$$

$$\Delta_\nu^\mu \dot{q}^\nu = -\frac{1}{\tau_q} \left[q_\mu + \lambda \frac{nT^2}{e+p} \nabla^\mu \frac{\nu}{T} + l_{q\pi} \nabla_\nu \pi^{\mu\nu} + l_{q\Pi} \nabla^\mu \Pi - \lambda T^2 q^\mu \partial_\mu \left(\frac{\tau_q u^\mu}{2\lambda T^2} \right) \right],$$

$$\Delta^{\mu\alpha} \Delta^{\nu\beta} \dot{\pi}_{\alpha\beta} = -\frac{1}{\tau_\pi} \left[\pi^{\mu\nu} - 2\eta \nabla^{\langle\mu} u^{\nu\rangle} - l_{\pi q} \nabla^{\langle\mu} q^{\nu\rangle} + \pi_{\mu\nu} \eta T \partial_\alpha \left(\frac{\tau_\pi u^\alpha}{2\eta T} \right) \right],$$

Input: “EOS” $\varepsilon = \varepsilon(p)$ initial and final conditions

Extract QGP viscosity with hydrodynamics



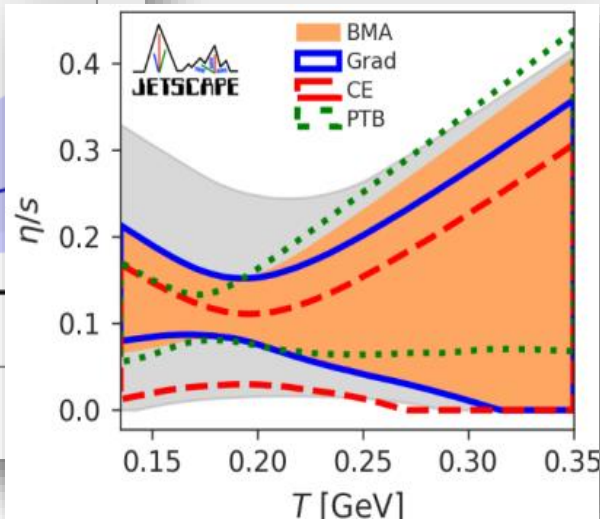
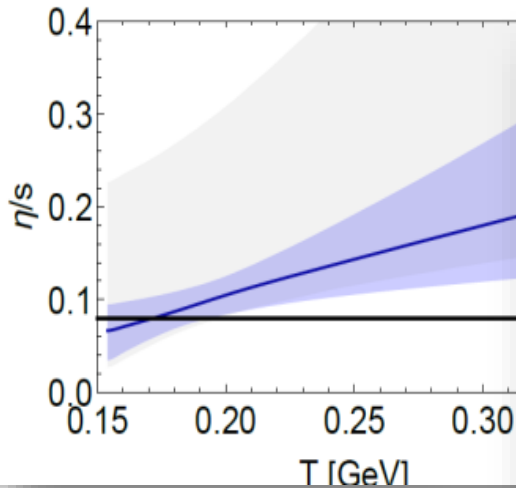
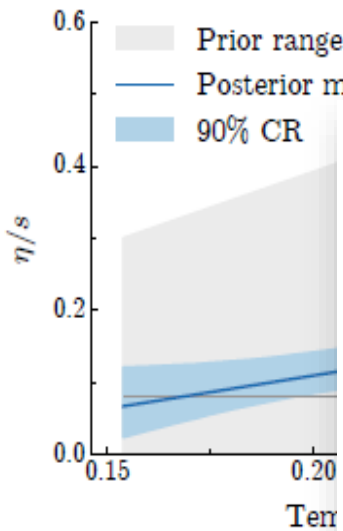
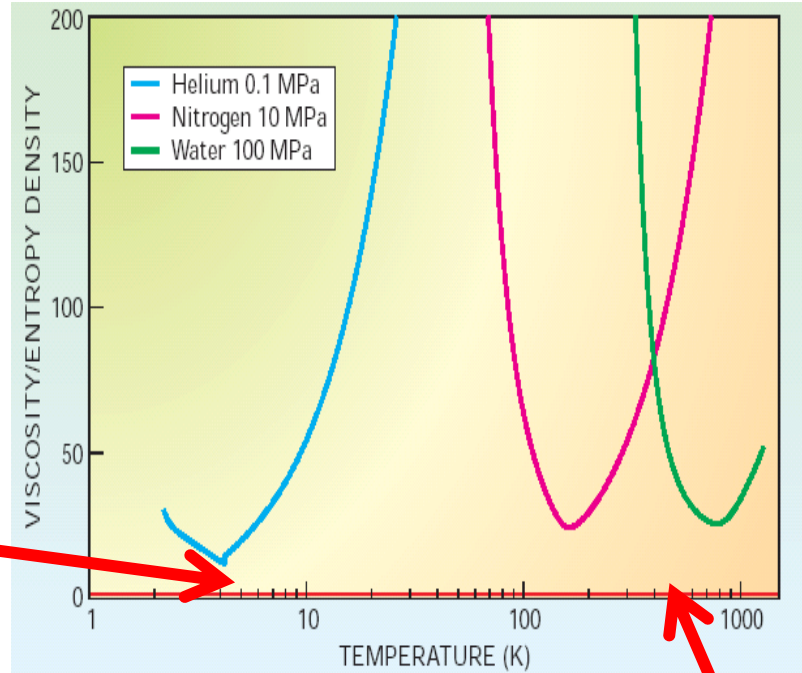
-An quantitative extraction of the QGP viscosity with iEBE-VISHNU and the massive data evaluation
 - $\eta/s(T)$ is very close to the KSS bound of $1/4\pi$

J. Bernhard, S. Moreland, S.A. Bass, J. Liu, U. Heinz, PRC 2015

QGP: most perfect liquid



AdS/CFT $\frac{\eta}{s} \geq \frac{h}{k_B}$



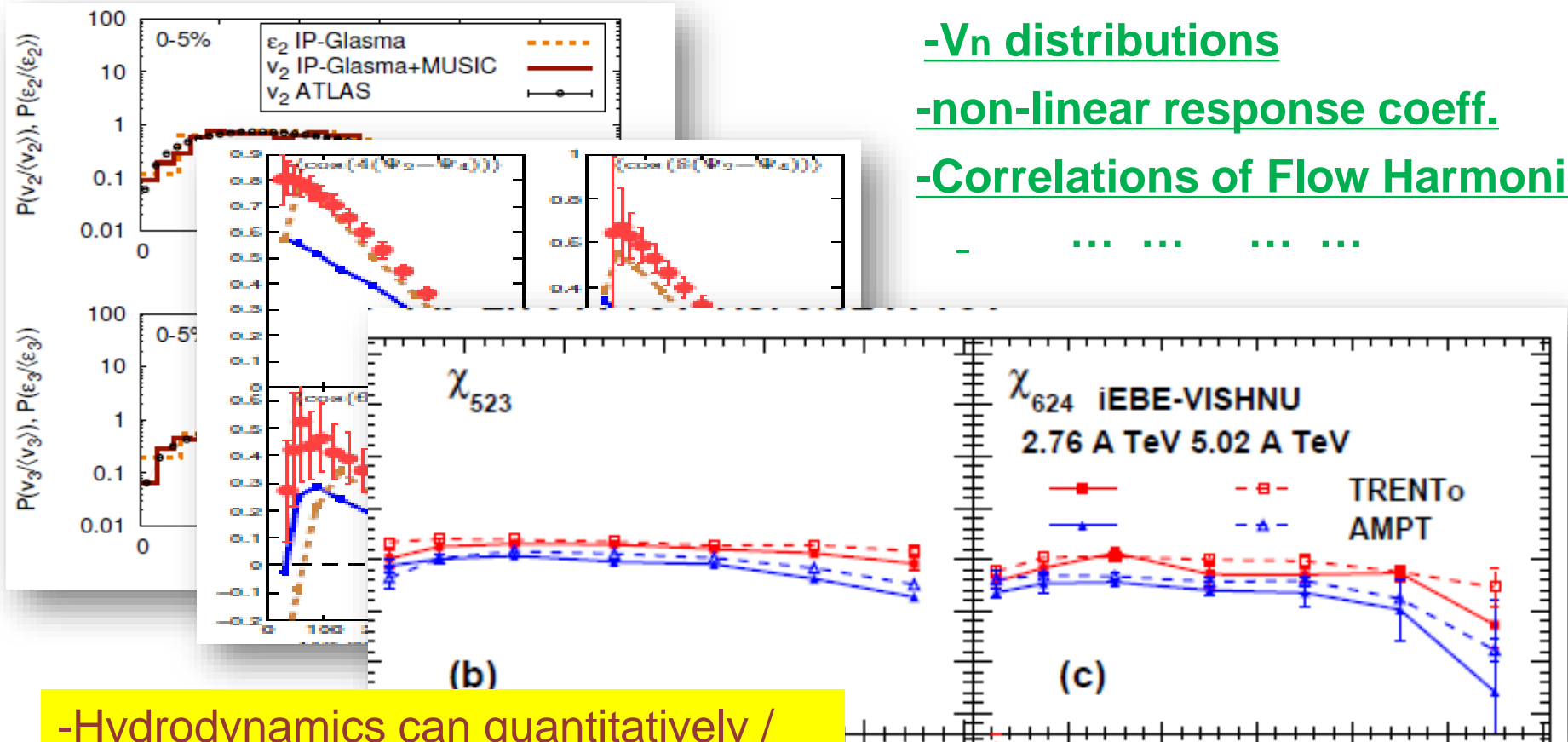
**QGP specific
shear viscosity
Extracted from
exp data**

Powerful predictions from hydrodynamics

- V_n distributions

-non-linear response coeff.

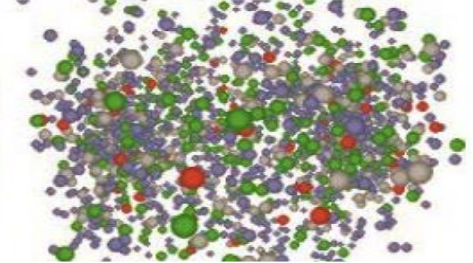
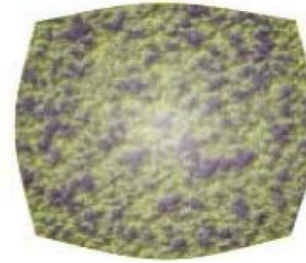
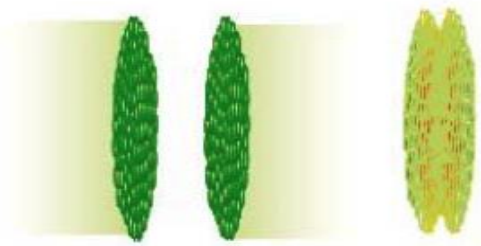
-Correlations of Flow Harmonics



-Hydrodynamics can quantitatively / qualitatively describe / predict various flow data

H. Xu, Z. Li and H. S*, Phys. Rev. C93, no. 6, 064905 (2016); W. Zhao, H. Xu and H. S*, Eur. Phys. J. C 77, no. 9, 645 (2017); X. Zhu, Y. Zhou, H. Xu and H. S*, Phys. Rev. C95, no. 4, 044902 (2017); W. Zhao, L. Zhu, H. Zheng, C. M. Ko and H. S*, Phys. Rev. C 98, no. 5, 054905 (2018); Li, Zhao, Zhou, H.S*, in preparation (2020)

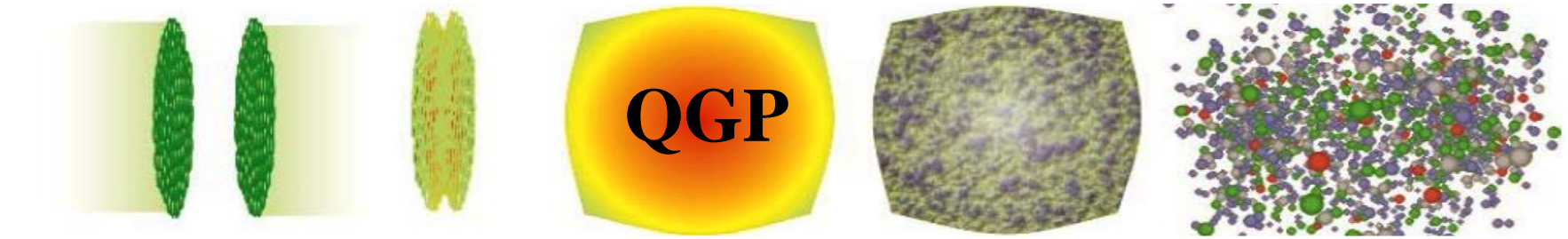
The QGP has been created in relativistic heavy ion collisions



Hottest Matter on Earth



The QGP has been created in relativistic heavy ion collisions



Hottest Matter on Earth



How tiny the QGP droplet could be?

Correlations & Flow in small systems

System size scan:

Pb+Pb Xe+Xe O+O p-p collisions ...

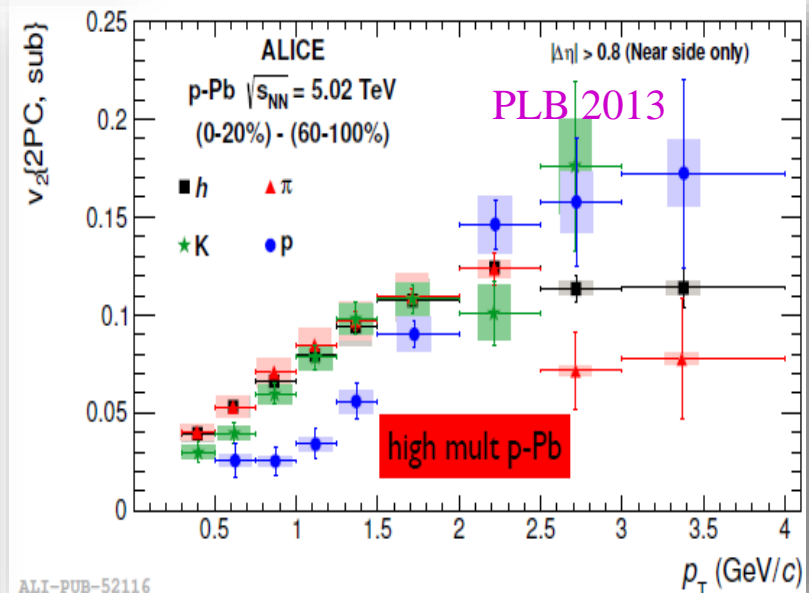
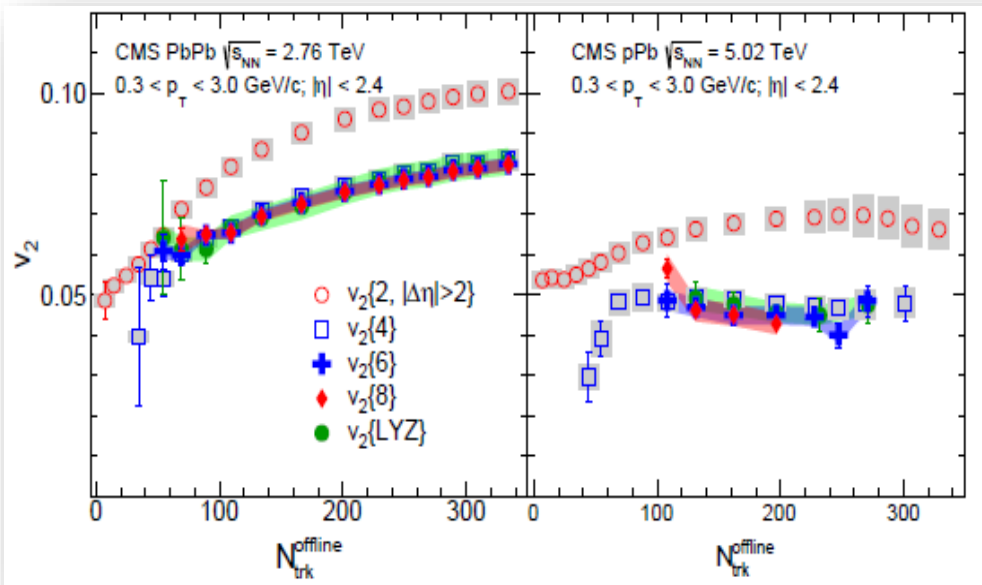
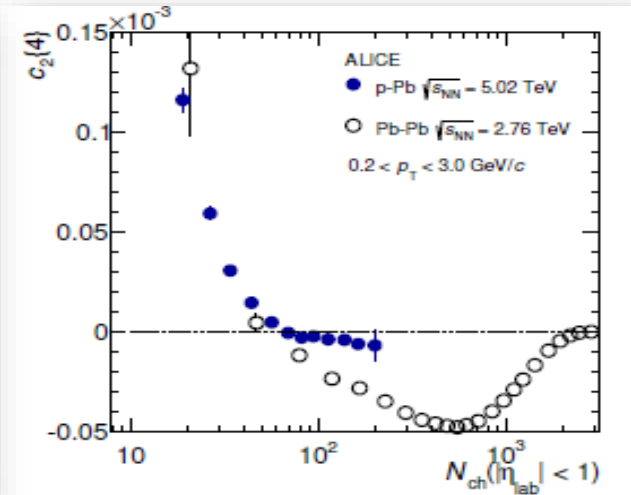
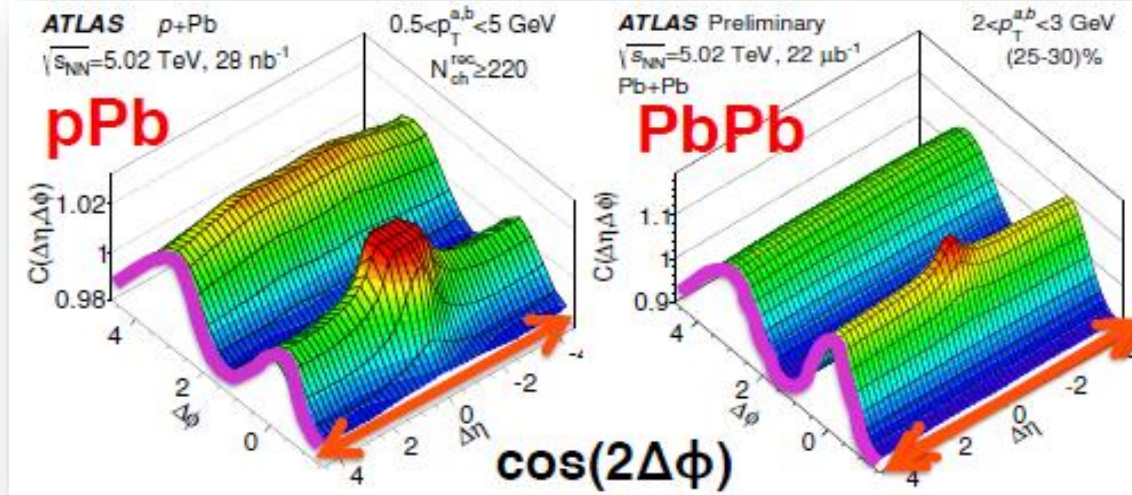
Geometry scan:

p-Au d+Au He-Au

Other collision systems:

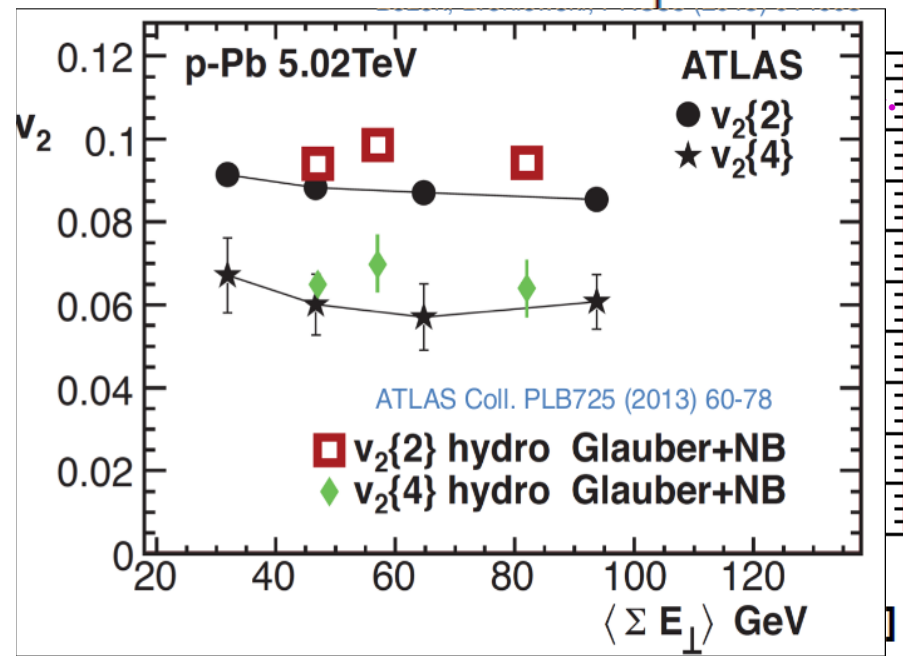
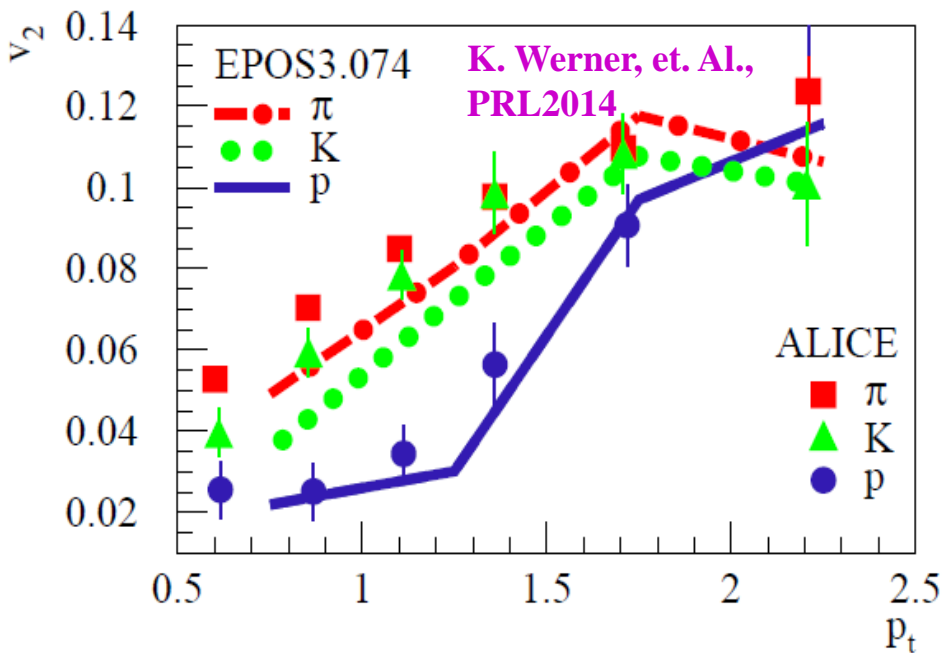
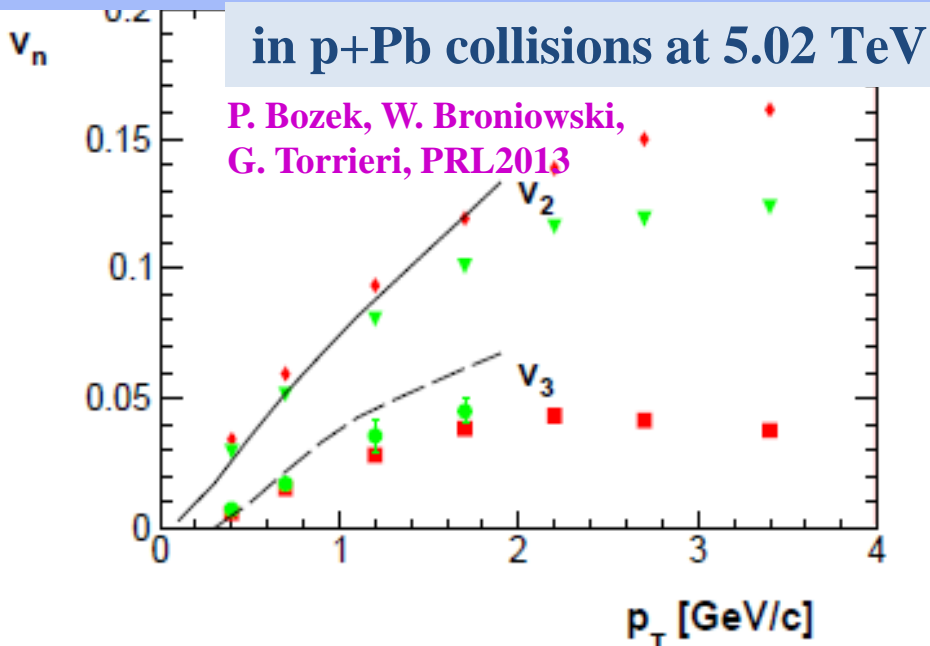
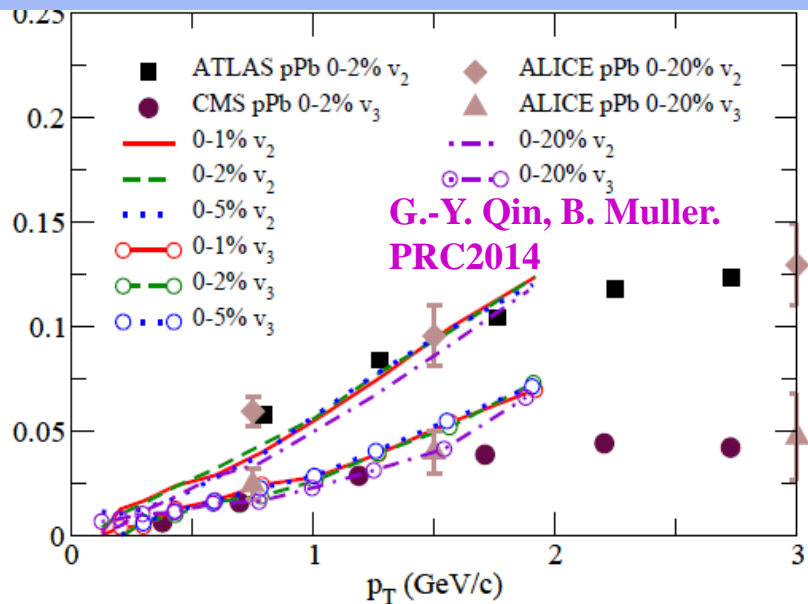
OBSERVABLES	A–A	p–A (high mult.)	pp (high mult.)	pp (low mult.)	UPC	ep	e ⁺ e ⁻ (high mult.)	e ⁺ e ⁻
Near-side ridge yield	✓ [1,2]	✓ [30,32,33]	✓ [30,31]	✓ [34]	—	✗ [74,75]	✓ [77]	✗ [76]
Anisotropic flow	✓ [3,4]	✓ [36,37,38,39]	✓ [35,37]	✓ [30]	✓ [72,73]	✗ [74,75]	✓ [77]	—
Multiparticle cumulants	✓ [5]	✓ [40-45]	✓ [40,41,45]	—	—	—	—	—
Mass ordering	✓ [6]	✓ [47-49]	✓ [46,48]	—	—	—	—	—

