Global polarization and spin alignment measurements in heavy-ion collisions at RHIC and LHC

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NICA-MPD Collaboration Meeting, 2024/10/14-16

Outline

- Brief introduction of the global polarization
- Measurements in heavy-ion collisions and prospects
 - Hyperon polarization
 - Vector meson spin alignment
- Summary

Global polarization in HIC



- The initial momentum gradient should result in a net angular momentum (shear) in this direction that will be transferred to quark spin via spin-orbit interaction, this effect may not be washed out during interaction and hadronization
- Spin-vorticity coupling Betz, Gyulassy, Torrieri Phys. Rev. C 76, 044901 (2007); Becattini, Piccinini, Rizzo Phys. Rev. C 77, 024906 (2008)
- Connection to classical world, the Barnett effect, a fraction of the *L* associated with the body rotation is transformed into the spin *L* of the electron

Experimental measurements: Λ



Experimental measurements: Λ (cont.)

STAR Col. Phys. Rev. C 76, 024915 (2007)



Experimental measurements: A (cont.2)



Measurements of Λ and Ξ, Ω





Experimental measurements: φ,K*

 $|1-1\rangle = |\downarrow\downarrow\rangle$





Experimental measurements: φ,K*(cont.)



New Measurements φ,K*0@non-central collisions



 New measurements extend the study to lower energies with high statistics, @200 GeV, a factor of ~50 more event statistics analyzed.

 We see that the signal for the φ meson occurs mainly within ~1.0-2.4 GeV/c; at larger p_T the results can be regarded as being consistent with 1/3 within ~2σ or less.

* 1st order EP: ZDC or BBC * 2nd order EP: TPC

STAR Col. Nature 614, 244 (2023)

New Measurements φ,K*0@non-central collisions



Study the fine structure vs. centrality

STAR Col. Nature 614, 244 (2023)



At high energies (\geq 62.4 GeV) for φ , and (\geq 39 GeV) for K^{*0} , ρ_{00} in central collisions tends to \leq 1/3. This might be caused by transerve local spin alignment and a contribution from the helicity polarization of quarks.

Results mid-central & averaged over p_T



- φ-meson is significantly above 1/3 for sqrt{s}≤
 62 GeV
- 2) K* is largely consistent with 1/3
- 3) Averaged over 62 GeV and below:
- 0.3541 ± 0.0017 (stat.) ± 0.0018 (sys.) for φ
- 0.3356 ± 0.0034 (stat.) ± 0.0043 (sys.) for K*

* Different approaches are used in the combinatorial bg. analysis

STAR Col. Nature **614**, 244 (2023)

Expectations ρ_{00} from theory

1

0.38

0.36

ළි 0.34

0.32

0.30

 $F^2 (m_\pi^2)$

$$d^{3}\boldsymbol{p}_{b} \exp \left(-\frac{\boldsymbol{p}_{b}^{2}}{a_{\phi}^{2}}\right) \frac{p_{b,x}^{2}}{E_{\boldsymbol{p}_{1}}E_{\boldsymbol{p}_{2}}}.$$

$$(42)^{10^{-5}}$$

$$c_{\varepsilon}: \text{ Vorticity tensor}^{[1]} < \frac{1/3}{(\text{Negative } \sim 10^{-4})}$$

$$c_{\varepsilon}: \text{ Electric field}^{[2]} > \frac{1/3}{(\text{Positive } \sim 10^{-5})}$$
Fragmentation}^{[3]} > \text{or, } < \frac{1}{3}
$$(1 \text{ Local spin alignment and } \rho_{001}^{\phi}(\mathbf{s}\boldsymbol{c}, \boldsymbol{p})$$

$$\text{Turbulent color field}^{[5]} < \frac{1}{3}$$

$$c_{\phi}: \text{ Vector meson strong} > \frac{1}{3}$$

$$\rho_{00}^{\phi} \approx \frac{1}{3} + c_{\omega} + c_{\varepsilon} + c_{\rm EM} + c_{\phi} + c_{\rm LV} + c_h + c_{\rm TC} + c_{\rm shear}$$

[1]. Yang et al., Phys. Rev. C 97, 034917 (2018) [2]. Sheng et al., Phys. Rev. D 101, 096005 (2020) [3]. Xia et al., Phys. Lett. B 817, 136325 (2021) [4]. Gao, Phys. Rev. D 104, 076016 (2021) [5]. Mulle (Yang, Phy Rev. D 105, L011901 (2022) [6]. Li, Liu, arXiv:2206.11890, Wagner, Weickgenannt, Speranza, arXiv:2207.01111

The local correlation or fluctuation of φ fields is the dominant mechanism for the observed φ -meson ρ_{00}

Sheng, et al., Phys. Rev. Lett. 131, 042304 (2023)





The spin-spin correlations

• Spin-spin correlations and hypernuclei polarization in HIC

Shen, Chen, Tang, arXiv:2407.21291

$$c'_{\Lambda\Lambda} = \frac{64}{\pi^2 \alpha_{\Lambda}^2} \frac{\langle \sin(\phi_i^* - \Psi_{\rm EP}) \sin(\phi_j^* - \Psi_{\rm EP}) \rangle}{\langle \cos^2(\Psi_{\rm EP} - \Psi_{\rm RP}) \rangle} - P_{\Lambda}^2,$$

$$B(c'_{OS}, c'_{SS}) = \frac{1}{2} \left(\frac{c'_{\Lambda\bar{\Lambda}} - c'_{\Lambda\Lambda}}{c'_{\Lambda\bar{\Lambda}} + c'_{\Lambda\Lambda}} + \frac{c'_{\bar{\Lambda}\Lambda} - c'_{\bar{\Lambda}\bar{\Lambda}}}{c'_{\bar{\Lambda}\Lambda} + c'_{\bar{\Lambda}\bar{\Lambda}}} \right)$$