Correlations & Flow in p-Pb collisions

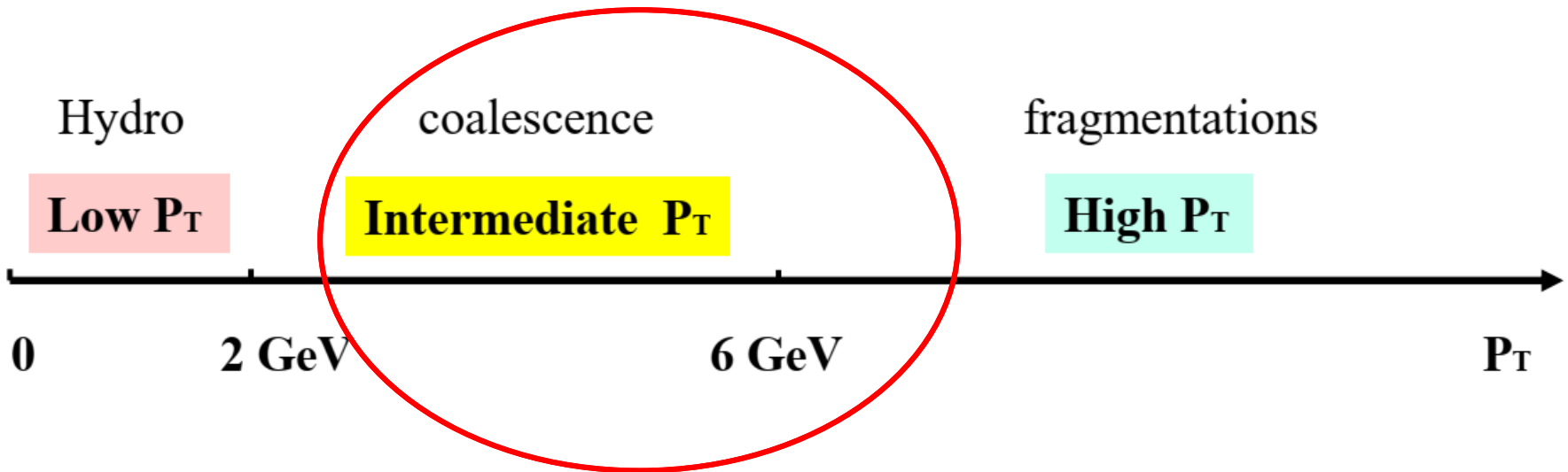


-Many flow-like signals have been observed in high multiplicity p-Pb collisions

Flow in p-Pb -- Hydrodynamics Simulations



Partonic flow and QGP signals in large & small systems



NCQ scaling of v_2 in p-Pb collisions

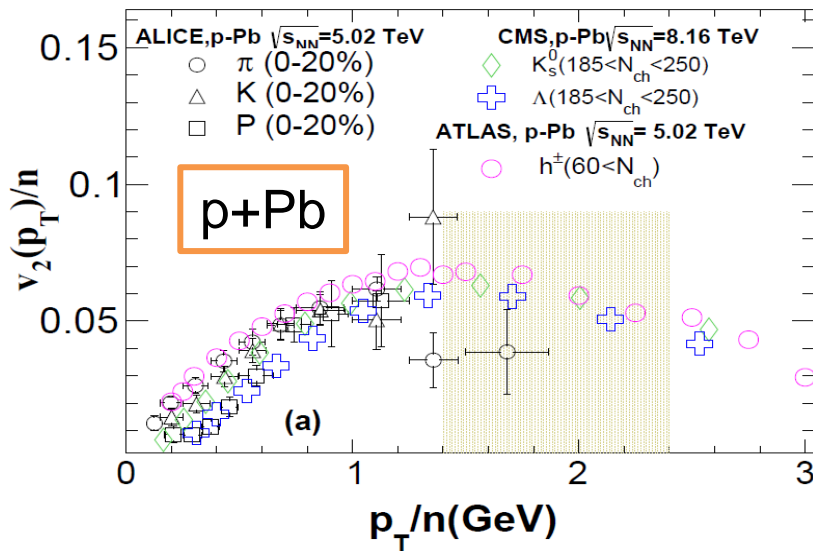
p+Pb collisions

- Where does such approximate NCQ scaling of v_2 come from?
- Is it an indication of partonic degree of freedom?

ALICE data: PLB,726,164 (2013).

CMS data: PRL, 121, 082301 (2018).

ATLAS data: PRC, 96, 024908 (2017).



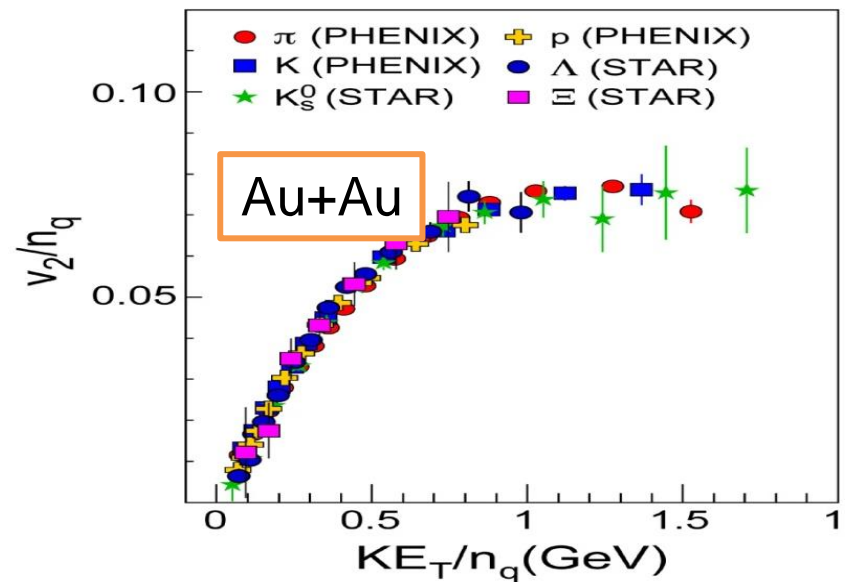
Au+Au collisions

QGP was discovered @RHIC

-strong elliptic flow

-jet quenching

-VCQ scaling of elliptic flow



Simple coalescence for large systems (Au+Au)

- thermal - thermal parton coalescence

$$\frac{dN_M}{d^3\mathbf{P}_M} = g_M \int d^3\mathbf{x}_1 d^3\mathbf{p}_1 d^3\mathbf{x}_2 d^3\mathbf{p}_2 f_q(\mathbf{x}_1, \mathbf{p}_1) f_{\bar{q}}(\mathbf{x}_2, \mathbf{p}_2) \times W_M(\mathbf{y}, \mathbf{k}) \delta^{(3)}(\mathbf{P}_M - \mathbf{p}_1 - \mathbf{p}_2)$$

$$\frac{dN_B}{d^3\mathbf{P}_B} = g_B \int d^3\mathbf{x}_1 d^3\mathbf{p}_1 d^3\mathbf{x}_2 d^3\mathbf{p}_2 d^3\mathbf{x}_3 d^3\mathbf{p}_3 f_{q_1}(\mathbf{x}_1, \mathbf{p}_1) \times f_{q_2}(\mathbf{x}_2, \mathbf{p}_2) f_{q_3}(\mathbf{x}_3, \mathbf{p}_3) W_B(\mathbf{y}_1, \mathbf{k}_1; \mathbf{y}_2, \mathbf{k}_2) \times \delta^{(3)}(\mathbf{P}_B - \mathbf{p}_1 - \mathbf{p}_2 - \mathbf{p}_3)$$

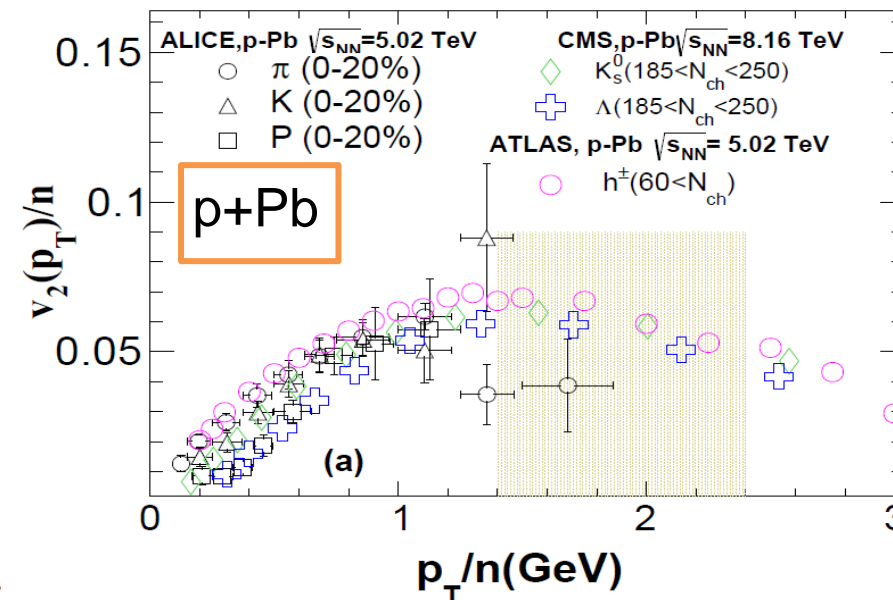
Complicated coalescence for small systems (p+Pb, p+p)

- the effects from jet/mini-jets
- fragmentations from high p_T become important

Complicated Coalescence processes:

- thermal - thermal parton coalescence
- thermal - hard parton coalescence
- hard - hard parton coalescence

Frag contributes more at intermediate p_T



Hydro-Coal-Frag Hybrid Model

Thermal hadrons (VISH2+1):

- generated by hydro.
- with Cooper-Frye.
- Meson: $P_T < 2P_1$; baryon: $P_T < 3P_1$.

Coalescence hadrons (Coal Model):

- generated by coalescences model including thermal-thermal, thermal-hard & hard-hard parton coalescence.

Fragmentation hadrons (LBT):

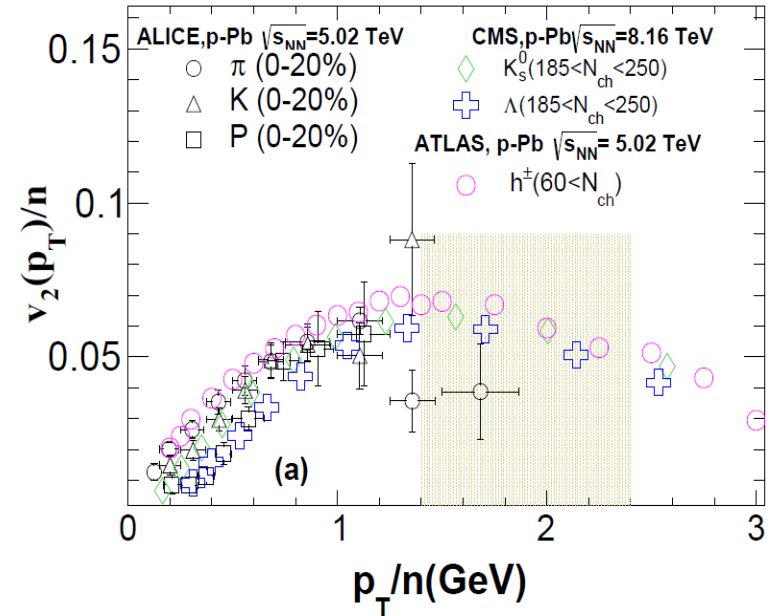
- Hard partons generated by PYTHIA8, then suffered energy loss by LBT

UrQMD afterburner:

- All hadrons are feed into UrQMD for hadronic evolution, scatterings and decays.

Zhao, Ko, Liu, Qin & Song.

Phys. Rev. Lett. 125 7 072301(2020)_{in}



Main Parameters:

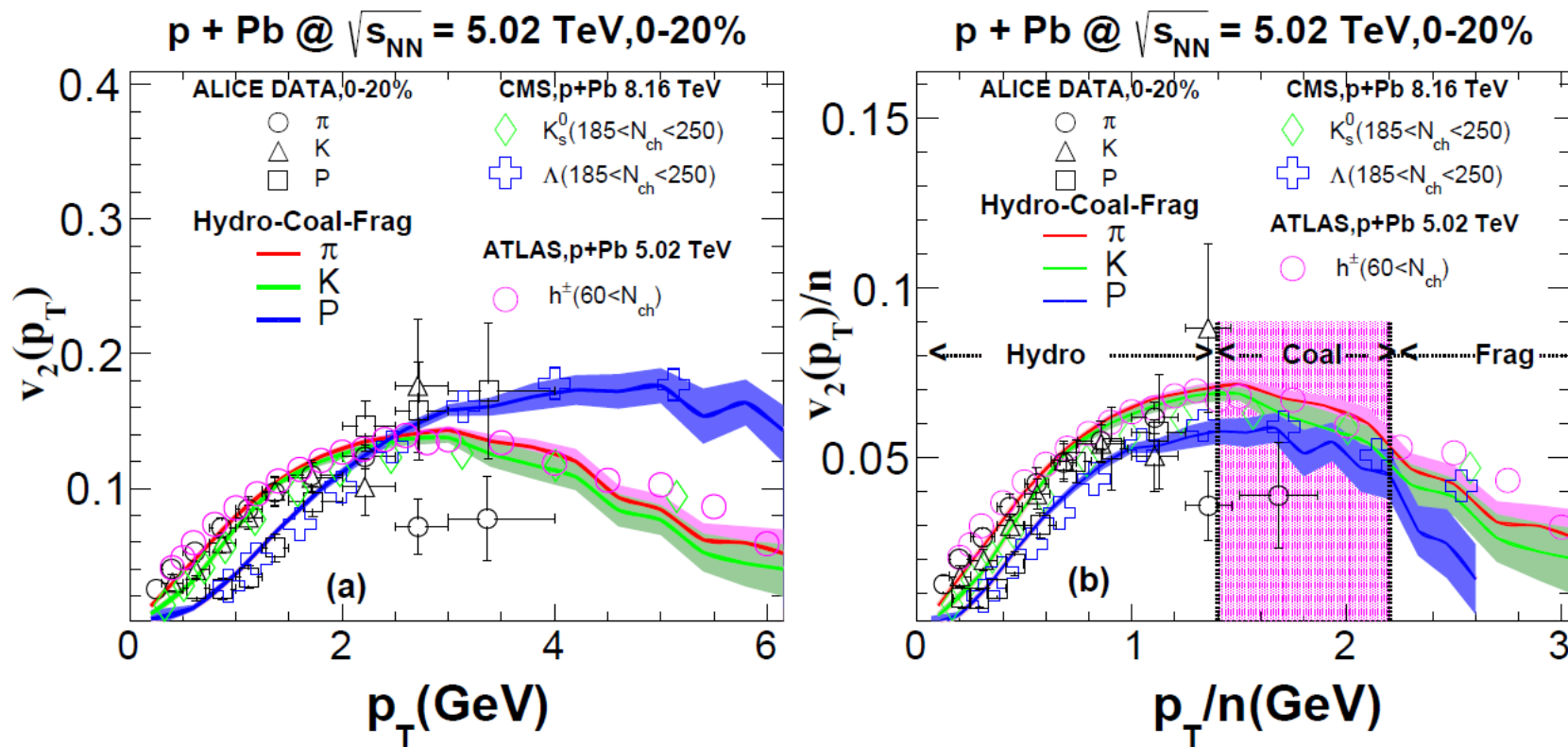
-Thermal hadrons from hydro with $P_T < P_1$.

-Hard partons from LBT with $P_T > P_2$.

Fixed by the p_T spectra

$P_{T1} = 1.6$ GeV and $P_{T2} = 2.6$ GeV

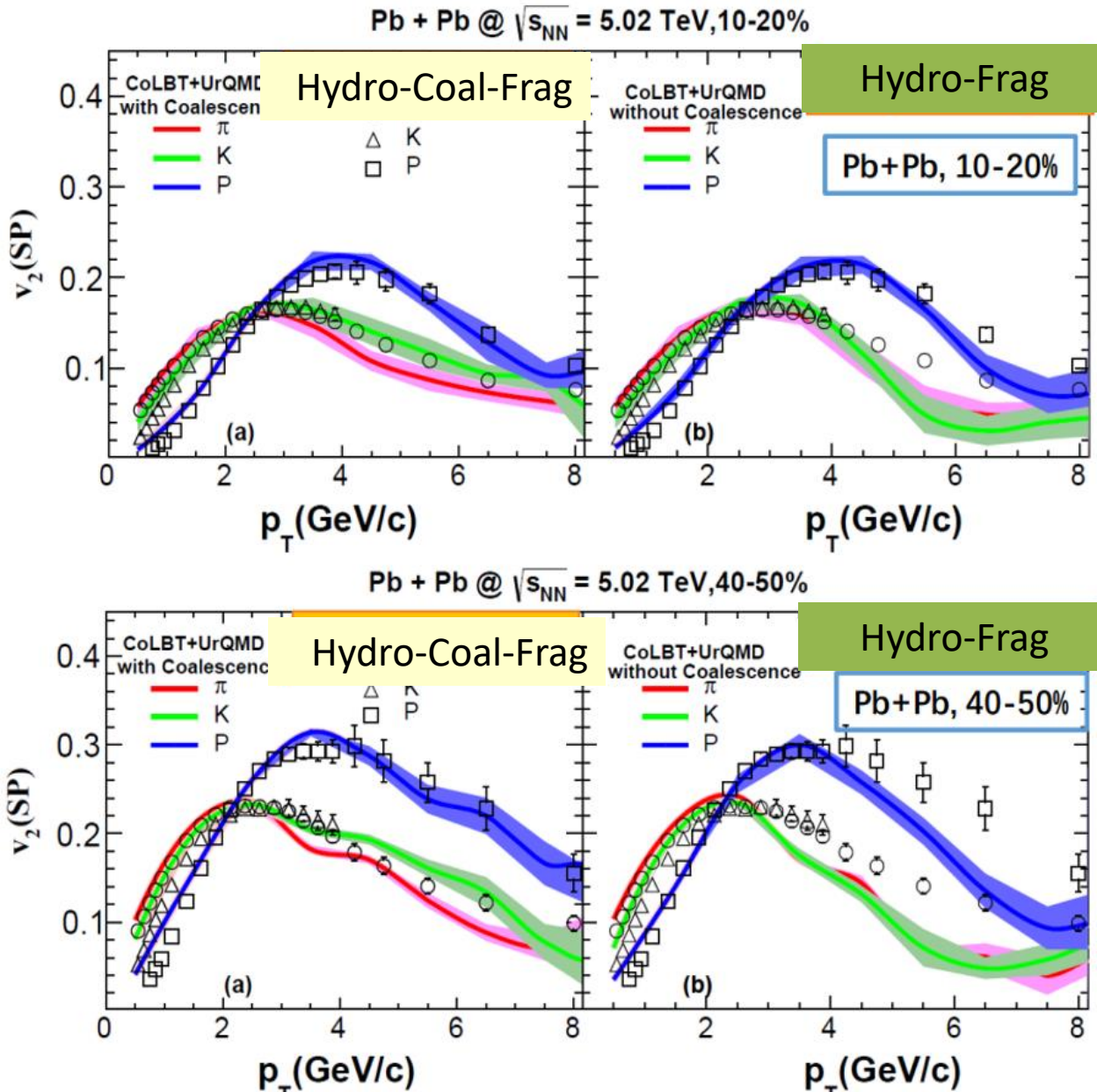
VCQ scaling of v_2 & hint partonic degree of freedom



-At intermediate p_T , Hydro-Coal-Frag model obtains an approximate NCQ scaling as shown by the data.

Strongly indication of partonic degree of freedom in small system !

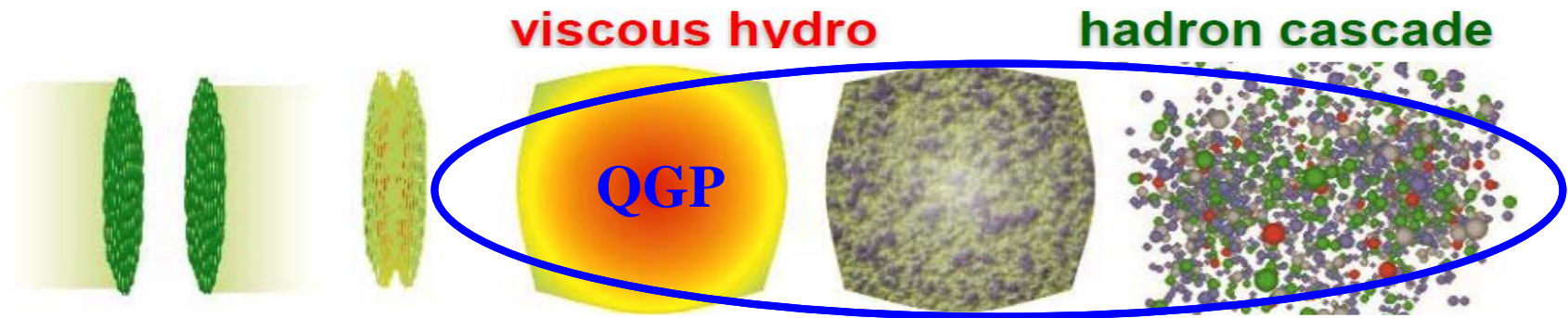
V_2 & the importance of quark coalescence



-CoLBT-hydro with coalescence works well for PID flow of Pb+Pb collisions from 0 to 8 GeV. **Quark coalescence is important at intermediate P_T**

thermal-hard parton Coalescence & Fragmentation Breaks up the NCQ scaling of v_2 in Pb+Pb collisions

Zhao, Chen, Luo, Ke & Wang. Phys. Rev. Lett. 128 2 022302(2022).



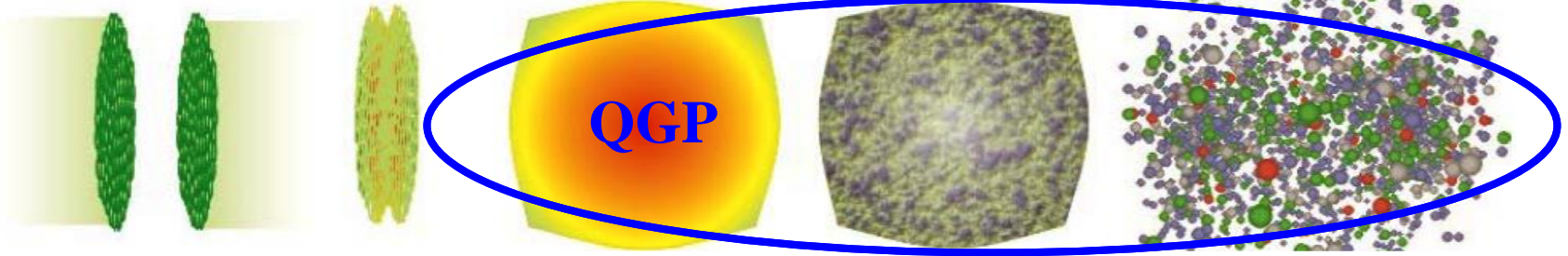
Theory: Hydrodynamics & hybrid approach are powerful tool to simulate the QGP fireball evolution and study its properties



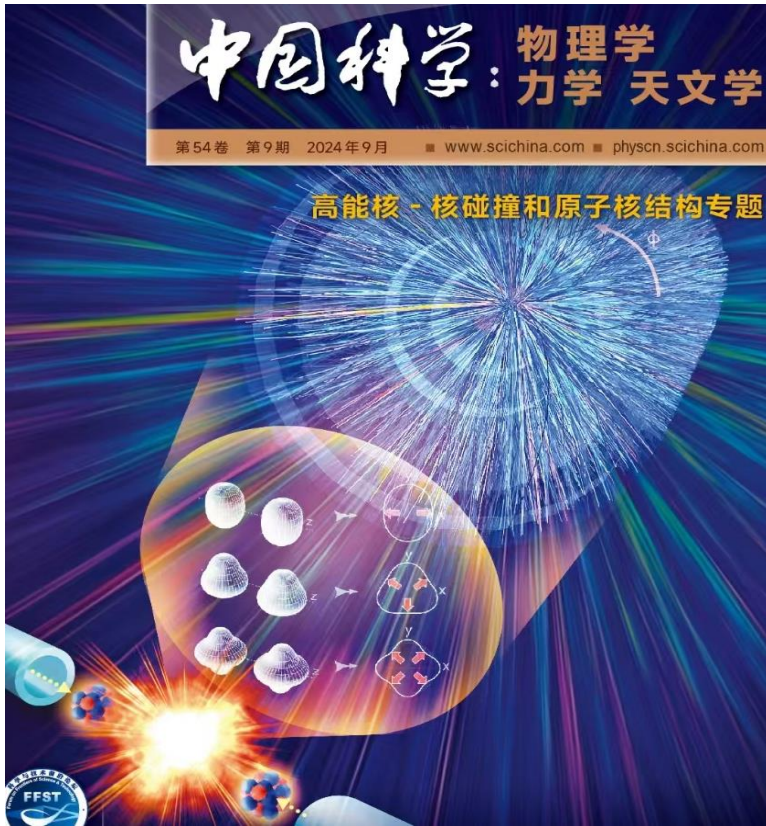
Experiment: various flow observable have been measured quantitatively described/predicted by hydro

viscous hydro

hadron cascade



Theory: Hydrodynamics & hybrid approach are powerful tool to simulate the QGP fireball evolution and study its properties



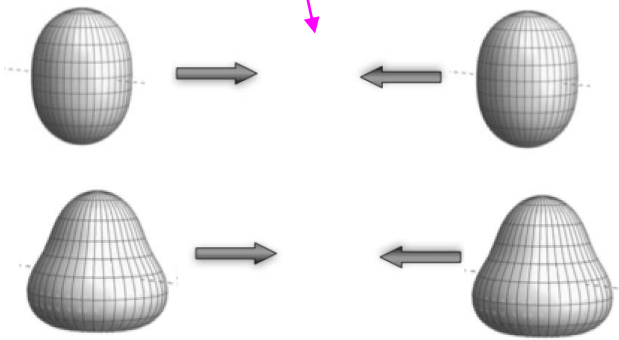
-We are ready to focus on the initial state of the QGP

nuclear structure of colliding nuclei

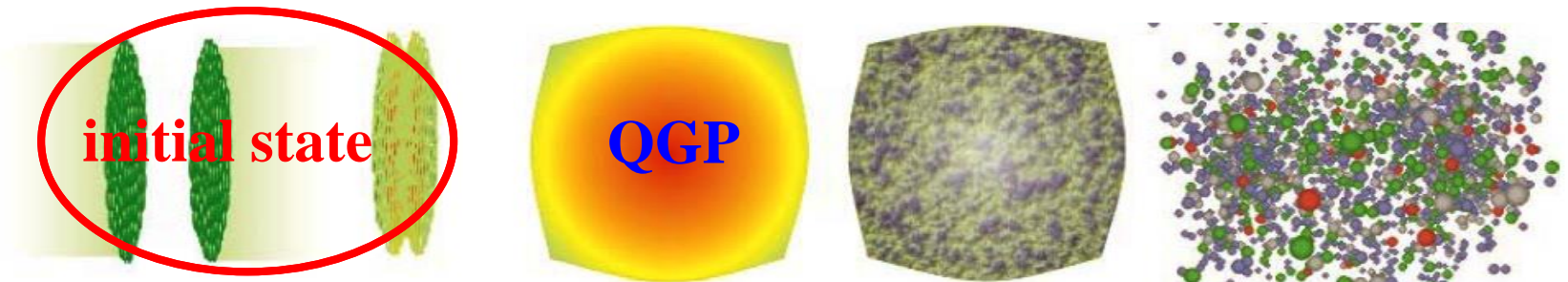


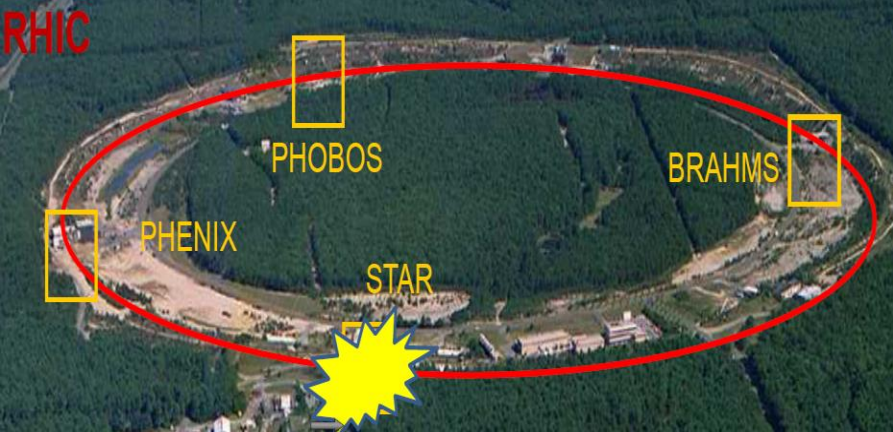
Probe nuclear structure with relativistic heavy ion collisions

- Relativistic heavy collisions start from nuclei



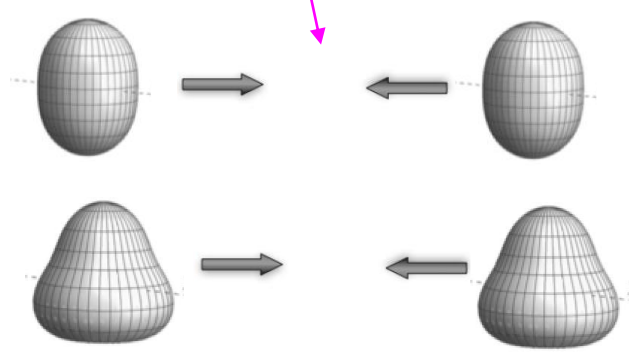
initial state with deformation



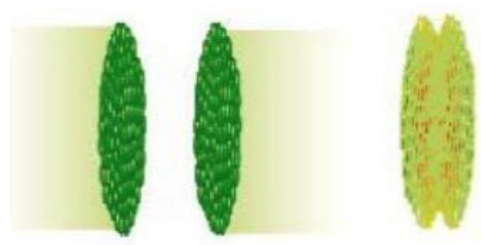


Probe nuclear structure with relativistic heavy ion collisions

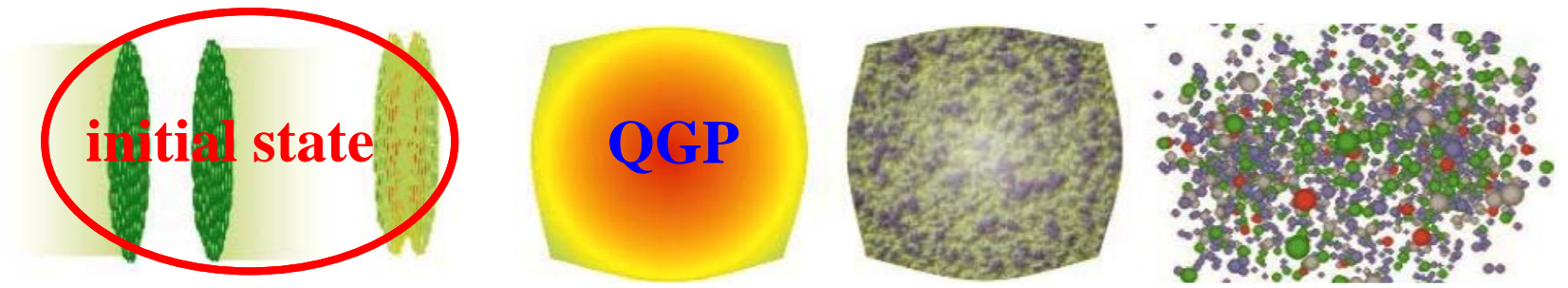
- Relativistic heavy collisions **start from nuclei**
- Collision time $< 10^{-24}$ s **directly probe the ground state of nuclei**

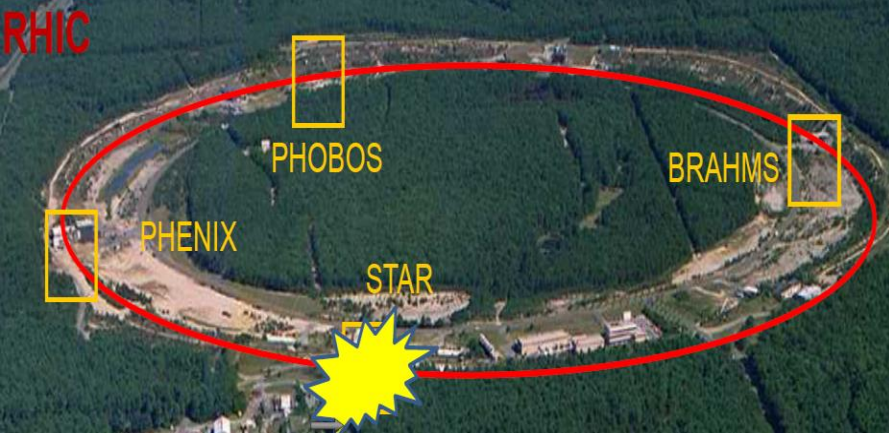


initial state with deformation



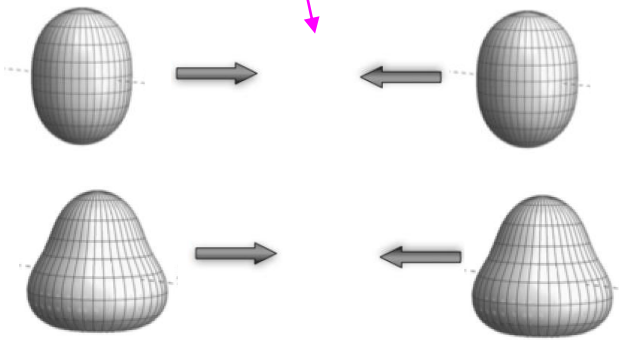
Collision time $< 10^{-24}$ s



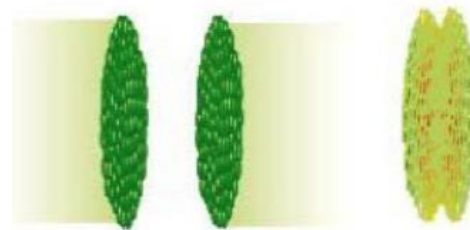


Probe nuclear structure with relativistic heavy ion collisions

- Relativistic heavy collisions **start from nuclei**
- Collision time $< 10^{-24}$ s **directly probe the ground state of nuclei**

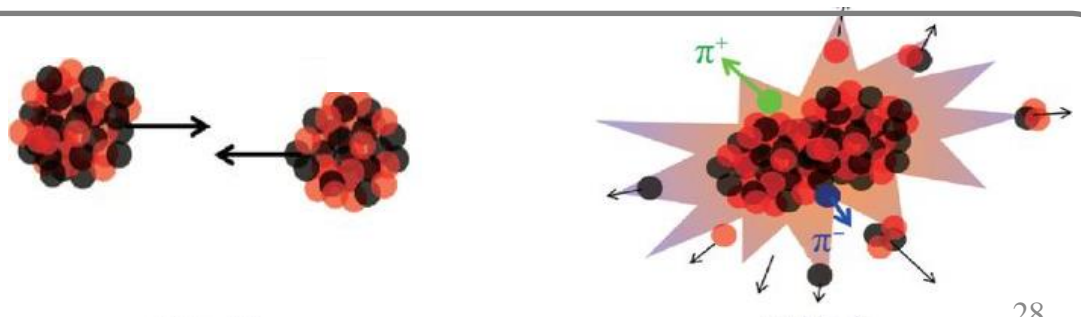


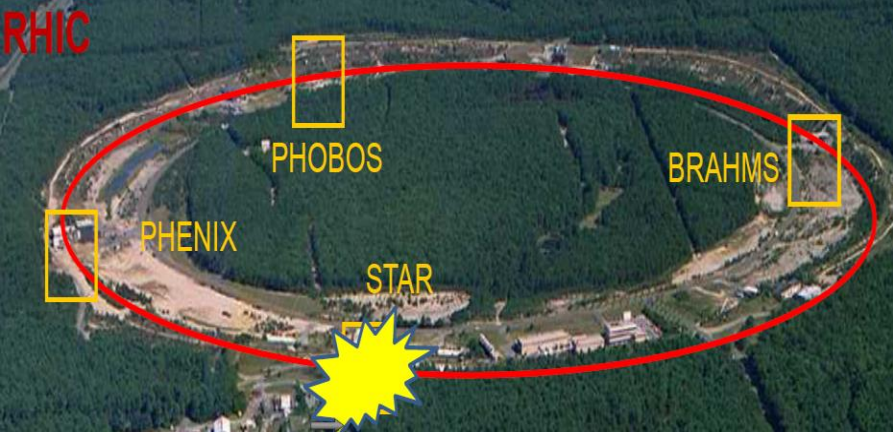
initial state with deformation



Collision time $< 10^{-24}$ s

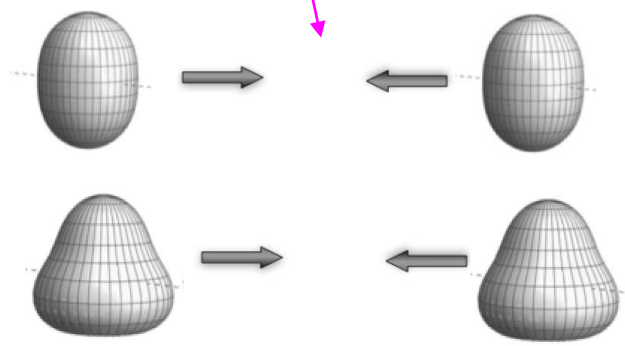
heavy ion collision at intermediate energies breaks up / excites nuclei during the collisions





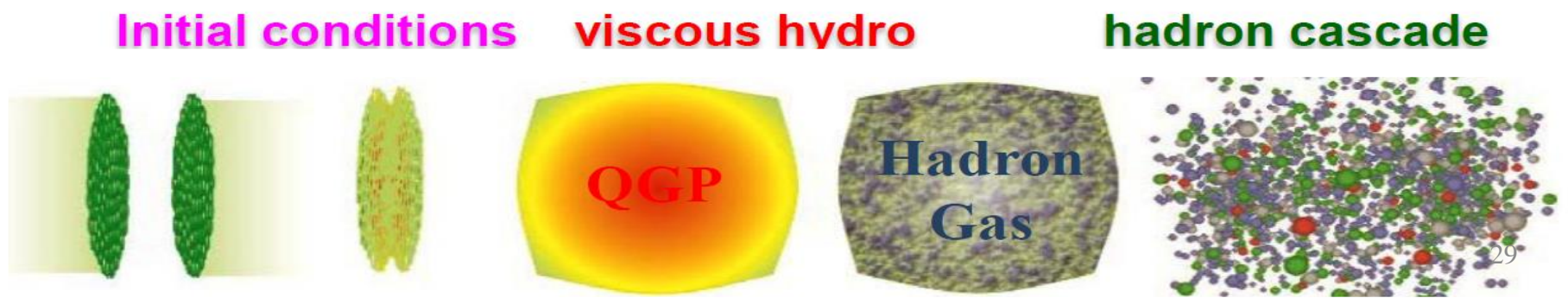
Probe nuclear structure with relativistic heavy ion collisions

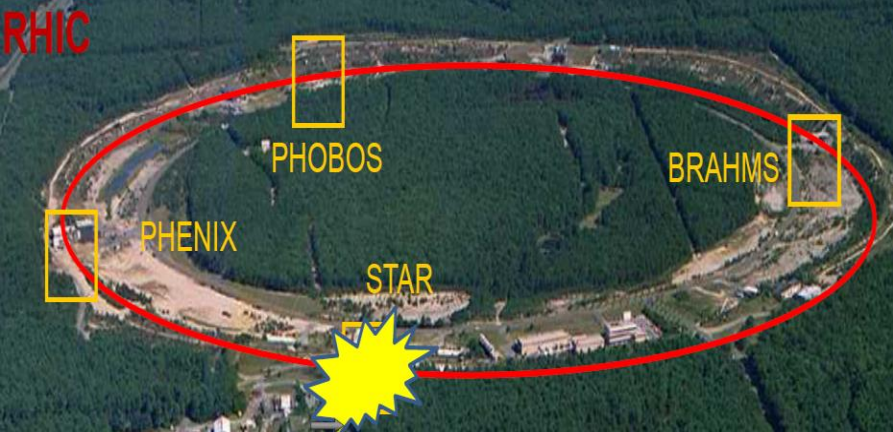
- Relativistic heavy collisions **start from nuclei**
- Collision time $< 10^{-24}$ s directly **probe the ground state of nuclei**
- **Well calibrated calculations for QGP evolution**; to focus on the initial state



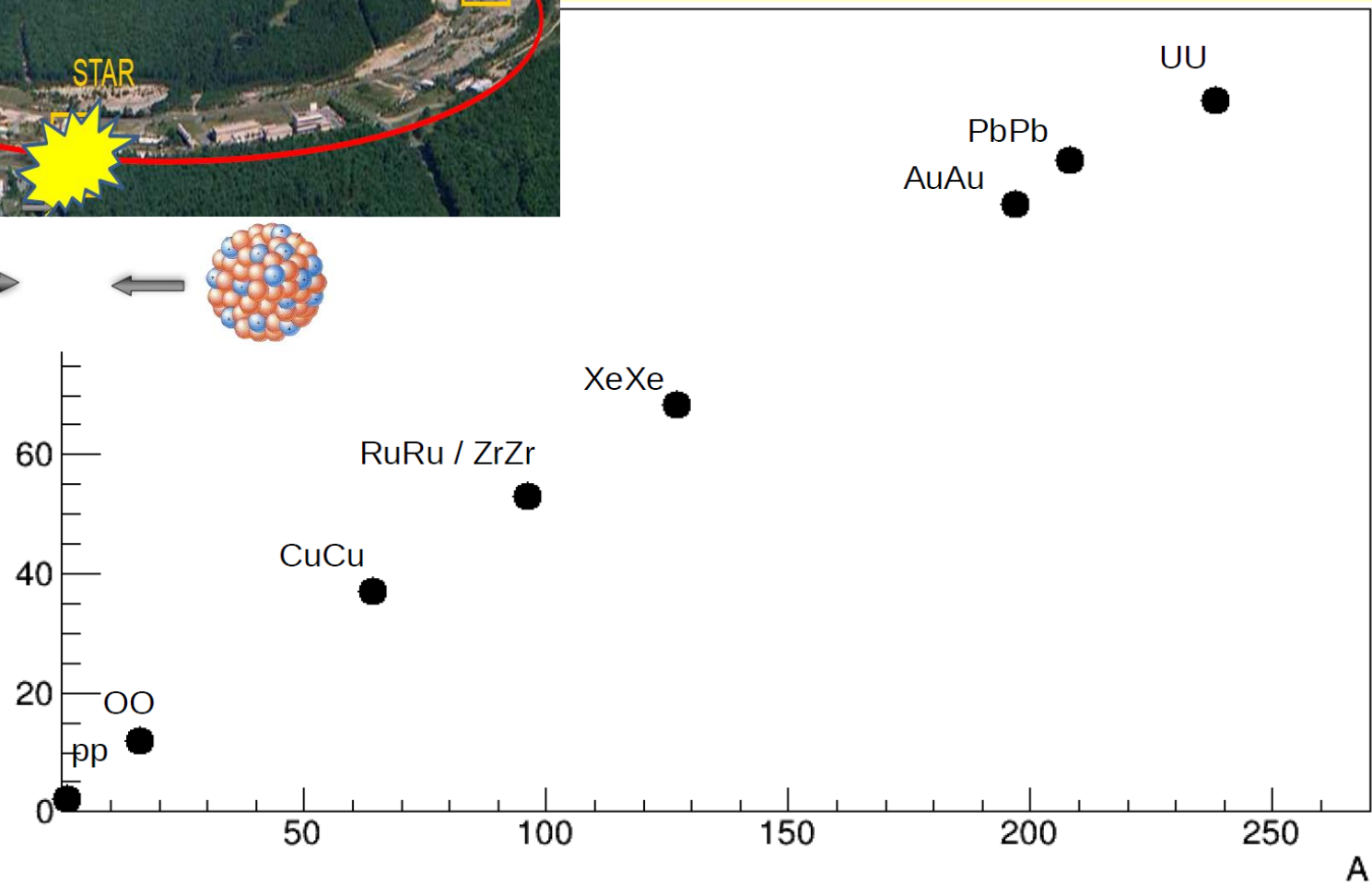
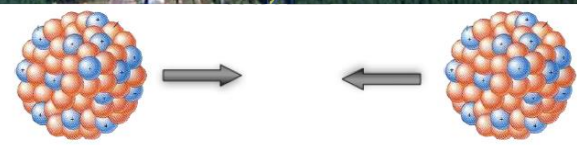
initial state with deformation

Well calibrated calculations



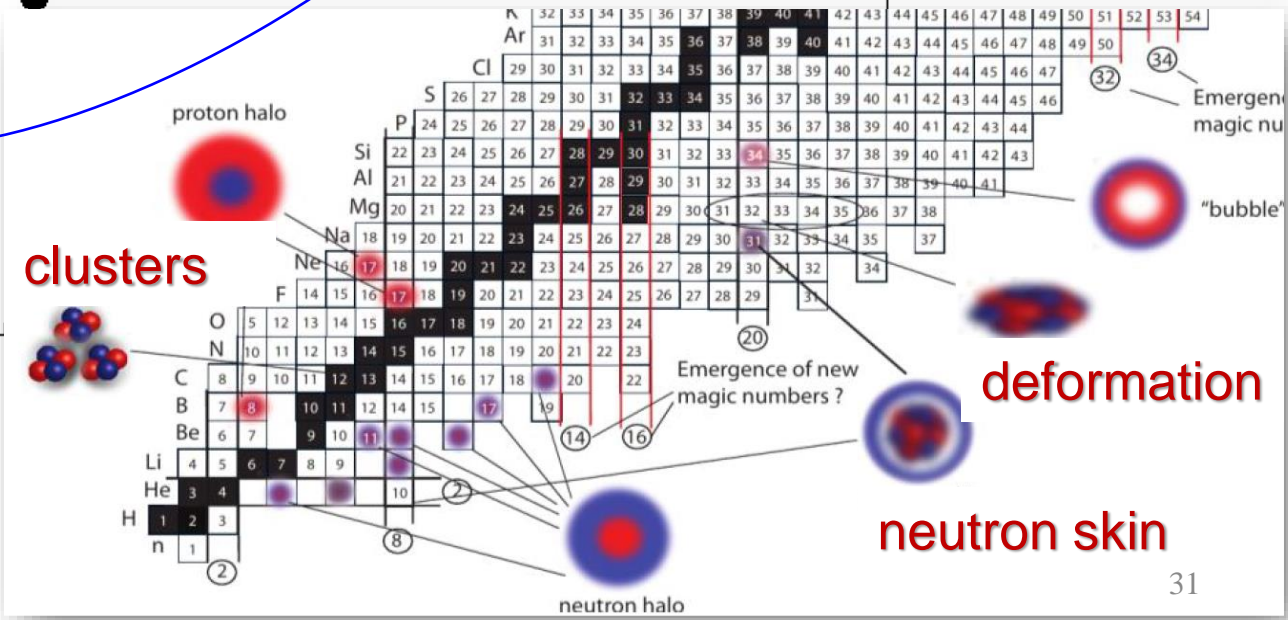
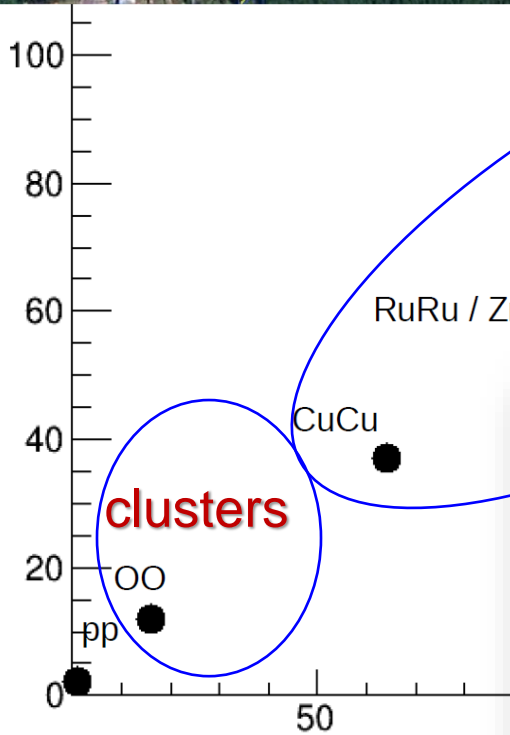
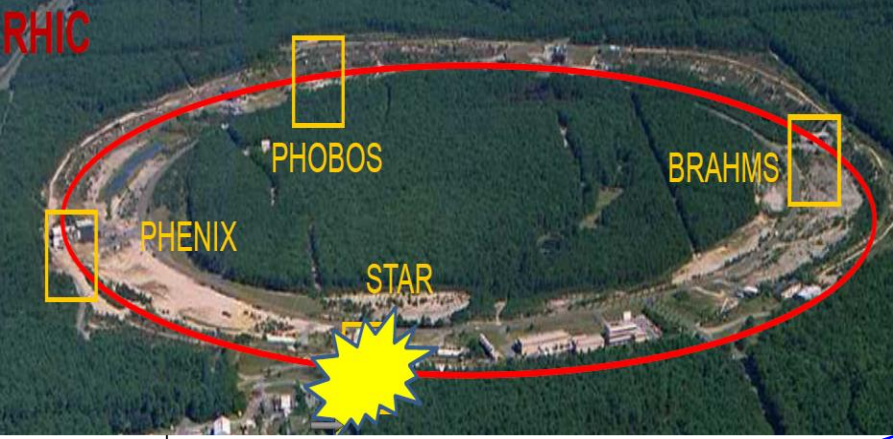


Rich collision systems at RHIC & the LHC

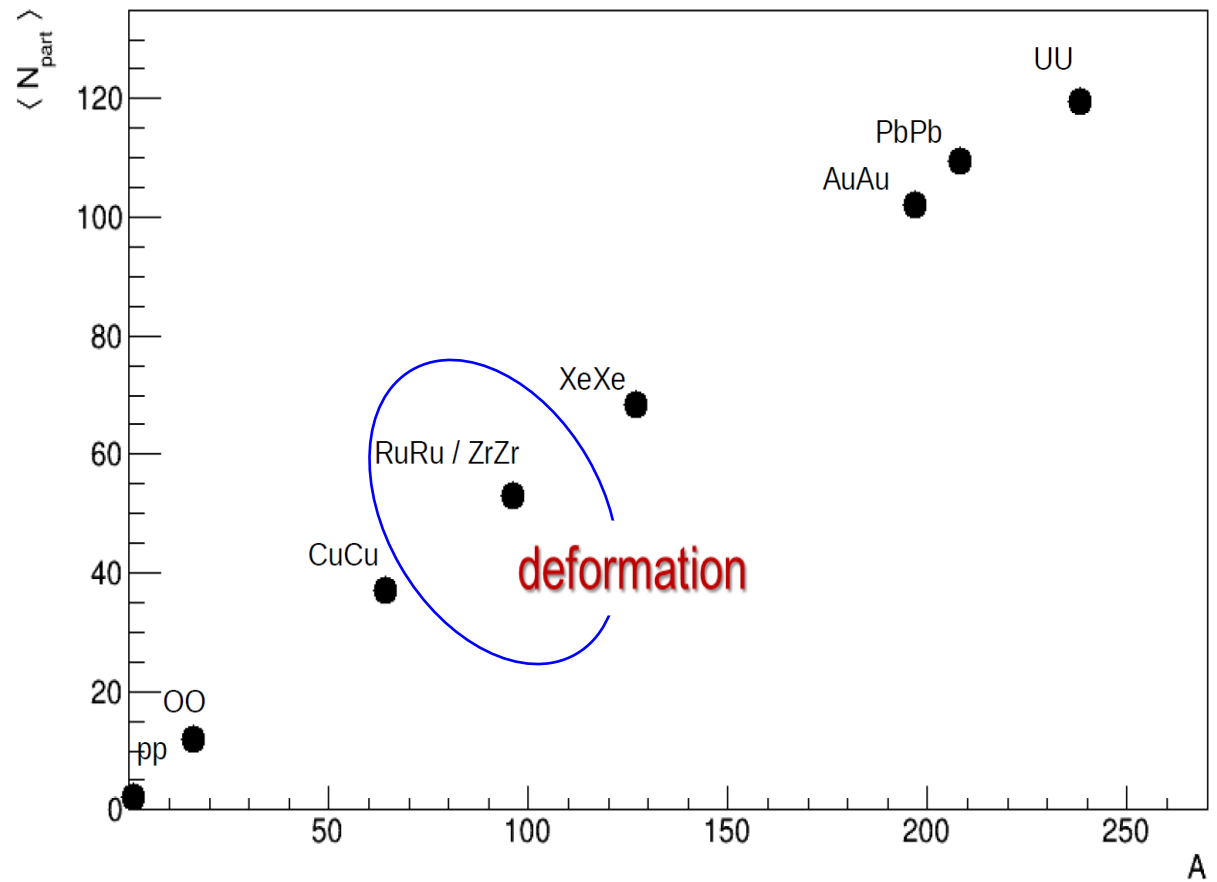
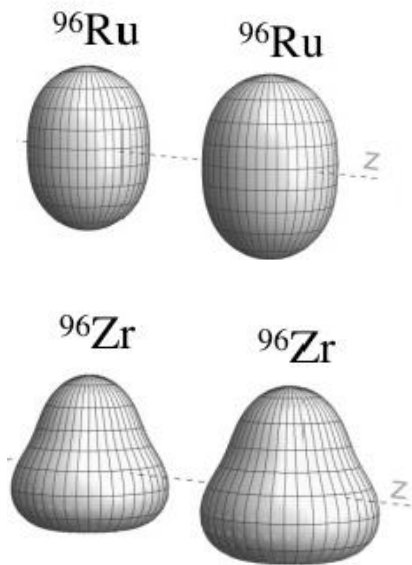