Sun, Liu, Zhen, Chen, Ko, Ma, arXiv:2405.12015

🌑 n 🕘 p 🌏 Λ

 ${}^3_{\Lambda}H(\frac{1}{2}^+)$

 ${}^{3}_{\Lambda}H(\frac{1}{2}^{+})$

J^P	structure	decay mode	$\frac{dN}{d\cos\theta^*}$
$\frac{1}{2}^{+}$	$\Lambda(\frac{1}{2}^+) - np(1^+)$	$^{3}_{\Lambda}\text{H} \rightarrow \pi^{-} + ^{3}\text{He}$	$\frac{1}{2}(1-\frac{1}{2.58}\alpha_{\Lambda}\mathcal{P}_{\Lambda}\cos\theta^{*})$
$\left \frac{1}{2}^+\right $	$\Lambda(\frac{1}{2}^+) - np(0^+)$	$^{3}_{\Lambda}\text{H} \rightarrow \pi^{-} + ^{3}\text{He}$	$\frac{1}{2}(1+\alpha_{\Lambda}\mathcal{P}_{\Lambda}\cos\theta^{*})$
$\frac{3}{2}^{+}$	$\Lambda(\frac{1}{2}^+) - np(1^+)$	$^{3}_{\Lambda}\text{H} \rightarrow \pi^{-} + ^{3}\text{He}$	$\frac{1}{2} \left(1 - \mathcal{P}_{\Lambda}^2 (3\cos^2 \theta^* - 1) \right)$
$\left[\frac{1}{2}\right]^{-}$	$\overline{\Lambda(\frac{1}{2}^{-})} - \overline{np}(1^{-})$	$\frac{3}{\overline{\Lambda}}\overline{\mathrm{H}} \to \pi^+ + {}^3\overline{\mathrm{He}}$	$\frac{1}{2}(1-\frac{1}{2.58}\alpha_{\bar{\Lambda}}\mathcal{P}_{\bar{\Lambda}}\cos\theta^*)$
$\frac{1}{2}^{-}$	$\bar{\Lambda}(\frac{1}{2}^{-}) - \overline{np}(0^{-})$	$ {}^3_{\overline{\Lambda}}\overline{\mathrm{H}} \to \pi^+ + {}^3\overline{\mathrm{He}}$	$\frac{1}{2}(1+\alpha_{\bar{\Lambda}}\mathcal{P}_{\bar{\Lambda}}\cos\theta^*)$
$\left[\frac{3}{2}\right]^{-}$	$\bar{\Lambda}(\frac{1}{2}^{-}) - \overline{np}(1^{-})$	$\frac{3}{\overline{\Lambda}}\overline{\mathrm{H}} \to \pi^+ + {}^3\overline{\mathrm{He}}$	$\frac{1}{2} \left(1 - \mathcal{P}_{\bar{\Lambda}}^2 (3\cos^2\theta^* - 1) \right)$

$$\mathcal{P}_{_{\Lambda}\mathrm{H}} \approx \frac{\frac{2}{3} \langle \mathcal{P}_n \rangle + \frac{2}{3} \langle \mathcal{P}_p \rangle - \frac{1}{3} \langle \mathcal{P}_\Lambda \rangle - \langle \mathcal{P}_n \mathcal{P}_p \mathcal{P}_\Lambda \rangle + C_-}{1 - \frac{2}{3} (\langle (\mathcal{P}_n + \mathcal{P}_p) \mathcal{P}_\Lambda \rangle) + \frac{1}{3} \langle \mathcal{P}_n \mathcal{P}_p \rangle + C_+}$$
(15)

$$C_{-} = -\frac{1}{4} \left(\langle c_{np}^{zz} \mathcal{P}_{\Lambda} \rangle + \langle c_{p\Lambda}^{zz} \mathcal{P}_{n} \rangle + \langle c_{n\Lambda}^{zz} \mathcal{P}_{p} \rangle \right) - \frac{1}{4} \langle c_{np\Lambda}^{zzz} \rangle,$$

$$C_{+} = \frac{1}{12} \left(\langle c_{np}^{zz} \rangle - 2 \langle c_{p\Lambda}^{zz} \rangle - 2 \langle c_{n\Lambda}^{zz} \rangle \right).$$
(16)

From φ to other mesons



- Forward rapidity J/ ψ ρ_{00} <1/3 at LHC
- Midrapidity J/ψ ρ₀₀ ~ 1/3 at RHIC

 D^{*+} shows a clear p_T dependence

 \rightarrow The underlying physics seems not converged?

Prospects at NICA-MPD for mesons

- The different species of vector meson spin alignment from RHIC/LHC seems not converged, independent measurements will be very helpful to understand the underlying physics
- NICA-MPD can identify the vector mesons well, thus will be excellent experiment to measure the light flavor spin alignment



Figs. from Zebo

Summary

- Spin polarization opens a new avenue to investigate heavy-ion collisions
- Global hyperon polarization is observed with the order of a few percent. It represents a measure of the average value of the global quark polarization in the system
- Global vector meson spin alignment is observed with a surprisingly large parttern for φ-meson. It represents a local fluctuation/correlation between quark and anti-quark polarization
- Measurements as a function of collision energies, different hadron species are on-going, rich physics to be explored, and the NICA-MPD will be very powerful to establish the feature of highbaryon density region