$^{197}\text{Au} + ^{197}\text{Au}$, $^{238}\text{U} + ^{238}\text{U}$, $^{208}\text{Pb} + ^{208}\text{Pb}$, $^{129}\text{Xe} + ^{129}\text{Xe}$, $^{96}\text{Zr} + ^{96}\text{Zr}$,
 $^{96}\text{Ru} + ^{96}\text{Ru}$, $^{64}\text{Cu} + ^{64}\text{Cu}$, $^{16}\text{O} + ^{16}\text{O}$, $p + ^{208}\text{Pb}$, $p + p$

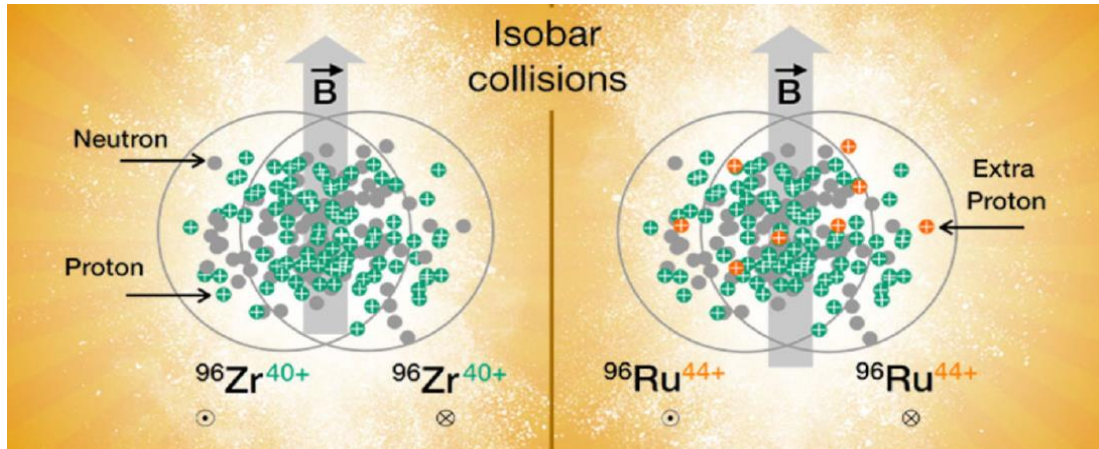
Rich collision systems at RHIC & the LHC



Study the deformation of ^{96}Ru and ^{96}Zr at RHIC isobar run



$^{96}\text{Ru}+^{96}\text{Ru}$ and $^{96}\text{Zr}+^{96}\text{Zr}$ Collisions @ RHIC isobar run



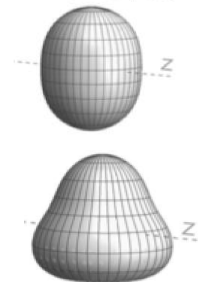
- To search the Chiral Magnetic Effect (CME)

- Obviously different early magnetic field for Ru+Ru and Zr+Zr collisions

Deformation of ^{96}Ru and ^{96}Zr

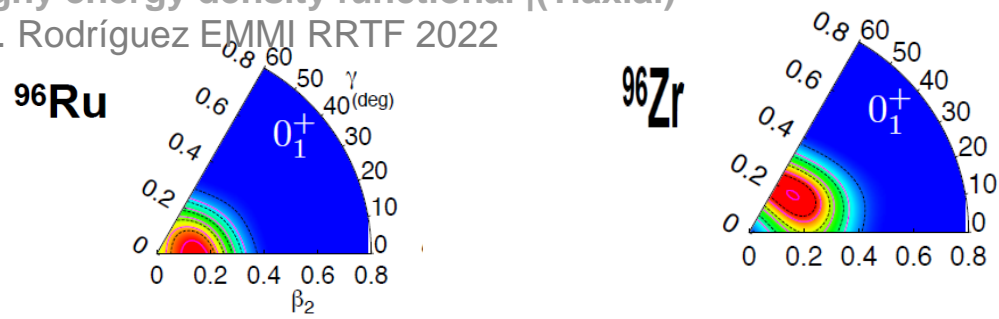
Conversion from $B(E_n)$ to β_n via: $\beta_2 = \frac{4\pi}{3ZR^2} \sqrt{\frac{B(E2)_{\uparrow}}{e^2}}$, $\beta_3 = \frac{4\pi}{3ZR_0^3} \sqrt{\frac{B(E3)_{\uparrow}}{e^2}}$

	β_2	$E_{2^+_1}$ (MeV)	β_3	$E_{3^-_1}$ (MeV)
^{96}Ru	0.154	0.83	-	3.08
^{96}Zr	0.062	1.75	0.202, 0.235, 0.27	1.90

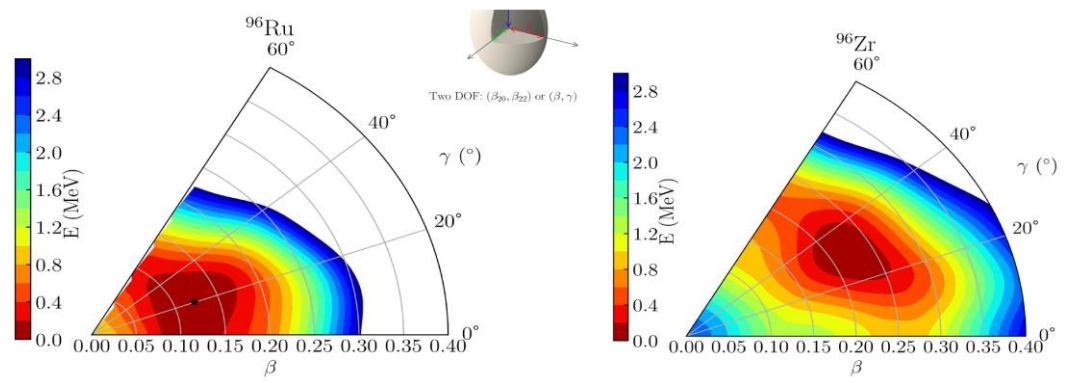


Deformation of ^{96}Ru & ^{96}Zr — personal comments

Gogny energy density functional |(Tiaxial)
T.R. Rodríguez EMMI RRTF 2022

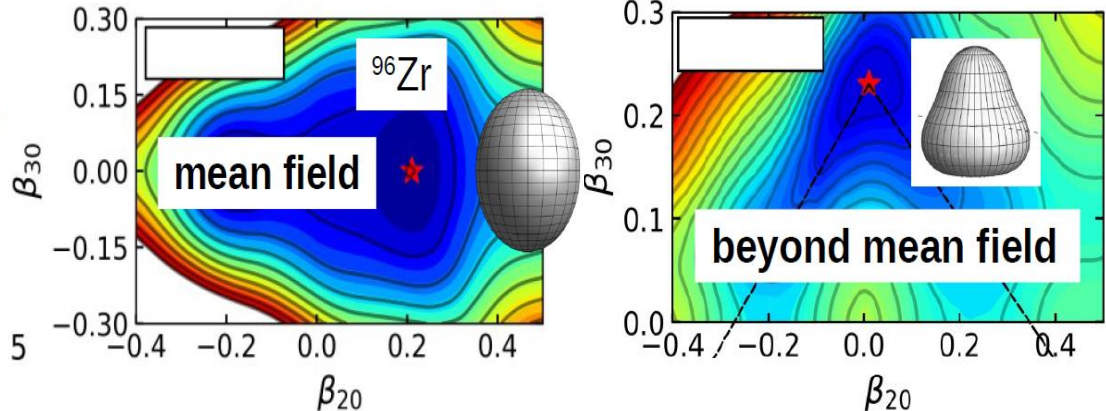


Skyrme EDF W Ryssens EMMI RRTF 2022



18 ⁺	8205.7	15 ⁻	6754.1
16 ⁺	6441.6	13 ⁻	5750.2
14 ⁺	5680.7	11 ⁻	4798.7
12 ⁺	4418.3	9 ⁻	3951.1
10 ⁺	3817.2	7 ⁻	3291.5
8 ⁺	2950.4	5 ⁻	2588.4
6 ⁺	2149.7		
4 ⁺	1518.1		
2 ⁺	832.6		
0 ⁺	0.0		

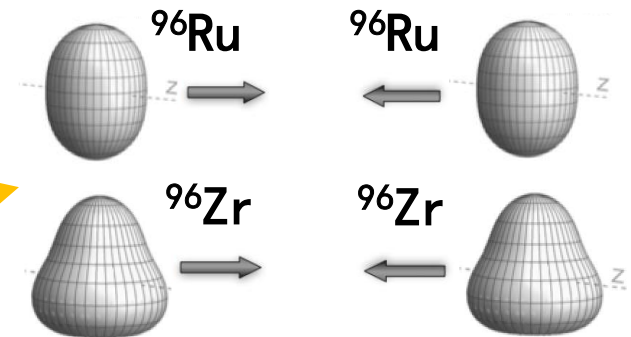
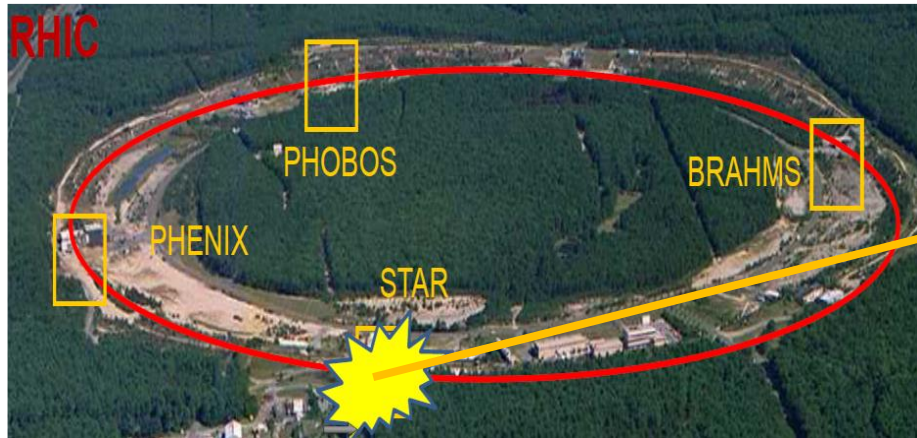
Beyond-mean-field Rong, Lu, arXiv:2201.02114



Nuclear structure physics obtain the deformation information from the spectrum with certain model calculations (not directly image the deformation in position space)

Probe the deformation (mass distributions) of ^{96}Ru & ^{96}Zr

Relativistic heavy ion collisions



**initial conditions:
(deformation / mass distributions)**

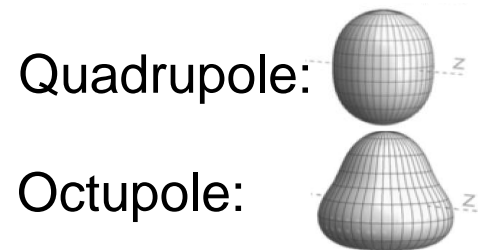
Initial conditions (TRENTO)

- Sample nucleon position in deformed nuclei with:

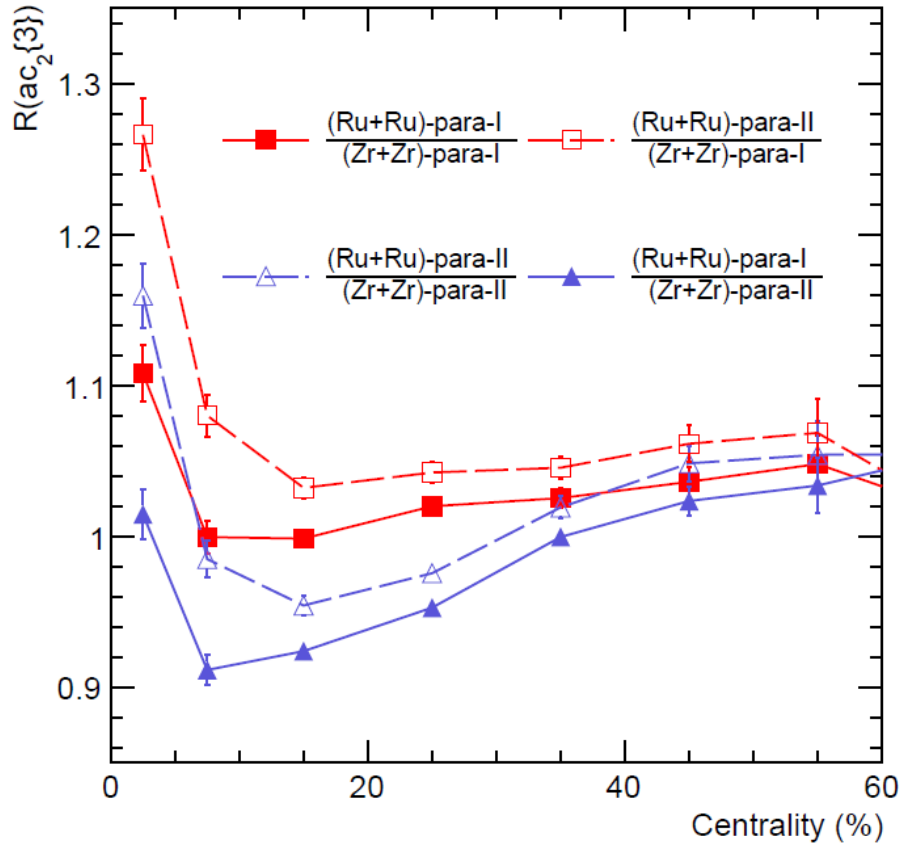
$$\rho(r, \theta, \phi) = \frac{\rho_0}{1 + e^{(r-R(\theta, \phi))/a_0}}$$

$$R(\theta, \phi) = R_0 \left(1 + \beta_2 [\cos \gamma Y_{2,0} + \sin \gamma Y_{2,2}] \right.$$

$$\left. + \beta_3 \sum_{m=-3}^3 \alpha_{3,m} Y_{3,m} + \beta_4 \sum_{m=-4}^4 \alpha_{4,m} Y_{4,m} \right)$$

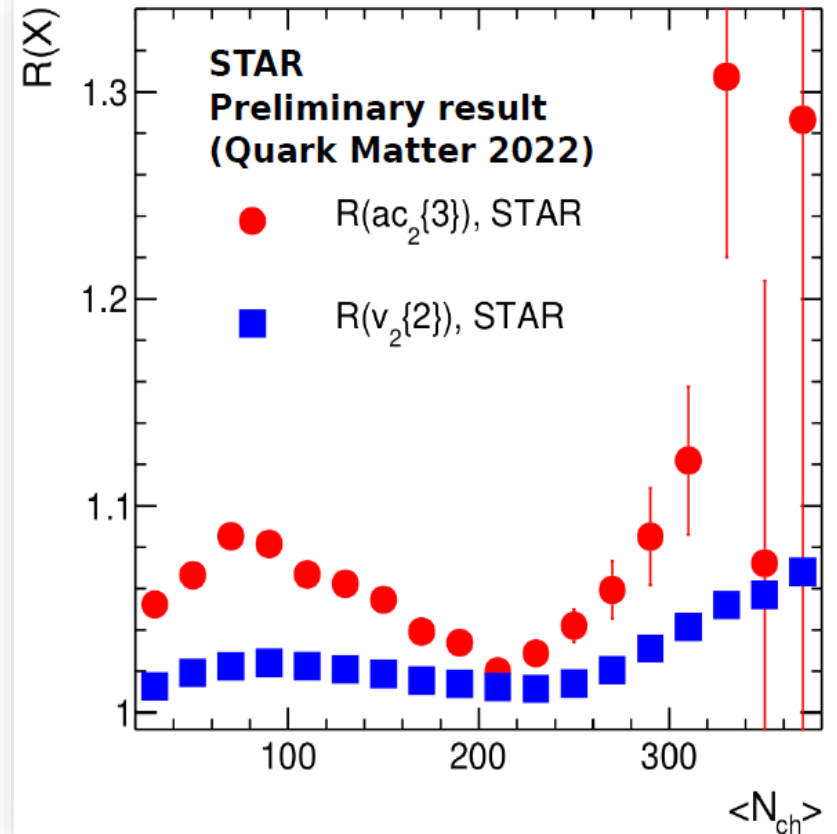


$ac_2\{3\}$ for Ru+Ru and Zr+Zr collisions

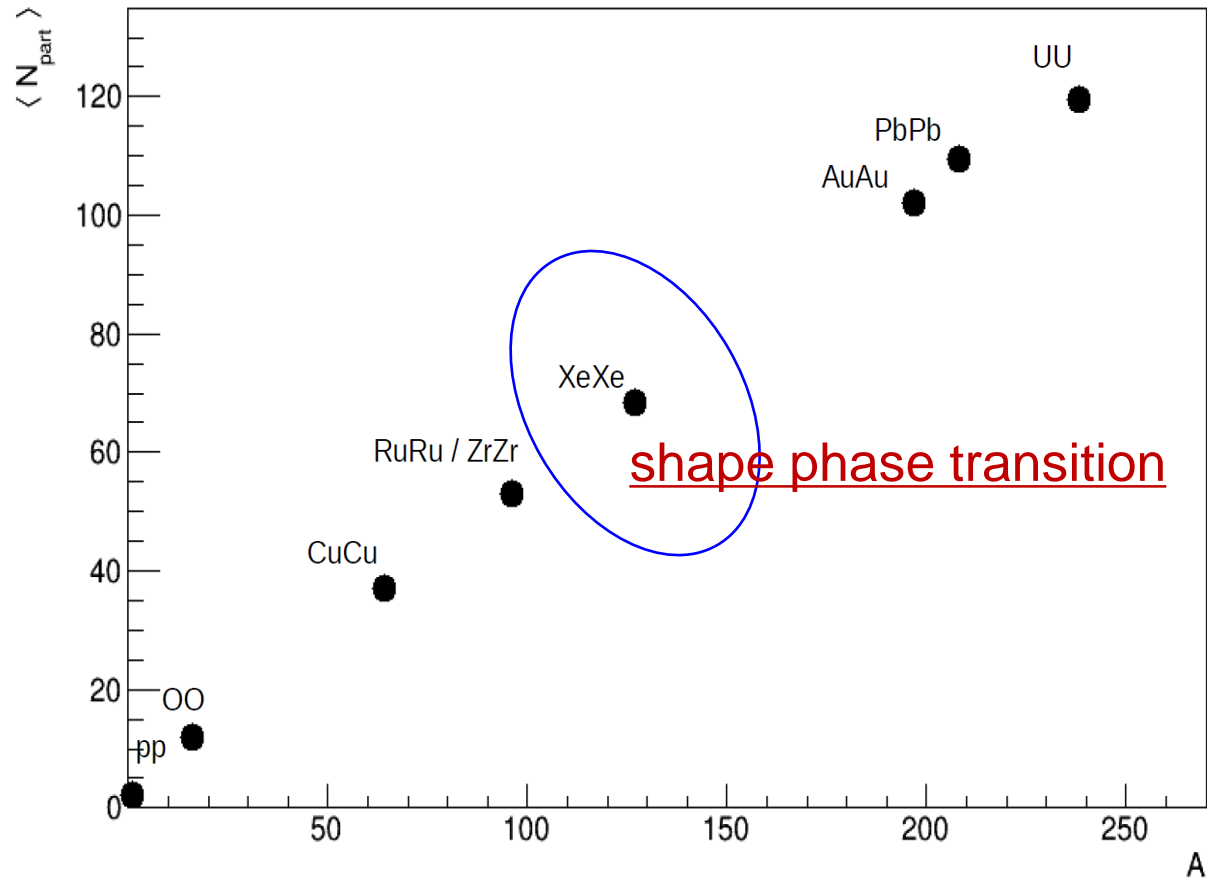


$ac_2\{3\}$ is sensitive to quadrupole and octupole deformations

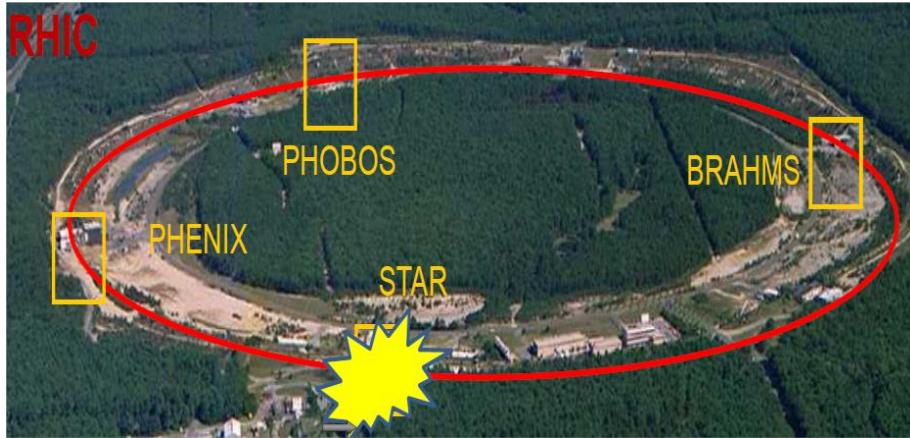
$$ac_2\{3\} = \langle v_2^2 v_4 \cos 4(\Phi_2 - \Phi_4) \rangle,$$



Probe the shape phase transition with Xe +Xe collisions



The Phase Transition



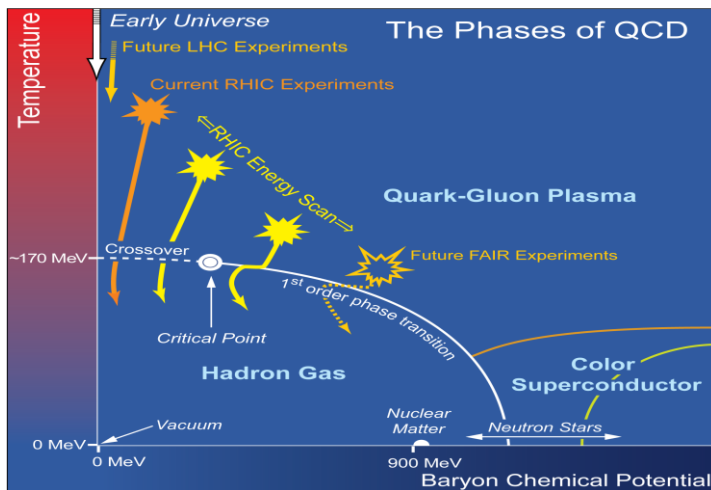
$^{129}\text{Xe} + ^{129}\text{Xe}$ collision

-explore the second-order shape phase transition occurring in the vicinity of $^{128-130}\text{Xe}$

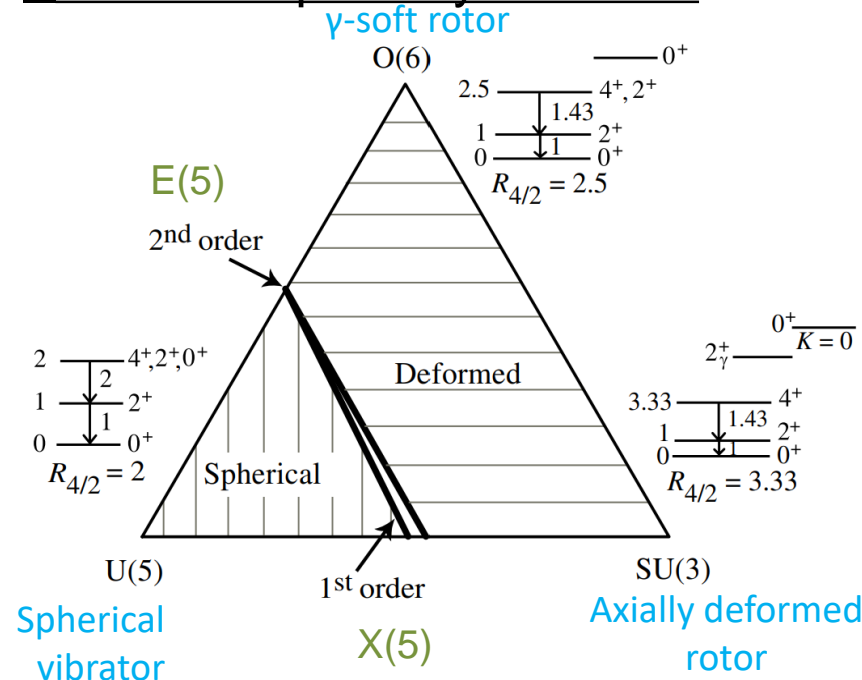
S. Zhao, H. Xu, Y. Zhou, Y. Liu, H. Song, arXiv: 2403.07441 [nucl-th]

Relativistic heavy ion collisions

-mainly aim to explore QCD Phase Transition



The critical point symmetries



Shape phase transition for Xe isotopes

The shape phase transition:

- rapid structural change along certain isotope or isotone chains
- the dynamic interplay between the spherical-driving pairing interaction and the deformation-driving proton-neutron interaction

The shape phase transition for the Xe isotopes:

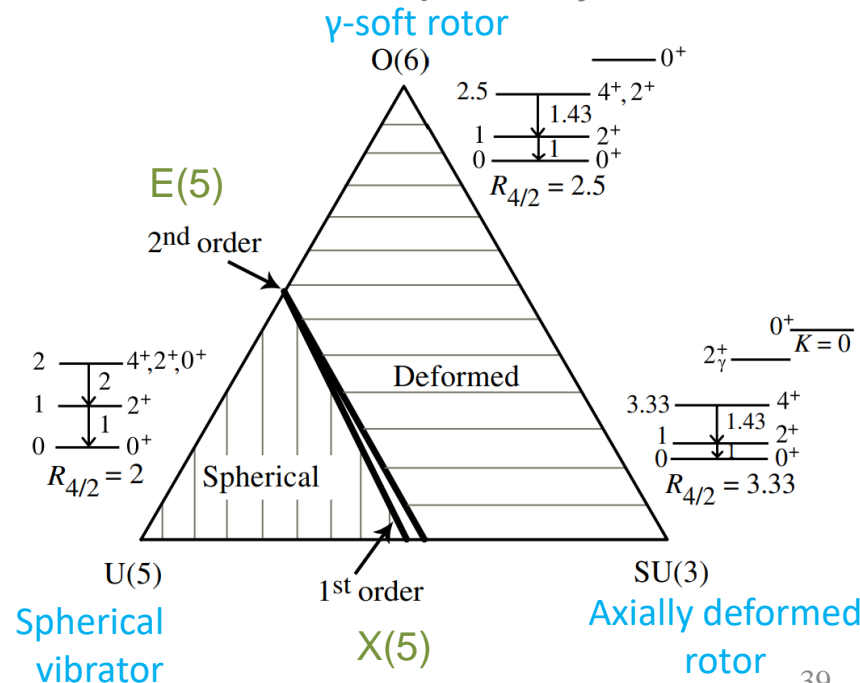
-Within the the framework of the interacting boson model (IBM), the Xe isotopes undergo a shape phase transition from a γ -soft rotor to a spherical vibrator

R. F. Casten, Nucl. Phys. A 439, 289 (1985). G. Puddu, O. Scholten, and T. Otsuka, Nucl. Phys. A 348, 109 (1980). R. F. Casten and P. Von Brentano, Phys. Lett. B 152, 22 (1985).

-the critical point is described by the $E(5)$ symmetry, associated with a 2nd order phase transition

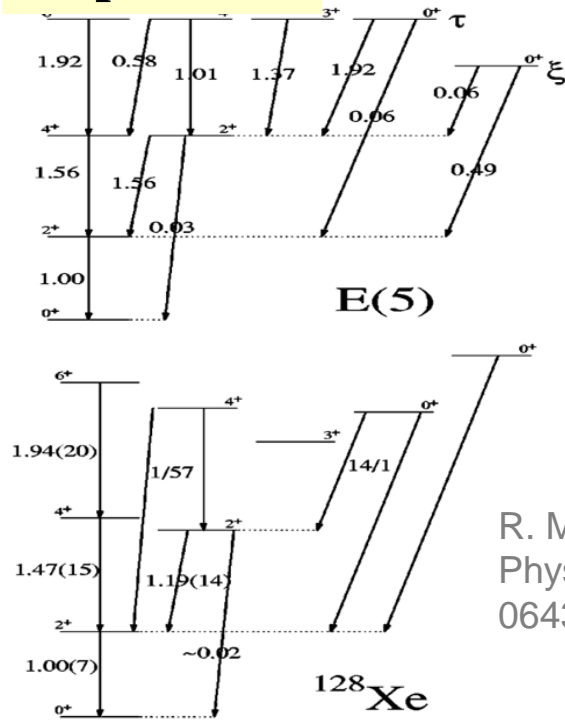
F. Iachello, Phys. Rev. Lett. 87, 052502 (2001).
F. Iachello, Phys. Rev. Lett. 85, 3580 (2000).

The critical point symmetries



E(5) symmetry near $^{128-130}\text{Xe}$

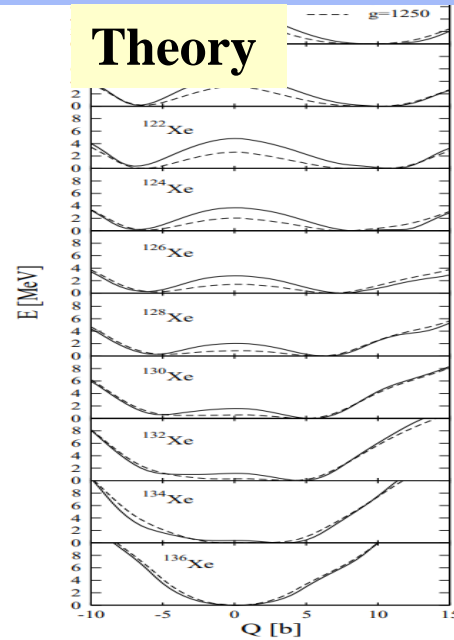
Experiment



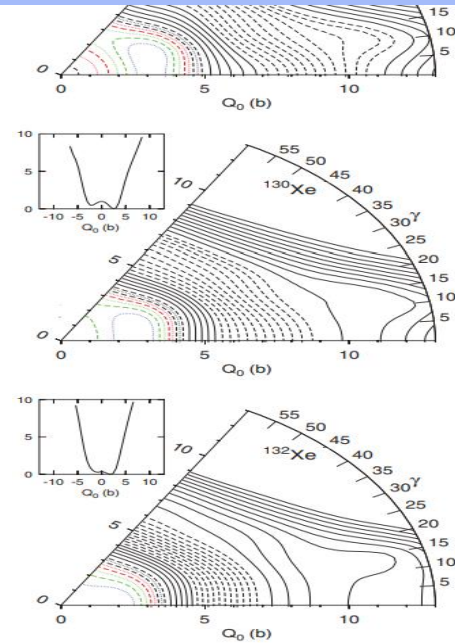
R. M. Clark, et. al.
Phys. Rev. C 69,
064322 (2004)

the measured energy spectroscopy of ^{128}Xe agrees well with the E(5) predic. (the normalized transition strengths, the branching ratios ...)

Theory



R. Rodriguez-Guzman, et.
al. Phys. Rev. C 76,
064303 (2007)



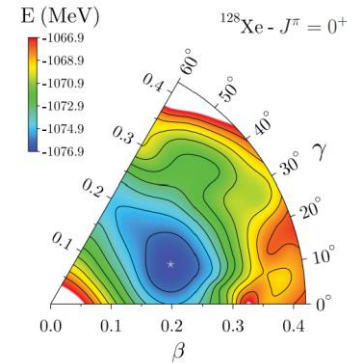
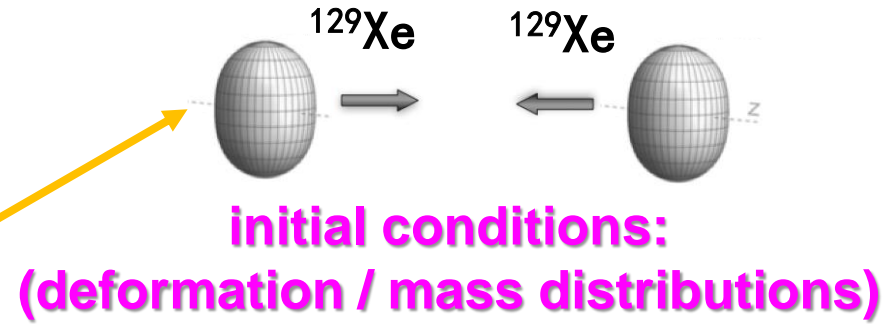
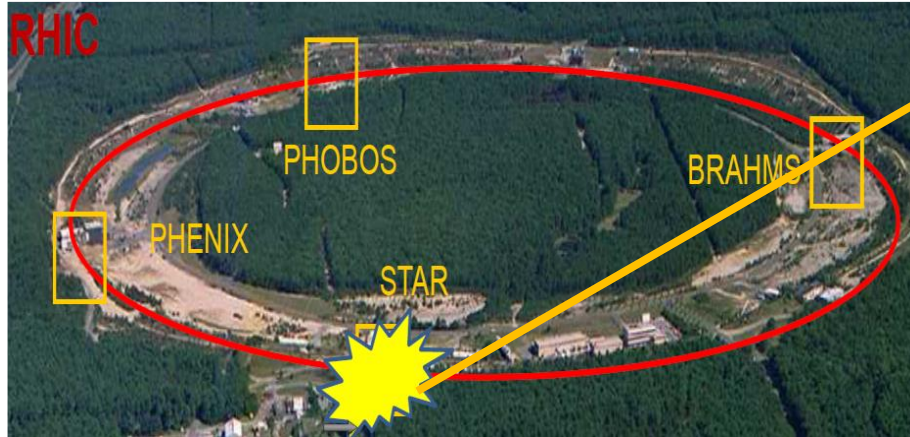
L.M.Robledo, et. al. Phys.
Rev.C 78 (2008) 034314

-Model calculations indicate a critical point of the second-order shape phase transition (E(5) symmetry) lies in the vicinity of $^{128-130}\text{Xe}$, associated with a γ -soft deformation

→ Exploring the 2nd order shape transition of Xe isotope with Xe+Xe collisions at the LHC and NICA

Probe the γ -soft deformation of ^{129}Xe

Relativistic heavy ion collisions



Rigid triaxial deformation ($\gamma=30^\circ$)

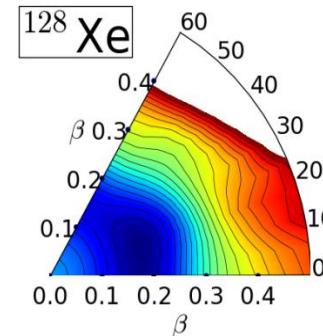
Bally et. al. Eur.Phys.J. A 58 (2022) 9, 187,

Initial conditions (TRENTO)

- Sample nucleon position in deformed nuclei with:

$$\rho(r, \theta, \phi) = \frac{\rho_0}{1 + e^{(r-R(\theta, \phi))/a_0}}$$

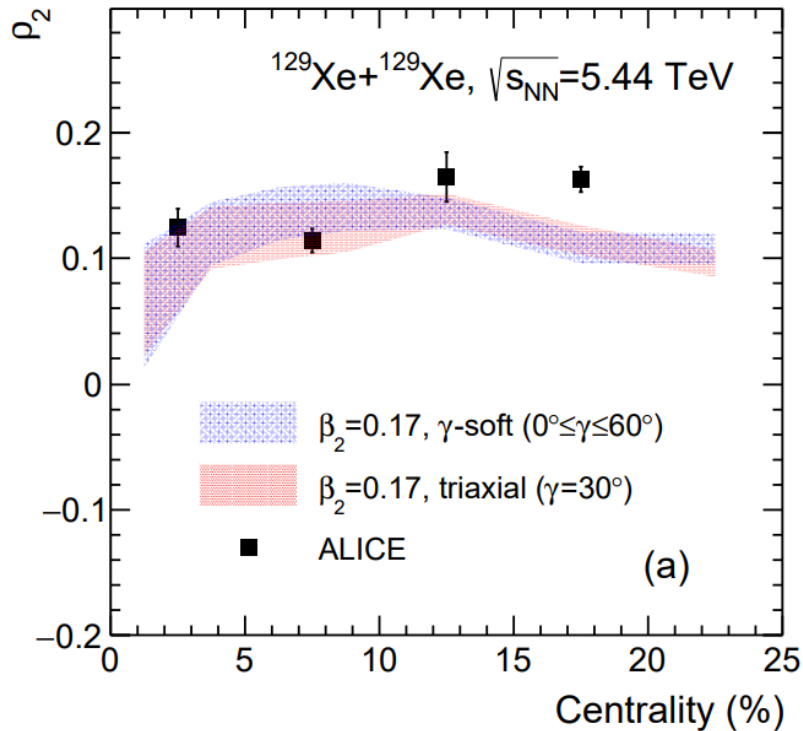
$$R(\theta, \phi) = R_0(1 + \beta_2[\cos \gamma Y_{2,0}(\theta, \phi) + \sin \gamma Y_{2,2}(\theta, \phi)]).$$



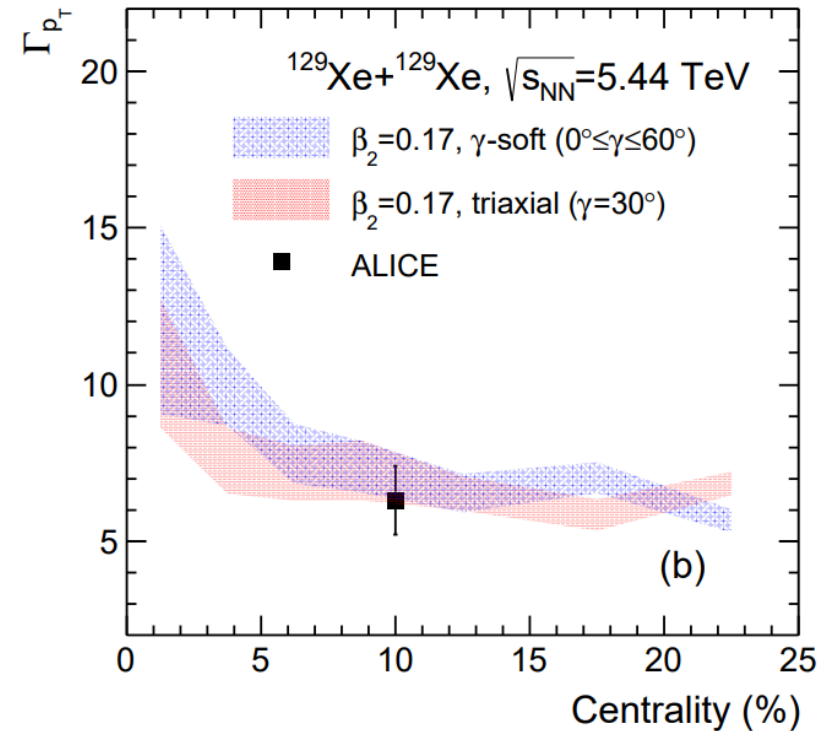
γ -soft (flat distribution in $0 \leq \gamma \leq 60^\circ$)

Z. P. Li, et. al. Phys. Rev. C 81, 034316 (2010),

3-particle correlation



$$\rho_2 \equiv \frac{\text{cov}(v_2\{2\}^2, [p_T])}{\sqrt{\text{var}(v_2\{2\}^2)}\sqrt{\text{var}([p_T])}}$$

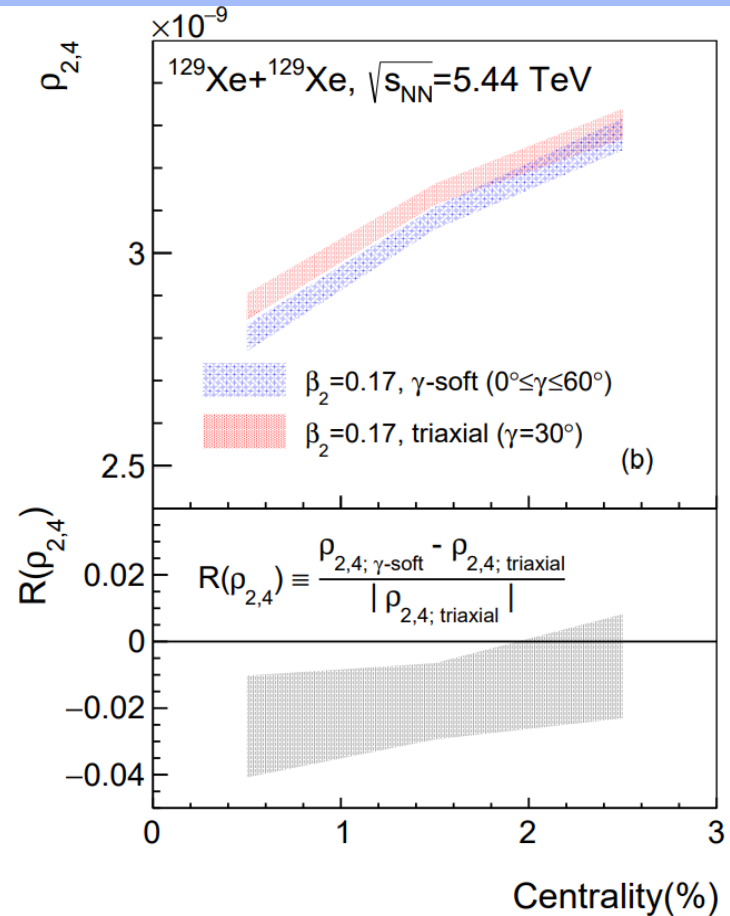
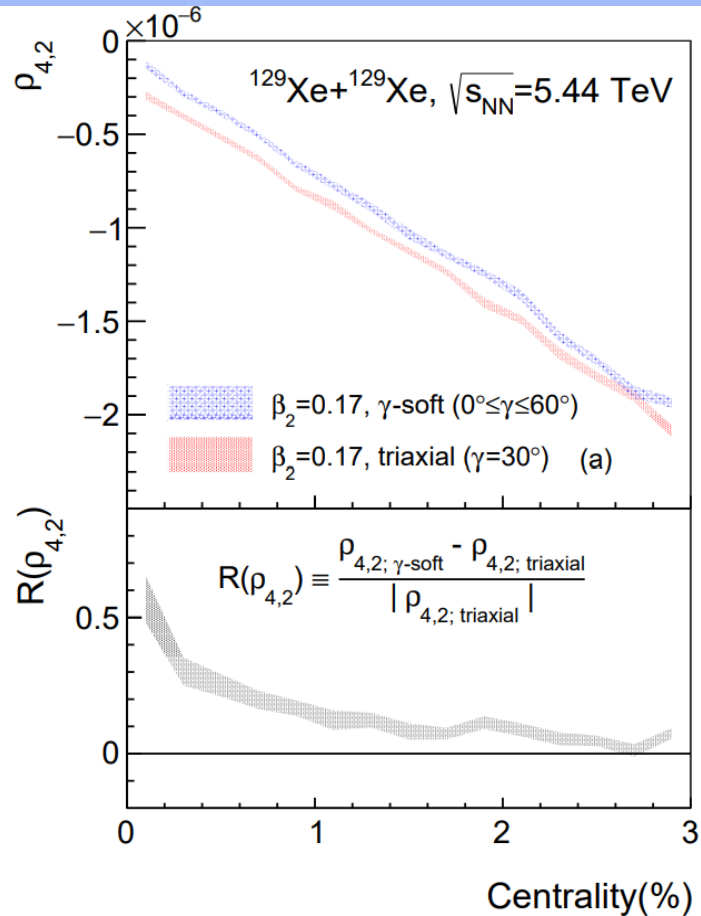


$$\Gamma_{p_T} = \frac{\langle \delta p_{T,i} \delta p_{T,j} \delta p_{T,k} \rangle \langle [p_T] \rangle}{\langle \delta p_{T,i} \delta p_{T,j} \rangle^2},$$

-Our calculations with rigid triaxial or γ -soft deformation of ^{129}Xe can describe the measured ρ_2 and Γ_{p_T} equally well.

$\rho_2, \Gamma_{p_T} \propto \beta_2^3 \cos(3\gamma)$ insensitive to triaxial deformation $\gamma=30^\circ$ and γ -soft $0 \leq \gamma \leq 60^\circ$

6-particle correlations

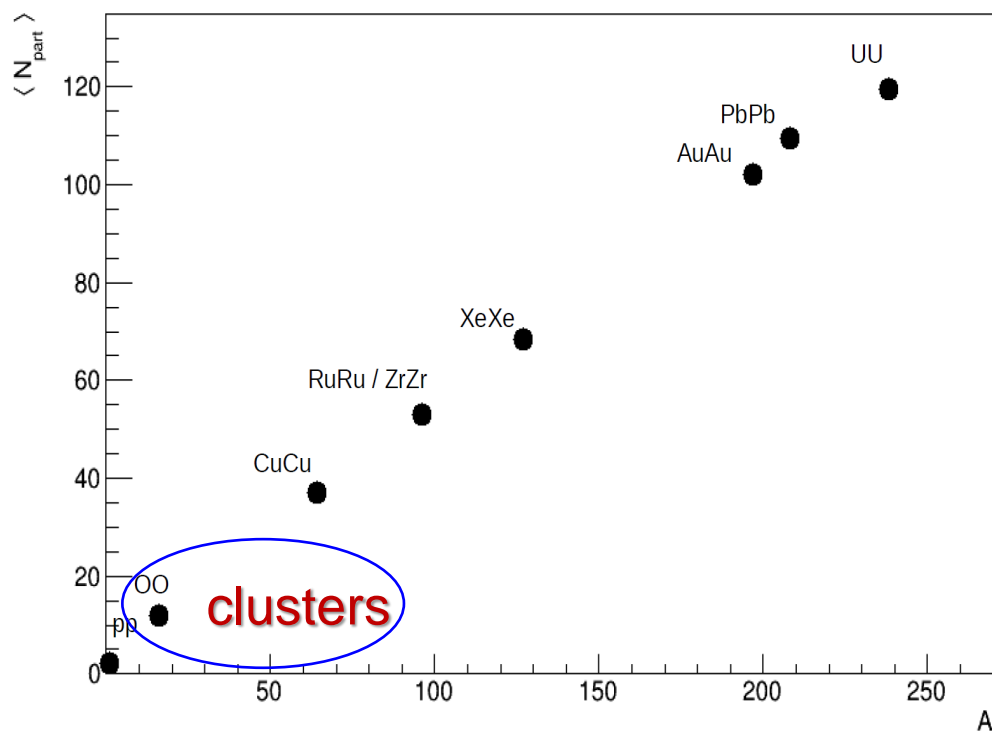


$$\rho_{4,2} \equiv \left(\frac{\langle \varepsilon_2^4 \delta d_\perp^2 \rangle}{\langle \varepsilon_2^4 \rangle \langle d_\perp \rangle^2} \right)_c$$

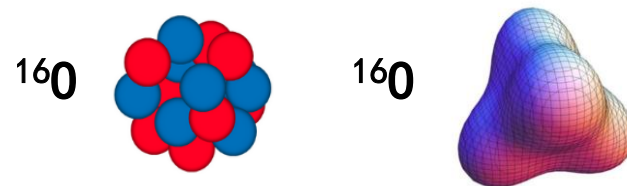
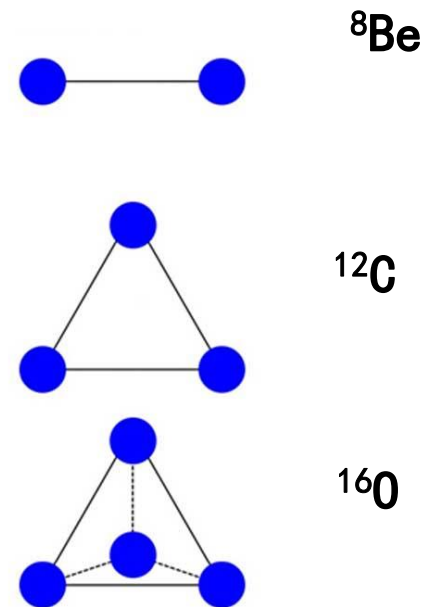
$$\rho_{2,4} \equiv \left(\frac{\langle \varepsilon_2^2 \delta d_\perp^4 \rangle}{\langle \varepsilon_2^2 \rangle \langle d_\perp \rangle^4} \right)_c \langle \cos(6\gamma) \rangle$$

The γ -soft deformation of ^{129}Xe lead to a clear enhancement of 6-particle correlations $\rho_{4,2}$ in ultra-central Xe+Xe collisions

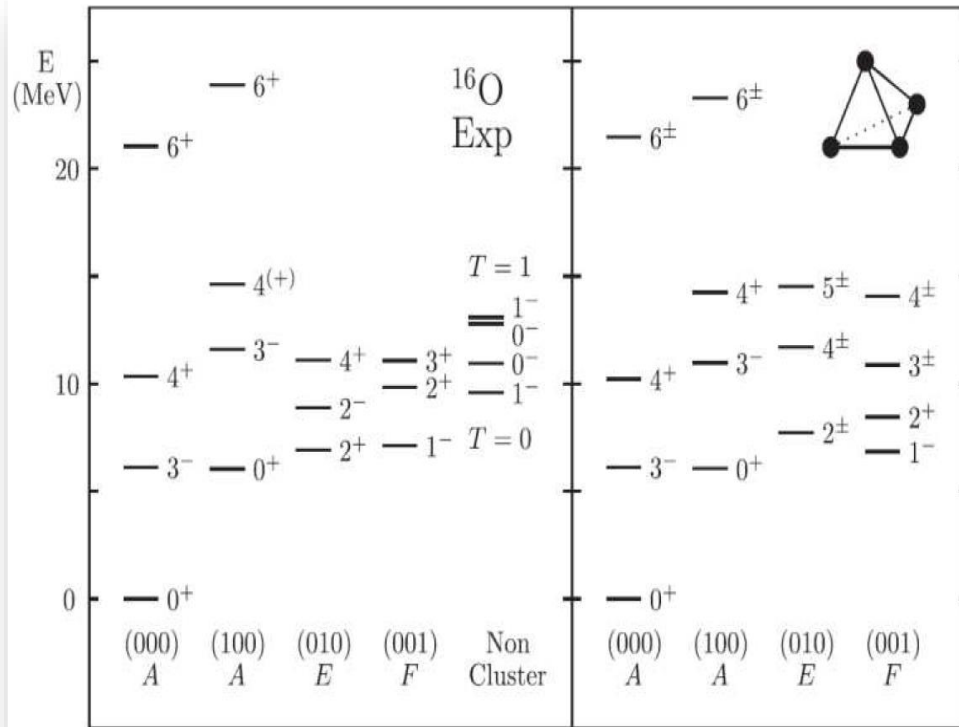
Probe the α -cluster of ^{16}O at RHIC and the LHC



$^{16}\text{O}+^{16}\text{O}$ collisions and $p+^{16}\text{O}$ collisions originally aim to study the possible formation of the QGP in small systems



α -cluster of ^{16}O from nuclear structure



-ACM calculations show that the low-lying states of ^{16}O can be described as rotation-vibration of a 4α cluster with tetrahedral symmetry.

R.Bijker and F.Iachello, Phys. Rev. Lett. 112, no.15, 152501 (2014)

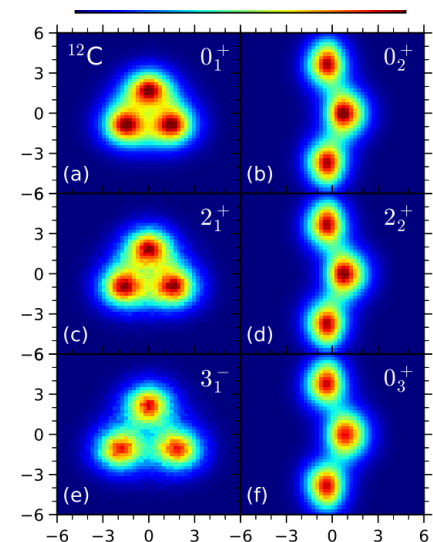
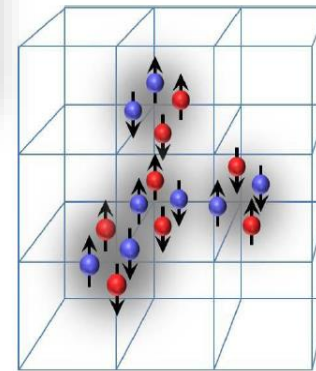
-ab initio lattice calculations demonstrate the nucleons are arranged in a tetrahedral alpha clusters in the ground state

E. Epelbaum, et al Phys. Rev. Lett.112, no.10, 102501 (2014)

Recent NLEFT calculations for light nuclei: Intrinsic shape composed of alpha clusters

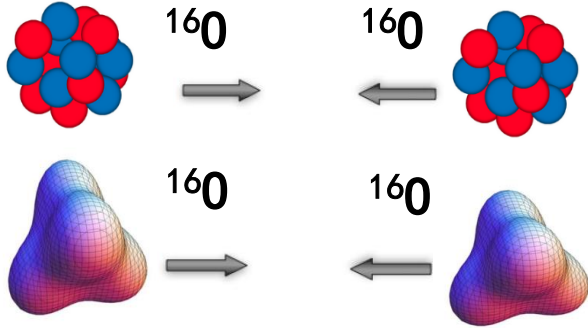
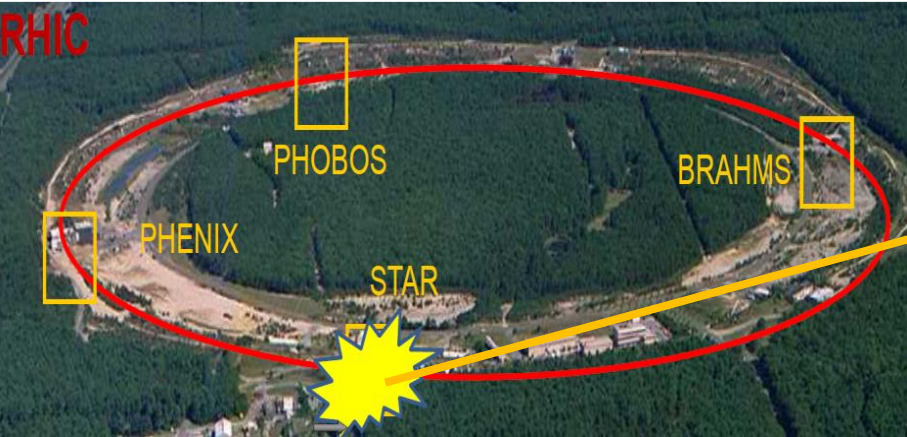
Shen, Elhatisari, Lahde, Lee, Lu, UGM, Nature Commun. **14** (2023) 2777

Dee Lee talk, today



Relativistic heavy ion collision to probe the structure of ^{16}O

Relativistic heavy ion collisions

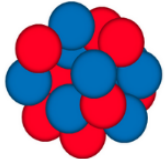


initial conditions:
(with or without α -cluster)

Initial conditions (TRENTO)

-Woods-Saxon:

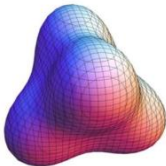
$$\rho(r, \theta, \phi) = \frac{\rho_0}{1 + e^{(r-R(\theta, \phi))/a_0}}$$



Spherical shape

-Alpha-Cluster:

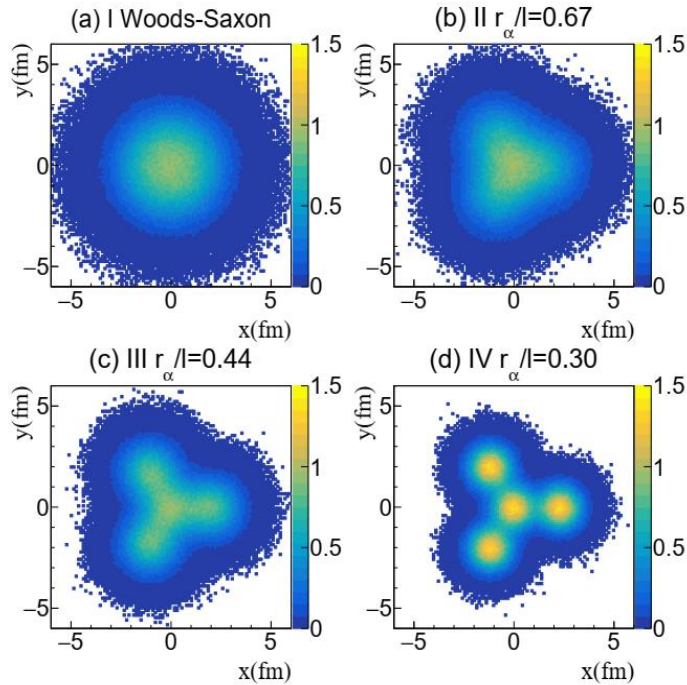
$$f_i(\mathbf{r}) = A \exp \left[-\frac{3(\mathbf{r} - \mathbf{r}_i)^2}{2r_\alpha^2} \right]$$



tetrahedral alpha clusters

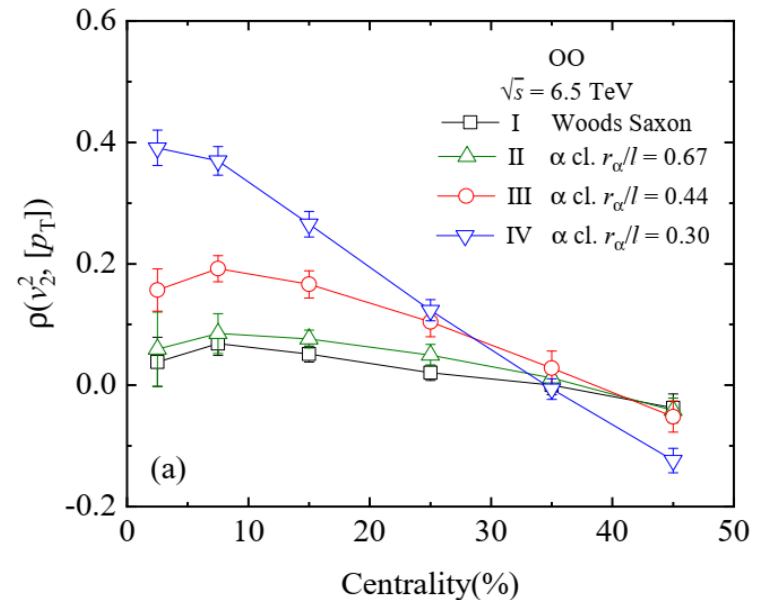
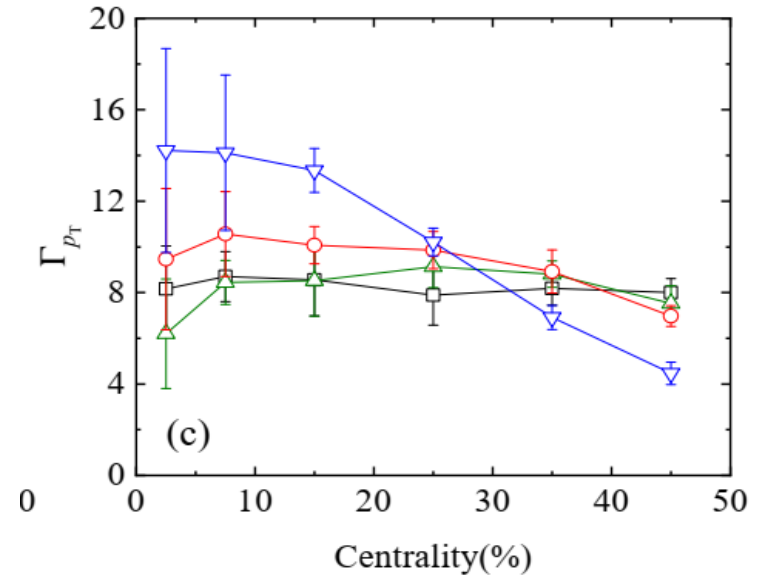
	distribution	l	r_α	r_α/l
I	Woods-Saxon			
II	α cluster	3.0	2.0	0.67
III	α cluster	3.6	1.6	0.44
IV	α cluster	4.0	1.2	0.30

Sensitive observables for α -clustering

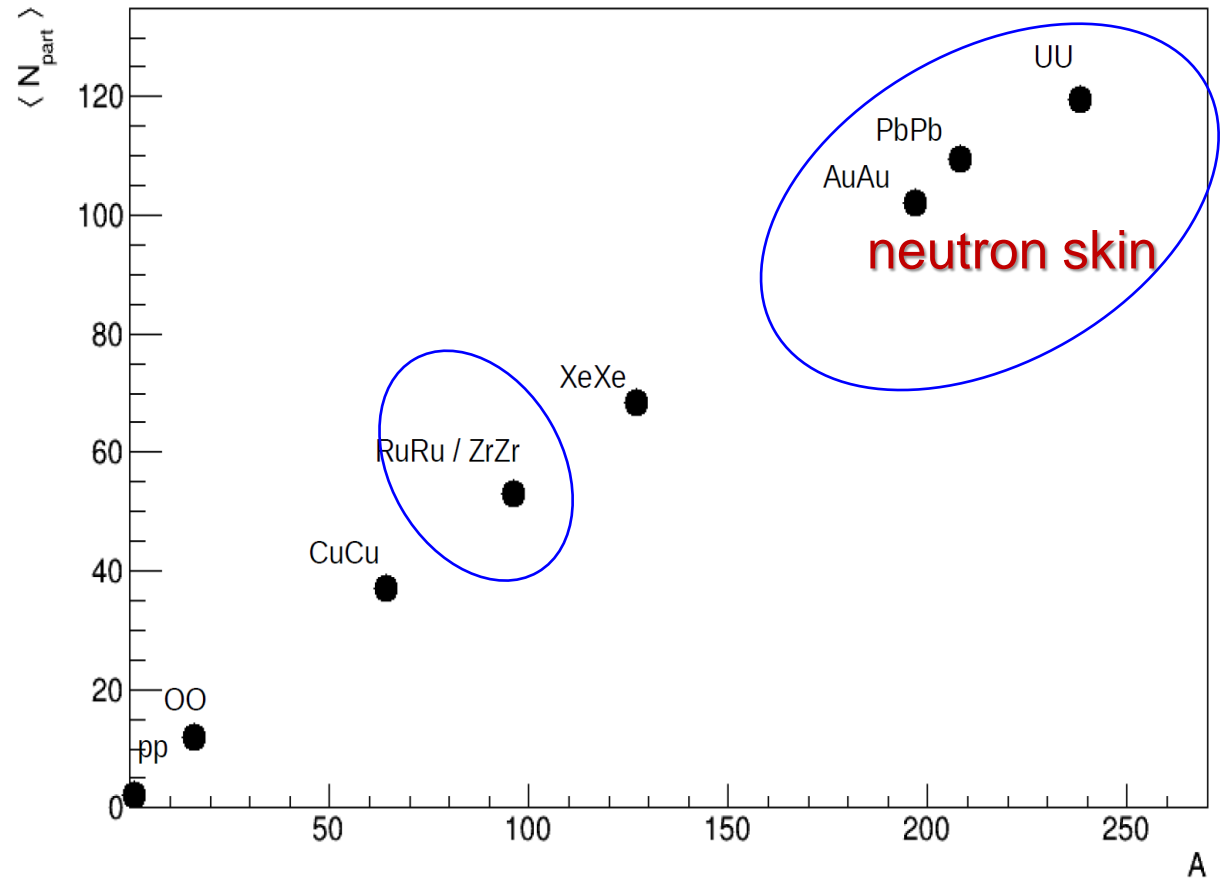
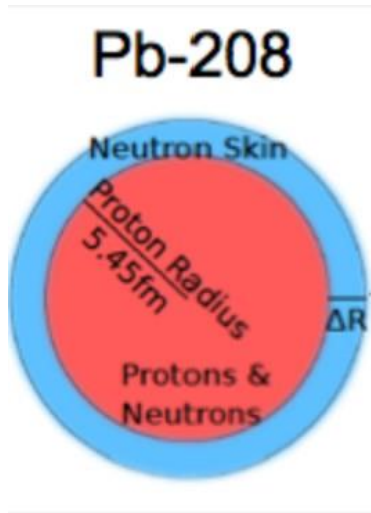


Several observables, such as the correlator Γ the $v_n - p_T$ correlations in $^{16}\text{O}+^{16}\text{O}$ collisions are sensitive to the compactness of the α cluster in the colliding nuclei, which can be used to constrain the detailed configurations of ^{16}O in the future.

Y. Wang, S. Zhao, B. Cao, H. Xu and H. Song. arXiv: 2401.15723 [nucl-th].



Probe neutron skin at RHIC and the LHC



Neutron skin & neutron star

EOS of nuclear matter

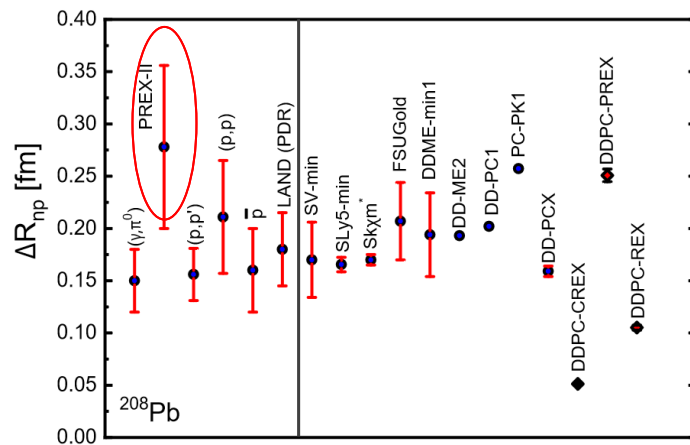
$$\epsilon(\rho, \alpha) = [\epsilon_{SNM}(\rho_0) + S(\rho_0)\alpha^2] + \alpha^2 L \frac{\rho - \rho_0}{3\rho_0} + \frac{1}{2}(K_0 + \alpha^2 K_{sym}) \left(\frac{\rho - \rho_0}{3\rho_0}\right)^2$$

L : the first order term in EOS; symmetry energy; Large L thick neutron skin

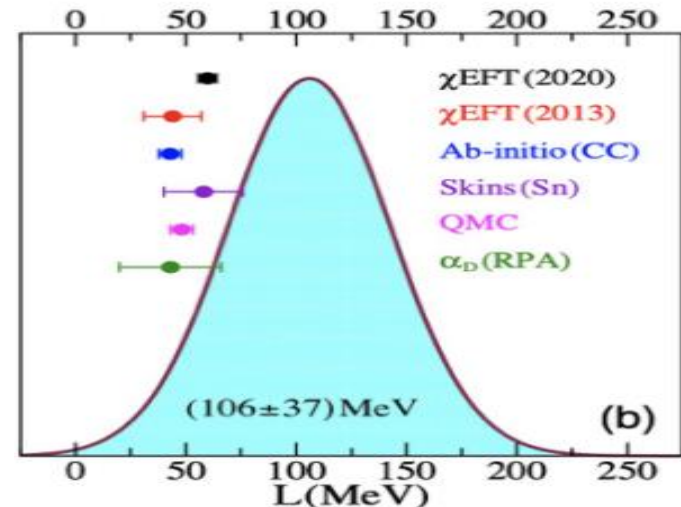
Probe the Neutron Skin at low energy nuclear physics

Parity-Violating Electron Scattering in Jefferson Lab

$$R_{skin}^{208} = 0.278_{-0.078}^{+0.078} fm \quad \text{Phys. Rev. Lett. 126, 172502, (2021)}$$



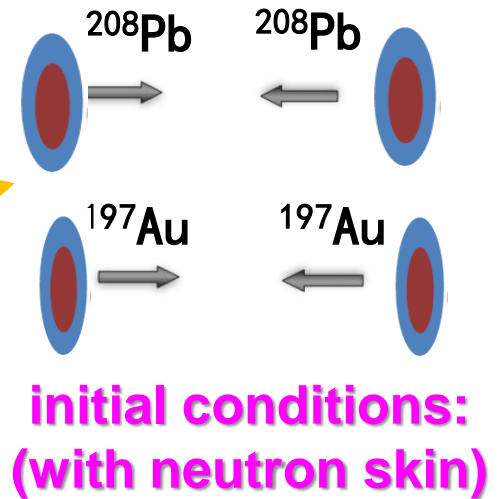
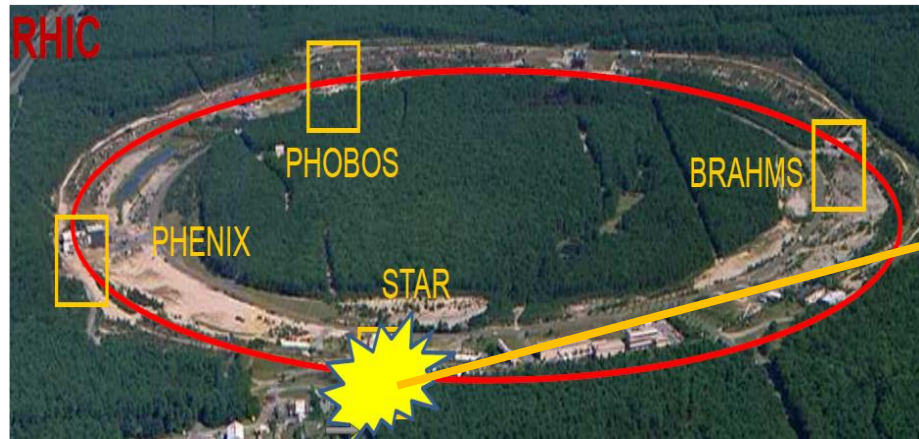
arXiv:2206.06527 Esra Yüksel and Nils Paar



Phys. Rev. Lett. 126, 172502 D. Adhikari et al

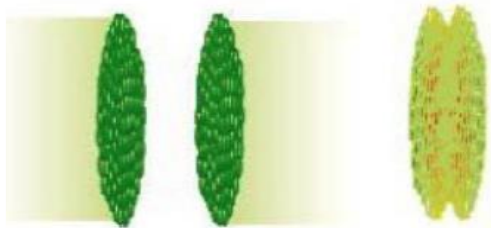
Relativistic heavy ion collision to probe the neutron skin

Relativistic heavy ion collisions



Well calibrated calculations

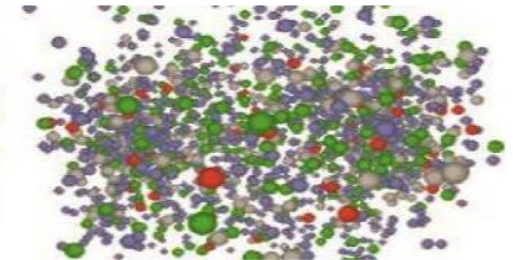
Initial conditions



viscous hydro

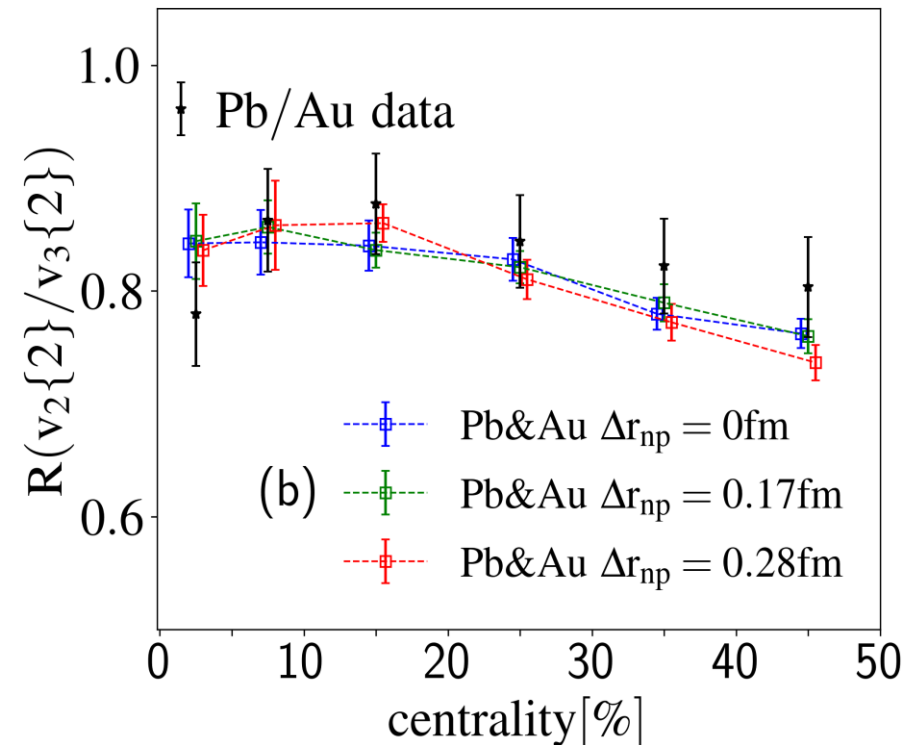
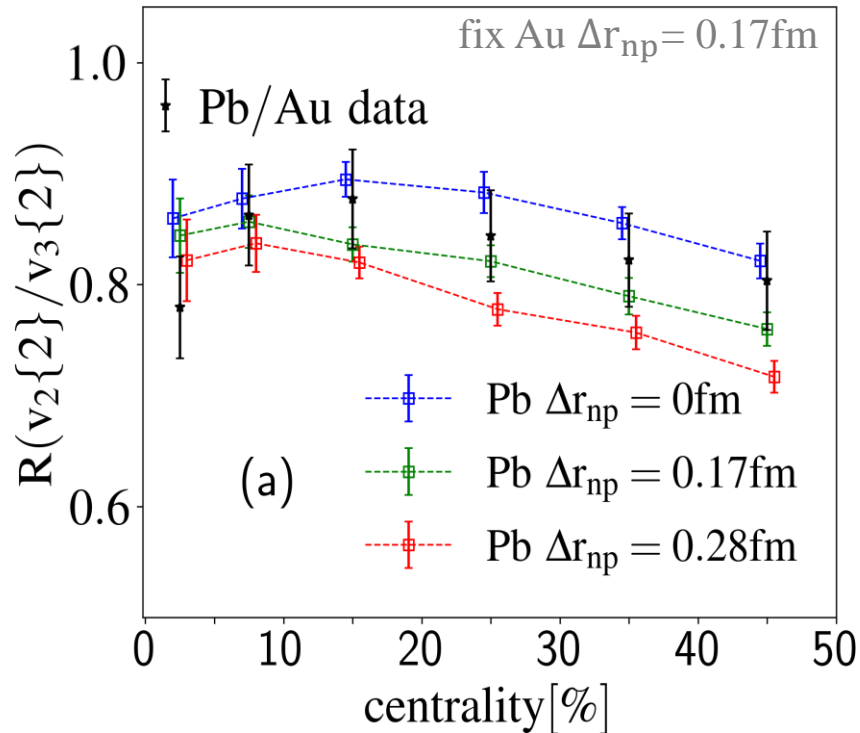


hadron cascade



Probing the neutron skin of ^{197}Au and ^{208}Pb

semi-isobaric double ratio



A scaling behavior was found in double ratio of $v_2\{2\}/v_3\{2\}$ when Au and Pb have the neutron skins of the same size,
The measured flow harmonics at various centrality suggest Au and Pb have similar neutron skin

Probing nuclear structure across energy scales

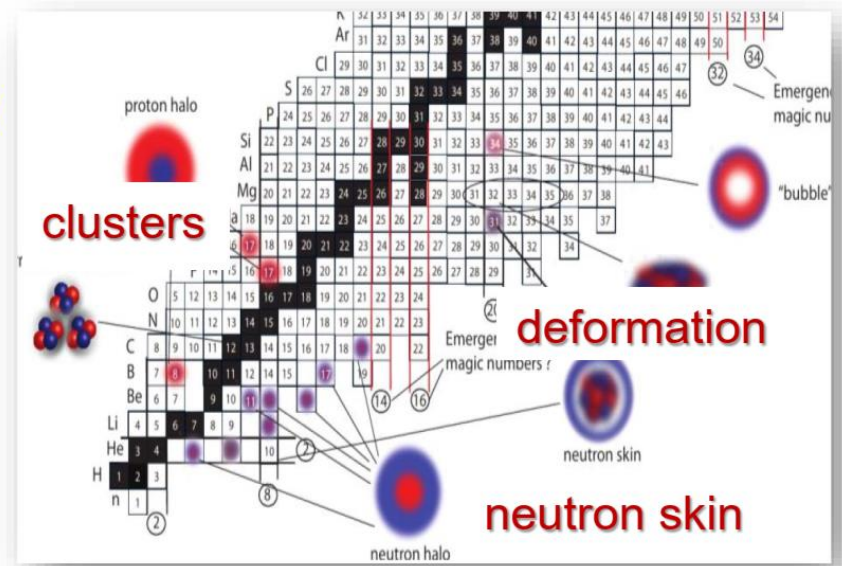
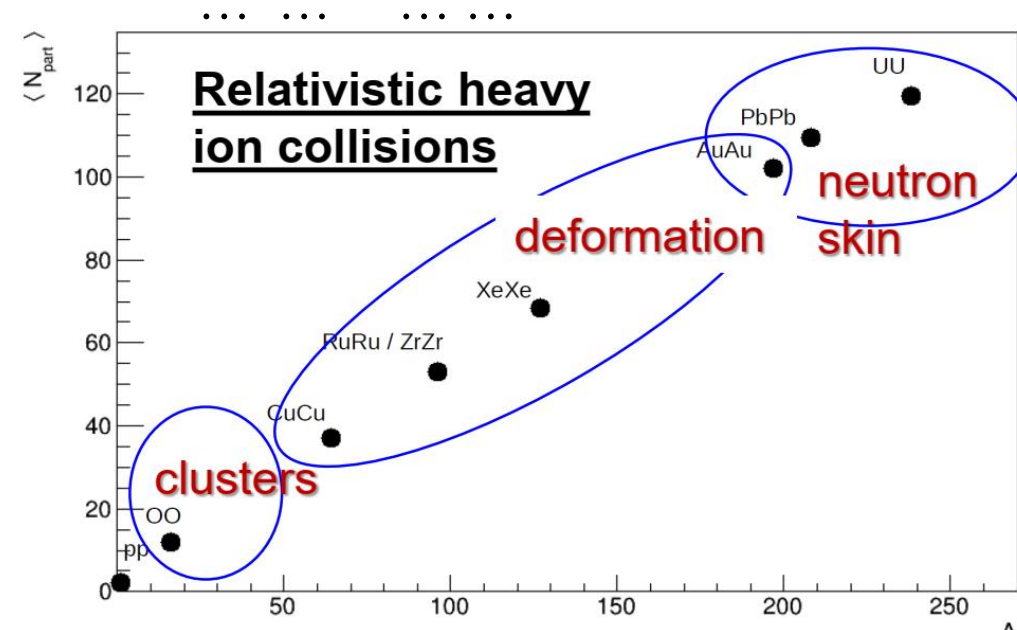
-Relativistic heavy ion collisions at RHIC and the LHC have provide rich collision systems to study various aspects in nuclear structure

-Past & future workshop, program for such intersection study

“Intersection of nuclear structure&high-energy nuclear collisions” INT, Jan.23-Feb.24 2023

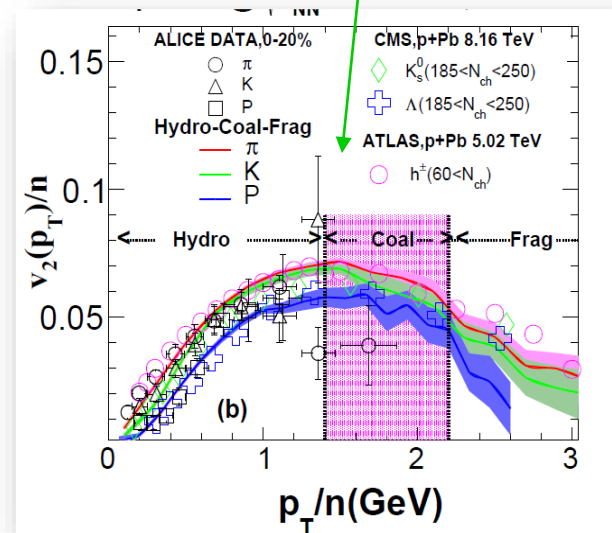
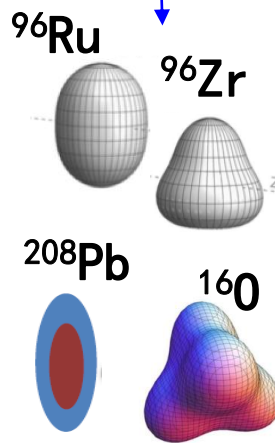
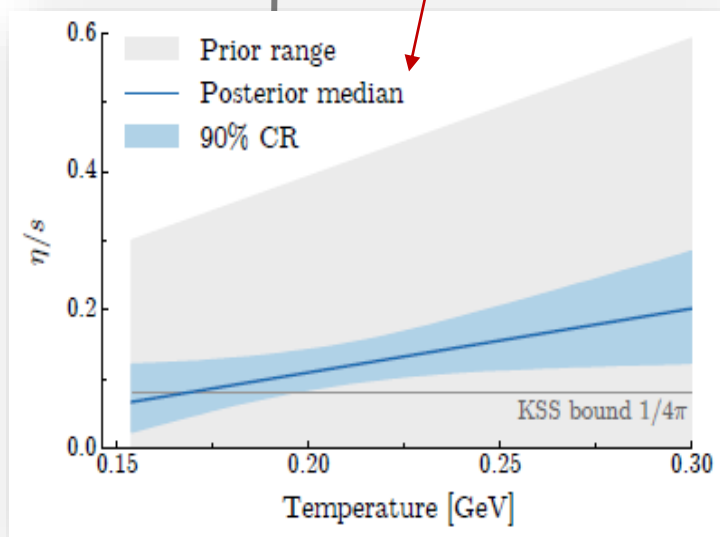
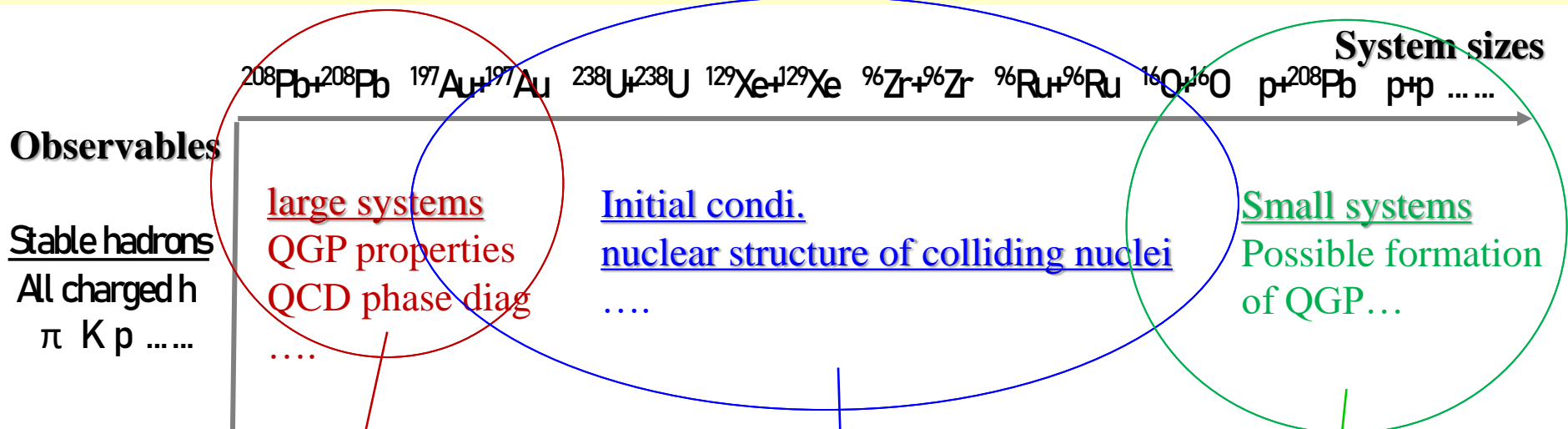
“Exploring Nuclear Physics across Energy Scales”, Beijing, April 15-17 2024

“Light ion collisions at the LHC” CERN, Oct 21-25 2024

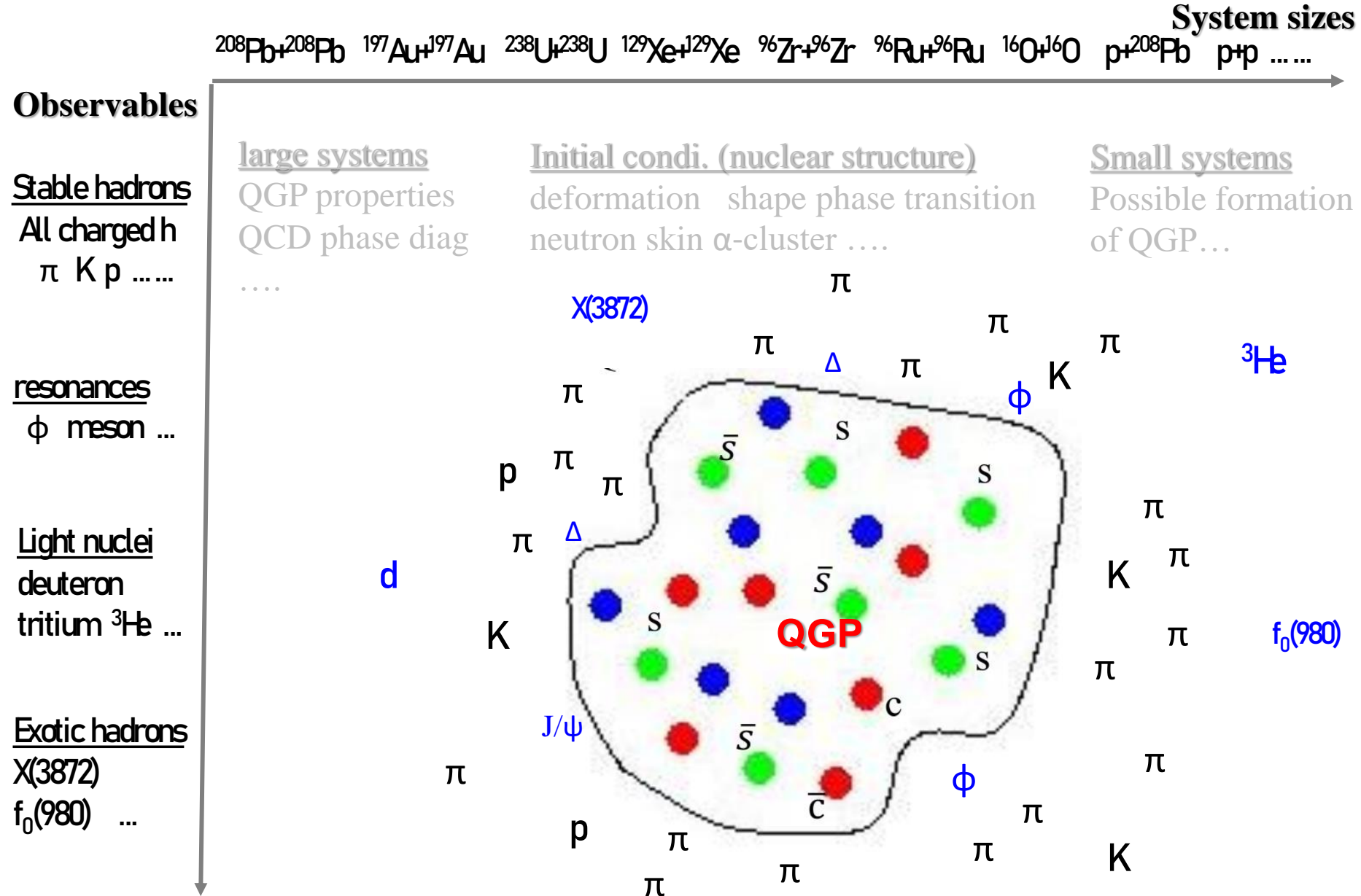


→ Probing nuclear structure at **NICA**, exploring the 2nd order shape transition of Xe isotope with Xe+Xe collisions

Soft Physics for high energy Nucleons-Nucleons Collisions



Soft Physics for high energy Nucleons-Nucleons Collisions



Soft Physics for high energy Nucleons-Nucleons Collisions

System sizes

$^{208}\text{Pb}+^{208}\text{Pb}$ $^{197}\text{Au}+^{197}\text{Au}$ $^{238}\text{U}+^{238}\text{U}$ $^{129}\text{Xe}+^{129}\text{Xe}$ $^{96}\text{Zr}+^{96}\text{Zr}$ $^{96}\text{Ru}+^{96}\text{Ru}$ $^{16}\text{O}+^{16}\text{O}$ $\text{p}+^{208}\text{Pb}$ $\text{p}+\text{p}$

Observables

Stable hadrons

All charged h
 π K p

large systems

QGP properties
QCD phase diag
.....

Initial condi. (nuclear structure)

deformation shape phase transition
neutron skin α -cluster

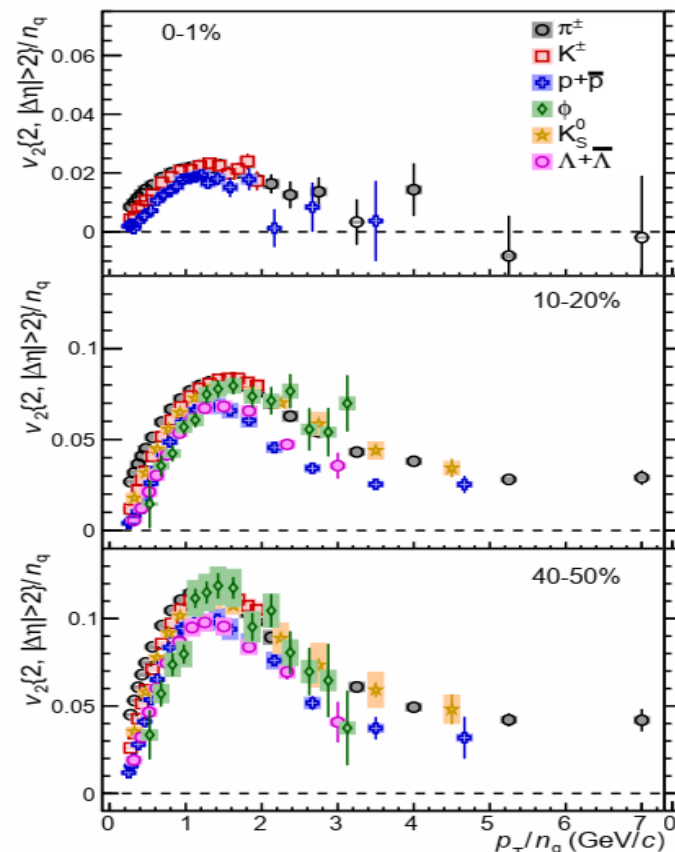
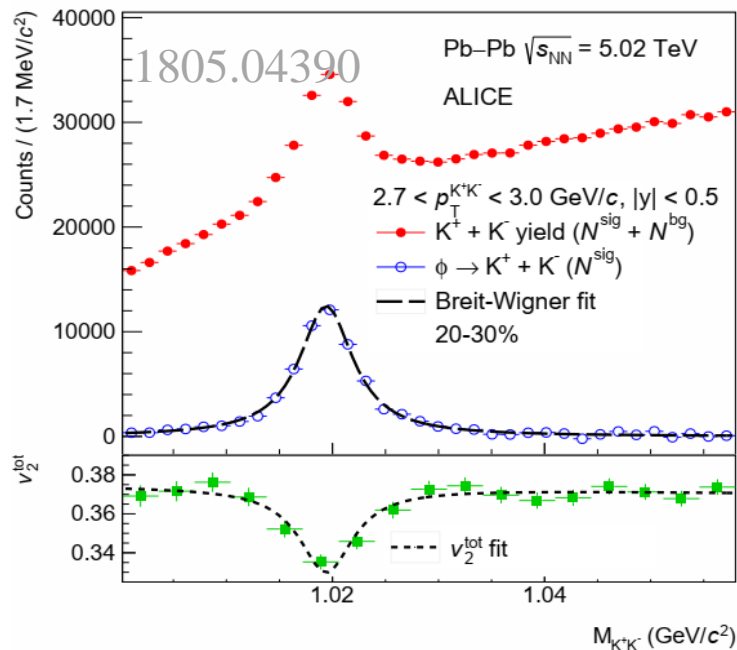
Small systems

Possible formation of QGP...

ϕ meson: (nice probe for the QGP)
long life time (45fm/c),
weakly interact with hadronic gas

resonances

ϕ meson ...



Soft Physics for high energy Nucleons-Nucleons Collisions

System sizes

$^{208}\text{Pb}+^{208}\text{Pb}$ $^{197}\text{Au}+^{197}\text{Au}$ $^{238}\text{U}+^{238}\text{U}$ $^{129}\text{Xe}+^{129}\text{Xe}$ $^{96}\text{Zr}+^{96}\text{Zr}$ $^{96}\text{Ru}+^{96}\text{Ru}$ $^{16}\text{O}+^{16}\text{O}$ $\text{p}+^{208}\text{Pb}$ $\text{p}+\text{p}$

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Light nuclei

deuteron
 tritium ^3He ...

large systems

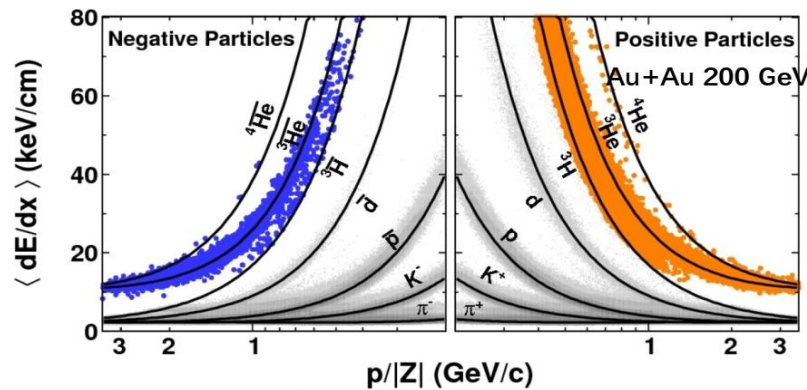
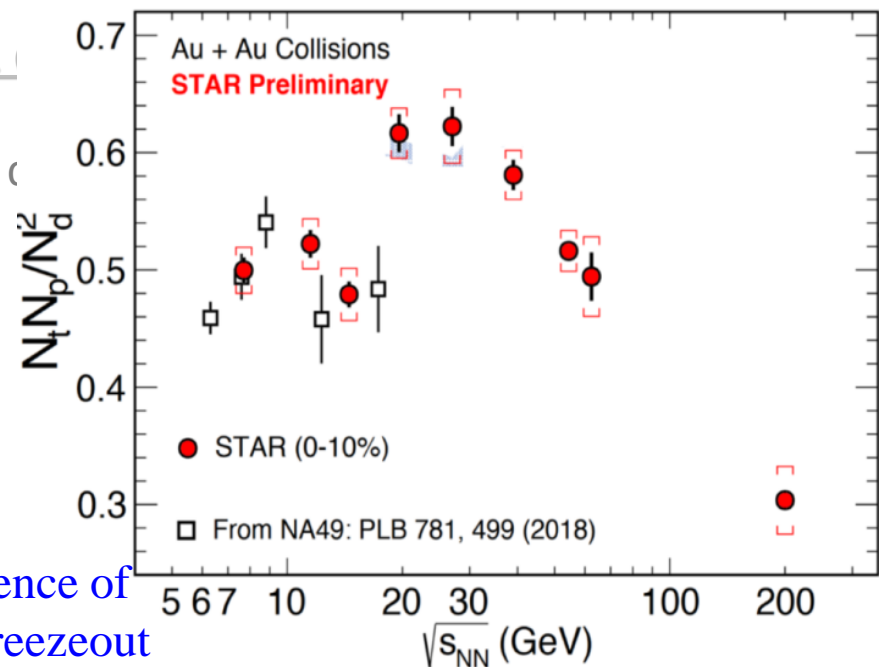
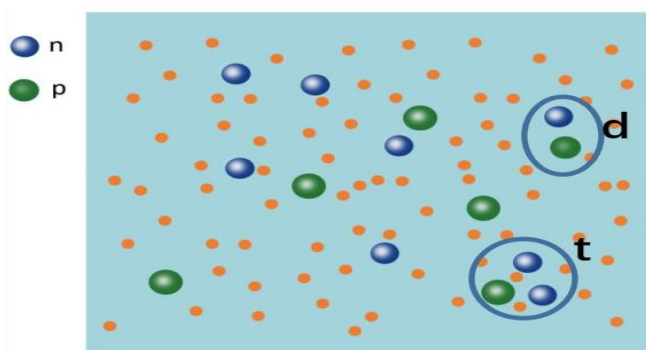
QGP properties
 QCD phase diag

Initial condi.

deformation
 neutron skin c

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 long life time (45fm/c),
 weakly interact with hadron gas
 reconstruct with $\phi \rightarrow K^+K^-$

Light nuclei: produced by coalescence of
 neutrons & protons at kinetic freezeout
 probe fluctuations & hadronic flow



STAR (Nature 473,353(2011))

Soft Physics for high energy Nucleons-Nucleons Collisions

System sizes

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Observables

Stable hadrons

All charged
 π K p

resonances

ϕ meson ...

Light nuclei

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Exotic hadrons

$X(3872)$
 $f_0(980)$...

large systems

QGP properties
QCD phase diag
.....

Initial condi. (nuclear structure)

deformation shape phase transition
neutron skin α -cluster

Small systems

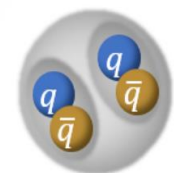
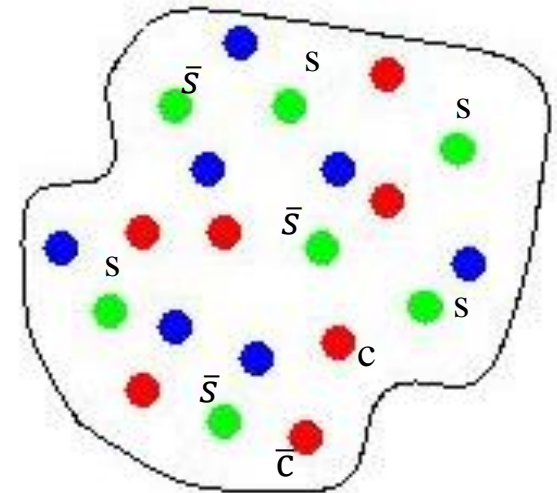
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probe fluctuations & hadronic flow

Exotic hadrons: easier to be produced with
QGP formation; identify compact or molecular
structure?

Quark Gluon Plasma



Probing exotic hadrons in relativistic heavy ion collisions

PRL **106**, 212001 (2011)

PHYSICAL REVIEW LETTERS

Identifying Multiquark Hadrons from Heavy Ion Collisions

Sungtae Cho,¹ Takenori Furumoto,^{2,3} Tetsuo Hyodo,⁴ Daisuke Jido,² Che Ming Ko,⁵ Marina Nielsen,⁶ Akira Ohnishi,² Takayasu Sekihara,^{2,7} Shigehiro Yasui,⁸ and

PHYSICAL REVIEW LETTERS **126**, 012301 (2021)

Deciphering the Nature of X(3872) in Heavy Ion Collisions

Hui Zhang,^{1,2,*} Jinfeng Liao,^{3,†} Enke Wang,^{1,2,‡} Qian Wang,^{1,2,4,§} and Hongxi Xie,^{1,2,5,¶}

using a **multiphase transport model (AMPT)** for describing such collisions and to study the production mechanism of either molecule or tetraquark picture, we compute observables for X(3872) in Pb-Pb collisions at the Large Hadron Collider. We find a crucial role, **leading to a 2-order-of-magnitude difference in the X(3872) yield centrality dependence between hadronic molecules and compact tetraquarks**, thus

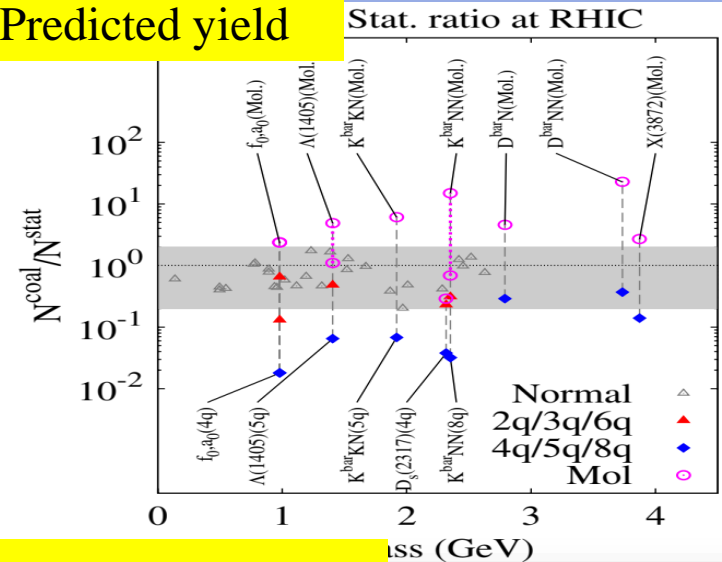
PHYSICAL REVIEW LETTERS **128**, 032001 (2022)

Evidence for X(3872) in Pb-Pb Collisions and Studies of its Prompt Production at $\sqrt{s_{NN}} = 5.02$ TeV

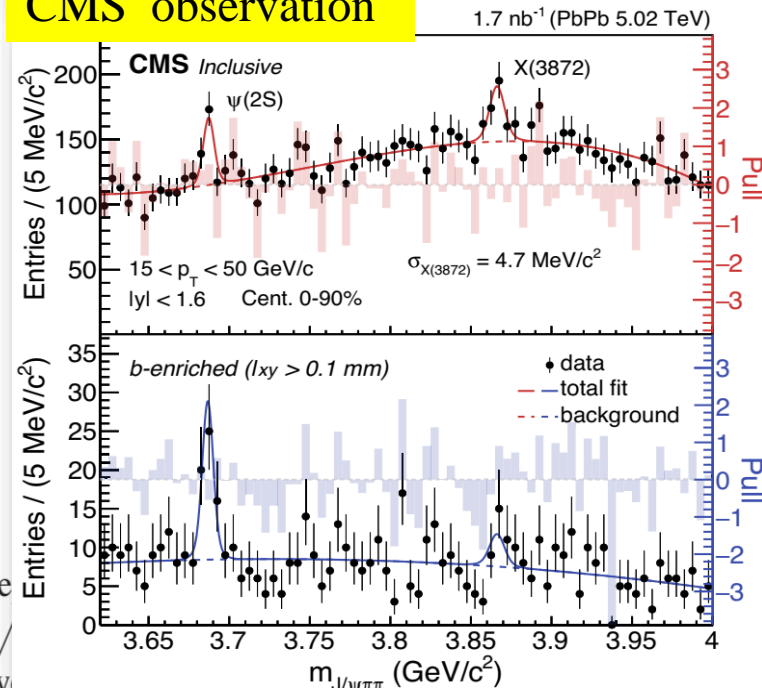
A. M. Sirunyan *et al.*^{*}
CMS Collaboration

The first evidence for X(3872) production in relativistic heavy ion collisions is reported. The production is studied in lead-lead (Pb-Pb) collisions at a center-of-mass energy of $\sqrt{s_{NN}} = 5.02$ TeV per nucleon pair, using the decay chain $X(3872) \rightarrow J/\psi \pi^+ \pi^- \rightarrow \mu^+ \mu^- \pi^+ \pi^-$. The data with

Predicted yield

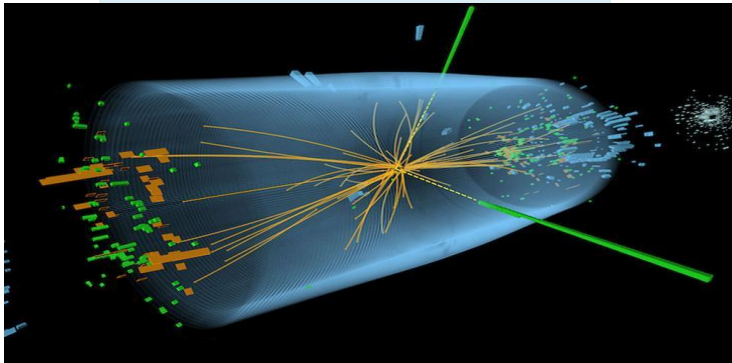


CMS observation



Advantage to study exotic hadrons in heavy ion collisions?

Particle physics



$f_0(980)$

$$I^G(J^{PC}) = 0^+(0^{++})$$

See the review on "Scalar Mesons below 1 GeV."

T-matrix pole $\sqrt{s} = (980-1010) - i(20-35)$ MeV [1]

Mass (Breit-Wigner) = 990 ± 20 MeV [1]

Full width (Breit-Wigner) = 10 to 100 MeV [1]

$f_0(980)$ DECAY MODES	Fraction (Γ_i/Γ)	p (MeV/c)
$\pi\pi$	seen	476
$K\bar{K}$	seen	36
$\gamma\gamma$	seen	495

$a_0(980)$

$$I^G(J^{PC}) = 1^-(0^{++})$$

See the review on "Scalar Mesons below 1 GeV."

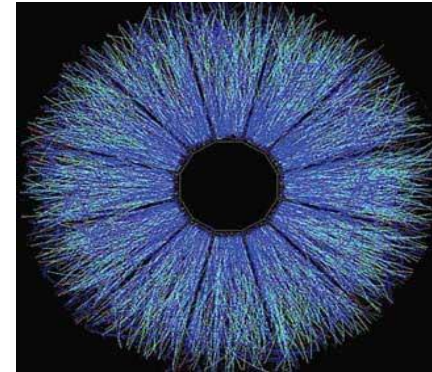
T-matrix pole $\sqrt{s} = (970-1020) - i(30-70)$ MeV [1]

Mass $m = 980 \pm 20$ MeV [1]

Full width $\Gamma = 50$ to 100 MeV [1]

$a_0(980)$ DECAY MODES	Fraction (Γ_i/Γ)	p (MeV/c)
$\eta\pi$	seen	319
$K\bar{K}$	seen	†

High energy nuclear physics



A large amount of particles produced
 → momentum distributions

- particle yield
- - p_T spectra
- flow anisotropy

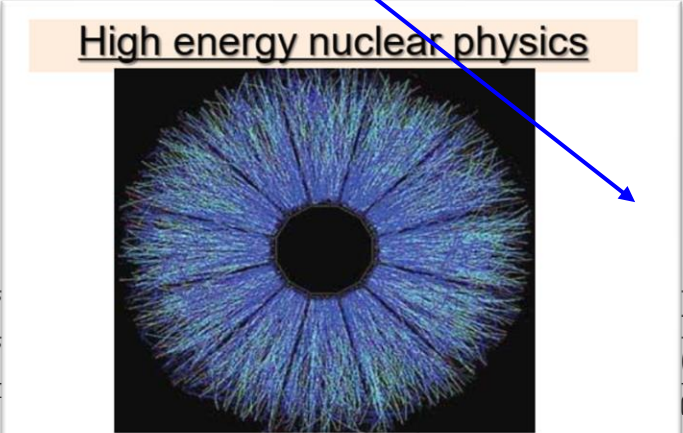
More info ↓

Advantage: provide complimentary information to constrain properties of hadrons

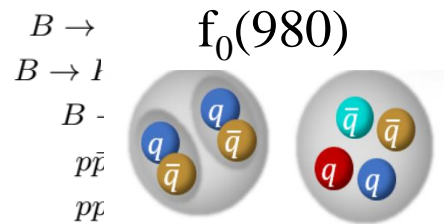
disadvantage: huge background

State	$I^G(J^{PC})$	M [MeV]	Γ [MeV]	S -wave threshold(s) [MeV]	Decay mode(s) [branching ratio(s)]
$f_0(500)$ (Peláez, 2016) ^a	$0^+(0^{++})$	449^{+22}_{-16}	550 ± 24	$\pi\pi(173^{+22}_{-16})$	$\pi\pi$ [dominant] $\gamma\gamma$
$\kappa(800)$	$\frac{1}{2}(0^+)$	682 ± 29	547 ± 24	$K\pi(48 \pm 29)$	πK
$f_0(980)$	$0^+(0^{++})$	990 ± 20	$10 \sim 100$	$K^+K^-(3 \pm 20)$ $K^0\bar{K}^0(-5 \pm 20)$	$\pi\pi$ [dominant] $K\bar{K}$ $\gamma\gamma$
$a_0(980)$	$1^-(0^{++})$	980 ± 20	$50 \sim 100$	$K\bar{K}(-11 \pm 20)$	$\eta\pi$ [dominant] $K\bar{K}$ $\gamma\gamma$
$f_1(1420)$	$0^+(1^{++})$	1426.4 ± 0.9	54.9 ± 2.6	$K\bar{K}^*(39.1 \pm 0.9)$	$K\bar{K}^*$ (dominant) $\eta\pi\pi$ [possibly seen] $\phi\gamma$
$a_1(1420)$	$1^-(1^{++})$	1414^{+15}_{-13}	153^{+8}_{-23}	$K\bar{K}^*(27^{+15}_{-13})$	$f_0(980)\pi$ [seen]

State	$I^G(J^{PC})$	M [MeV]	Γ [MeV]	S -wave threshold(s) [MeV]	Observed mode(s) (branching ratios)
$X(3872)$	$0^+(1^{++})$	3871.69 ± 0.17	< 1.2	$D^{*+}D^- + c.c.(-8.15 \pm 0.20)$ $D^{*0}\bar{D}^0 + c.c.(0.00 \pm 0.18)$	$B \rightarrow K[\bar{D}^{*0}D^0] \sim 21\%$ $B \rightarrow$ $B \rightarrow l$ $B \rightarrow p\bar{l}$ $B \rightarrow p\bar{l}$ $B \rightarrow K[J/\psi\omega](> 1.9\%)$ $B \rightarrow [J/\psi\gamma](> 6 \times 10^{-3})$ $B \rightarrow [\psi(2S)\gamma](> 3.0\%)$
$X(3940)$	$?^?(?^{??})$			-75.1 ± 9	$e^+e^- \rightarrow J/\psi[D\bar{D}^*]$
$X(4160)$	$?^?(?^{??})$			(139^{+29}_{-25})	$e^+e^- \rightarrow J/\psi[D^*\bar{D}^*]$
$Z_c(3900)$	$1^+(1^{++})$			$0.8 + 2.4$	$e^+e^- \rightarrow \pi[D\bar{D}^* + c.c.]$



High energy nuclear physics



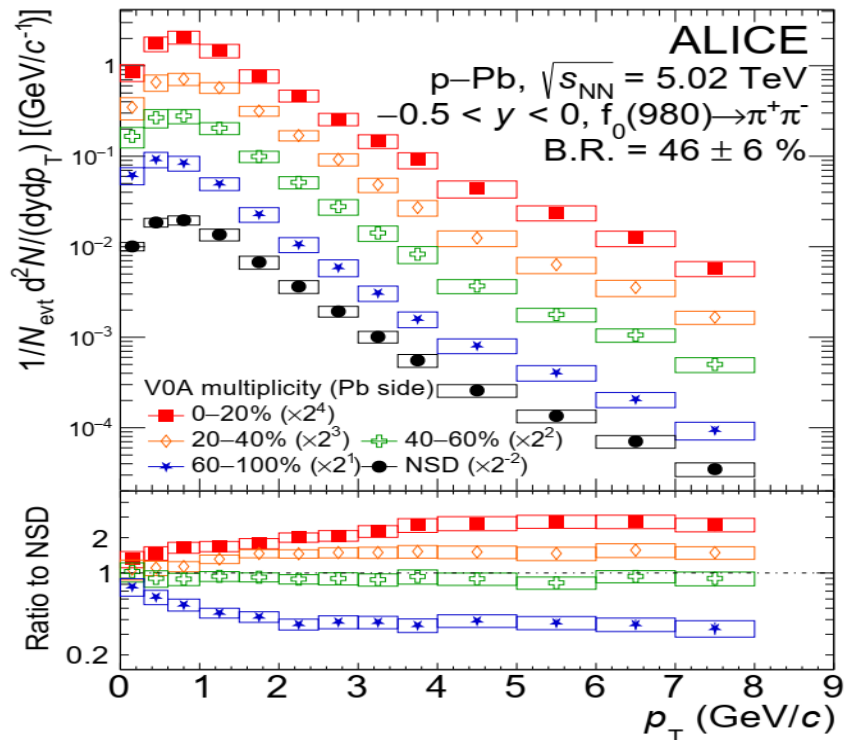
A large amount of particles produced
→ momentum distributions

Probing $f_0(980)$ in p-Pb collisions

particle yield

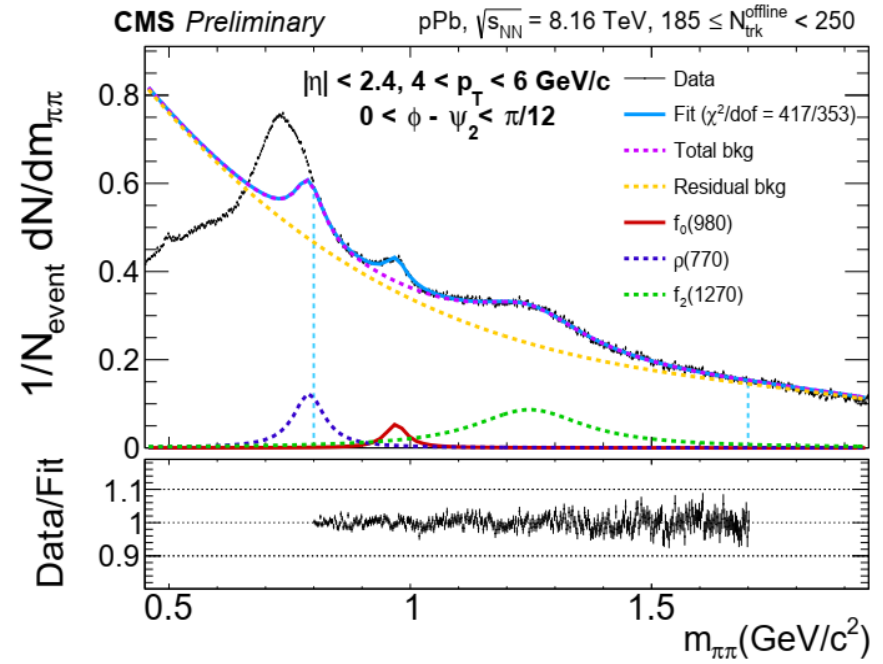
Multiplicity class (VOA)	dN/dy
0–20%	$0.206 \pm 0.005 \pm 0.014$
20–40%	$0.153 \pm 0.004 \pm 0.010$
40–60%	$0.113 \pm 0.002 \pm 0.008$
60–100%	$0.064 \pm 0.001 \pm 0.005$

p_T spectra

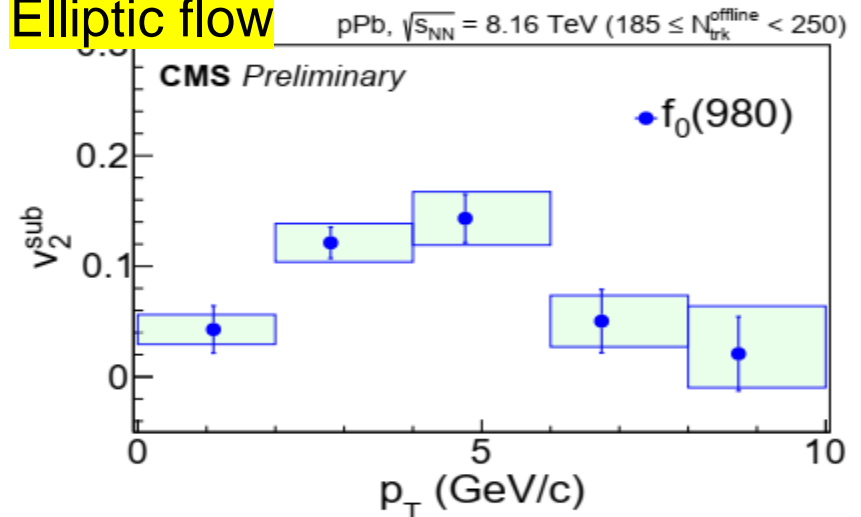


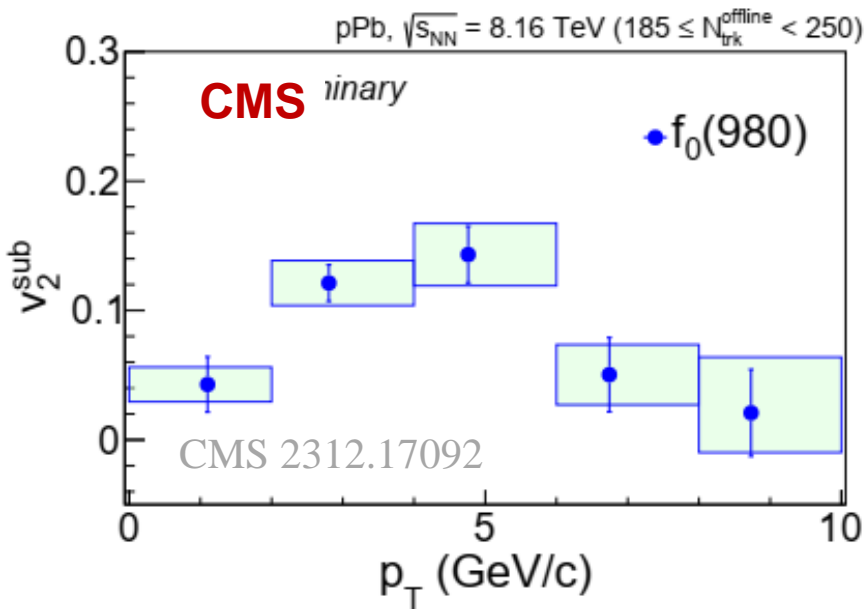
ALICE Phys.Lett.B 853 138665 (2024)

Reconstruction of $f_0(980)$ CMS 2312.17092



Elliptic flow





flow:

- anisotropy for particle momentum distribution
- contains rich information to constrain the formation /properties of an exotic hadron

Comments:

- Dominant decay channel: $f_0(980) \rightarrow \pi^+ \pi^-$; no PID in CMS, use all charged tracks
- p-Pb collisions**, less background contamination for $f_0(980)$ reconstruction
 - flow is not easy to measure in small systems
 - CMS: event plane method for flow measurement
 - non flow subtraction
 -

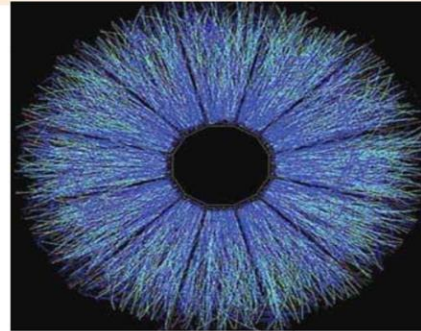
Flow for $f_0(980)$: need efforts from ALICE, STAR collaboration
efforts from the theory community

TABLE I Mesons that contain at most one heavy quark that cannot be easily accommodated in the $q\bar{q}$ quark model. Their quantum numbers $I^G(J^{PC})$, masses, widths, the nearby S -wave thresholds, $m_{\text{threshold}}$, where we add in brackets $M - m_{\text{threshold}}$, and the observed decay modes are listed in order. The data without references are taken from the 2016 edition of the Review of Particle Physics (Patrignani *et al.*, 2016).

State	$I^G(J^{PC})$	M [MeV]	Γ [MeV]	S -wave threshold(s) [MeV]	Decay mode(s) [branching ratio(s)]
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$a_1(1420)$	$1^-(1^{++})$	1414^{+15}_{-13}	153^{+8}_{-23}	$K\bar{K}^*(27^{+15}_{-13})$	$f_0(980)\pi$ [seen]
$X(1835)$	$?^?(0^{-+})$	$1835.8^{+4.0}_{-3.2}$	112 ± 40	$p\bar{p}(-40.7^{+4.0}_{-3.2})$	$p\bar{p}$ $\eta'\pi\pi$ $K_S^0 K_S^0 \eta$
$D_{s0}^*(2317)^+$	$0(0^+)$	2317.7 ± 0.6	<		$D_s^+ \pi^0$ [(48 ± 11)% [(18 ± 4)% π^- [(4 ± 1)% $7)^+ \gamma$ [(4 ± 5)%]
$D_{s1}(2460)^+$	$0(1^+)$	2459.5 ± 0.6	<		DK D^*K
$D_{s1}^*(2860)^+$	$0(1^-)$	2859 ± 27	159		

→ With similar PID as STAR explore light exotic hadrons at NICA,

High energy nuclear physics



A large amount of particles produced
→ momentum distributions

-particle yield

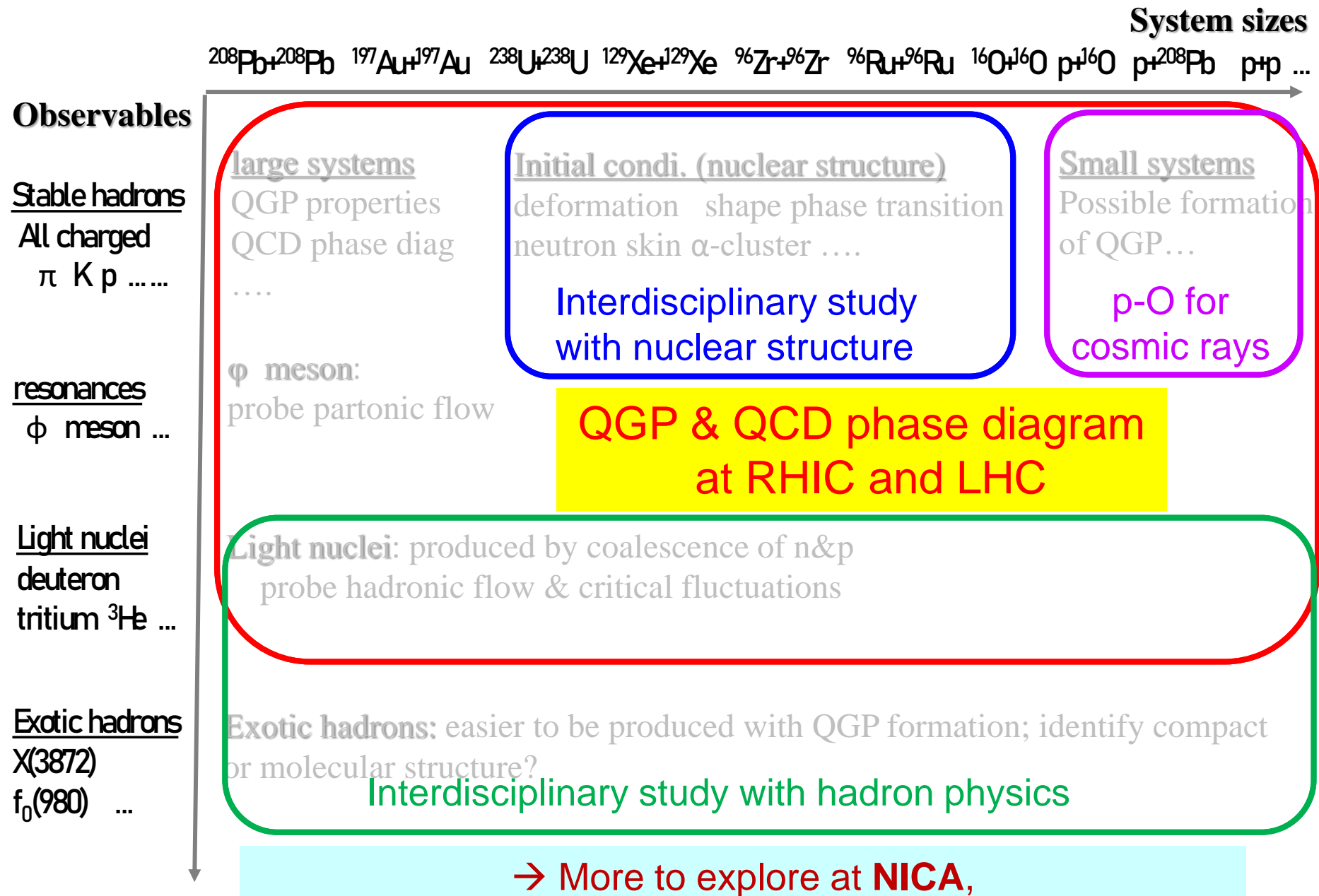
→ - pT spectra

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Soft Physics for high energy Nucleons-Nucleons Collisions

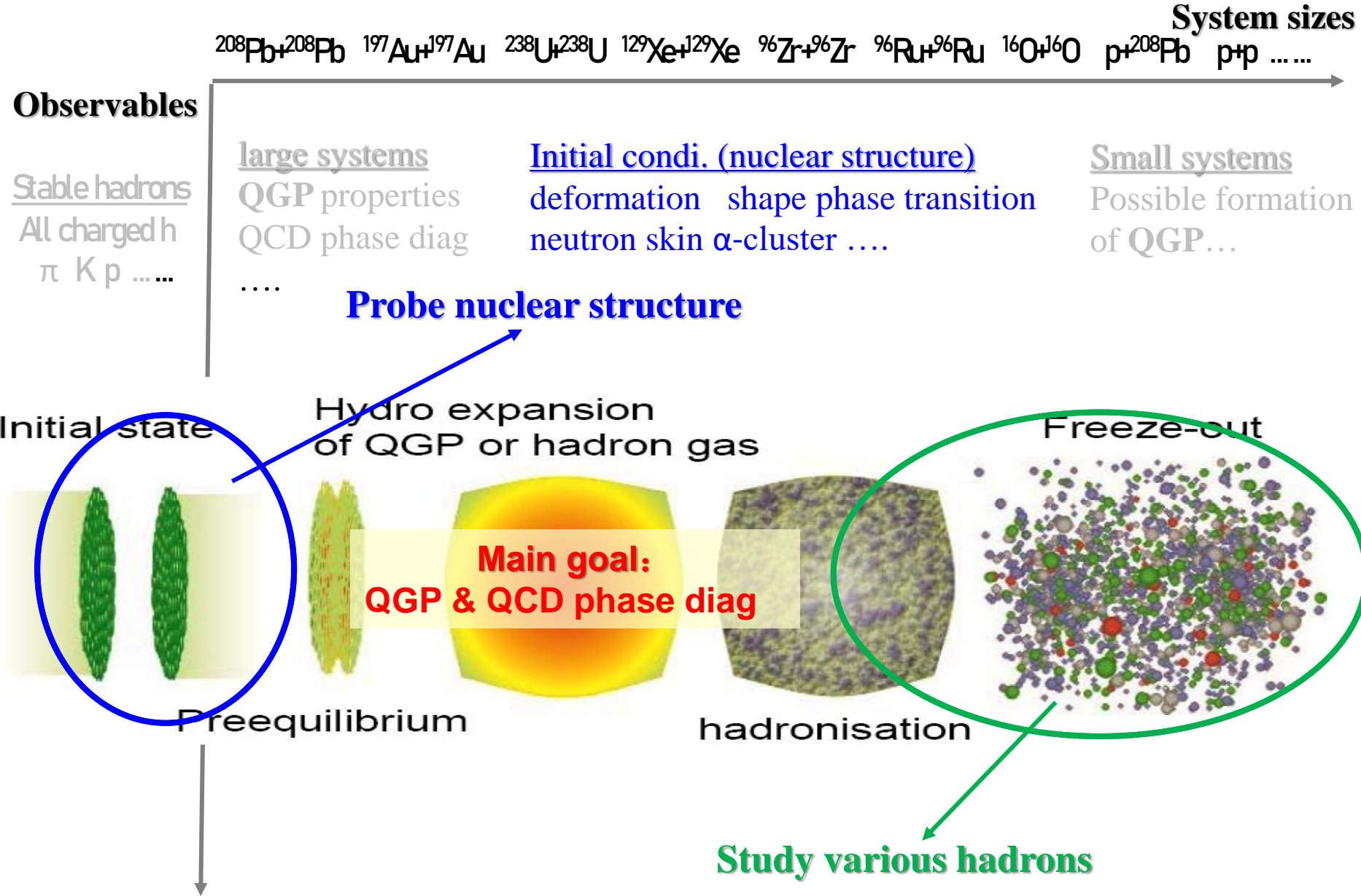
	System sizes									
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Observables										
<u>Stable hadrons</u>	<u>large systems</u>			<u>Initial condi. (nuclear structure)</u>			<u>Small systems</u>			
All charged	QGP properties			deformation shape phase transition			Possible formation			
π K p	QCD phase diag			neutron skin α -cluster			of QGP...			
									
<u>resonances</u>										
ϕ meson ...	ϕ meson: probe partonic flow									
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deuteron	Light nuclei: produced by coalescence of n&p probe hadronic flow & critical fluctuations									
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Soft Physics for high energy Nucleons-Nucleons Collisions

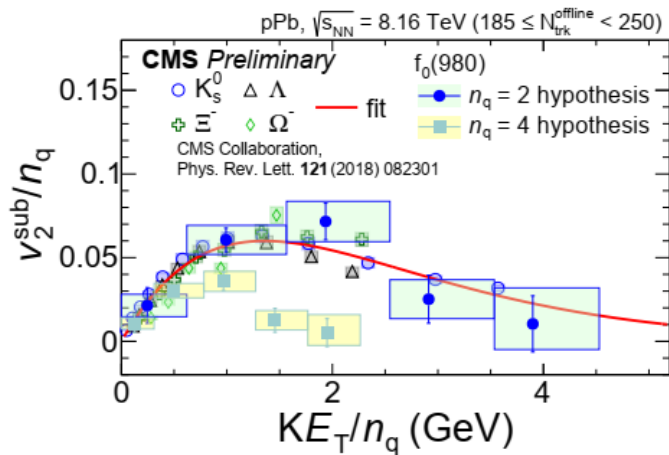


Many Thanks

Soft Physics for high energy Nucleons-Nucleons Collisions



(Gu An for CMS)



- ▶ v_2 of $f_0(980)$ measured as a function of p_T up to 10 GeV/c
- ▶ Assuming NCQ scaling, n_q of $f_0(980)$ is consistent with 2.
- ▶ $n_q = 4$ (tetra-quark state or $K\bar{K}$ molecule) excluded with 7.7σ .
- ▶ $n_q = 3$ ($q\bar{q}g$ hybrid) excluded with 3.5σ .
- ▶ Our data favor $q\bar{q}$ normal meson state for $f_0(980)$.